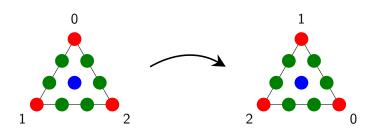
## Symmetry in Finite Element Exterior Calculus

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## Symmetry of Scalar Elements



$$\mathcal{P}_3\Lambda^0(\mathcal{T}^2) = \left\langle \lambda_0^3, \lambda_1^3, \lambda_2^3, \lambda_1^2\lambda_2, \lambda_2^2\lambda_1, \lambda_2^2\lambda_0, \lambda_0^2\lambda_2, \lambda_0^2\lambda_1, \lambda_1^2\lambda_0, \lambda_0\lambda_1\lambda_2 \right\rangle.$$

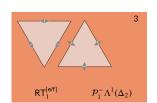
• When computing matrix of, e.g.,  $a(u, v) = \int_{T^2} \nabla u \cdot \nabla v$ , can exploit sixfold symmetry of  $T^2$  to compute fewer entries.

$$a(\lambda_0^3, \lambda_1^2 \lambda_2) = a(\lambda_0^3, \lambda_2^2 \lambda_1)$$

$$= a(\lambda_1^3, \lambda_2^2 \lambda_0) = a(\lambda_1^3, \lambda_0^2 \lambda_2)$$

$$= a(\lambda_2^3, \lambda_0^2 \lambda_1) = a(\lambda_2^3, \lambda_1^2 \lambda_0)$$

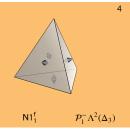
# Symmetry of Vector Elements Whitney Elements



$$\langle \lambda_1 d\lambda_2 - \lambda_2 d\lambda_1, \lambda_2 d\lambda_0 - \lambda_0 d\lambda_2, \lambda_0 d\lambda_1 - \lambda_1 d\lambda_0 \rangle.$$



$$\begin{split} & \left\langle \lambda_1 \, d\lambda_2 - \lambda_2 \, d\lambda_1, \right. \\ & \left. \lambda_2 \, d\lambda_0 - \lambda_0 \, d\lambda_2, \right. \\ & \left. \lambda_0 \, d\lambda_1 - \lambda_1 \, d\lambda_0, \right. \\ & \left. \lambda_0 \, d\lambda_3 - \lambda_3 \, d\lambda_0, \right. \\ & \left. \lambda_1 \, d\lambda_3 - \lambda_3 \, d\lambda_1, \right. \\ & \left. \lambda_2 \, d\lambda_3 - \lambda_3 \, d\lambda_2 \right\rangle. \end{split}$$

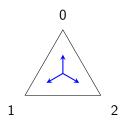


$$\langle \lambda_1 d\lambda_2 \wedge d\lambda_3 \\ + \lambda_2 d\lambda_3 \wedge d\lambda_1 \\ + \lambda_3 d\lambda_1 \wedge d\lambda_2, \\ \dots, \\ \lambda_0 d\lambda_1 \wedge d\lambda_2 \\ + \lambda_1 d\lambda_2 \wedge d\lambda_0 \\ + \lambda_2 d\lambda_0 \wedge d\lambda_1 \rangle$$

Geometric symmetry  $\Rightarrow$  basis symmetry (up to sign).

## Symmetry of Vector Elements

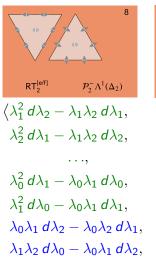
Lack of Symmetric Bases

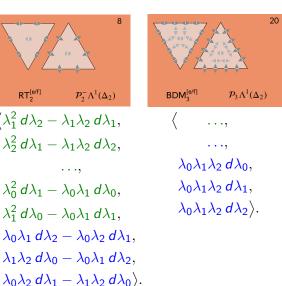


$$\mathcal{P}_0 \Lambda^1(T^2)$$

$$= \langle d\lambda_0, d\lambda_1, d\lambda_2 \rangle,$$

$$d\lambda_0 + d\lambda_1 + d\lambda_2 = 0$$





#### Results

### Theorem (if: Licht, 2019; only if: YBK, 2021)

The following spaces have symmetry-invariant bases up to sign if and only if the corresponding condition holds.

$$\mathcal{P}_r\Lambda^1(T^2)$$
 if and only if  $r \notin 3\mathbb{N}_0$ ,  $\mathcal{P}_r^-\Lambda^1(T^2)$  if and only if  $r \notin 3\mathbb{N}_0 + 2$ .

