VE320 – Summer 2024

Introduction to Semiconductor Devices

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Chapter 3 Introduction to the Quantum Theory of Solids

Outline

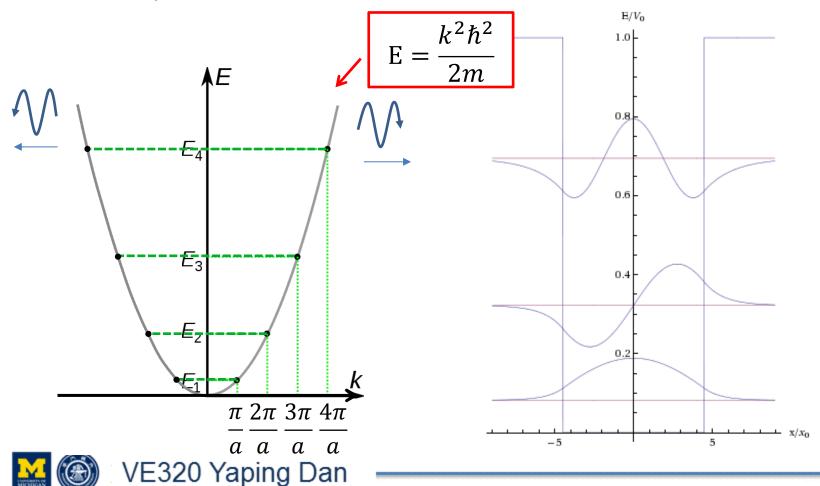
- 3.1 Allowed and Forbidden Energy Bands
- 3.2 Electrical Conduction in Solids
- 3.3 Extension to Three Dimensions
- 3.4 Effective Mass
- 3.5 Density of States Function
- 3.6 Statistical Mechanics

Outline

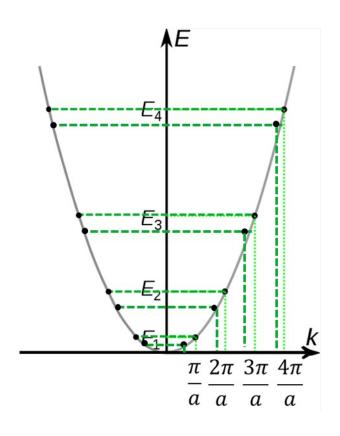
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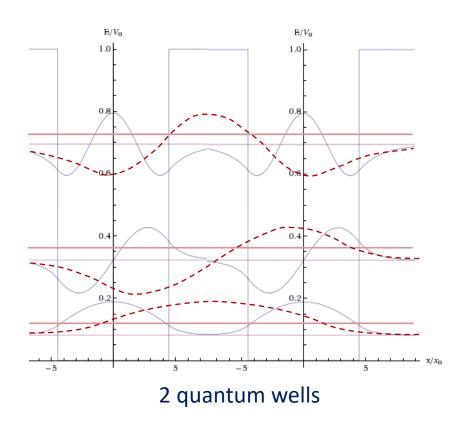
Forming energy bands: analytical

Previously: Electrons in Finite Quantum Well

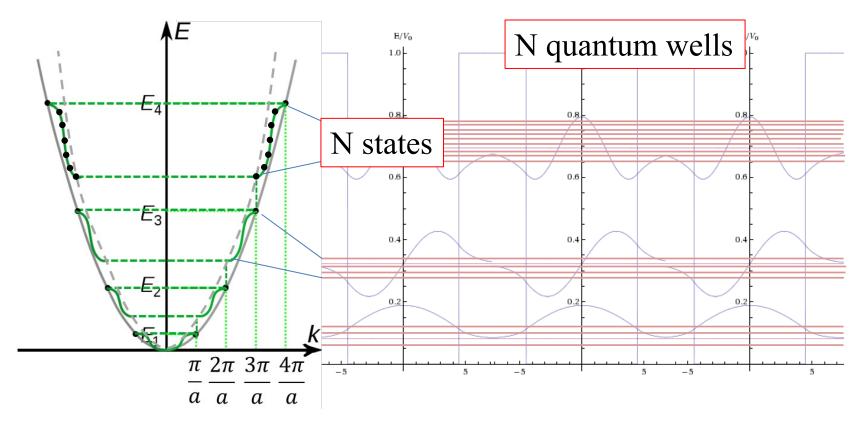


Forming energy bands: analytical

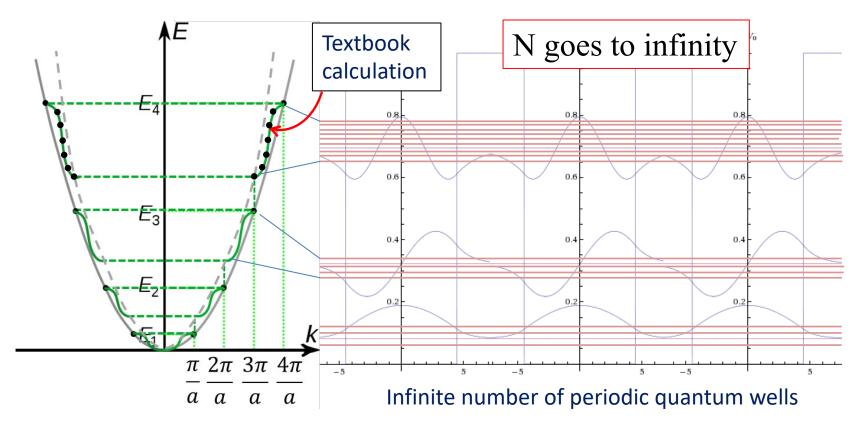




Forming energy bands: analytical



Forming energy bands: analytical



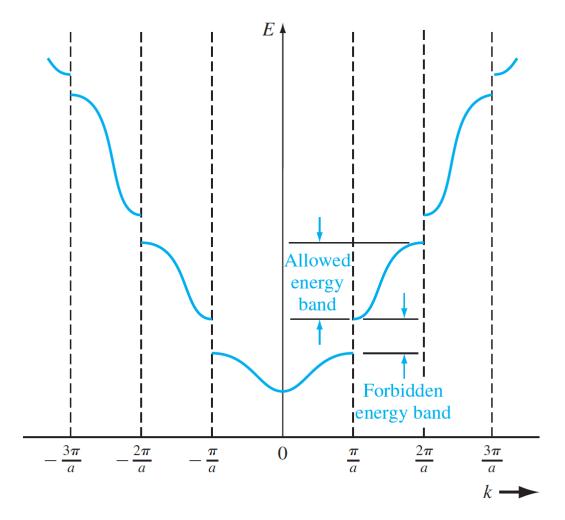
On P.67

eq.(3.22)

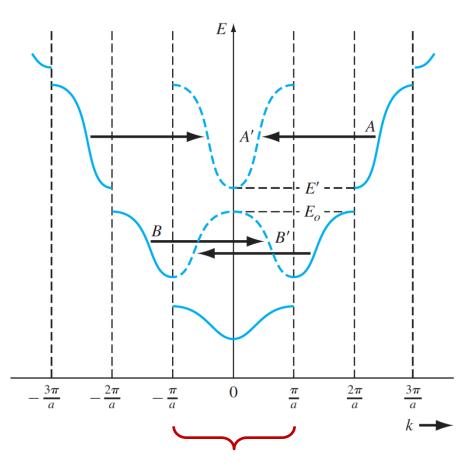
$$\frac{mV_0ba}{\hbar^2}\frac{\sin(\alpha a)}{\alpha a} + \cos(\alpha a) = \cos(ka)$$

