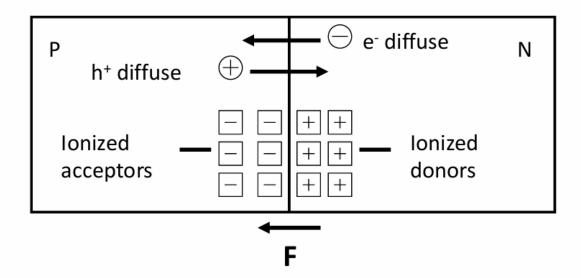
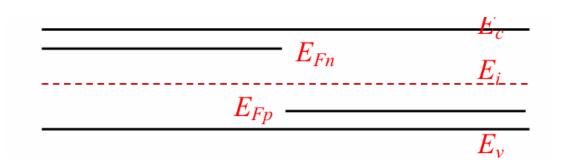
RC6

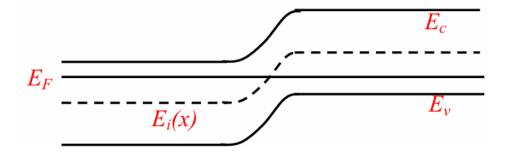
What happens when you join p-type and n-type material?



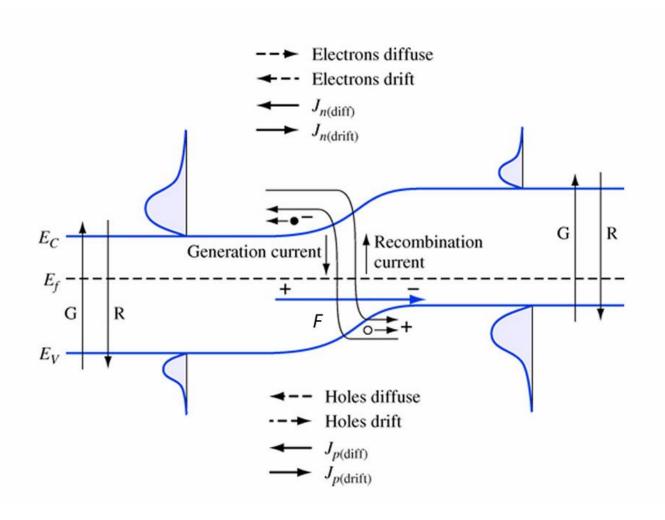
- Diffusion drives e's and h's to lower concentration regions
- Ionized dopants remain, creating an electric field
- Under equilibrium, diffusion and drift processes are balanced

Band diagram before and after equilibrium





Equilibrium current flow in P-N junction



Built-in potential

$$E_{C} = \sum_{V_{bi}} V_{bi}$$

$$qV_{bi} = E_{G} - (E_{C} - E_{F})_{n} - (E_{F} - E_{V})_{p}$$

$$(E_{C} - E_{F})_{n} = -kT \ln\left(\frac{N_{D}}{N_{C}}\right)$$

$$(E_{F} - E_{V})_{p} = -kT \ln\left(\frac{N_{A}}{N_{V}}\right)$$

$$n_{i}^{2} = N_{C}N_{V} \exp\left(-\frac{E_{G}}{kT}\right)$$

$$\Rightarrow E_{G} = kT \log\left(\frac{N_{C}N_{V}}{n_{i}^{2}}\right)$$

The depletion approximation

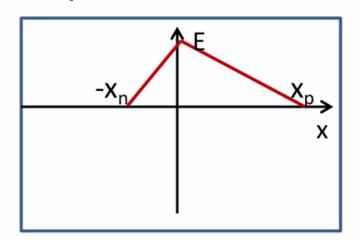
- Analyze p-n junction using electrostatics (Poisson's equation)
- Depletion approximation no mobile charge in depletion region, only fixed charge from the ionized impurities
- Depletion region is necessary to generate the electric field that cancels out the diffusion current.

Zero applied bias-charge density and electric field

$$\frac{d^2V(x)}{dx^2} = \begin{cases} \frac{q}{\varepsilon} N_a & 0 \le x \le x_p \\ -\frac{q}{\varepsilon} N_d & -x_n \le x < 0 \end{cases}$$

$$\begin{array}{c|c} & & & & \\ & + & & \\ & -x_n & N_a & - & \\ & & & \times \end{array}$$

$$E = -\frac{dV(x)}{dx} = \begin{cases} -\frac{q}{\varepsilon} N_a \ x + \frac{q}{\varepsilon} N_a \ x_p & 0 \le x \le x_p \\ \frac{q}{\varepsilon} N_d \ x + \frac{q}{\varepsilon} N_d \ x_n & -x_n \le x < 0 \end{cases}$$



Zero applied bias-potential and space charge width

After applying two restrictions to the potential, we can get the space charge width.

