

When quantum memory is useful for dense coding

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Why am I talking about this at resource workshop?

Question: In dense coding, when does receiver want to possess quantum memory?

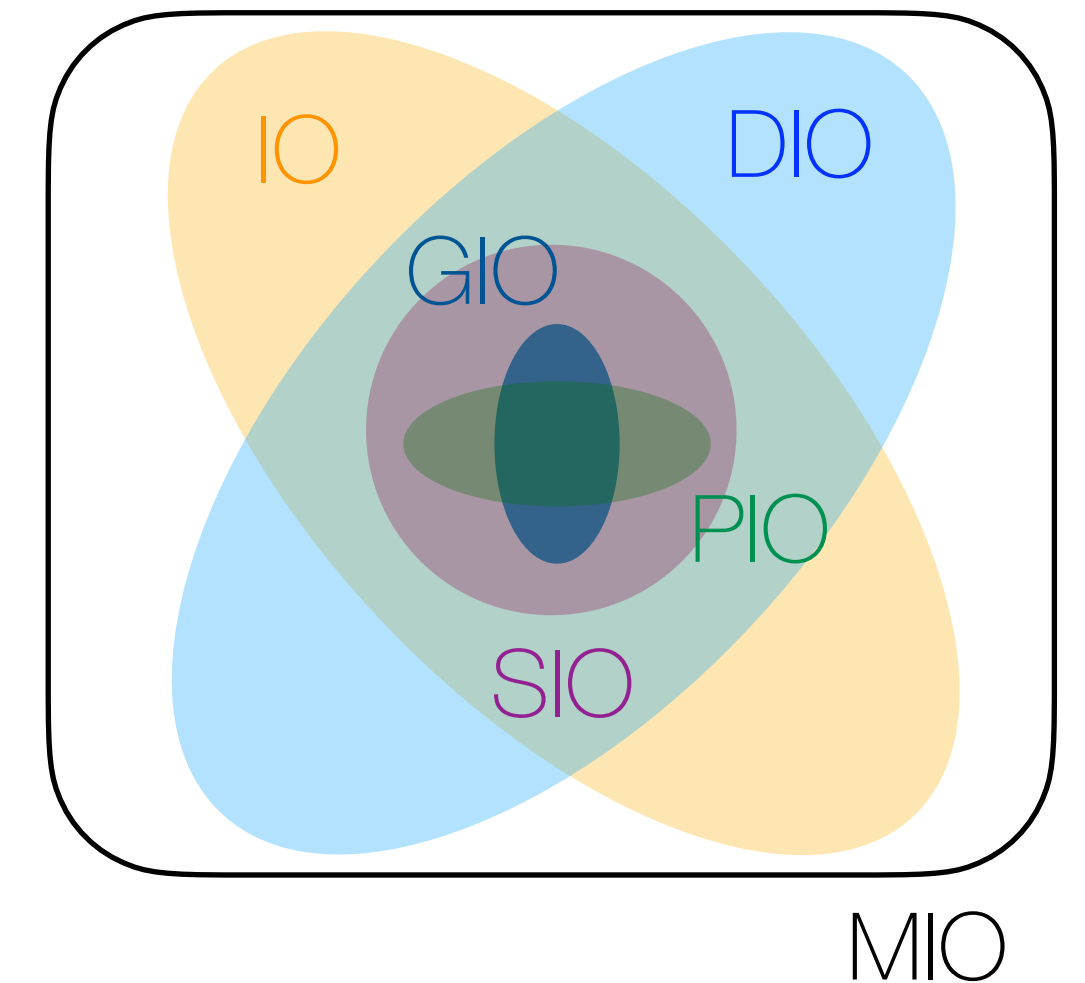
- Operational significance of entanglement from the perspective of quantum memory

- Application of resource theory of (speakeable) coherence

Started by a fundamental motivation: how to quantify superposition?

Operational significance unclear: zoo of free operations

Physical significance unclear: artificial preferred basis (c.f. asymmetry)