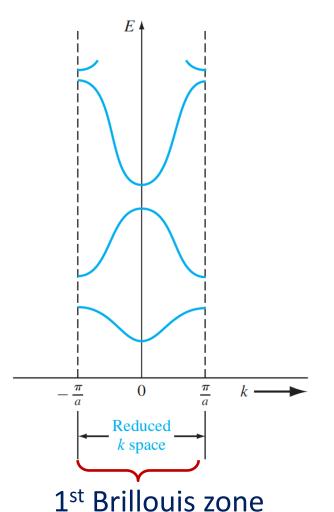


Forming energy bands: analytical



Band structure in physical and k space for 1D periodic quantum wells



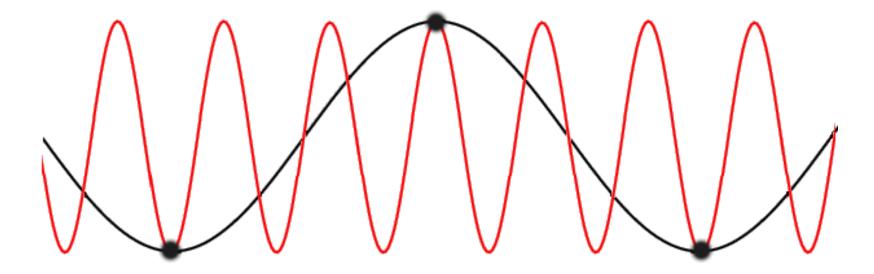


1st Brillouis zone



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- Black wave with a smaller k (longer wavelength) is in the 1st Brillouis zone.
- Red wave with a larger k (short wavelength) is outside of 1st Brillouis zene.
- Both waves have the same frequency (same energy).
- Both waves can describe the exact same information of a particle.

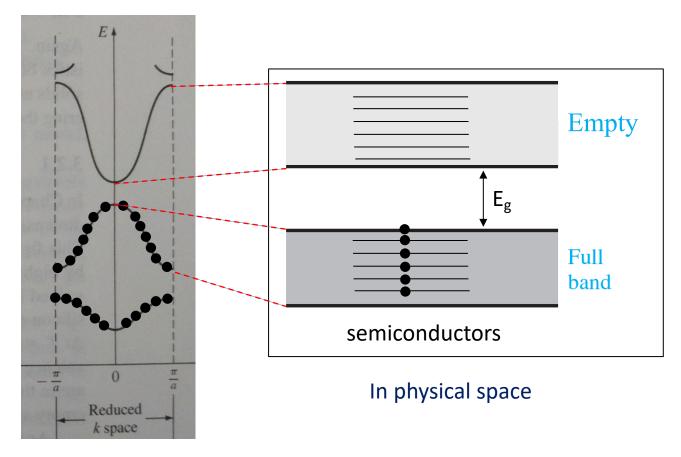


http://en.wikipedia.org/wiki/Phonon#/media/File:Phonon k 3k.gif

Outline

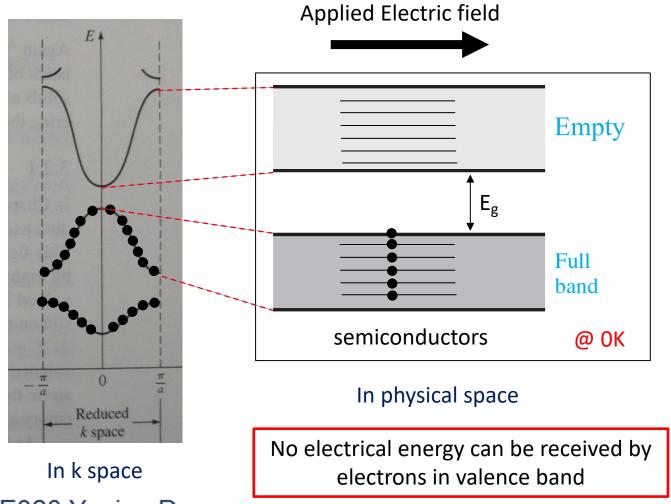
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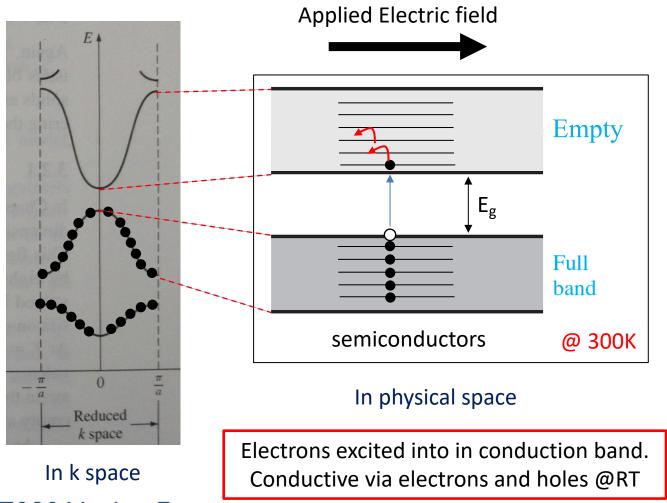
Energy band of semiconductors

