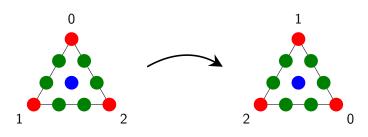
# Duality and Symmetry in Finite Element Exterior Calculus

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Pennsylvania State University

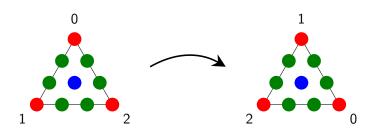
June 19-25, 2022

## Symmetry of Scalar Elements



$$\mathcal{P}_3\Lambda^0(\mathcal{T}^2) = \left\langle \lambda_0^3, \lambda_1^3, \lambda_2^3, \lambda_1^2\lambda_2, \lambda_2^2\lambda_1, \lambda_2^2\lambda_0, \lambda_0^2\lambda_2, \lambda_0^2\lambda_1, \lambda_1^2\lambda_0, \lambda_0\lambda_1\lambda_2 \right\rangle.$$

## Symmetry of Scalar Elements



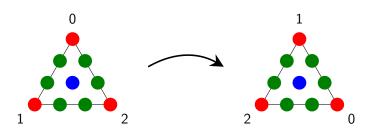
$$\mathcal{P}_3\Lambda^0(T^2) = \left\langle \lambda_0^3, \lambda_1^3, \lambda_2^3, \lambda_1^2\lambda_2, \lambda_2^2\lambda_1, \lambda_2^2\lambda_0, \lambda_0^2\lambda_2, \lambda_0^2\lambda_1, \lambda_1^2\lambda_0, \lambda_0\lambda_1\lambda_2 \right\rangle.$$

• When computing matrix of, e.g.,  $a(u, v) = \int_{\mathcal{T}^2} \nabla u \cdot \nabla v$ , can exploit sixfold symmetry of  $\mathcal{T}^2$  to compute fewer entries.

$$\begin{split} a\left(\lambda_0^3, \lambda_1^2 \lambda_2\right) &= a\left(\lambda_1^3, \lambda_2^2 \lambda_0\right) = a\left(\lambda_2^3, \lambda_0^2 \lambda_1\right) \\ &= a\left(\lambda_0^3, \lambda_2^2 \lambda_1\right) = a\left(\lambda_1^3, \lambda_0^2 \lambda_2\right) = a\left(\lambda_2^3, \lambda_1^2 \lambda_0\right) \end{split}$$



## Symmetry of Scalar Elements



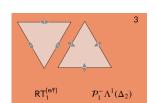
$$\mathcal{P}_3\Lambda^0(T^2) = \left\langle \lambda_0^3, \lambda_1^3, \lambda_2^3, \lambda_1^2\lambda_2, \lambda_2^2\lambda_1, \lambda_2^2\lambda_0, \lambda_0^2\lambda_2, \lambda_0^2\lambda_1, \lambda_1^2\lambda_0, \lambda_0\lambda_1\lambda_2 \right\rangle.$$

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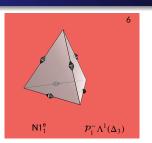
$$a\left(\lambda_0^3, \lambda_1^2 \lambda_2\right) = a\left(\lambda_1^3, \lambda_2^2 \lambda_0\right) = a\left(\lambda_2^3, \lambda_0^2 \lambda_1\right)$$
$$= a\left(\lambda_0^3, \lambda_2^2 \lambda_1\right) = a\left(\lambda_1^3, \lambda_0^2 \lambda_2\right) = a\left(\lambda_2^3, \lambda_1^2 \lambda_0\right)$$

• More generally,  $\int_{\mathcal{T}^2} g^{-1} (du \otimes dv) \sqrt{\det g} = \sqrt{\det g} g^{-1} \left( \int_{\mathcal{T}^2} du \otimes dv \right).$ 