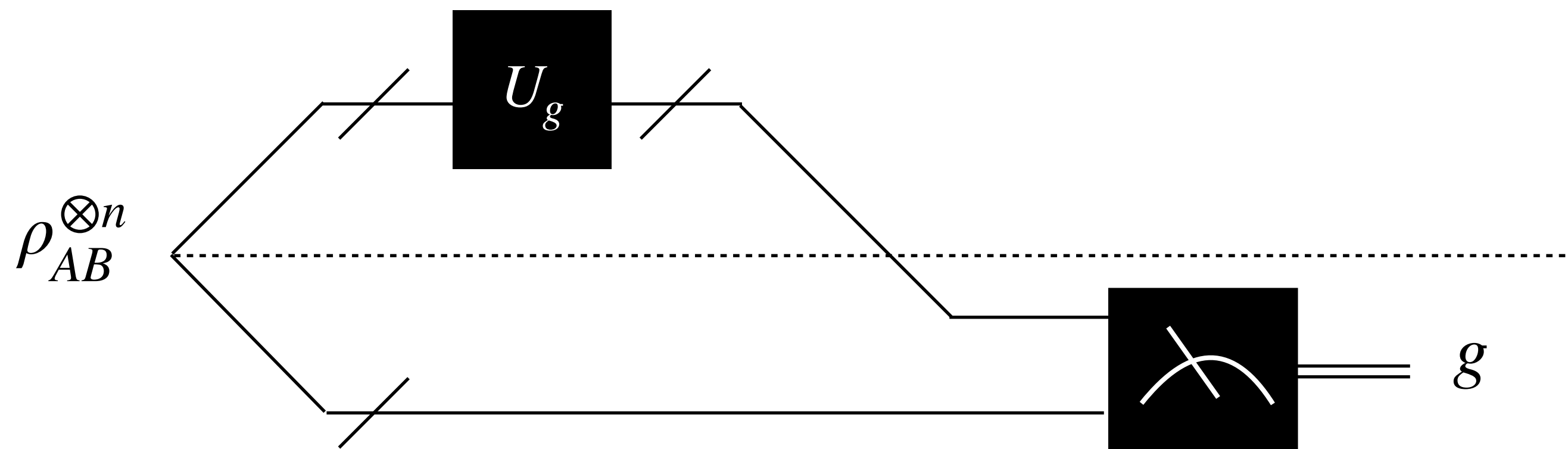


We find that the analysis of free operations (GIO!) is closely related to dense coding

Standard setting of dense coding

Classical message $k \in \mathcal{K}_n$ encoded in an group element $g \in G \times \cdots \times G = G^n$

Encode $g \in G^n$ onto a state $\rho_{AB}^{\otimes n}$ by applying a group representation $\{U_g\}_{g \in G^n}$



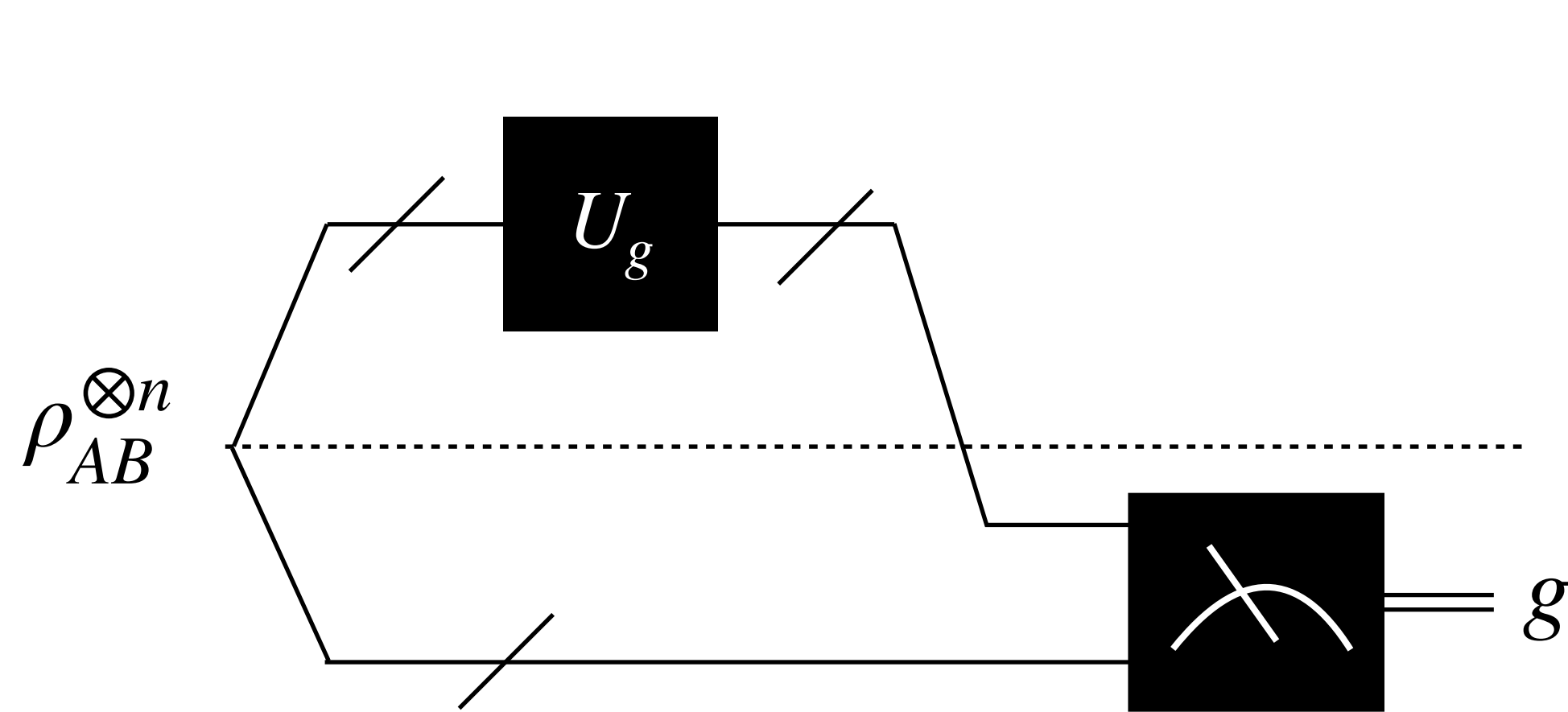
Capacity: rate of classical bits that can be reliably sent

$$C_c(\rho_{AB}) = \sup_p S \left(\sum_{g \in G} p_g U_g \rho_{AB} U_g^\dagger \right) - \sum_{g \in G} p_g S(U_g \rho_{AB} U_g^\dagger) = D(\rho_{AB} \| \mathcal{G}_A(\rho_{AB}))$$