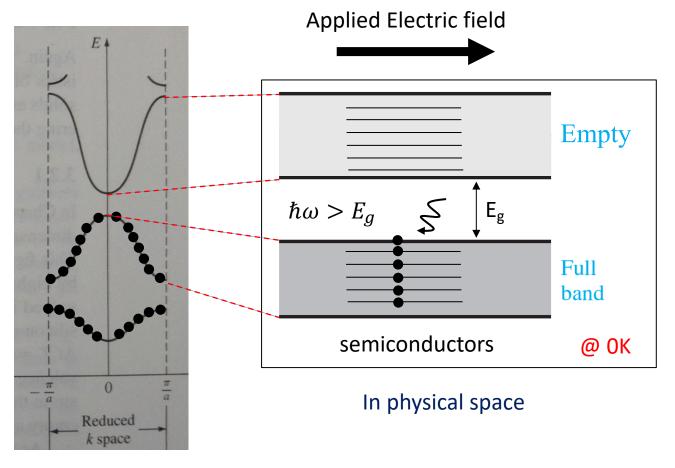


In k space



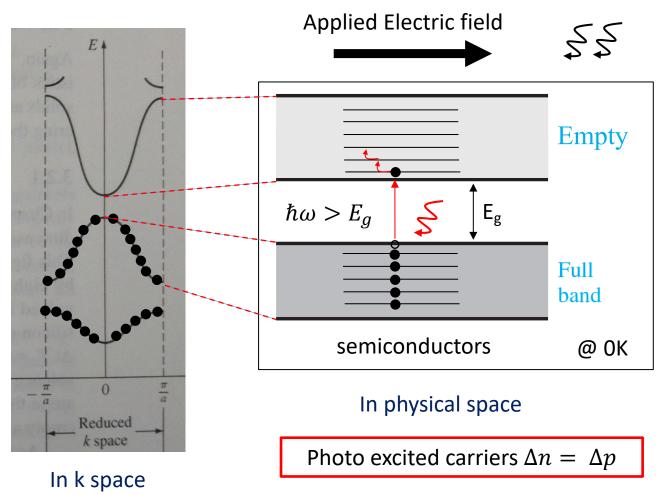




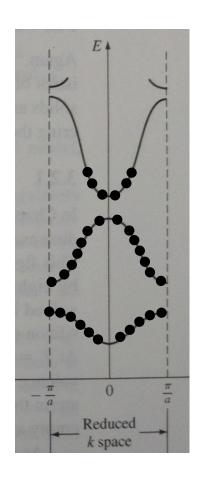




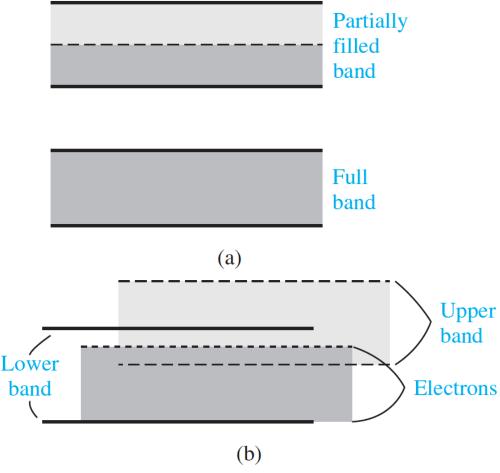




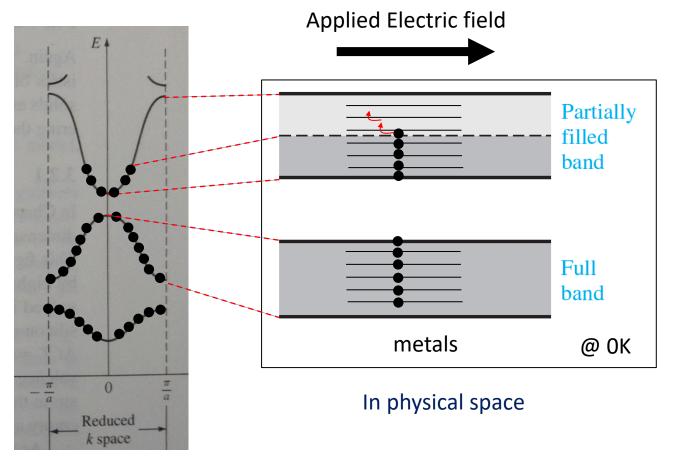
Energy band of metals



Forming energy bands is complicated.



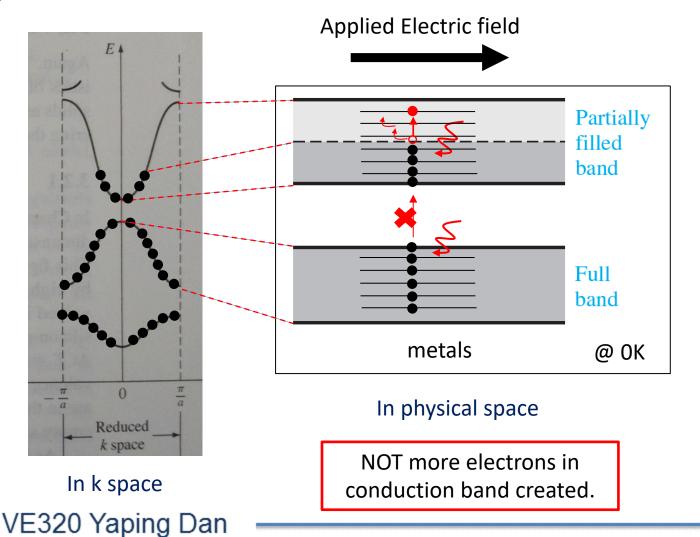
Energy band of metals







Energy band of metals



Metals, semiconductors and insulators

Insulators are wide bandgap semiconductors!

Physicists call semiconductors as Insulators.

Check your understanding

When the temperature decreases to 0K, how does the

conductivity change for metals and semiconductors?

Check your understanding

What's the difference between metals, semiconductors and

insulators in terms of energy bands?

Doping in semiconductors

Intrinsic semiconductors:

pure semiconductor, no doping, no defects

<u>n-type semiconductors</u>:

Charge carriers are negative, i.e. electrons

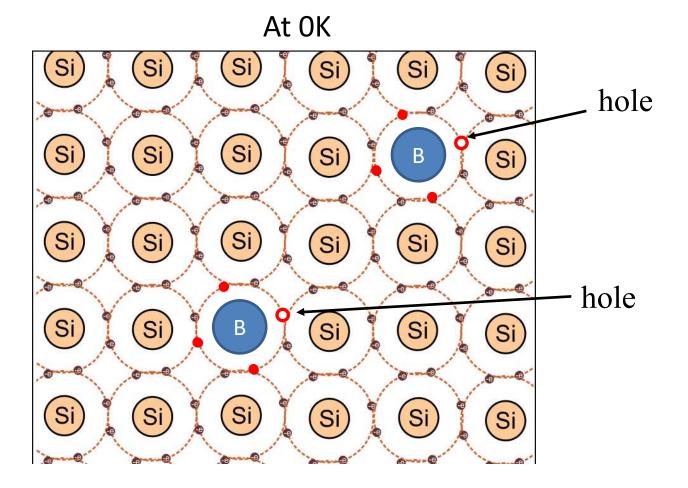
Doped by donor-type of dopants (impurities)

p-type semiconductors:

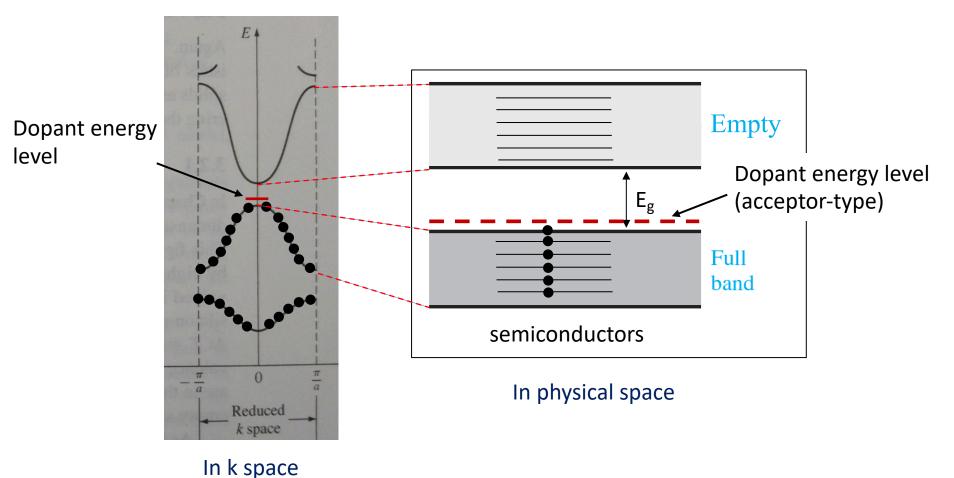
Charge carriers are positive, i.e. holes

Doped by acceptor-type of dopants (impurities)

