Professor Rio
EN.585.615.81.SP21 Mathematical Methods
Mid-term Exam
Johns Hopkins University
Student: Yves Greatti

# Question 1

a. Graph of the function attached as a separate pdf.

b. Since we have made the function $f(x)$ even using an even extension, all the $b_k$ coefficients in its Fourier series are zero. With a period $L = 4$, we determine the remaining coefficients $a_k$:

$$a_k = \frac{2}{4} \int_{-2}^{2} x \cos{(\frac{2k\pi x}{4})} dx$$

And since $f$ is even now

$$\begin{aligned} a_k &= \frac{4}{4} \int_{0}^{2} x \cos{(\frac{2k\pi x}{4})} dx \\ &= \int_{0}^{2} x \cos{(\frac{k\pi x}{2})} dx \end{aligned}$$

Using integration by parts, for $k > 0$:

$$\begin{aligned} a_k &= \frac{2}{k\pi}[x \sin(\frac{k\pi x}{2})]_0^2 - \frac{2}{k\pi} \int_0^2 \sin(\frac{k\pi x}{2}) dx \\ &= 0 - \frac{2}{k\pi}(-\frac{2}{k\pi})[\cos\frac{k\pi x}{2})]_0^2 \\ &= \frac{4}{(k\pi)^2)}[\cos(k\pi) - \cos(0)] \\ &= \frac{4}{(k\pi)^2}[(-1)^k - 1] \end{aligned}$$

Then

$$a_k = \begin{cases} -\dfrac{8}{(k\pi)^2} \text{ for odd } k \\ 0 \text{ for even } k \end{cases}$$

And $a_0 = \frac{2}{4} \int_{-2}^{2} x dx = \frac{4}{4} \int_0^2 x dx = \frac{1}{2}[x^2]_0^2 = 2$. With the coefficients $a_k$ determined, we obtain the Fourier series for $f(x)$:

$$\begin{aligned} f(x) &= \frac{2}{2} - \sum_{k=1}^{\infty} \frac{8}{(k\pi)^2} \cos(\frac{2k\pi x}{4}) \; k \text{ odd} \\ x &= 1 - \frac{8}{\pi^2} \sum_{k=0}^{\infty} \frac{1}{(2k+1)^2} \cos(\frac{(2k+1)\pi x}{2}) \end{aligned}$$

c. Applying Parseval's identity for Fourier series and using the result of part b.:

$$
\begin{aligned}
\frac{1}{4}\int_{-2}^{2} x^2 dx &= \frac{2^2}{4} + \frac{1}{2}\sum_{k=1}^{\infty}(a_k^2 + 0)\ k \text{ odd} \\
\frac{2}{4}\int_{0}^{2} x^2 dx &= 1 + \frac{1}{2}\sum_{k=0}^{\infty}\Big(\frac{8}{(2k+1)^2\pi^2}\Big)^2 \\
\frac{1}{2}[\frac{x^3}{3}]_0^2 &= 1 + \frac{1}{2}\cdot\frac{64}{\pi^4}\sum_{k=0}^{\infty}\frac{1}{(2k+1)^4} \\
\frac{4}{3} - 1 &= \frac{32}{\pi^4}\sum_{k=0}^{\infty}\frac{1}{(2k+1)^4} \\
\sum_{k=0}^{\infty}\frac{1}{(2k+1)^4} &= \frac{\pi^4}{32}\cdot\frac{1}{3}
\end{aligned}
$$

Therefore

$$\sum_{n=0}^{\infty}\frac{1}{(2n+1)^4} = \frac{\pi^4}{96}$$

# Question 2

a. Graph of the function attached as a separate pdf.

b.

$$f(t) = A\Big[H(t) - H(t-\tau)\Big]$$

c.

$$\tilde{f}(w) = F\{f(t)\} = \frac{1}{\sqrt{2\pi}}\int_{-\infty}^{\infty} f(t)e^{-iwt}\,dt$$

Since $f(t) = 0$ for $t \geq 0$ or $t \leq \tau$:

$$
\begin{aligned}
\tilde{f}(w) &= \frac{1}{\sqrt{2\pi}}\int_0^{\tau} A\cdot e^{-iwt}\,dt \\
&= \frac{A}{\sqrt{2\pi}}\Big(\frac{1}{-iw}\Big)[e^{-iwt}]_0^{\tau} \\
&= \frac{iA}{w\sqrt{2\pi}}(e^{-iw\tau} - 1) \\
&= \frac{iA}{w\sqrt{2\pi}}e^{-iw\frac{\tau}{2}}(e^{-iw\frac{\tau}{2}} - e^{iw\frac{\tau}{2}})
\end{aligned}
$$