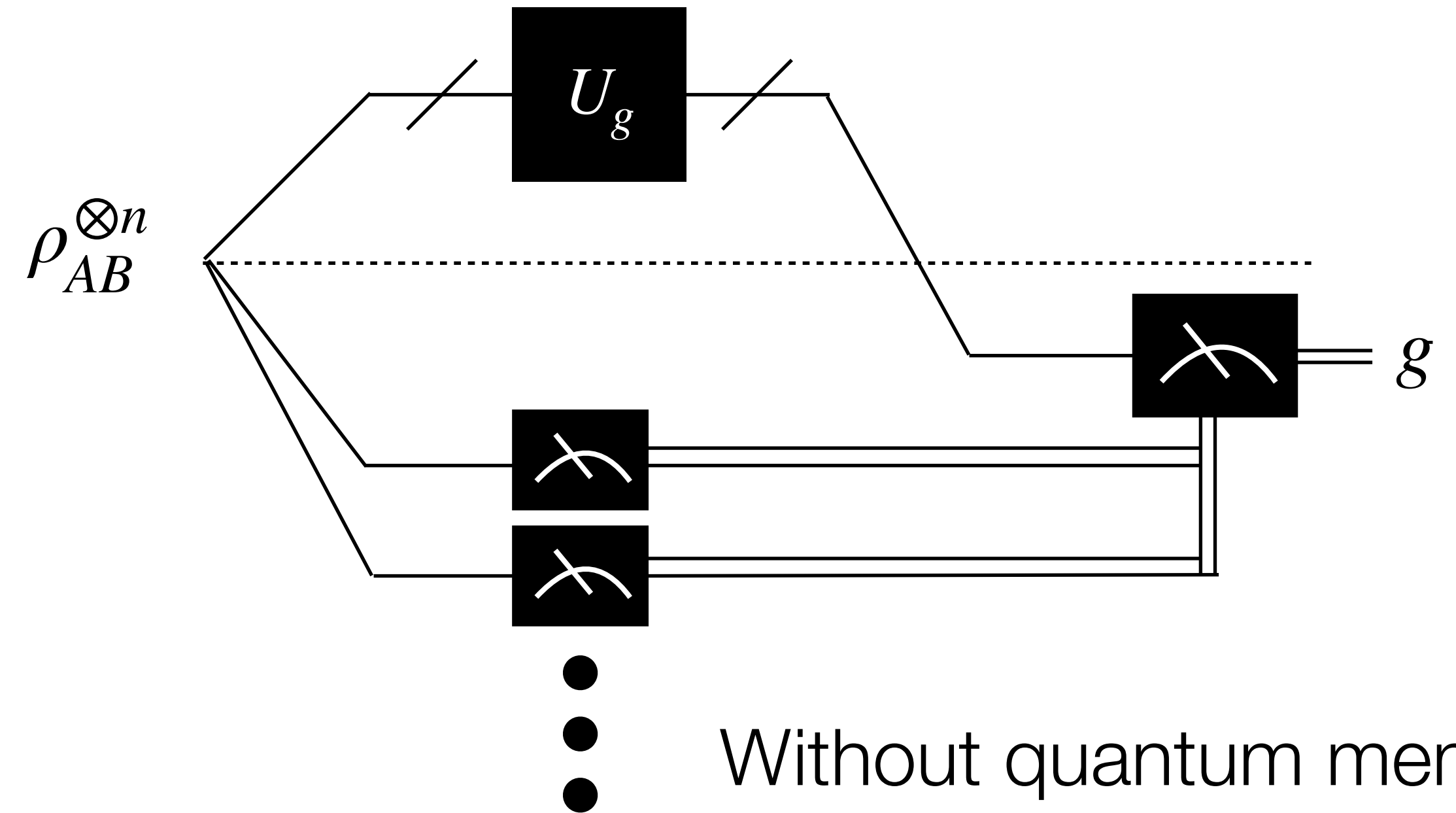
$$\mathcal{G}(\cdot) = \frac{1}{|G|} \sum_{g \in G} U_g \cdot U_g^\dagger$$

“Group twirling”

Dense coding without quantum memory



With full quantum memory



Without quantum memory

Bob measures his side immediately after ρ_{AB} is distributed.

c.f. [Hayashi, Wang, PRX Quantum '22] for the case of LOCC decoder

If Bob measures with basis $\mathcal{B} = \{ |k\rangle \}_k$, they share

$$\mathcal{B}_B(\rho_{AB}) = \sum_k \langle k | \rho_{AB} | k \rangle \otimes |k\rangle \langle k| = \sum_k p_k \rho_{AB|k}$$

$$p_k = \text{Tr}_A \langle k | \rho_{AB} | k \rangle$$

$$\rho_{AB|k} = \frac{1}{p_k} \langle k | \rho_{AB} | k \rangle \otimes |k\rangle \langle k|$$