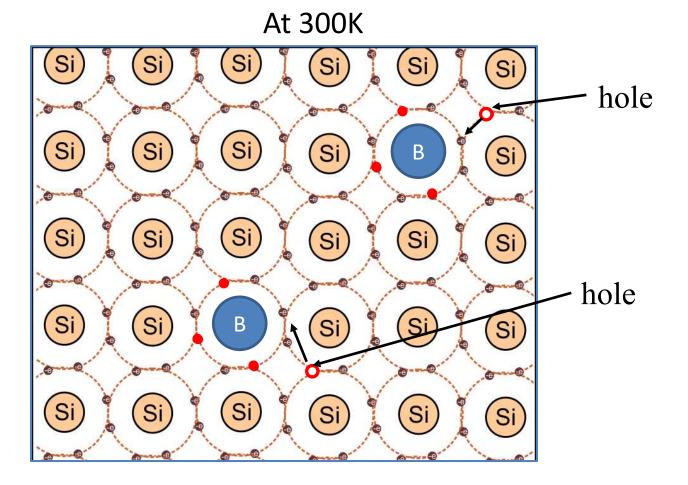
Acceptor-type or **p-type** doping (from the view of physical world)



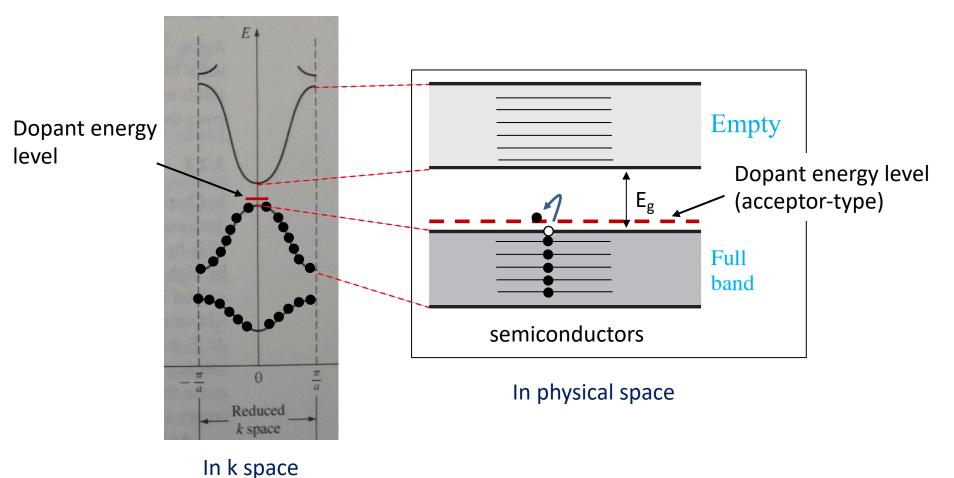
Acceptor-type or **p-type** doping (from the view of energy band)



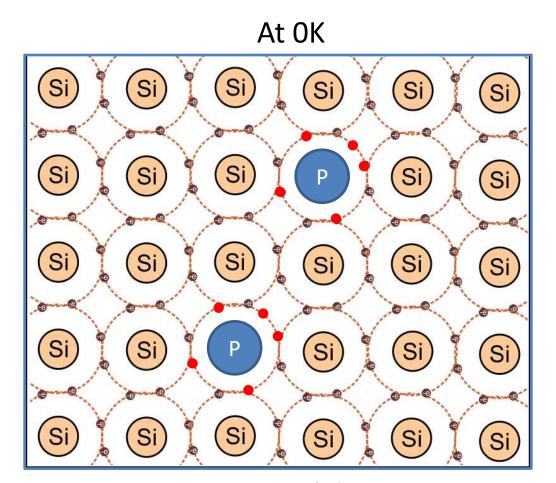
Acceptor-type or **p-type** doping (from the view of physical world)



Acceptor-type or **p-type** doping (from the view of energy band)

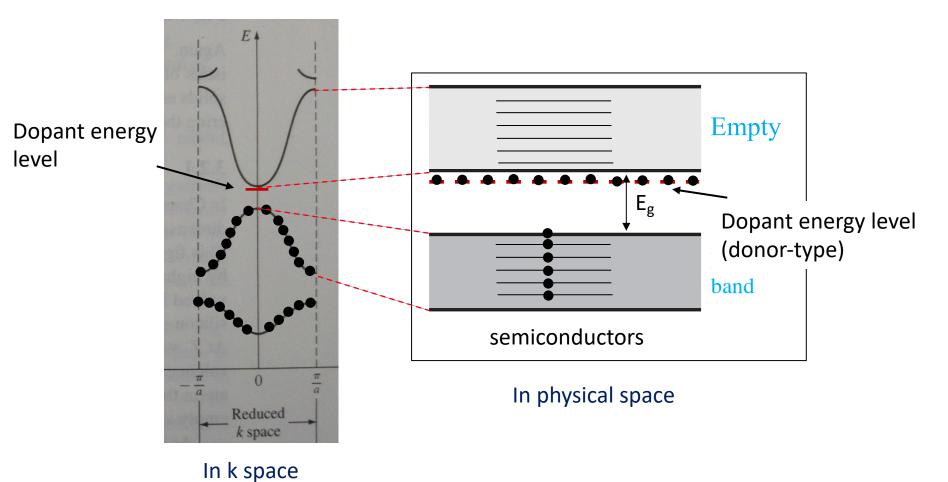


Donor-type or **n-type** doping (from the view of physical world)



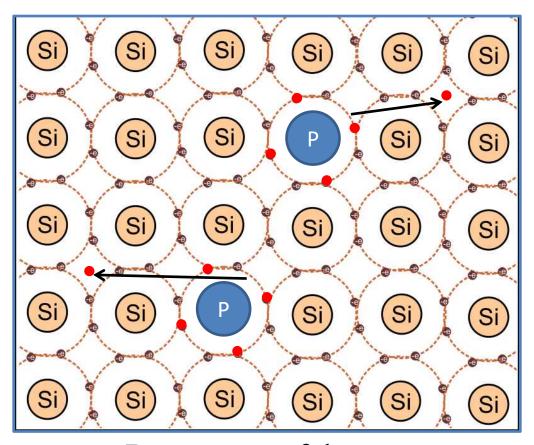
Donor-type of dopants

Donor-type or **n-type** doping (from the view of energy band)



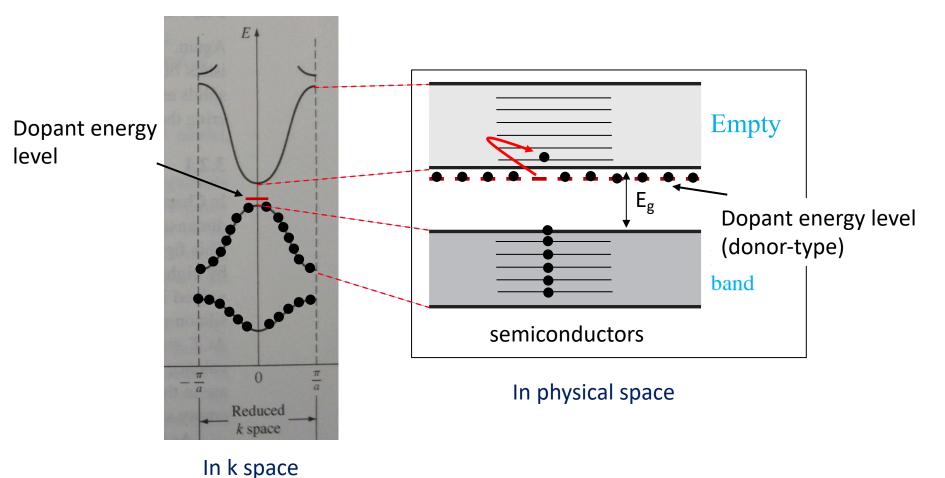
Donor-type or **n-type** doping (from the view of physical world)

