

PHYS414 Final Project Report

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Newton Part

Part(A)

$$\frac{dm(r)}{dr} = 4\pi r^2 \rho(r)$$

$$\frac{dP(r)}{dr} = - \frac{Gm(r)\rho(r)}{r^2}$$

$$P = K \rho^{1+\frac{1}{n}}$$

$$\Rightarrow \frac{dP(r)}{dr} \cdot \left(- \frac{r^2}{G\rho(r)} \right) = m(r) \Rightarrow \frac{d}{dr} \left(\frac{dP(r)}{dr} \cdot \left(- \frac{r^2}{G\rho(r)} \right) \right) = \frac{dm(r)}{dr}$$

by plugging in the given relation for the first derivative of $m(r)$:

$$\Rightarrow \frac{d}{dr} \left(\frac{dP(r)}{dr} \cdot \left(- \frac{r^2}{G\rho(r)} \right) \right) = 4\pi r^2 \rho(r)$$

Using the relation:

$$\Rightarrow \frac{dP}{dr} = \frac{d}{dr} (K \rho^{1+\frac{1}{n}}) = \left(1 + \frac{1}{n}\right) K \rho^{\frac{1}{n}} \frac{d\rho}{dr} \quad \text{by substituting}$$

$$\Rightarrow \frac{d}{dr} \left(\left(1 + \frac{1}{n}\right) K \rho^{\frac{1}{n}} \frac{d\rho}{dr} \cdot \left(- \frac{r^2}{G\rho(r)} \right) \right) = 4\pi r^2 \rho(r)$$

$$\Rightarrow 4\pi r^2 \rho(r) + \frac{d}{dr} \left(\left(1 + \frac{1}{n}\right) \frac{K r^2}{G} \cdot \rho^{\frac{1}{n}-1} \frac{d\rho}{dr} \right) = 0$$

$$\Rightarrow \frac{1}{4\pi r^2} \frac{d}{dr} \left(\left(1 + \frac{1}{n}\right) \frac{K r^2}{G} \rho^{\frac{1}{n}-1} \frac{d\rho}{dr} \right) + \rho(r) = 0$$

$$\text{let } r = \alpha \xi \quad \text{Where } \alpha' = \frac{(n+1)K}{4\pi G} \rho_c^{\frac{1}{n}-1}$$

$$\text{let } \Theta = \rho_c \rho^{\frac{1}{n}}$$

$$\Rightarrow \frac{dp}{dr} = \frac{dp}{d\theta} \frac{d\theta}{d\xi} \cdot \frac{d\xi}{dr}$$

$$\rho = \rho_c \theta^n \Rightarrow \frac{dp}{d\theta} = \rho_c n \theta^{n-1}; \quad \xi = \frac{r}{\alpha} \Rightarrow \frac{d\xi}{d\theta} = \frac{1}{\alpha}$$

$$\Rightarrow \frac{dp}{dr} = \rho_c n \theta^{n-1} \cdot \frac{1}{\alpha} \cdot \frac{d\theta}{d\xi} \quad \text{by substituting these relations to the equation we obtain}$$

$$\Rightarrow \frac{1}{\alpha^2 \xi^2} \cdot \frac{d}{d\xi} \left(\left(1 + \frac{1}{n}\right) \frac{K r^2}{4\pi G} \left(\rho_c \theta^n\right)^{\frac{1}{n}-1} \cdot \cancel{\rho_c n \theta^{n-1}} \cdot \frac{1}{\alpha} \frac{d\theta}{d\xi} + \cancel{\rho_c \theta^n} \right) = 0$$

$$\Rightarrow \frac{1}{\alpha^2 \xi^2} \cdot \frac{d}{d\xi} \left(\frac{(n+1)K}{4\pi G} \cdot \underbrace{\rho_c^{\frac{1}{n}-1}}_{\alpha^2} \cdot \frac{r^2}{\alpha} \frac{d\theta}{d\xi} \right) + \theta^n = 0$$

$$\Rightarrow \frac{1}{\xi^2} \cdot \frac{d}{d\xi} \left(\underbrace{\frac{\alpha^2 r^2}{\alpha^4}}_{\xi^2} \cdot \frac{d\theta}{d\xi} \right) + \theta^n = 0 \Rightarrow \boxed{\frac{1}{\xi^2} \cdot \frac{d}{d\xi} \left(\xi^2 \frac{d\theta}{d\xi} \right) + \theta^n = 0}$$

In order to find the mass of the star we need to integrate the mass density over its volume. We know that the surface of the star is at ξ . We need to find