From Euler identity:

$$e^{-iw\frac{\tau}{2}} - e^{iw\frac{\tau}{2}} = -2i\sin w\frac{\tau}{2}$$

Therefore

$$\begin{aligned}\tilde{f}(w) &= \frac{2A}{w\sqrt{2\pi}} e^{-iw\frac{\tau}{2}} \sin w\frac{\tau}{2} \\ &= \sqrt{\frac{2}{\pi}}\frac{A}{w} e^{-iw\frac{\tau}{2}} \sin w\frac{\tau}{2} \\ &= A\sqrt{\frac{2}{\pi}} e^{-iw\frac{\tau}{2}} \frac{\tau}{2}\frac{\sin(w\frac{\tau}{2})}{w\frac{\tau}{2}} \\ &= \frac{A}{\sqrt{2\pi}}\tau e^{-iw\frac{\tau}{2}} \text{sinc}(w\frac{\tau}{2})\end{aligned}$$

d. Let $A = \frac{1}{\tau}$ then substituting in $f(t)$ from part c., gives:

$$F\{\lim_{\tau\to 0} f(t)\} = \lim_{\tau\to 0} F\{f(t)\} = \lim_{\tau\to 0} \frac{1}{\sqrt{2\pi}} e^{-iw\frac{\tau}{2}} \frac{\sin(w\frac{\tau}{2})}{w\frac{\tau}{2}}$$

$$\lim_{\theta\to 0} \frac{\sin(\theta)}{\theta} = 1 \text{ by Hospitals rule}$$
$$\lim_{\tau\to 0} e^{-iw\frac{\tau}{2}} = \lim_{\tau\to 0} e^0 = 1$$

Therefore

$$F\{\lim_{\tau\to 0} f(t)\} = \frac{1}{\sqrt{2\pi}}$$

e. The Fourier transform of $f(t)$ as $\tau \to 0$ is the Fourier transform of a $\delta$-function as we can expect as we ”transform” the rectangular function $f(t)$ to a Dirac impulse.

# Question 7

$$x^2\frac{d^2y}{dx^2} + x\frac{dy}{d} - y = x\,, y(e) = 0\,, y'(e) = 2$$

a. This is Euler differential equation, and we make the change of variable $x = e^t$ or $t = \ln(x)$. Then

$$\begin{aligned}\frac{dy}{dx} &= \frac{dy}{dt}\frac{dt}{dx} = \frac{dy}{dt}\frac{d\ln x}{dx} = \frac{dy}{dt}\frac{1}{x} = \frac{1}{x}\frac{dy}{dt} \\ x\frac{dy}{dx} &= \frac{dy}{dt}\end{aligned}$$

And since this is a Legendre ODE with $\alpha = 1$ and $\beta = 0$, we can use the expression for the second derivative $(\alpha x + \beta)^2\frac{d^2y}{dx^2} = \alpha^2\frac{d}{dt}[\frac{d}{dt} - 1]y$. With $\alpha = 1$ and $\beta = 0$, we have: $\frac{d^2y}{dx^2} = \frac{d^2y}{t^2} - \frac{dy}{dt}$.

Substitute into the above equation yields:

$$\begin{aligned}(\frac{d^2y}{dt^2} - \frac{dy}{dt}) + \frac{dy}{dt} - y &= e^t \\ \frac{d^2y}{dt^2} - y &= e^t\end{aligned}$$

b. The homogeneous equation is

$$\frac{d^2y}{dt^2} - y = 0$$

Assume a solution of the form $y(t) = Ae^{\lambda t}$ gives the characteristic equation $\lambda^2 - 1 = 0$ which has for roots $\lambda = \pm 1$ and gives for solution $y(t) = c_1 e^t + c_2 e^{-t}$.

c. The ODE to solve is:

$$\frac{d^2y}{dt^2} - y = 0$$

It is in standard form and it is defined at any point $t$, it is analytic, thus we take as solution $y(t) = \sum_{t=0}^{\infty} a_n t^n$. So:

$$y'(t) = \sum_{t=0}^{\infty} n a_n t^{n-1}$$

$$y''(t) = \sum_{t=0}^{\infty} n(n-1) a_n t^{n-2}$$

by reindexing

$$y''(t) = \sum_{t=-2}^{\infty} (n+2)(n+1) a_{n+2} t^n$$

$$y''(t) = \sum_{t=0}^{\infty} (n+2)(n+1) a_{n+2} t^n$$

Substitute into the ODE gives:

$$\sum_{t=0}^{\infty} (n+2)(n+1) a_{n+2} t^n - \sum_{t=0}^{\infty} a_n t^n = 0$$

$$\sum_{t=0}^{\infty} [(n+2)(n+1) a_{n+2} - a_n] t^n = 0$$

or

$$a_{n+2} = \frac{1}{(n+2)(n+1)} a_n$$

$$a_n = \frac{1}{n(n-1)} a_{n-2}$$

Take $a_0 = a_1 = 1$ and we generate the coefficients:

. $n = 2$ then $a_2 = \frac{1}{2 \cdot 1} a_0 = \frac{1}{2 \cdot 1} = \frac{1}{2!}$

. $n = 3$ then $a_3 = \frac{1}{3 \cdot 2} a_1 = \frac{1}{3 \cdot 2} = \frac{1}{3!}$

. $n = 4$ then $a_4 = \frac{1}{4 \cdot 3} a_2 = \frac{1}{4 \cdot 3 \cdot 2 \cdot 1} = \frac{1}{4!}$