> 0K



Donor-type of dopants

Donor-type or **n-type** doping (from the view of energy band)





Doping in semiconductors

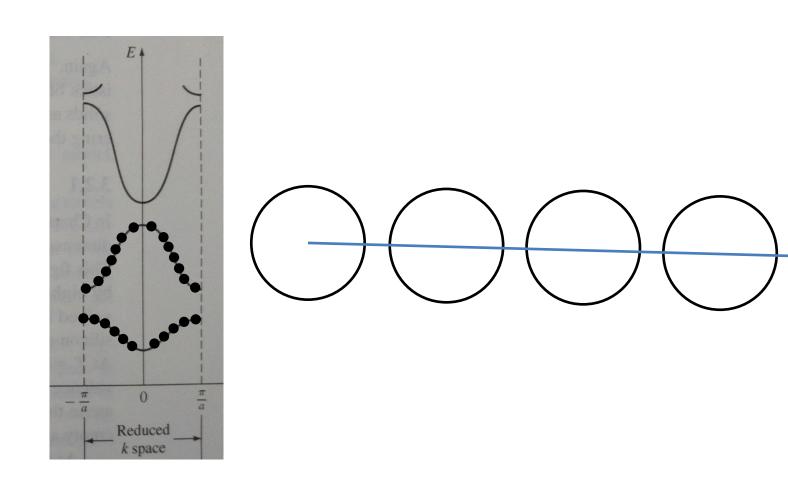
Si atomic concentration: 5 x 10²² cm⁻³

	Low concentration of doping	Medium concentration doping	High concentration of doping
Concentration (cm ⁻³)	< 10 ¹⁶	10 ¹⁶ -10 ¹⁸	10 ¹⁸ - 10 ²⁰
Relative concentration	1ppm	1 -100 ppm	100 ppm – 1%

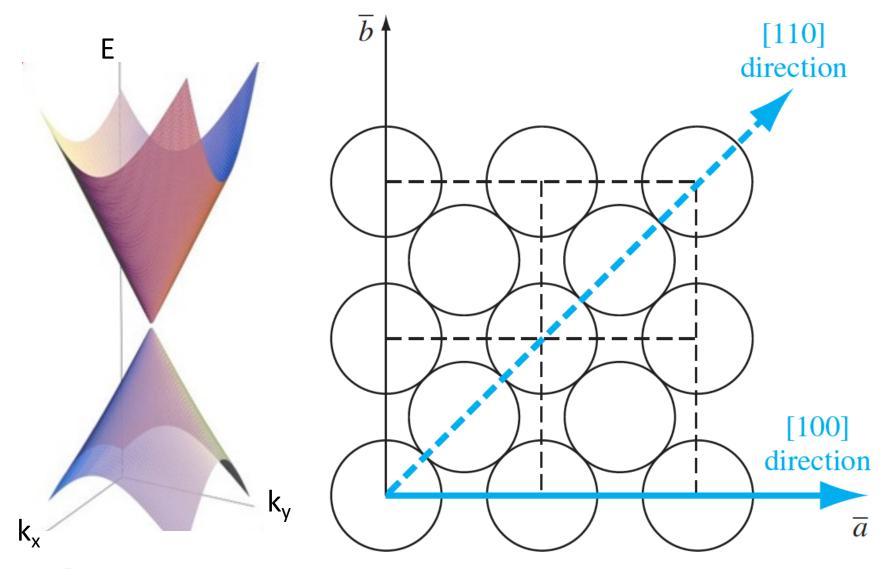
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3.3 Extension to Three Dimensions

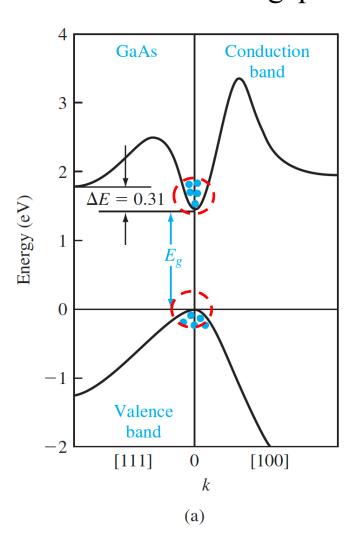


3.3 Extension to Three Dimensions

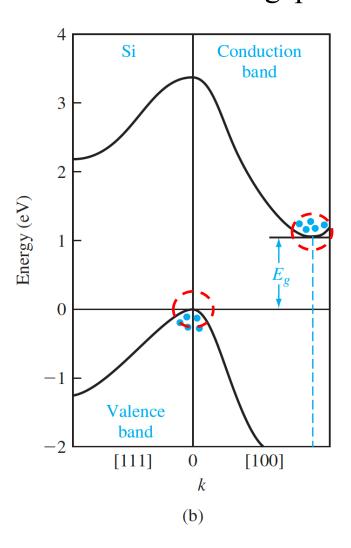


3.3 Extension to Three Dimensions

Direct bandgap



Indirect bandgap

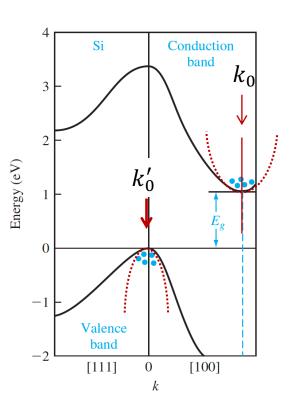


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3.4 Effective Mass

- So far the energy band structure is theoretically calculated.
- How to experimentally find it?





3.4 Effective Mass (For Electrons in the conduction band)

Electrons in free space have:

$$E_f(k) = \frac{\hbar^2 k^2}{2m} \Rightarrow \frac{d^2 E_f(k)}{dk^2} = \frac{\hbar^2}{m}$$

Electrons in crystalline semiconductors have:

$$E(k) = E(k = k_0) + \frac{d^2 E}{2dk^2} |_{k=k_0} (k - k_0)^2$$

$$\frac{d^2E}{dk^2}|_{k_0} = \frac{\hbar^2}{m^*}$$

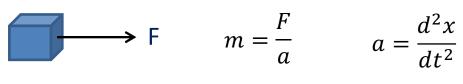
- $\frac{d^2E}{dk^2}|_{k_0} = \frac{\hbar^2}{m^*} \qquad \qquad \bullet \text{ m* has a unit of mass}$ We call it the effective mass of electrons in the crystal

$$E(k) = E(k = k_0) + \frac{\hbar^2}{2m^*} (k - k_0)^2$$

3.4 Effective Mass

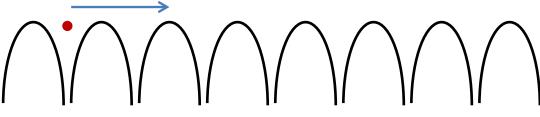
How to understand effective mass

Example: use Newton's law to find mass of an object

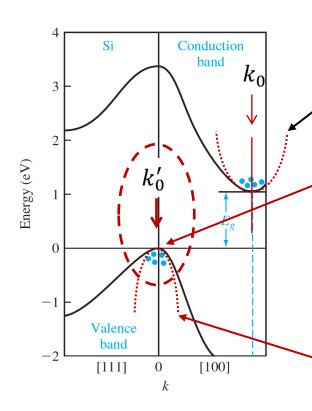


In the air

Modulated by Electric potential of ions



3.4 Effective Mass (For Electrons in the valence band)



Electrons in conduction band

$$E(k) = E(k_0) + \frac{d^2E}{2dk^2}|_{k_0}(k - k_0)^2$$

$$\frac{d^2E}{dk^2}|_{k_0'} = \frac{\hbar^2}{m^*} < 0$$

Negative mass would make no sense.