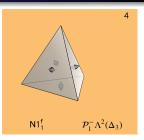
## Symmetry of Vector Elements Whitney Elements



$$\langle \lambda_1 d\lambda_2 - \lambda_2 d\lambda_1,$$
  
 $\lambda_2 d\lambda_0 - \lambda_0 d\lambda_2,$   
 $\lambda_0 d\lambda_1 - \lambda_1 d\lambda_0 \rangle.$ 



$$\begin{split} & \left\langle \lambda_1 \, d\lambda_2 - \lambda_2 \, d\lambda_1, \right. \\ & \left. \lambda_2 \, d\lambda_0 - \lambda_0 \, d\lambda_2, \right. \\ & \left. \lambda_0 \, d\lambda_1 - \lambda_1 \, d\lambda_0, \right. \\ & \left. \lambda_0 \, d\lambda_3 - \lambda_3 \, d\lambda_0, \right. \\ & \left. \lambda_1 \, d\lambda_3 - \lambda_3 \, d\lambda_1, \right. \\ & \left. \lambda_2 \, d\lambda_3 - \lambda_3 \, d\lambda_2 \right\rangle. \end{split}$$

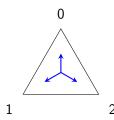


$$\langle \lambda_1 d\lambda_2 \wedge d\lambda_3 \\ + \lambda_2 d\lambda_3 \wedge d\lambda_1 \\ + \lambda_3 d\lambda_1 \wedge d\lambda_2, \\ \dots, \\ \lambda_0 d\lambda_1 \wedge d\lambda_2 \\ + \lambda_1 d\lambda_2 \wedge d\lambda_0 \\ + \lambda_2 d\lambda_0 \wedge d\lambda_1 \rangle$$

Geometric symmetry  $\Rightarrow$  basis symmetry (up to sign).

## Symmetry of Vector Elements

Lack of Symmetric Bases



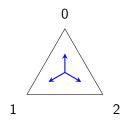
$$\mathcal{P}_0 \Lambda^1(T^2)$$

$$= \langle d\lambda_0, d\lambda_1, d\lambda_2 \rangle,$$

$$d\lambda_0 + d\lambda_1 + d\lambda_2 = 0$$

## Symmetry of Vector Elements

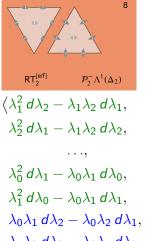
Lack of Symmetric Bases

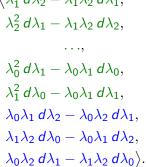


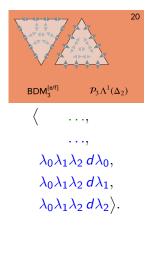
$$\mathcal{P}_0 \Lambda^1(T^2)$$

$$= \langle d\lambda_0, d\lambda_1, d\lambda_2 \rangle,$$

$$d\lambda_0 + d\lambda_1 + d\lambda_2 = 0$$







#### Results

#### Theorem (if: Licht, 2019; only if: YBK, 2021)

The following spaces have symmetry-invariant bases up to sign if and only if the corresponding condition holds.

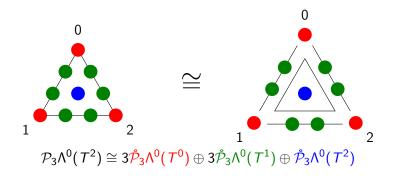
$$\mathcal{P}_r\Lambda^1(T^2)$$
 if and only if  $r \notin 3\mathbb{N}_0$ ,  $\mathcal{P}_r^-\Lambda^1(T^2)$  if and only if  $r \notin 3\mathbb{N}_0 + 2$ .

#### Theorem (YBK, 2021)

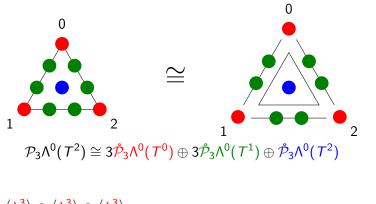
The following spaces have symmetry-invariant bases up to sign if and only if the corresponding condition holds.

$$\mathcal{P}_r\Lambda^1(T^3)$$
 always,  
 $\mathcal{P}_r^-\Lambda^1(T^3)$  if and only if  $r \notin 3\mathbb{N}_0 + 2$ ,  
 $\mathcal{P}_r\Lambda^2(T^3)$  always,  
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#### Methods Recursion

