Modern 2016 PS 23 Key

Q12D.4 In this problem, we will consider the covalent bond, which illustrates how "sharing" a quanton can lead to an attractive force that has no easy classical explanation. Consider an H₂⁺ molecule, where a pair of protons share a single electron. In order to make the problem tractable, we will ignore the repulsion between protons and model their interaction with the shared electron as if the protons were two square wells separated by a small distance, each about as wide as a hydrogen atom (0.1 nm). (This model may seem absurdly idealized, but more realistic models simply increase the complexity without adding any new insight.)

Run SchroSolver and choose the "Symmetric Well" potential energy function and a fixed well width of 0.5 nm. The ground state EEF should (by symmetry) predict that the electron is equally likely to be in either well, since the wells are identical. However, the eigenfunction might achieve this by having "bumps" in each well that have the same sign (a "symmetric" state) or opposite signs (an "antisymmetric" state). You can choose which is displayed by selecting the type from a pop-up menu in the window.

(a) Start with a center-to-center well separation of 0.75 nm. Show that the symmetric EEF has a slightly lower energy than the antisymmetric one (hand in a printout that provides supporting evidence).

(b) Explain qualitatively (physically) why this must be so. (Hint: Compare SchroSolver plots of the symmetric and antisymmetric functions, and explain why the wavelike "bump" in each well has gentler curvature in the symmetric case than in the antisymmetric case. How is that curvature linked to the energy of the state?)

(c) So the symmetric EEF is the true ground state for this system. Select this type of function and vary the wells' center-to-center separation (without changing anything else) from 0.55 nm in steps of 0.05 nm up to 1.10 nm, and make a chart of ground state energies (to at least 6 decimal places). You should see that the system's energy increases as we pull the wells apart.

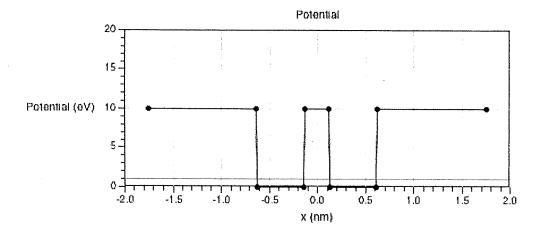
(d) Explain why this must be. (Hint: Again, first explain why the "bump" in each well has gentler curvature if

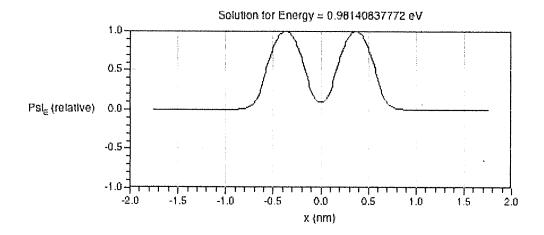
the wells are closer together.)

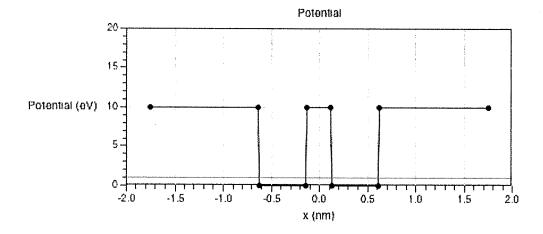
(e) Explain why this means that sharing the electron creates an attractive force between the wells. (Hint: Imagine pulling the wells apart. This will cause the system's energy to increase, as we have just seen. Then think about the work you must do on the system to do this.)

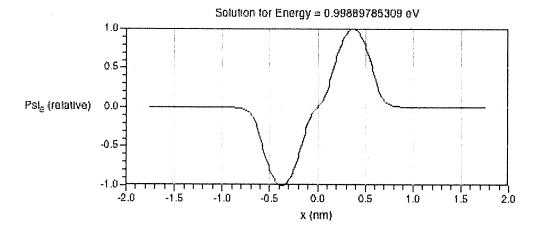
(f) Show that if the electron is in the antisymmetric state, sharing an electron leads to a repulsive force between the wells, and explain why this must be so.

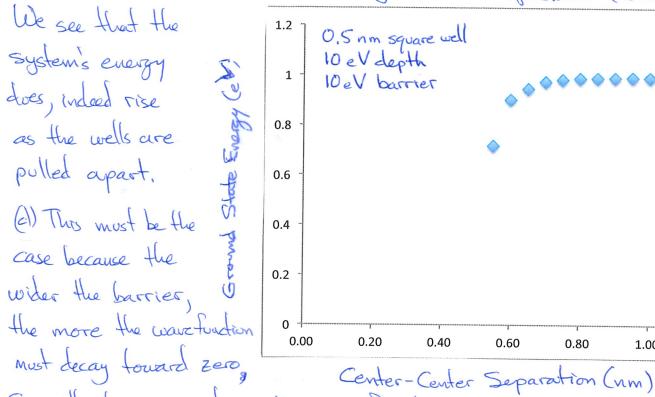
	3						
(a) _	I am runn	ing this on So	chrosolver. We see	e the symmetric			
(a) I am running this on Schrosolver. We see the symmetric solution has lower energy than the centisymmetric solution.							
Lsy	jm = 0,981 e	V Earl	G = 0.999eV (s	ee next page).			
Esym = 0.981 eV Earti = 0.999 eV (see next page). (b) This must be so because the antisymmetric solution has a node							
(zero crossing) in the barrier. This requires V_{anti} to drop a bit lower in value near the barrier edge, meaning the λ is slightly							
lower in value near the barrier edge, meaning the I is slightly							
shorter in the antisymmetric case,							
(1) I am systematically varying center-center separation and							
•	seeing the	effect on grow	d state energy	1			
	r _{cc}	Esym	rac .	Esym			
	0,55	0.716308	0,85	0,988889			
	0.60	0.900654	0.90	0.990111			
	0.65	0.944628	0.95	0.990481			
		5 972 163	1.00	0-,990744			
	0.70	0.973 153	_	500077			
	0.70	0.781408	1.05	0.990880			