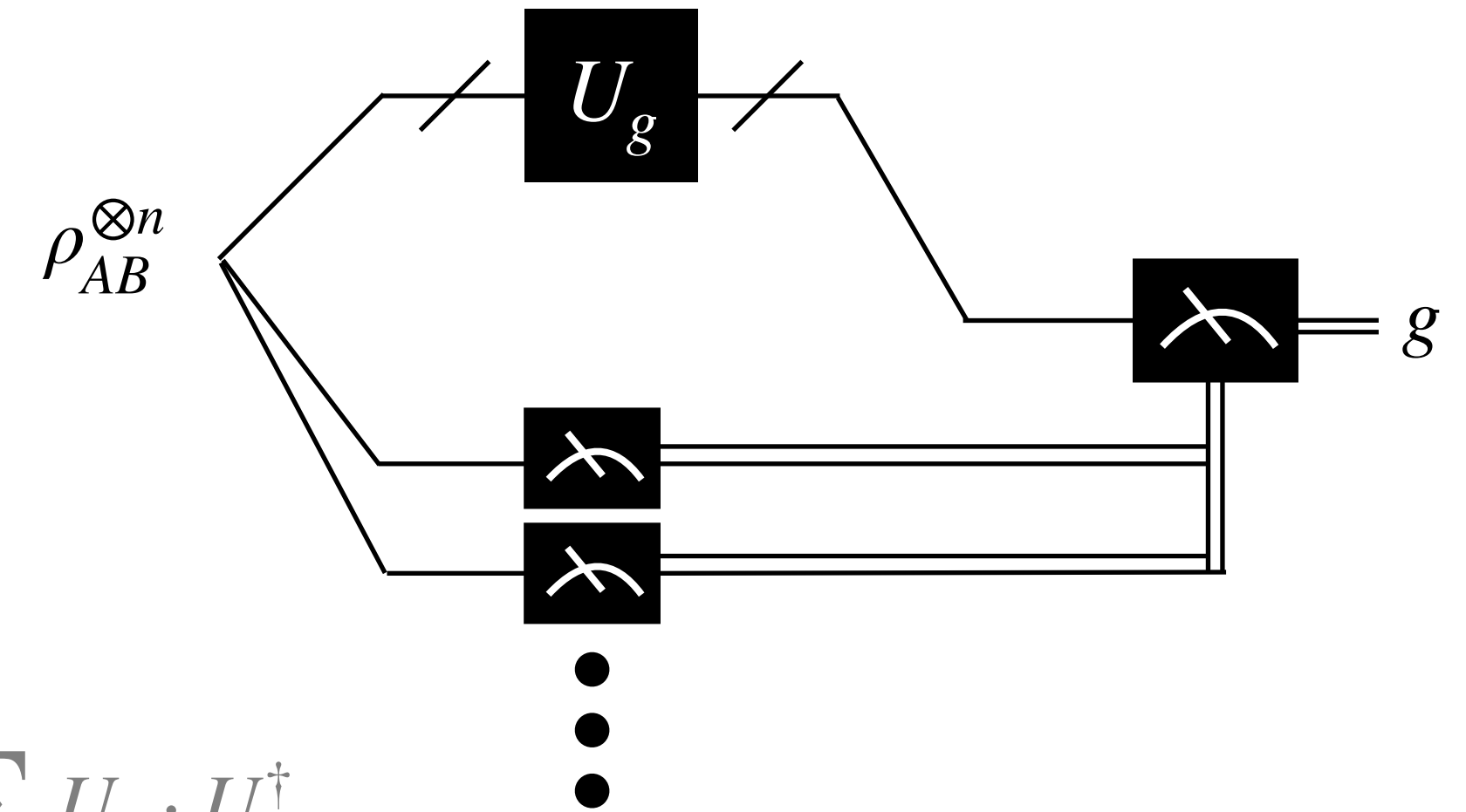
Capacity with and without quantum memory

$$\mathcal{B}_B(\rho_{AB}) = \sum_k \langle k | \rho_{AB} | k \rangle \otimes |k\rangle\langle k| = \sum_k p_k \rho_{AB|k}$$

Capacity with this shared state

$$\begin{aligned} C_c(\mathcal{B}_B(\rho_{AB})) &= D(\mathcal{B}_B(\rho_{AB}) \| \mathcal{G}_A \circ \mathcal{B}_B(\rho_{AB})) \\ &= D(\mathcal{B}_B(\rho_{AB}) \| \mathcal{B}_B \circ \mathcal{G}_A(\rho_{AB})) \end{aligned}$$

$$\mathcal{G}(\cdot) = \frac{1}{|G|} \sum_{g \in G} U_g \cdot U_g^\dagger$$



$$C_c(\rho_{AB}) = D(\rho_{AB} \| \mathcal{G}_A(\rho_{AB})) \geq D(\mathcal{B}_B(\rho_{AB}) \| \mathcal{B}_B \circ \mathcal{G}_A(\rho_{AB})) = C_c(\mathcal{B}_B(\rho_{AB}))$$

with memory

data processing

without memory

When the equality holds, Bob does not need to hold a quantum memory.

Q-memory-uselessness

Example: $G = \mathbb{Z}_l$, $U_g = \sum_i |g + i\rangle\langle i|$, $\rho_{AB} = |\Phi_{AB}\rangle\langle\Phi_{AB}|$, $|\Phi_{AB}\rangle = \frac{1}{\sqrt{l}} \sum_{i=1}^l |ii\rangle$

Capacity with memory $C_c(\rho_{AB}) = S(\mathcal{E}_A(\Phi_{AB})) = \log l$

This capacity can be achieved without quantum memory

$$U_g \otimes \mathbb{I} |\Phi_{AB}\rangle = \frac{1}{\sqrt{l}} \sum_i |i + g\rangle |i\rangle$$

- Bob measures with computational basis and get i
- Bob measures the received qubit with computational basis and get $i + g$
—————→ Bob can reconstruct g

If $C_c(\rho_{AB}) = C_c(\mathcal{B}_B(\rho_{AB}))$, we say that ρ_{AB} is **\mathcal{B} -q-memory useless**.

If ρ_{AB} is \mathcal{B} -q-memory useless for some basis \mathcal{B} , ρ_{AB} is **q-memory useless**.

Can we characterize q-memory uselessness?