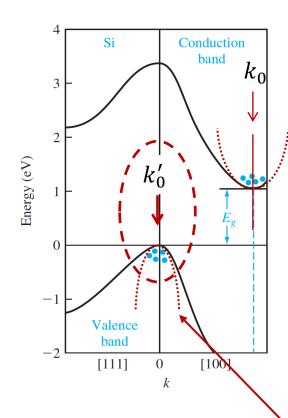


Electrons in valence band

$$E(k) = E(k'_0) - \frac{\hbar^2}{2m_p^*}|_{k'_0}(k - k'_0)^2$$

Positive mass

3.4 Effective Mass (A new particle defined: holes)



Electrons in conduction band

$$E(k) = E(k_0) + \frac{\hbar^2}{2m_n^*} |_{k_0} (k - k_0)^2$$

Positive mass

- Equivalent to a positive charge carrier
- Different effective mass
- Electrons and holes can come from dopants separately

Holes in valence band

$$E(k) = E(k'_0) - \frac{\hbar^2}{2m_p^*} |_{k'_0} (k - k'_0)^2$$

Positive mass

3.4 Effective Mass

Effective mass can be experimentally measured

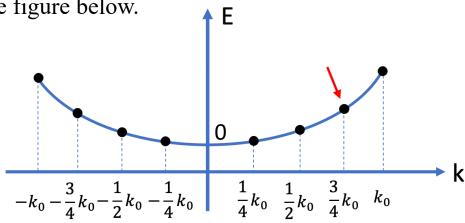
	Symbol	Germanium	Silicon	Gallium Arsenide
Bandgap	E _g (eV)	0.66	1.12	1.424
Electrons	m_e^*/m_0	0.55	1.08	0.067
Holes	$m_{\rm h}^{\ *}/m_0$	0.37	0.56	0.48

It means that energy band structure can be experimentally found.

(at least near conduction band bottom and valence band top)

Problem Example #1

- 1. In an infinite quantum well, the wavenumber k and energy E is quantized as shown in Figure 1. Please answer the following questions:
- a) Write the static wavefunction of the dot $(k = \frac{3}{4}k_0)$ that the red arrow is pointing to. Find the wavelength of this wavefunction.
- b) If all the states are filled with electrons, how many electrons can be filled in the figure below? The electron spin is not considered.
- c) If the effective mass of this quantum system is m_n^* , can this system allow an electron to have a quantized energy of $\frac{\hbar^2 k_0^2}{9m_n^*}$? If yes, mark the state of this electron in the figure below. How about a quantized energy of $\frac{\hbar^2 k_0^2}{8m_n^*}$? If yes, mark the state of this electron in the figure below.



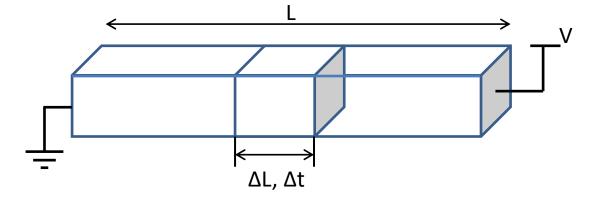
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n type semiconductor

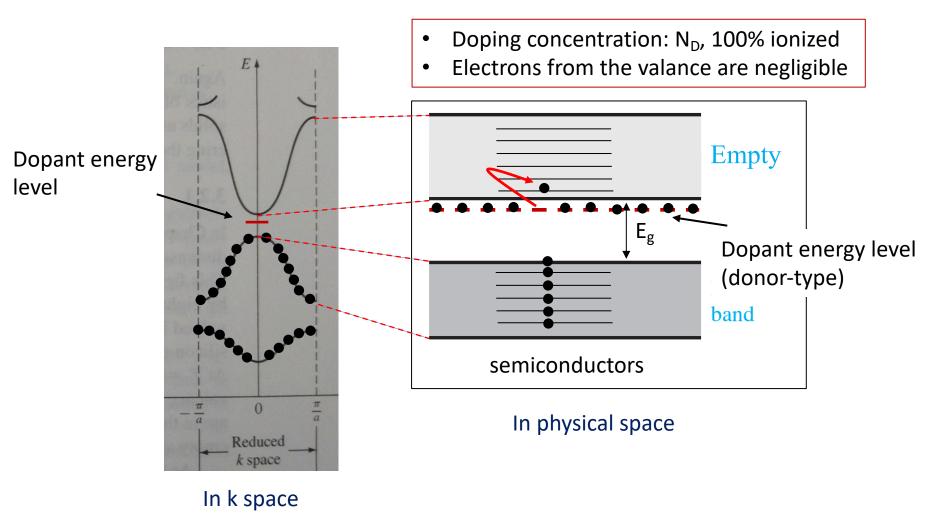
$$I = \frac{\Delta Q}{\Delta t} = \frac{nqA_c\Delta L}{\Delta t} = nqA_cv$$



$$v = \mu E = \mu V/L$$

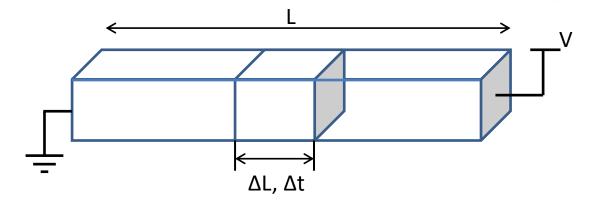
$$I = \frac{\Delta Q}{\Delta t} = \frac{nqA_c\Delta L}{\Delta t} = nqA_c\mu V/L \qquad \Rightarrow \quad \sigma = \frac{I}{V} = \frac{nqA_c\mu}{L}$$

Donor-type or **n-type** doping (from the view of energy band)



n type semiconductor

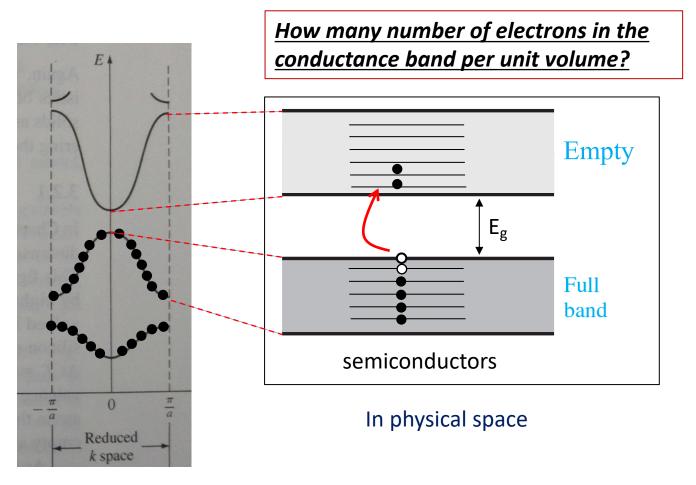
$$I = \frac{\Delta Q}{\Delta t} = \frac{nqA_c\Delta L}{\Delta t} = nqA_cv$$



$$v = \mu E = \mu V/L$$

$$I = \frac{\Delta Q}{\Delta t} = \frac{nqA_c\Delta L}{\Delta t} = nqA_c\mu V/L \qquad \Rightarrow \quad \sigma = \frac{I}{V} = \frac{N_D qA_c\mu}{L}$$