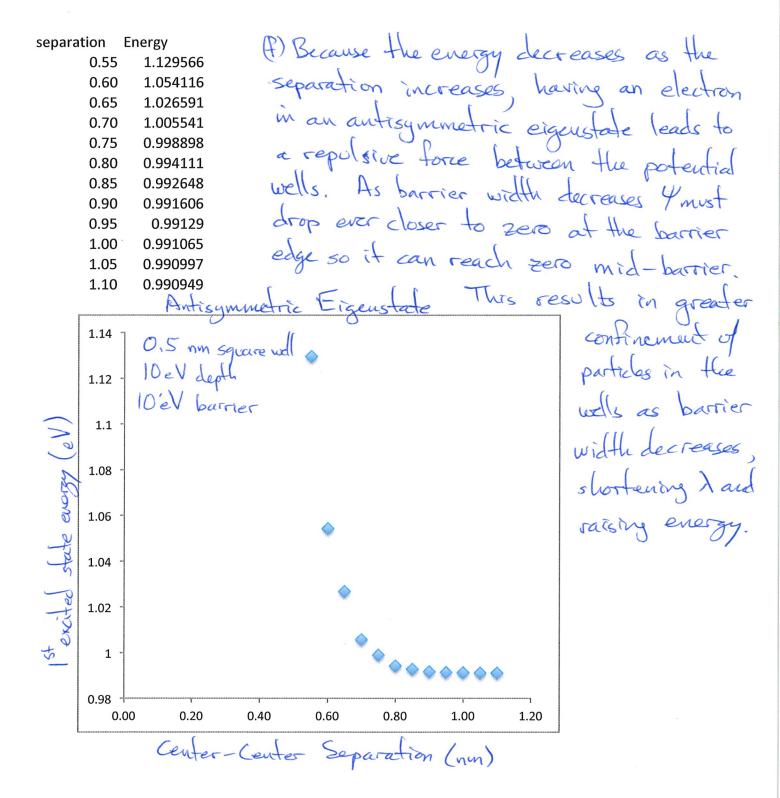


So with larger separation, the wavefunction is more confined in its square well, leading to shorter it and higher energy.

(e) Pulling the wells apart raises the energy of the symmetric eigenstate. So it requires work to pull them apart. Since $dW = \vec{F} \cdot d\vec{r}$, this means there is a force pulling the wells together (an attractive force).

(4) Now for the energy trent in the lowest autisymmetric state

Ccc	Earti	Vcc	Eanti
0.55 0.60 0.65 0.70 0.75	1.129566 1.054116 1.026591 1.005541 0.398898 0.994111	0.85 0.90 0.95 1.00 1.05	0,9926A8 0,991606 0,991290 0,991065 0,990997 0,996949



- Q13D.3 (Requires integral calculus.) One can calculate the gravitational binding energy of a spherical object (such as the earth) as follows. Assume that the earth's density is a constant $\rho = M/\left[\frac{4}{3}\pi R^3\right]$, where M is the mass of the earth and R is its radius. Imagine disassembling the earth by removing a thin shell of material of thickness dr then another shell of thickness dr and so on.
 - (a) Show that the energy required to move a shell of thickness dr away from the earth when its radius has already been pared down to r is then

$$E = \frac{3GM^2r^4dr}{R^6}$$
 (Q13.23)

- **(b)** Do an appropriate integral to find the gravitational binding energy of the sphere.
- (c) Use this to find the earth's gravitational mass deficit (in kilograms), assuming that the system's constituent parts are its atoms.

(a) The mass of the spherical shell is
$$dM = \rho A dr = \left(\frac{M}{3\pi}R^3\right)(A\pi r^2) dr = \frac{3Mr^2dr}{R^3}$$
The mass inside a radius r is
$$M(r) = \rho V = \left(\frac{M}{\frac{4}{3}\pi}R^3\right)\left(\frac{4}{3}\pi r^3\right) = \frac{Mr^3}{R^3}$$
The gravitational potential energy in place is then
$$U = \frac{GM(r)dM}{r} = -\frac{G}{r}\left(\frac{Mr^3}{R^3}\right)\left(\frac{3Mr^2}{R^3}\right) dr = \frac{3GM^2r^4dr}{R^6}$$
So the energy required to take this wass to infinity is
$$E = 0 - U = 3GM^2r^4dr$$

(b) To get the total gravitational binding energy we integrate $E_b = \int dr E(r) = \int \frac{3GM^2}{R^6} r dr = \frac{3GM^2}{R^6} \frac{r^5}{5} \Big|_0^2 = \frac{3GM^2}{5} \frac{3GM^2}{5} \Big|_0^2 = \frac{$