Group representation

Consider the group G and projective unitary representation $\{U_g\}_{g \in G}$ acting on \mathcal{H}_A .

$$\mathcal{H}_A = \bigoplus_{\lambda \in \hat{G}} \mathcal{H}_\lambda \otimes \mathcal{M}_\lambda$$

\hat{G} : subset of irreducible representations

\mathcal{H}_λ : representation space \mathcal{M}_λ : multiplicity space

Let Π_λ be a projection onto $\mathcal{H}_\lambda \otimes \mathcal{M}_\lambda$, and $q_\lambda := \text{Tr}[(\Pi_\lambda \otimes \mathbb{I}_B) \rho_{AB}]$
 $\rho_\lambda := \frac{1}{q_\lambda} (\Pi_\lambda \otimes \mathbb{I}_B) \rho_{AB} (\Pi_\lambda \otimes \mathbb{I}_B)$

Then, we get $\mathcal{G}_A(\rho_{AB}) = \bigoplus_{\lambda \in \hat{G}} q_\lambda \frac{\mathbb{I}_\lambda}{d_\lambda} \otimes \text{Tr}_{\mathcal{H}_\lambda}(\rho_\lambda)$

$$C_c(\rho_{AB}) = H(q) + \sum_{\lambda \in \hat{G}} q_\lambda \left(\log d_\lambda + S(\text{Tr}_{\mathcal{H}_\lambda} \rho_\lambda) \right) - S(\rho_{AB})$$

$$C_c(\rho_{AB}) = D(\rho_{AB} \| \mathcal{G}_A(\rho_{AB}))$$

Characterization of q-memory-uselessness

We here focus on pure resource state $\rho_{AB} = |\psi_{AB}\rangle\langle\psi_{AB}|$

Suppose that unitary representation is multiplicity-free, i.e., $\mathcal{H}_A = \bigoplus_{\lambda \in \hat{G}} \mathcal{H}_\lambda$
e.g., cyclic group $G = \mathbb{Z}_l$ $U_g = \sum_{j=0}^{l-1} |j+g\rangle\langle j|$

Each irrep is one-dimensional and gains different phase under group action.

$|\psi_{AB}\rangle$ is q-memory-useless if and only if $|\psi_{AB}\rangle$ has the form

$$|\psi_{AB}\rangle = \sum_{\lambda \in \hat{G}} \sqrt{q_\lambda} |\psi_\lambda\rangle \otimes \sum_{k=1}^{d_B} \sqrt{p_k} e^{i\theta_{\lambda,k}} |k\rangle \quad |\psi_\lambda\rangle \in \mathcal{H}_\lambda$$

for some probability distributions $\{q_\lambda\}_\lambda$ and $\{p_k\}_k$.

Characterization of q-memory-uselessness

Equivalently, let us write the resource state as

$$|\psi_{AB}\rangle = \sum_{\lambda \in \hat{G}} \sqrt{q_\lambda} |\psi_\lambda\rangle |v_\lambda\rangle$$

$$|\psi_{AB}\rangle = \sum_{\lambda \in \hat{G}} \sqrt{q_\lambda} |\psi_\lambda\rangle \otimes \sum_{k=1}^{d_B} \sqrt{p_k} e^{i\theta_{\lambda,k}} |k\rangle$$

Then, we have the following characterization.

$|\psi_{AB}\rangle$ is q-memory-useless if and only if $|v_\lambda\rangle$ can be written in the form

$$|v_\lambda\rangle = \sum_{k=1}^{d_B} \sqrt{p_k} e^{i\theta_{\lambda,k}} |k\rangle \quad \lambda \in \hat{G}$$

for some basis $\{|k\rangle\}_k$ and phases $\{e^{i\theta_{\lambda,k}}\}_{\lambda,k}$ for all $\lambda \in \hat{G}$.

The problem is whether $|v_1\rangle, \dots, |v_l\rangle$ with $l := |\hat{G}|$ admits this form.

Characterization of q-memory-uselessness

$$|v_1\rangle, \dots, |v_l\rangle \text{ with } l := |\hat{G}| \text{ admits } |v_\lambda\rangle = \sum_{k=1}^{d_B} \sqrt{p_k} e^{i\theta_{\lambda,k}} |k\rangle \quad \bullet \bullet \bullet \star$$

When does \star happen? One case: when $\{|v_\lambda\rangle\}_\lambda$ are orthogonal. Take

$$|k\rangle = \frac{1}{\sqrt{l}} \sum_{s=1}^l e^{i\frac{2\pi ks}{l}} |v_s\rangle \quad k = 1, \dots, l$$

and $|l+1\rangle, \dots, |d_B\rangle$ to be orthogonal to $|v_1\rangle, \dots, |v_l\rangle$.

Another case: when $l = 2$. Choose $\{|1\rangle, |2\rangle\}$ satisfying

$$|v_1\rangle = \frac{1}{\sqrt{2}}(|1\rangle + |2\rangle) \quad |v_2\rangle = \frac{1}{\sqrt{2}}(e^{i(\theta'+\theta)}|1\rangle + e^{i(\theta'-\theta)}|2\rangle) \quad e^{i\theta'} \cos \theta = \langle v_1 | v_2 \rangle$$

Can we get a general characterization of when \star happens?