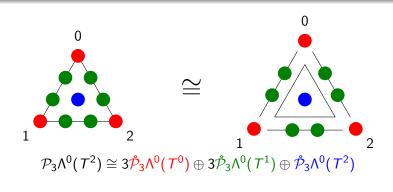


#### Methods Recursion



$$\begin{split} \left\langle \lambda_0^3 \right\rangle \oplus \left\langle \lambda_1^3 \right\rangle \oplus \left\langle \lambda_2^3 \right\rangle \\ \oplus \left\langle \lambda_1^2 \lambda_2, \lambda_2^2 \lambda_1 \right\rangle \oplus \left\langle \lambda_2^2 \lambda_0, \lambda_0^2 \lambda_2 \right\rangle \oplus \left\langle \lambda_0^2 \lambda_1, \lambda_1^2 \lambda_0 \right\rangle \oplus \left\langle \lambda_0 \lambda_1 \lambda_2 \right\rangle \end{split}$$

#### Methods Recursion

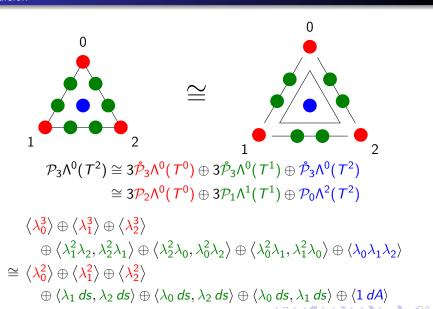


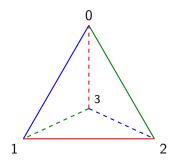
$$\langle \lambda_0^3 \rangle \oplus \langle \lambda_1^3 \rangle \oplus \langle \lambda_2^3 \rangle$$

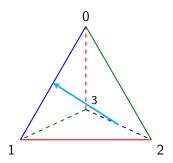
$$\oplus \langle \lambda_1^2 \lambda_2, \lambda_2^2 \lambda_1 \rangle \oplus \langle \lambda_2^2 \lambda_0, \lambda_0^2 \lambda_2 \rangle \oplus \langle \lambda_0^2 \lambda_1, \lambda_1^2 \lambda_0 \rangle \oplus \langle \lambda_0 \lambda_1 \lambda_2 \rangle$$

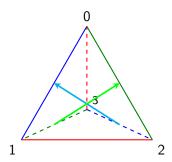
$$\cong \langle \lambda_0^2 \rangle \oplus \langle \lambda_1^2 \rangle \oplus \langle \lambda_2^2 \rangle$$

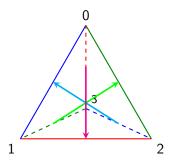
$$\oplus \langle \lambda_1 \, ds, \lambda_2 \, ds \rangle \oplus \langle \lambda_0 \, ds, \lambda_2 \, ds \rangle \oplus \langle \lambda_0 \, ds, \lambda_1 \, ds \rangle \oplus \langle 1 \, dA \rangle$$



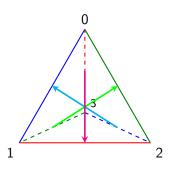








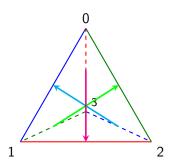
#### Tetrahedron Basis



$$\mathcal{P}_0 \Lambda^1 (T^3)$$

$$= \langle d\lambda_0 + d\lambda_1 - d\lambda_2 - d\lambda_3, d\lambda_0 + d\lambda_2 - d\lambda_1 - d\lambda_3, d\lambda_1 + d\lambda_2 - d\lambda_0 - d\lambda_3 \rangle$$

$$=: \langle \alpha, \beta, \gamma \rangle.$$



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$$=: \langle \alpha, \beta, \gamma \rangle.$$

$$\mathcal{P}_{2}\Lambda^{1}(T^{3})$$

$$= \mathcal{P}_{2}\Lambda^{0}(T^{3}) \otimes \mathcal{P}_{0}\Lambda^{1}(T^{3})$$