#### Theorem (YBK, 2021)

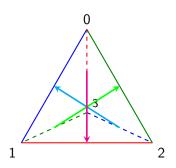
The following spaces have symmetry-invariant bases up to sign if and only if the corresponding condition holds.

$$\begin{array}{lll} \mathcal{P}_r \Lambda^1(T^3) & \text{always}, \\ \mathcal{P}_r^- \Lambda^1(T^3) & \text{if and only if} & r \notin 3\mathbb{N}_0 + 2, \\ \mathcal{P}_r \Lambda^2(T^3) & \text{always}, \\ \mathcal{P}_r^- \Lambda^2(T^3) & \text{always}. \end{array}$$

#### Methods Recursion

$$\begin{array}{c}
0 \\
1 \\
2 \\
\mathcal{P}_{3}\Lambda^{0}(T^{2}) \cong 3\mathring{\mathcal{P}}_{3}\Lambda^{0}(T^{0}) \oplus 3\mathring{\mathcal{P}}_{3}\Lambda^{0}(T^{1}) \oplus \mathring{\mathcal{P}}_{3}\Lambda^{0}(T^{2}) \\
\cong 3\mathcal{P}_{2}\Lambda^{0}(T^{0}) \oplus 3\mathcal{P}_{1}\Lambda^{1}(T^{1}) \oplus \mathcal{P}_{0}\Lambda^{2}(T^{2}) \\
\langle \lambda_{0}^{3} \rangle \oplus \langle \lambda_{1}^{3} \rangle \oplus \langle \lambda_{2}^{3} \rangle \\
\oplus \langle \lambda_{1}^{2}\lambda_{2}, \lambda_{2}^{2}\lambda_{1} \rangle \oplus \langle \lambda_{2}^{2}\lambda_{0}, \lambda_{0}^{2}\lambda_{2} \rangle \oplus \langle \lambda_{0}^{2}\lambda_{1}, \lambda_{1}^{2}\lambda_{0} \rangle \oplus \langle \lambda_{0}\lambda_{1}\lambda_{2} \rangle \\
\cong \langle \lambda_{0}^{2} \rangle \oplus \langle \lambda_{1}^{2} \rangle \oplus \langle \lambda_{2}^{2} \rangle \\
\oplus \langle \lambda_{1} ds, \lambda_{2} ds \rangle \oplus \langle \lambda_{0} ds, \lambda_{2} ds \rangle \oplus \langle \lambda_{0} ds, \lambda_{1} ds \rangle \oplus \langle 1 dA \rangle
\end{array}$$

#### Methods Tetrahedron Basis



$$\mathcal{P}_0 \Lambda^1(T^3)$$

$$= \langle d\lambda_0 + d\lambda_1 - d\lambda_2 - d\lambda_3, d\lambda_0 + d\lambda_2 - d\lambda_1 - d\lambda_3, d\lambda_1 + d\lambda_2 - d\lambda_0 - d\lambda_3 \rangle$$

$$=: \langle \alpha, \beta, \gamma \rangle.$$

$$\mathcal{P}_{2}\Lambda^{1}(T^{3})$$

$$= \mathcal{P}_{2}\Lambda^{0}(T^{3}) \otimes \mathcal{P}_{0}\Lambda^{1}(T^{3})$$

$$= \langle \lambda_{0}^{2}\alpha, \lambda_{0}^{2}\beta, \lambda_{0}^{2}\gamma, \lambda_{1}^{2}\alpha, \lambda_{1}^{2}\beta, \lambda_{1}^{2}\gamma, \lambda_{2}^{2}\alpha, \lambda_{2}^{2}\beta, \lambda_{2}^{2}\gamma, \lambda_{3}^{2}\alpha, \lambda_{3}^{2}\beta, \lambda_{3}^{2}\gamma, \lambda_{0}\lambda_{1}\alpha, \lambda_{0}\lambda_{1}\beta, \lambda_{0}\lambda_{1}\gamma, \lambda_{0}\lambda_{2}\alpha, \lambda_{0}\lambda_{2}\beta, \lambda_{0}\lambda_{2}\gamma, \lambda_{0}\lambda_{3}\alpha, \lambda_{0}\lambda_{3}\beta, \lambda_{0}\lambda_{3}\gamma, \lambda_{1}\lambda_{2}\alpha, \lambda_{1}\lambda_{2}\beta, \lambda_{1}\lambda_{2}\gamma, \lambda_{1}\lambda_{3}\alpha, \lambda_{1}\lambda_{3}\beta, \lambda_{1}\lambda_{3}\gamma, \lambda_{2}\lambda_{3}\alpha, \lambda_{2}\lambda_{3}\beta, \lambda_{2}\lambda_{3}\gamma \rangle.$$

#### Methods Obstructions

#### Representations of $\mathbb{Z}/3$

- The 1D representation **1** where  $\mathbb{Z}/3$  acts trivially.
- The 2D representation **2** where  $\mathbb{Z}/3$  acts by  $120^{\circ}$  rotations.
- The 3D representation **3** where  $\mathbb{Z}/3$  acts by permuting the coordinates.
  - ${f 3}\cong {f 1}\oplus {f 2}$  because  $\langle (1,1,1) \rangle$  is an invariant subspace.

#### Invariant bases

1 and 3 have symmetry-invariant bases, but 2 does not.



#### Proposition

A representation  $V \cong m\mathbf{1} \oplus n\mathbf{2}$  has a  $\mathbb{Z}/3$ -invariant basis up to sign if and only if  $m \geq n$ .

#### Thank You and References



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