Characterization with Gram matrix

$$|v_1\rangle, \dots, |v_l\rangle \text{ with } l := |\hat{G}| \text{ admits } |v_\lambda\rangle = \sum_{k=1}^{d_B} \sqrt{p_k} e^{i\theta_{\lambda,k}} |k\rangle \quad \bullet \bullet \bullet \star$$

$V = (|v_1\rangle, \dots, |v_l\rangle)$ set of vectors $\longrightarrow J(V)_{ij} := \langle v_i | v_j \rangle$ Gram matrix of V

\star holds iff there exists a set $U = (|u_1\rangle, \dots, |u_l\rangle)$ of vectors such that

$$|u_\lambda\rangle = \sum_{k=1}^{d_B} \sqrt{p_k} e^{i\theta_{\lambda,k}} |\tilde{k}\rangle \quad \lambda = 1, \dots, l \text{ for the computational basis } \{|\tilde{k}\rangle\}_k, \text{ and}$$

$$J(V) = J(U) \quad \text{“Vectors with the same Gram matrix are connected by a unitary.”}$$

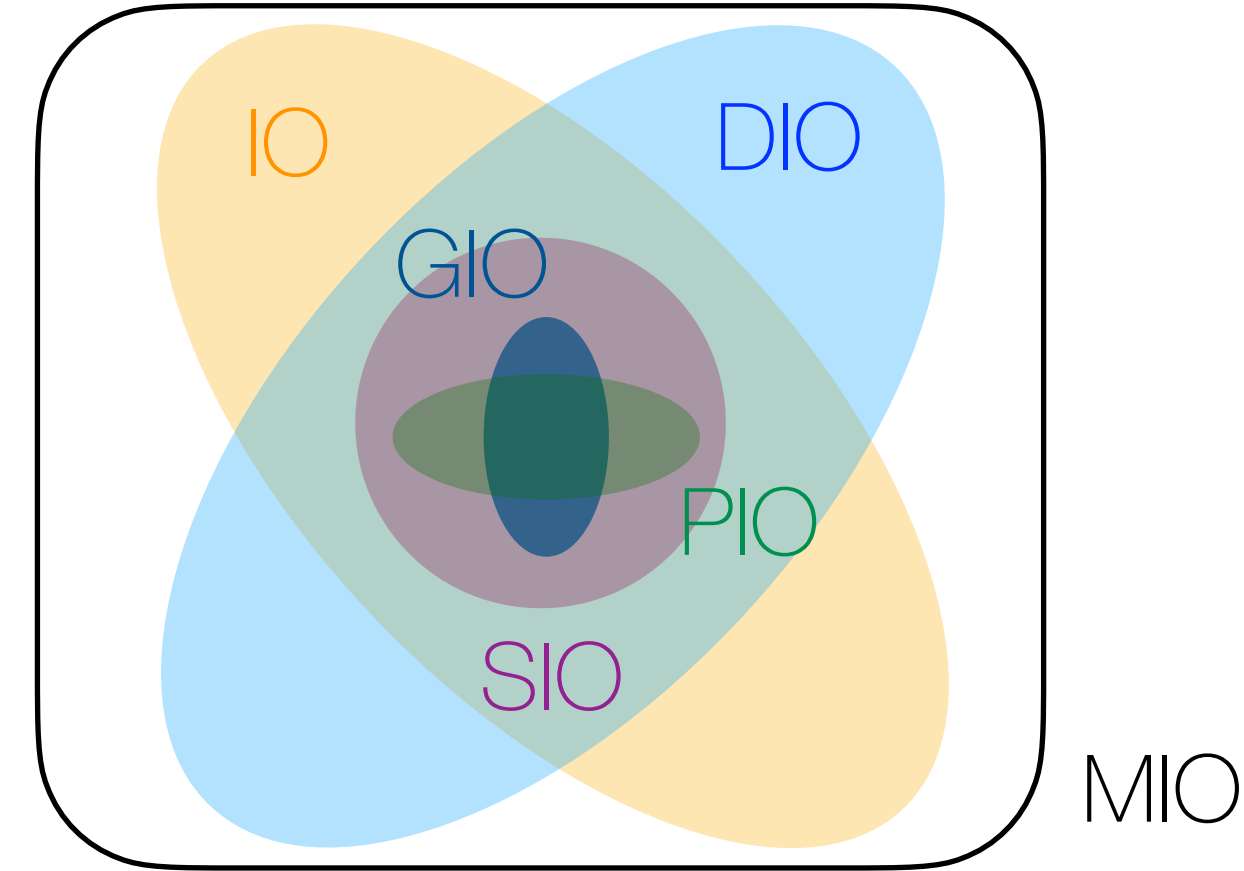
Observation $J(U) = l \sum_{k=1}^{d_B} p_k Z_k |+\rangle\langle+| Z_k^\dagger \quad |+\rangle = \frac{1}{\sqrt{l}} \sum_{\lambda=1}^l |\lambda\rangle \quad Z_k = \sum_{\lambda=1}^l e^{i\theta_{\lambda,k}} |\lambda\rangle\langle\lambda|$

$\star \iff |+\rangle$ can be transformed to the state $J(V)/l$ by a probabilistic incoherent unitary!