Mathematica Part

Newton Part A

```
In[ ]:= DSolve[{((1/\xi)^2)*D[\xi^2*\theta'[\xi],\xi]+ \theta[\xi]^n == 0}]
```

```
DSolve[{ \theta[\xi]^n + \frac{2 \xi \theta'[\xi] + \xi^2 \theta''[\xi]}{\xi^2} == 0}] :
```

```
In[ ]:= AsymptoticDSolveValue[{((1/\xi)^2)*D[\xi^2*\theta'[\xi],\xi]+ \theta[\xi]^n == 0, \theta[0] == 1} \theta[\xi], {\xi, 0, 6}, n]
```

```
AsymptoticDSolveValue[{ \theta[\xi]^n + \frac{2 \xi \theta'[\xi] + \xi^2 \theta''[\xi]}{\xi^2} == 0, \theta[0] == 1}, \theta[\xi], {\xi, 0, 6}]
```

```
In[ ]:= AsymptoticDSolveValue[{ \xi*\theta[\xi]^n + 2 \theta'[\xi] + \xi \theta''[\xi] == 0, \theta[0] == 1}, \theta[\xi], {\xi, 0, 4}]
```

```
Out[ ]:= 1 - \frac{\xi^2}{6} + \frac{n \xi^4}{120}
```

```
In[ ]:= 1 - \frac{\xi^2}{6} + \frac{n \xi^4}{120}
```

```
DSolve[{ \xi*\theta[\xi]^n + 2 \theta'[\xi] + \xi \theta''[\xi] == 0, \theta'[0] == 0, \theta[0] == 1}, \theta[\xi], {\xi, 0, \infty}]
```

```
1 - \frac{\xi^2}{6} + \frac{n \xi^4}{120}
```

In[]:= FullSimplify[DSolve[{ $\xi * \theta[\xi]^4 + 2 * \theta'[\xi] + \xi \theta''[\xi] = 0$, $\theta'[0] = 0$, $\theta[0] = 1$ }, $\theta[\xi]$,
{ $\xi, 0, \infty$ }]]

... Solve: Inconsistent or redundant transcendental equation. After reduction, the bad equation is $i(2 i c_1 + c_2) == 0$.

... Solve: Inconsistent or redundant transcendental equation. After reduction, the bad equation is $2 i c_1 + c_2 == 0$.

... Solve: Inverse functions are being used by Solve, so some solutions may not be found; use Reduce for complete solution information.

$$\left\{ \left\{ \theta[\xi] \rightarrow \frac{\sin[\xi]}{\xi} \right\} \right\}$$

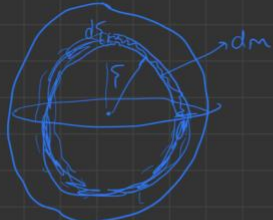
In order to find the mass of the star we need to integrate the mass density over its volume. We know that the surface of the star is at ξ . We need to find

$$M_S = \int_{V_S} dm \quad \text{and relate } dm \text{ to } d\xi$$

Given relations:

$$\frac{dm(r)}{dr} = 4\pi r^2 \rho(r) \quad \rho = \rho_c \theta^n \quad \& \quad r = \alpha \xi$$

$$\frac{dm}{d(\alpha \xi)} = 4\pi (\alpha \xi)^2 \rho_c \theta^n \Rightarrow dm = 4\pi \alpha^3 \xi^2 \rho_c \theta^n d\xi$$



$$\Rightarrow M_S = \int_{R_S} dm = \int_0^{\xi} 4\pi \alpha^3 \xi'^2 \rho_c \theta^n(\xi') d\xi'$$

$$= 4\pi \alpha^3 \rho_c \int_0^{\xi} \xi'^2 \theta^n(\xi') d\xi'$$

from the equation derived above:

$$\frac{1}{\xi^2} \frac{d}{d\xi} \left(\xi^2 \frac{d\theta}{d\xi} \right) + \theta^n = 0 \Rightarrow \underline{\xi^2 \theta^n} = - \frac{d}{d\xi} \left(\xi^2 \frac{d\theta}{d\xi} \right)$$

by substituting:

$$\Rightarrow M_S = 4\pi \alpha^3 \rho_c \int_0^{\xi} - \frac{d}{d\xi} \left(\xi^2 \frac{d\theta}{d\xi} \right) = -4\pi \alpha^3 \rho_c \left[\xi^2 \frac{d\theta}{d\xi} \right]_{\xi_n} \quad \xi_n = \frac{R}{\alpha}$$