$$= \langle \lambda_{0}^{2}\alpha, \lambda_{0}^{2}\beta, \lambda_{0}^{2}\gamma, \lambda_{1}^{2}\alpha, \lambda_{1}^{2}\beta, \lambda_{1}^{2}\gamma, \lambda_{2}^{2}\alpha, \lambda_{2}^{2}\beta, \lambda_{2}^{2}\gamma, \lambda_{3}^{2}\alpha, \lambda_{3}^{2}\beta, \lambda_{3}^{2}\gamma, \lambda_{0}\lambda_{1}\alpha, \lambda_{0}\lambda_{1}\beta, \lambda_{0}\lambda_{1}\gamma, \lambda_{0}\lambda_{2}\alpha, \lambda_{0}\lambda_{2}\beta, \lambda_{0}\lambda_{2}\gamma, \lambda_{0}\lambda_{3}\alpha, \lambda_{0}\lambda_{3}\beta, \lambda_{0}\lambda_{3}\gamma, \lambda_{1}\lambda_{2}\alpha, \lambda_{1}\lambda_{2}\beta, \lambda_{1}\lambda_{2}\gamma, \lambda_{1}\lambda_{3}\alpha, \lambda_{1}\lambda_{3}\beta, \lambda_{1}\lambda_{3}\gamma, \lambda_{2}\lambda_{3}\alpha, \lambda_{2}\lambda_{3}\beta, \lambda_{2}\lambda_{3}\gamma, \lambda_{2}\lambda_{3}\alpha, \lambda_{2}\lambda_{3}\beta, \lambda_{2}\lambda_{3}\gamma \rangle.$$

#### Representations of $\mathbb{Z}/3$

- The 1D representation  ${\bf 1}$  where  $\mathbb{Z}/3$  acts trivially.
- The 2D representation **2** where  $\mathbb{Z}/3$  acts by  $120^{\circ}$  rotations.
- The 3D representation **3** where  $\mathbb{Z}/3$  acts by permuting the coordinates.
  - ${f 3}\cong {f 1}\oplus {f 2}$  because  $\langle (1,1,1) \rangle$  is an invariant subspace.

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#### Invariant bases

1 and 3 have symmetry-invariant bases, but 2 does not.



#### Proposition

A representation  $V \cong m\mathbf{1} \oplus n\mathbf{2}$  has a  $\mathbb{Z}/3$ -invariant basis up to sign if and only if  $m \geq n$ .

#### References



Martin Licht.

Symmetry and invariant bases in finite element exterior calculus.

https://arxiv.org/abs/1912.11002.



Yakov Berchenko-Kogan.

Symmetric bases for finite element exterior calculus spaces.

https://arxiv.org/abs/2112.06065.



Douglas N. Arnold, Richard S. Falk, and Ragnar Winther. Finite element exterior calculus, homological techniques, and applications. *Acta Numer.*, 15:1–155, 2006.



D. N. Arnold and A. Logg. Periodic Table of the Finite Elements. SIAM News, 47(9), 2014.

#### References



Martin Licht.

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On basis constructions in finite element exterior calculus.

Adv. Comput. Math., 48(2), 2022.



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Duality in finite element exterior calculus and Hodge duality on the sphere.

Found. Comput. Math., 21(5):1153-1180, 2021.



Previously...

Previously...

#### Previously...

0 — 1

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$$\langle \lambda_0 ds, \lambda_1 ds \rangle \cong \langle \lambda_0^2 \lambda_1, \lambda_1^2 \lambda_0 \rangle$$
  
 $\mathcal{P}_1 \Lambda^1(T^1) \cong \mathring{\mathcal{P}}_3 \Lambda^0(T^1)$ 

#### Previously...

$$\begin{split} \langle \lambda_0 \; \textit{ds}, \lambda_1 \; \textit{ds} \rangle &\cong \left\langle \lambda_0^2 \lambda_1, \lambda_1^2 \lambda_0 \right\rangle \\ \mathcal{P}_1 \Lambda^1(\mathcal{T}^1) &\cong \mathring{\mathcal{P}}_3 \Lambda^0(\mathcal{T}^1) \end{split}$$

#### FEEC Duality (Arnold, Falk, and Winther, 2006)

$$\mathcal{P}_r \Lambda^k(T^n) \cong \mathring{\mathcal{P}}_{r+k+1}^- \Lambda^{n-k}(T^n),$$
  
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#### An explicit map (Licht, 2018)

$$\begin{split} \mathcal{P}_1 \Lambda^1(T^2) &\to \mathring{\mathcal{P}}_3^- \Lambda^1(T^2), \qquad \quad \mathcal{P}_1^- \Lambda^1(T^2) \to \mathring{\mathcal{P}}_2 \Lambda^1(T^2), \\ \lambda_1 \, d\lambda_1 &\mapsto \lambda_0 \lambda_1^2 \, d\lambda_2 - \lambda_1^2 \lambda_2 \, d\lambda_0, \quad \lambda_0 \, d\lambda_1 - \lambda_1 \, d\lambda_0 \mapsto \lambda_0 \lambda_1 \, d\lambda_2. \end{split}$$