Genuinely Incoherent Operations (GIO)

$\mathcal{I} = \text{conv}\{ |i\rangle\langle i| \} : \text{incoherent states}$

Λ is GIO iff $\Lambda(\sigma) = \sigma, \quad \forall \sigma \in \mathcal{I}$



Characterization of GIO

Λ is GIO iff $\Lambda(\rho) = A \odot \rho$ for some positive semidefinite matrix A with $A_{ii} = 1 \quad \forall i$

$(X \odot Y)_{ij} = X_{ij}Y_{ij} : \text{Hadamard product}$

[de Vicente, Streltsov, J. Phys. A '17]

J is a Gram matrix if and only if $J_{ii} = 1 \quad \forall i$ and $J \geq 0$

→ Every Gram matrix J has one-to-one correspondence with GIO $\Lambda_J(\rho) = J \odot \rho$

Q-memory uselessness with GIO

$$|v_1\rangle, \dots, |v_l\rangle \text{ with } l := |\hat{G}| \text{ admits } |v_\lambda\rangle = \sum_{k=1}^{d_B} \sqrt{p_k} e^{i\theta_{\lambda,k}} |k\rangle \quad \bullet \bullet \bullet \star$$

$\star \iff |+\rangle$ can be transformed to the state $J(V)/l$ by a probabilistic incoherent unitary

Observe that $J(V)/l = J(V) \odot |+\rangle\langle +| = \Lambda_{J(V)}(|+\rangle\langle +|)$ always holds.

i.e., $|+\rangle\langle +|$ can be always transformed to the state $J(V)/l$ by GIO.

$\star \iff \text{GIO } \Lambda_{J(V)}$ can be implemented by a probabilistic incoherent unitary

When does this happen?

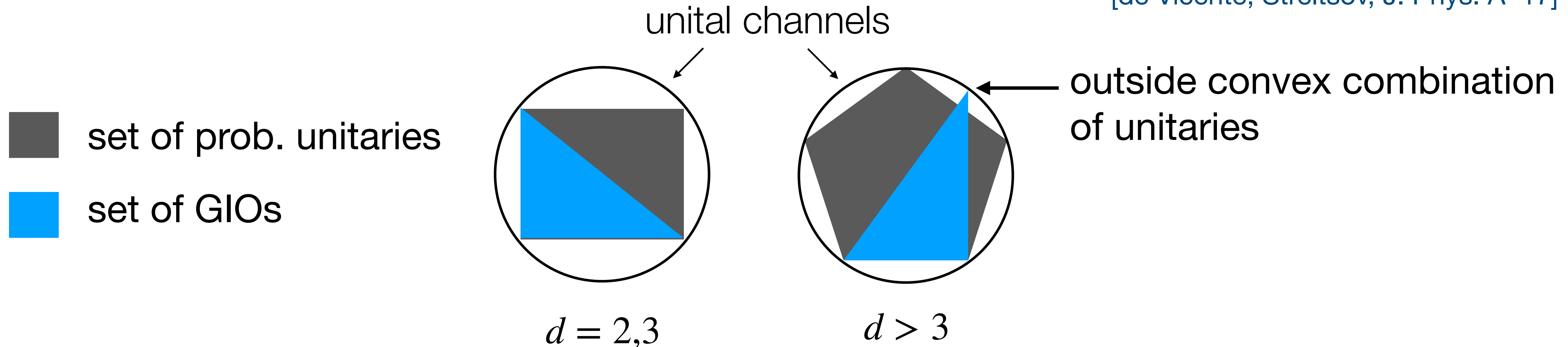
GIO implementable by probabilistic incoherent unitary

Let d be a dimension of the system that GIO operations act on.

When $d = 2, 3$, every $\Lambda \in \mathcal{O}_{\text{GIO}}$ can be implemented by a probabilistic application of incoherent unitaries.

When $d > 3$, there exists $\Lambda \in \mathcal{O}_{\text{GIO}}$ that cannot be implemented by a probabilistic applications of incoherent unitaries.

[de Vicente, Streltsov, J. Phys. A '17]



Application to dense coding

Combining these arguments lead to the following characterization.

If $l = |\hat{G}|$ takes $l = 2, 3$, $|\psi_{AB}\rangle$ is always q-memory-useless.

i.e., $C_c(|\psi_{AB}\rangle\langle\psi_{AB}|) = C_c(\mathcal{B}_B(|\psi_{AB}\rangle\langle\psi_{AB}|))$ for every basis \mathcal{B}_B on Bob.

If $l > 3$, there exists a state $|\psi_{AB}\rangle$ such that Bob's quantum memory is useful.

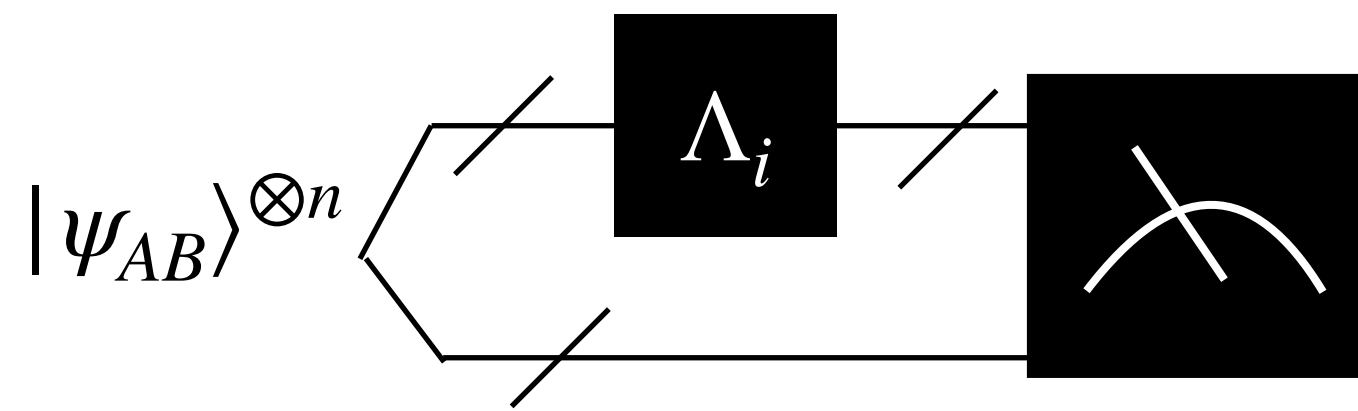
i.e., $C_c(|\psi_{AB}\rangle\langle\psi_{AB}|) > C_c(\mathcal{B}_B(|\psi_{AB}\rangle\langle\psi_{AB}|))$ for some basis \mathcal{B}_B on Bob.