$\xi \text{ at the surface}$

$$\Rightarrow M_S = -4\pi \rho_c \alpha^3 \cdot \xi_n^3 \cdot \frac{1}{\xi_n} \theta'(\xi_n) = -4\pi \rho_c \alpha^3 \cdot \frac{R^3}{\alpha^3} \cdot \left(- \frac{\theta'(\xi_n)}{\xi_n} \right)$$

$$\Rightarrow M_S = 4\pi \rho_c R^3 \left(- \frac{\theta'(\xi_n)}{\xi_n} \right)$$

Using the relation, we can write p_c as:

$$\alpha^1 = \frac{(n+1)K}{2\pi G} p_c^{\frac{1}{n}-1} \Rightarrow p_c = \left(\frac{4\pi G}{(n+1)K} \cdot \alpha^2 \right)^{\frac{n}{n-1}} \Rightarrow p_c = \left(\frac{4\pi G}{(n+1)K} \right)^{\frac{n}{n-1}} \cdot \alpha^{\frac{2n}{n-1}}$$

by substituting into the equation and using the definition of ξ_n at R :

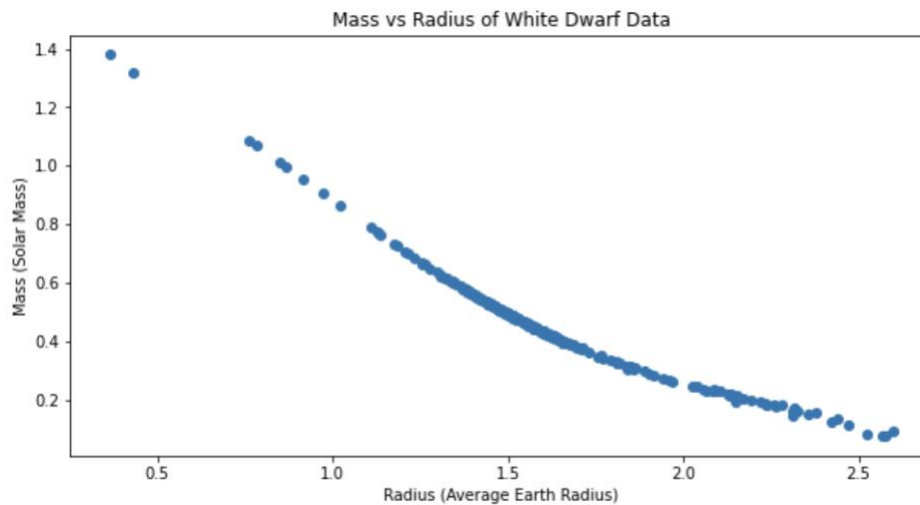
$$M_S = 4\pi \left(\frac{4\pi G}{(n+1)K} \right)^{\frac{n}{n-1}} \left(\frac{\xi_n}{R} \right)^{\frac{2n}{n-1}} R^3 \left(-\frac{\theta'(\xi_n)}{\xi_n} \right)$$

$$\Rightarrow M_S \propto R^{\frac{3n-3-2n}{n-1}} = R^{\frac{n-3}{n-1}}$$

$$\Rightarrow M_S \propto R^{\frac{3-n}{1-n}}$$

55 of 159

Part(B)



Part(C)

```
In[ ]:= Assuming[{x -> 0}, Series[C*(x*(2*x^2 - 3)*(x^2 + 1)^(1/2) + 3*ArcSinh[x]), {x, 0, 10}]]
```

$$\text{Out[]} = \frac{8 C x^5}{5} - \frac{4 C x^7}{7} + \frac{C x^9}{3} + O[x]^{11}$$

Using mathematica; Series for the first 3 terms by assuming $x \rightarrow 0$ ($x \ll 1$)
I obtained:

$$P = \frac{8C x^5}{5} - \frac{4C x^7}{7} + \frac{C x^9}{3} + O(x^{11}) \quad ; \quad x = \left(\frac{P}{D}\right)^{1/9}$$

Further simplification for $x \ll 1$ - get the leading term:

$$P = \frac{8C}{5} \cdot \left(\frac{P}{D}\right)^{5/9} = \frac{8C}{5} \frac{1}{D^{5/9}} P^{\frac{5+9-9}{9}}$$

$$\Rightarrow P = \underbrace{\frac{8C}{5} \frac{1}{D^{5/9}}}_{K_*} \cdot P^{\underbrace{1 - \frac{5-9}{9}}_{1/n_*}}$$

$$\Rightarrow \boxed{K = \frac{8C}{5D^{5/9}} \quad , \quad n_* = \frac{9}{5-9}}$$