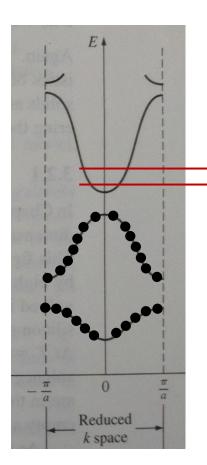
If the semiconductor is intrinsic:







If the semiconductor is intrinsic:



How many number of electrons in the conductance band per unit volume?

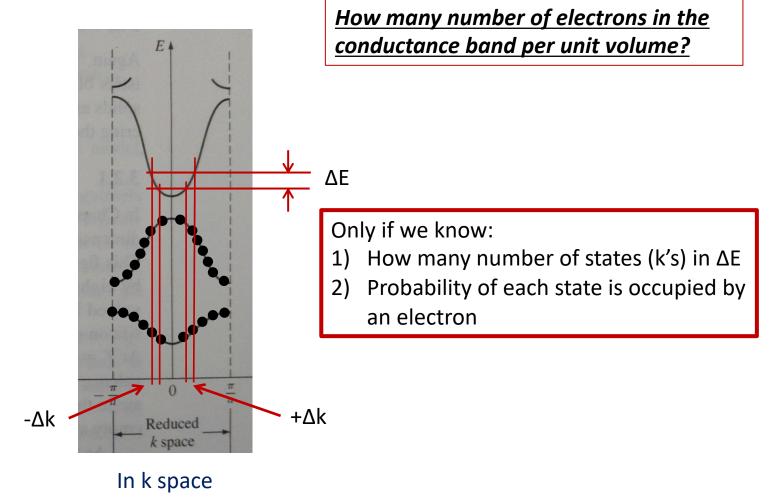
ΔΕ

Only if we know:

- 1) How many number of states in ΔE
- 2) Probability of each state is occupied by an electron

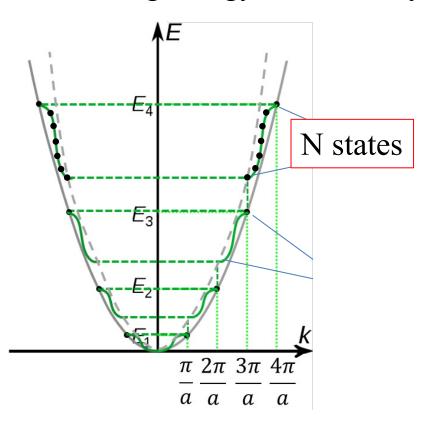


If the semiconductor is intrinsic:



3.1 Allowed and Forbidden Energy Bands

Forming energy bands: analytical



The total number of states in the whole crystal within $(0, \pi/a)$: N



The "density" of states in the whole crystal within $(0, \pi/a)$: $\frac{N}{\pi/a}$



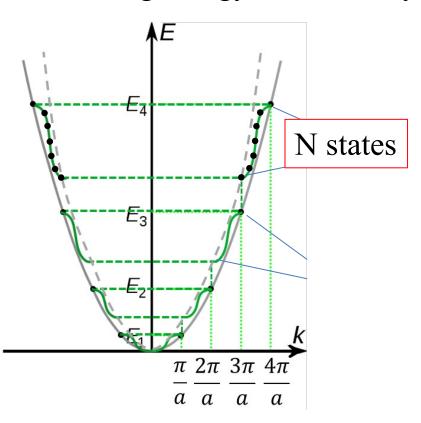
The number of states of the whole crystal within Δk : $\frac{N}{\pi/a} \times \Delta k$



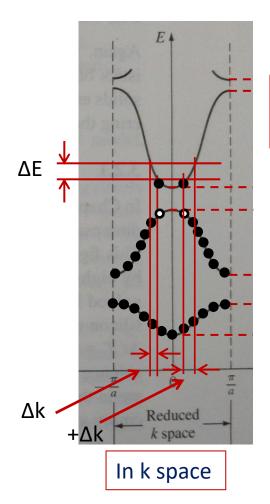
The number of states per unit volume within Δk : $\frac{N}{\pi/a} \times \Delta k \frac{1}{Na} = \frac{\Delta k}{\pi}$

3.1 Allowed and Forbidden Energy Bands

Forming energy bands: analytical

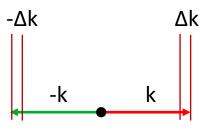


One-dimensional

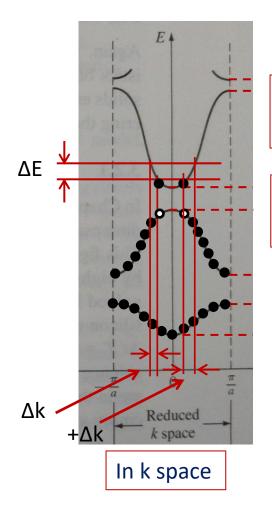


$$E = E(k) = E_c + \frac{\hbar^2}{2m_n^*}k^2$$

$$k = \mp \frac{\sqrt{2m_n^*(E - E_c)}}{\hbar}$$

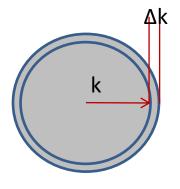


Two-dimensional



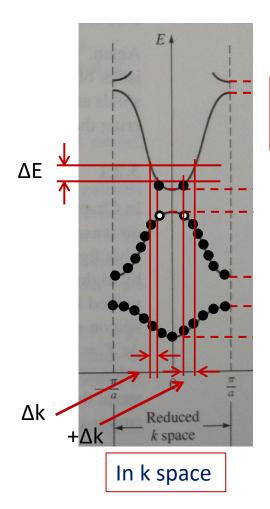
$$E = E(k) = E_c + \frac{\hbar^2}{2m_n^*}k^2$$

$$k = \mp \frac{\sqrt{2m_n^*(E - E_c)}}{\hbar}$$



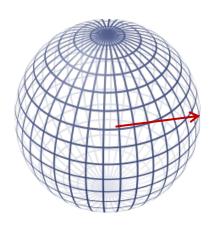


Three-dimensional



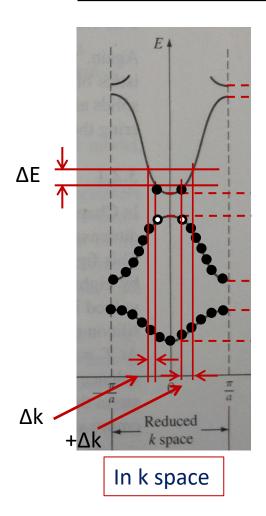
$$E = E(k) = E_c + \frac{\hbar^2}{2m_n^*}k^2$$

$$k = \mp \frac{\sqrt{2m_n^*(E - E_c)}}{\hbar}$$





Three-dimensional



spin
$$g(E) = \frac{dV_k}{dE} = \frac{2}{2} \frac{2\pi (2m^*)^{3/2}}{h^3} \sqrt{E - E_c}$$

The concept of electron spin was developed later, which is out of the scope of Schrodinger Equation.