#### Previously...

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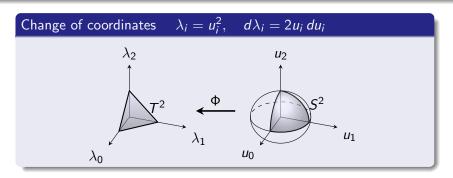
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#### The Hodge star (YBK, 2019)

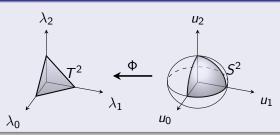
The two maps are the same; have formula using Hodge star on  $S^n$ .

## The sphere



## The sphere

### Change of coordinates $\lambda_i = u_i^2$ , $d\lambda_i = 2u_i du_i$



#### The duality map

- **①** Change coordinates to the sphere  $\Phi^* : \Lambda^k(T^n) \to \Lambda^k(S^n)$ .
- 2 Apply the Hodge star on the sphere.
- **3** Multiply by the bubble function  $u_N := u_0 \cdots u_n$ .
- Ohange coordinates back to the simplex.

$$(\Phi^*)^{-1} \circ u_N *_{S^n} \circ \Phi^*$$

$$a=\lambda_1\,d\lambda_1\in \mathcal{P}_1 \Lambda^1(\mathcal{T}^2)$$

$$a = \lambda_1 d\lambda_1 \in \mathcal{P}_1 \Lambda^1(T^2)$$

**1** 
$$\alpha = \Phi^* a = 2u_1^3 du_1 \in \mathcal{P}_3 \Lambda^1(S^2).$$

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- $\bullet b = (\Phi^*)^{-1} a = \lambda_0 \lambda_1^2 d\lambda_2 \lambda_1^2 \lambda_2 d\lambda_0 \in \mathring{\mathcal{P}}_3^- \Lambda^1(T^2).$

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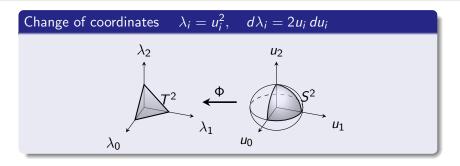
## $a = \lambda_1 d\lambda_1 \in \mathcal{P}_1 \Lambda^1(T^2)$

- $\bullet b = (\Phi^*)^{-1} a = \lambda_0 \lambda_1^2 d\lambda_2 \lambda_1^2 \lambda_2 d\lambda_0 \in \mathring{\mathcal{P}}_3^- \Lambda^1(T^2).$

## $a = \lambda_0 d\lambda_1 - \lambda_1 d\lambda_0 \in \mathcal{P}_1^- \Lambda^1(T^2)$

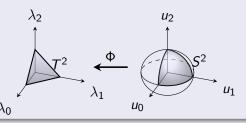
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## Polynomial forms on the simplex and the sphere



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#### Theorem

The map  $\Phi^* : \Lambda^k(T^n) \to \Lambda^k(S^n)$  gives isomorphisms:

$$\mathcal{P}_{r}\Lambda^{k}(T^{n}) \xrightarrow{\simeq} \mathcal{P}_{2r+k}\Lambda_{e}^{k}(S^{n}),$$

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## $\mathcal{P}_r \Lambda^k(T^n) \cong \mathcal{P}_{2r+k} \Lambda^k_e(S^n)$

$$\lambda_0 \lambda_1^2 \, d\lambda_2 \wedge d\lambda_3 \mapsto u_0^2 u_1^4 (2u_2 \, du_2) \wedge (2u_3 \, du_3)$$

$$= 4u_0^2 u_1^4 u_2 u_3 \, du_2 \wedge du_3$$

$$\mathcal{P}_3 \Lambda^2 (T^3) \to \mathcal{P}_8 \Lambda_e^2 (S^3).$$

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#### Example

$$\begin{split} \lambda_0 \lambda_1^2 \, d\lambda_2 \wedge d\lambda_3 &\mapsto u_0^2 u_1^4 (2u_2 \, du_2) \wedge (2u_3 \, du_3) \\ &= 4u_0^2 u_1^4 u_2 u_3 \, du_2 \wedge du_3 \\ \mathcal{P}_3 \Lambda^2 (\mathcal{T}^3) &\to \mathcal{P}_8 \Lambda_e^2 (\mathcal{S}^3). \end{split}$$