q-memory-uselessness of

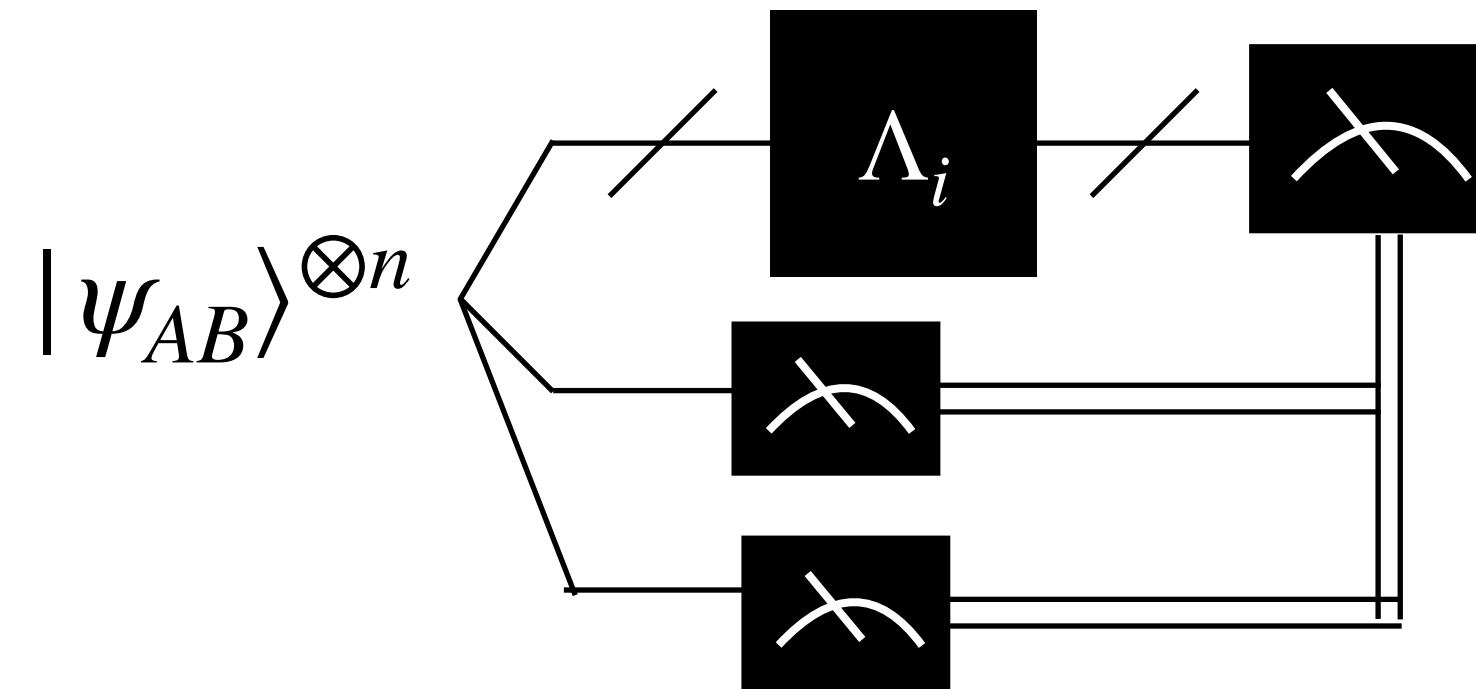
$$|\psi_{AB}\rangle = \sum_{\lambda \in \hat{G}} \sqrt{q_\lambda} |\psi_\lambda\rangle |v_\lambda\rangle \iff \text{Implementability of } \Lambda_{J(V)} \in \mathcal{O}_{\text{GIO}}$$

Implication to channel discrimination

Discriminate whether Λ_1 or Λ_2 applied.



with memory



without memory

e.g., quantum illumination

$$\Lambda_1 = (1 - p) \text{id} + p \mathcal{R}_\tau \quad \Lambda_2 = \mathcal{R}_\tau \quad \mathcal{R}_\tau: \text{preparation of thermal state } \tau$$

Here, we consider a variant of this, where $\Lambda_1 = \text{id}$ $\Lambda_2 = \mathcal{G}$ $\mathcal{G}(\cdot) = \frac{1}{|G|} \sum_{g \in G} U_g \cdot U_g^\dagger$

Error exponent of Type-II error in asymmetric channel discrimination

$$D(\psi_{AB} \| \mathcal{G}_A(\psi_{AB})) \geq D(\mathcal{B}_B(\psi_{AB}) \| \mathcal{G}_A \otimes \mathcal{B}_B(\psi_{AB}))$$

Same analysis can carry over!

with memory

without memory

Summary

- Introduced dense coding where receiver immediately measures out the shared entangled resource, representing the scenario that the receiver does not possess quantum memory.
- Characterized when the receiver's quantum memory becomes useful in terms of the given resource state and group representation.
- It turns out that the problem reduces to the analysis of GIO—whether the given GIO admits the implementation of probabilistic incoherent unitaries.
- Resource theory of coherence could find further applications as an analytical tool?

Thank you!