Space charge width

$$x = 0 \Rightarrow N_d \ x_n = N_a \ x_p$$

$$V(x = x_n) = V_{bi} = \frac{kT}{q} ln(\frac{N_d \ N_a}{n_i^2})$$

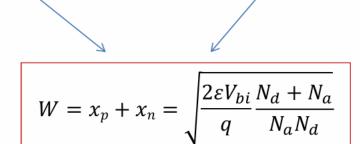
$$V(x) = \begin{cases} \frac{q}{\varepsilon} N_a \ (\frac{1}{2}x^2 - x_p x + \frac{x_p^2}{2}) & 0 \le x \le x_p \\ -\frac{q}{\varepsilon} N_d \ (\frac{1}{2}x^2 + x_n x - \frac{N_a}{N_d} \frac{x_p^2}{2}) & -x_n \le x < 0 \end{cases}$$

Zero applied bias-space charge width

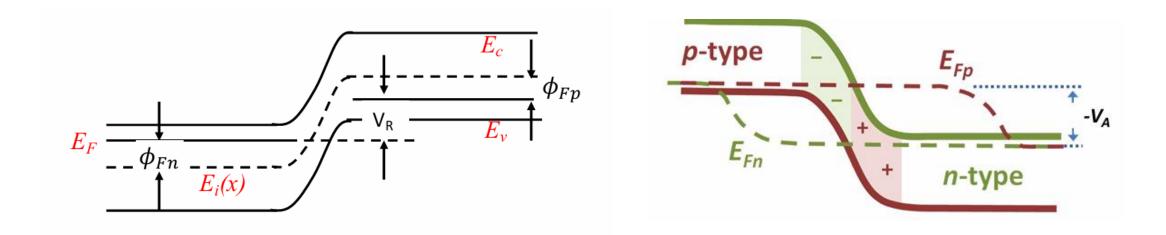
Space charge width

$$x_p = \sqrt{\frac{2\varepsilon V_{bi}}{q} \frac{N_d}{N_a} \frac{1}{N_a + N_d}}$$

$$x_p = \sqrt{\frac{2\varepsilon V_{bi}}{q} \frac{N_d}{N_a} \frac{1}{N_a + N_d}} \qquad N_d \ x_n = N_a \ x_p \Rightarrow x_n = \sqrt{\frac{2\varepsilon V_{bi}}{q} \frac{N_a}{N_d} \frac{1}{N_a + N_d}}$$



Reversed applied bias-band diagram



Reversed applied bias

Just replace all the built in voltage with built in voltage plus bias voltage.

Space charge width and electric field

$$x_{p} = \sqrt{\frac{2\varepsilon(V_{bi} + V_{R})}{q} \frac{N_{d}}{N_{a}} \frac{1}{N_{a} + N_{d}}}$$

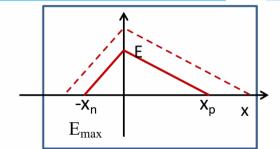
$$x_{p} = \sqrt{\frac{2\varepsilon(V_{bi} + V_{R})}{q} \frac{N_{d}}{N_{a}} \frac{1}{N_{a} + N_{d}}} \qquad N_{a}^{-} x_{n} = N_{d}^{+} x_{p} \Rightarrow x_{n} = \sqrt{\frac{2\varepsilon(V_{bi} + V_{R})}{q} \frac{N_{a}}{N_{d}} \frac{1}{N_{a} + N_{d}}}$$

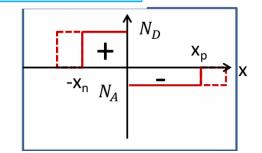
$$W = \sqrt{\frac{2\varepsilon(V_{bi} + V_R)}{q} \frac{N_d + N_a}{N_a N_d}}$$

$$W = \sqrt{\frac{2\varepsilon(V_{bi} + V_R)}{q} \frac{N_d + N_a}{N_a N_d}} \qquad E = -\frac{dV(x)}{dx} = \begin{cases} -\frac{q}{\varepsilon} N_a x + \frac{q}{\varepsilon} N_a x_p & 0 \le x \le x_p \\ \frac{q}{\varepsilon} N_d x + \frac{q}{\varepsilon} N_d x_n & -x_n \le x < 0 \end{cases}$$

$$E_{\text{max}} = \left\{ \frac{2e(V_{bi} + V_R)}{\epsilon_s} \left(\frac{N_a N_d}{N_a + N_d} \right) \right\}^{1/2} \qquad E_{\text{max}} = \frac{2(V_{bi} + V_R)}{W}$$

$$E_{\text{max}} = \frac{2(V_{bi} + V_R)}{W}$$





Reversed applied bias-junction capacitance

$$C' = \frac{dQ}{dV_{\rm R}}|_{V_R = V_{R_0}} = qN_d \frac{db}{dV_R}|_{V_R = V_{R_0}} = \sqrt{\frac{q\varepsilon}{2(V_{bi} + V_{R_0})} \frac{N_d N_a}{N_a + N_d}} = \frac{\varepsilon}{W}$$

One-sided junction

Either high concentration of donor or high concentration of acceptor