Problem Example #2

Determine the number of quantum states (per unit volume) in silicon between (E_v-kT) and E_v at 300K.

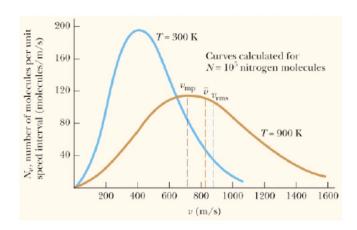
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Maxwell-Boltzmann probability function:

- distinguishable
- no limit on the particle number in each state
- Example: gas molecules in a container



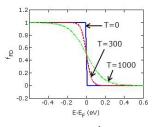
Bose-Einstein probability function:

- indistinguishable,
- no limit on the particle number in each state
- Example: photons

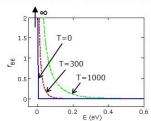
Fermi-Dirac probability function:

- indistinguishable
- one particle limit in each state
- Example: electrons in solids

Fermi-Dirac vs. Bose-Einstein Statistics



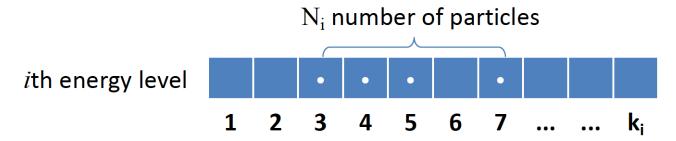
$$f_{\mathit{FD}}(E) = \frac{1}{\exp\!\left(\frac{E - E_{\mathit{F}}}{k_{\mathit{B}}T}\right) + 1}$$



$$f_{BB}(E) = \frac{1}{\exp\left(\frac{E}{k_B T}\right) - 1}$$

Fermi-Dirac probability function:

- indistinguishable
- one particle limit in each state
- Example: electrons in solids



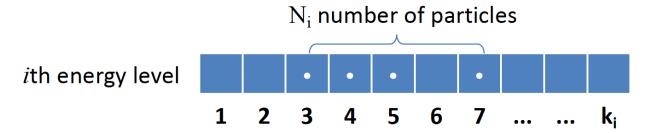
The totoal number of ways of arranging $N_{\rm i}$ particles in each ith energy level

$$k_i(k_i-1)\cdots(k_i-(N-1))=\frac{k_i!}{(k_i-N_i)}$$

(Particles are distinguishable)

Fermi-Dirac probability function:

- indistinguishable
- one particle limit in each state
- Example: electrons in solids



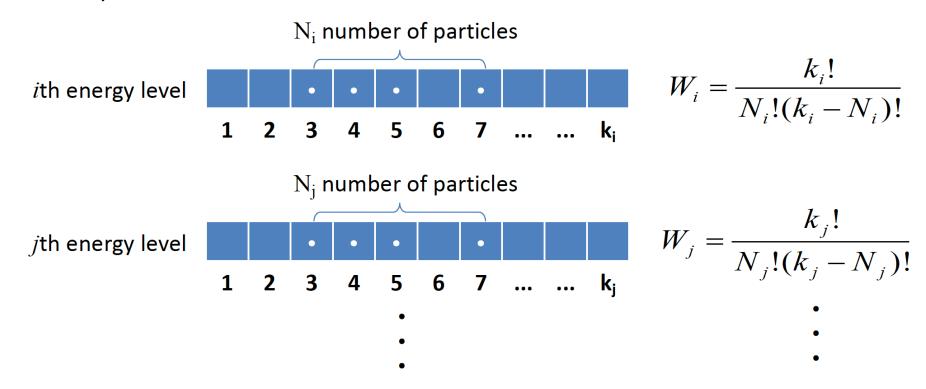
The totoal number of ways of arranging $N_{\rm i}$ indistinguishable particles in each ith energy level

$$W_i = \frac{k_i!}{N_i!(k_i - N_i)!}$$

(Particles are indistinguishable)

Fermi-Dirac probability function:

- indistinguishable
- one particle limit in each state
- Example: electrons in solids



For a given total number (N) of particles, the total number of ways of arranging indistiguishable particles among n energy levels is

$$W = \prod_{i=1}^{n} \frac{k_{i}!}{N_{i}!(k_{i} - N_{i})!}$$

$$f_{F}(E)$$

The highest probable distribution at following given constraints:

$$N = \sum_{i=1}^{n} N_i$$
 constant

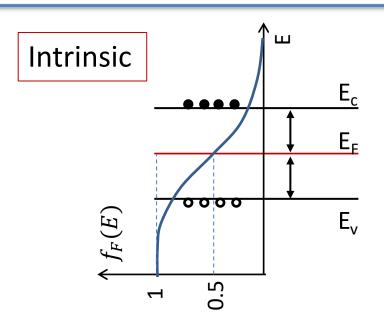
$$E_{total} = \sum_{i=1}^{n} E_i N_i$$
 constant

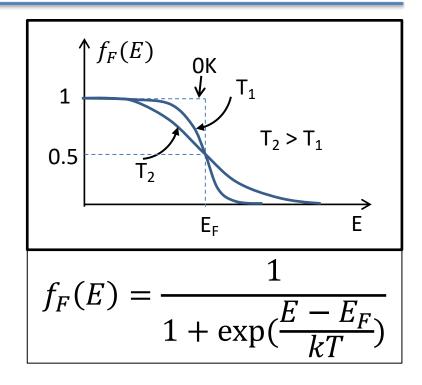
The probability of a state at energy E being occupied by an electron:

$$f_F(E) = \frac{1}{1 + \exp\left(\frac{E - E_F}{kT}\right)}$$

E is the energy level; E_F is the Fermi energy level; k is the Boltzmann constant; T is the absolute temperature.

3.6 Fermi distribution and Fermi level





Physical meaning of Fermi energy level:

At equilibrium, when an electron is added the system, the change of the system energy

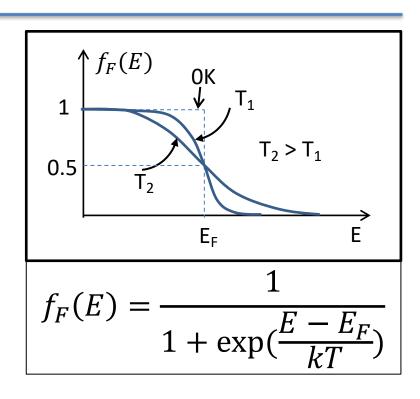
3.6 Boltzmann distribution

when
$$\exp\left(\frac{E-E_F}{kT}\right) \gg 1 \Rightarrow E-E_F > 3kT$$

$$f_F(E) = \frac{1}{1 + \exp(\frac{E - E_F}{kT})}$$

$$f_F(E) \approx \exp(-\frac{E - E_F}{kT})$$

Boltzmann distribution



3.6 Boltzmann distribution

Problem Example #3

Assume that the Fermi energy level is 0.35eV above the valence band energy. Let T=300K. Determine the probability of a state being empty of an electron at $E=E_{\rm v}-kT/2$.