$\vdots$

. $a_n = \frac{1}{n(n-1)} a_{n-2} = \cdots = \frac{1}{n!}$

The first solution we obtain is: $y_1(t) = \sum_{t=0}^{\infty} a_n t^n = \sum_{t=0}^{\infty} \frac{t^n}{n!} = e^t$. Secondly, if we set $a_0 = 1$ and choose $a_1 = -1$, then we obtain a second independent solution:

. $n = 2$ then $a_2 = \frac{1}{2\cdot 1} a_0 = \frac{1}{2\cdot 1} = \frac{1}{2!}$
. $n = 3$ then $a_3 = \frac{1}{3\cdot 2} a_1 = -\frac{1}{3\cdot 2} = \frac{-1}{3!}$
. $n = 4$ then $a_4 = \frac{1}{4\cdot 3} a_2 = \frac{1}{4\cdot 3\cdot 2\cdot 1} = \frac{1}{4!}$
. $n = 5$ then $a_5 = \frac{1}{5\cdot 4} a_3 = \frac{-1}{5\cdot 4\cdot 3\cdot 2\cdot 1} = \frac{-1}{5!}$

$\vdots$

. $a_n = \frac{1}{n(n-1)} a_{n-2} = \cdots = \frac{(-1)^n}{n!}$

We have the second solution: $y_2(t) = \sum_{t=0}^{\infty} a_n t^n = \sum_{t=0}^{\infty} \frac{(-t)^n}{n!}$, recognizing the last series as $e^{-t}$, we can write the general solution of the homogeneous equation as

$$y_H(t) = c_1 e^t + c_2 e^{-t}$$

which is the solution we found in question b.

d. The differential equation to solve is

$$\frac{d^2 y}{dt^2} - y = e^t$$

Next we use the variation of parameters method, we are looking for a solution $y_p(t) = k_1(t)e^t + k_2(t)e^{-t}$. We solve for derivatives of k's a system of two equations:

$$\begin{cases} k_1' e^t + k_2' e^{-t} & = 0 \\ k_1' e^t - k_2' e^{-t} & = e^t \end{cases}$$

Multiplying through by $e^t$ gives:

$$\begin{cases} k_1' e^{2t} + k_2' & = 0 \\ k_1' e^{2t} - k_2' & = e^{2t} \end{cases}$$

Adding first equation to second yields $2k_1' e^{2t} = e^{2t}$ or $k_1' = \frac{1}{2}$ and $k_1 = \frac{t}{2}$. Substitute

$$\begin{aligned} k_2' &= -k_1' e^{2t} \\ &= -\frac{1}{2} e^{2t} \end{aligned}$$

integrating

$$k_2 = -\frac{e^{2t}}{4}$$

Therefore:

$$\begin{aligned} y_p(t) &= k_1(t)e^t + k_2(t)e^{-t} \\ &= \frac{t}{2} e^t - \frac{e^{2t}}{4} e^{-t} \\ &= \frac{t}{2} e^t - \frac{e^t}{4} \\ &= \frac{e^t}{2}(t - \frac{1}{2}) \end{aligned}$$

e. The general solution is: $y(t) = y_H(t) + y_p(t) = c_1 e^t + c_2 e^{-t} + \frac{e^t}{2}(t - \frac{1}{2})$, simplifying the constants, we can rewrite the general solution as $y(t) = c_1 e^t + c_2 e^{-t} + \frac{t}{2} e^t$. Plugging back $x = e^t$ or $t = \ln(x)$ gives

$$y(x) = c_1 x + \frac{c_2}{x} + \frac{x \ln x}{2}$$

f. The total solution is

$$y(x) = c_1 x + \frac{c_2}{x} + \frac{x \ln x}{2}$$
$$y'(x) = c_1 x - \frac{c_2}{x^2} + \frac{1}{2}(1 + \ln x)$$

And the initial conditions are $y(e) = 0\,, y'(e) = 2$, plugging back these into the previous equations gives

$$\begin{cases} y(e) = c_1 e + \frac{c_2}{e} + \frac{e \ln e}{2} = 0 \\ y'(e) = c_1 - \frac{c_2}{e^2} + \frac{1}{2}(1 + \ln e) = 2 \end{cases}$$

$$\Rightarrow \begin{cases} c_1 e + c_2 e^{-1} = -\frac{e}{2} \\ c_1 - c_2 e^{-2} = 1 \end{cases}$$

$$\Rightarrow \begin{cases} c_1 e^2 + c_2 = -\frac{e^2}{2} \\ c_1 - c_2 e^{-2} = 1 \end{cases}$$

Adding equation (1) to equation (2) leads to $2c_1 = e^2 - \frac{e^2}{2} = \frac{e^2}{2}$, $c_1 = \frac{1}{4}$,$c_2 = e^2(c_1 - 1) = \frac{3}{4}e^2$. Reporting these constants into the expression of the total solution gives:

$$y(x) = \frac{1}{4}x - \frac{3}{4}e^2 \frac{1}{x} + \frac{x \ln x}{2}$$

$$y(x) = \frac{x^2 + 2x^2 \ln(x) - 3e^2}{4x}$$