#### Definition

- A form is even if it is invariant under all coordinate reflections.
  - e.g.  $R_2: (u_0, u_1, u_2, u_3) \mapsto (u_0, u_1, -u_2, u_3)$ .
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#### The image of $\Phi^*$ is even



$$\mathcal{P}_r^- \Lambda^k(T^n) \cong \mathcal{P}_{2r+k}^- \Lambda_e^k(S^n)$$

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• Let X denote the radial vector field

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#### A new definition of $\mathcal{P}_r^-$

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• Key fact:  $\Phi_* X = 2X$ .

$$\sum_{i=0}^{n} u_{i} \frac{\partial}{\partial u_{i}} = \sum_{i=0}^{n} u_{i} \frac{\partial \lambda_{i}}{\partial u_{i}} \frac{\partial}{\partial \lambda_{i}} = \sum_{i=0}^{n} u_{i} (2u_{i}) \frac{\partial}{\partial \lambda_{i}} = 2 \sum_{i=0}^{n} \lambda_{i} \frac{\partial}{\partial \lambda_{i}}.$$

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- $D\Phi$  maps vectors normal to  $S_i^{n-1}$  to zero.  $(\frac{\partial u_i^2}{\partial u_i} = 0 \text{ on } S_i^{n-1}.)$

### Recap

#### Theorem

The map  $\Phi^*: \Lambda^k(T^n) \to \Lambda^k(S^n)$  gives isomorphisms:  $\mathcal{P}_r \Lambda^k(T^n) \xrightarrow{\simeq} \mathcal{P}_{2r+k} \Lambda^k_e(S^n),$   $\mathcal{P}_r^- \Lambda^k(T^n) \xrightarrow{\simeq} \mathcal{P}_{2r+k}^- \Lambda^k_e(S^n),$   $\mathcal{P}_r \Lambda^k(T^n) \xrightarrow{\simeq} \mathcal{P}_{2r+k}^- \Lambda^k_e(S^n),$ 

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$$(\Phi^*)^{-1} \circ u_{N} *_{S^n} \circ \Phi^* : \qquad \mathcal{P}_r \Lambda^k(T^n) \cong \mathring{\mathcal{P}}_{r+k+1}^- \Lambda^{n-k}(T^n),$$
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#### $\mathsf{Theorem}$

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### Proposition

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$$*_{S^2}(*_{S^2} \alpha) = u_1^3 u_0(u_1 du_0 - u_0 du_1) - u_1^3 u_2(u_2 du_1 - u_1 du_2)$$

$$= -u_1^3(u_0^2 + u_1^2 + u_2^2) du_1 + \frac{1}{2}u_1^4 d(u_0^2 + u_1^2 + u_2^2)$$

$$= -\alpha \in \mathcal{P}_3 \Lambda^1(S^2).$$

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### It's okay

- These are the only counterexamples.
- Proposition still holds if we add span $\{1\}$  or span $\{vol_{S^n}\}$  to the left-hand side when necessary.

### The duality map

$$u_{N}*_{S^{n}}: \begin{array}{c} \mathcal{P}_{2r+k}\Lambda_{e}^{k}(S^{n}) \cong \tilde{\mathcal{P}}_{2r+n+k+2}^{-}\Lambda_{e}^{n-k}(S^{n}), \\ \mathcal{P}_{2r+k}^{-}\Lambda_{e}^{k}(S^{n}) \cong \tilde{\mathcal{P}}_{2r+n+k}\Lambda_{e}^{n-k}(S^{n}). \end{array}$$

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### Proposition

Restricting to  $T^n$ ,

$$\hat{\mathcal{P}}_r^- \Lambda^k(T^n) = \mathcal{P}_r^- \Lambda^k(T^n)$$

except 
$$\hat{\mathcal{P}}_0^- \Lambda^0(T^n) = \operatorname{span}\{1\}.$$

### While we're at it

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#### The volume form

$$\operatorname{vol}_{S^n} = u_0^{-1} du_1 \wedge \cdots du_n.$$

Perhaps it should be in  $\mathcal{P}_{-1}\Lambda^n(S^n)$  after all?



The cohomology of smooth closed manifolds

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### Duality

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## Thank you



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