Using the generic relation between M and K and plugging K_* and n_* :

$$M = \underbrace{-4\pi \left(\frac{4\pi G}{(n_*+1)K_*} \right)^{\frac{n_*}{1-n_*}}}_{A} \sum_{n_*}^{\frac{n_*+1}{n_*-1}} \theta'(\xi_{n_*}) R^{\underbrace{\frac{3-n_*}{1-n_*}}_B}$$

Introduce A & B constants to simplify the relationship between M & R

$$\Rightarrow M = A R^B$$

Where:

$$A = -4\pi \left(\frac{4\pi G}{(n_*+1)K_*} \right)^{\frac{n_*}{1-n_*}} \sum_{n_*}^{\frac{n_*+1}{n_*-1}} \theta'(\xi_{n_*}) \quad ; \quad B = \frac{3-n_*}{1-n_*}$$

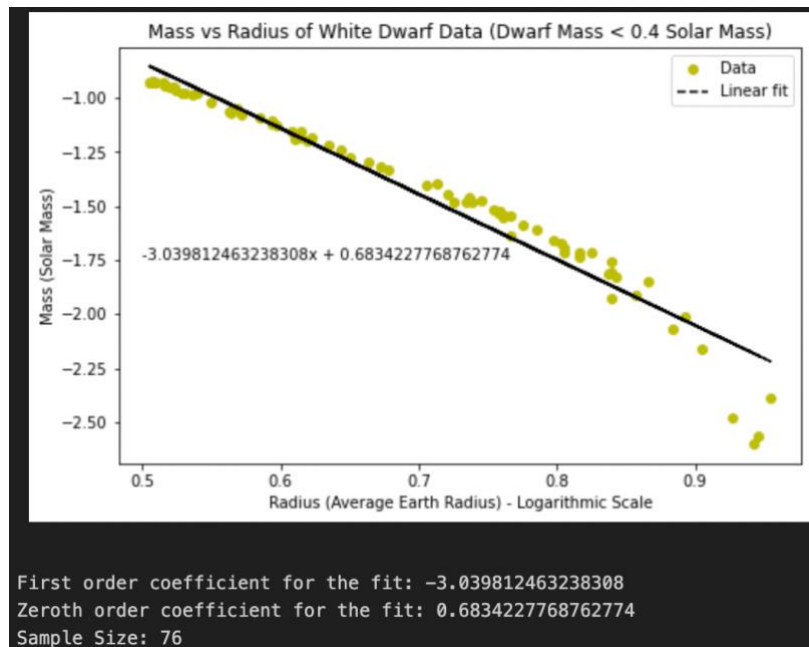
Given a power relation, as in our case $M = A R^B$ where one of the parameters appear in the exponent; it is intuitive to do a fit to the logarithm of the relation:

$$\ln M = \ln(A R^B) = \ln A + \ln R^B = \ln A + B \ln R$$

We have a linear relation between $\ln M$ and $\ln R$

$$\Rightarrow \ln M = B \ln R + \ln A$$

\Rightarrow linear fit to the \ln of the data points



From the linear fit of the form $mx+b$, the obtained coefficients are:

$$m = -3.0398 \quad b = 0.6834$$

n_* can be found directly.

$$B = \frac{3-n_*}{1-n_*} = -3.0398 \Rightarrow n_* \rightarrow 1.48507 \approx 1.5$$

For K_* we have the relation:

$$\ln A = \ln \left[-4\pi \left(\frac{4\pi G}{(n_*+1)K_*} \right)^{\frac{n_*}{1-n_*}} \sum_{n_*}^{\frac{n_*+1}{n_*-1}} \theta'(\xi_{n_*}) \right] = 0.6834$$

where $n_* = 1.5$

So we need the values for ξ_{n_*} and $\theta'(\xi_{n_*})$

Given the definition $\theta(\xi_{n_*}) = 0$: which corresponds to the surface of the star we obtain a root-finding problem to find ξ_{n_*} . But first we need to find the behavior of θ given the differential equation.

We have the Lane-Emden equation:

$$\frac{1}{\xi^2} \frac{d}{d\xi} \left(\xi^2 \frac{d\theta}{d\xi} \right) + \theta^n = 0$$

Can be written:

$$\frac{1}{\xi^2} \left(2\xi \frac{d\theta}{d\xi} + \xi^2 \frac{d^2\theta}{d\xi^2} \right) + \theta^n = 0$$

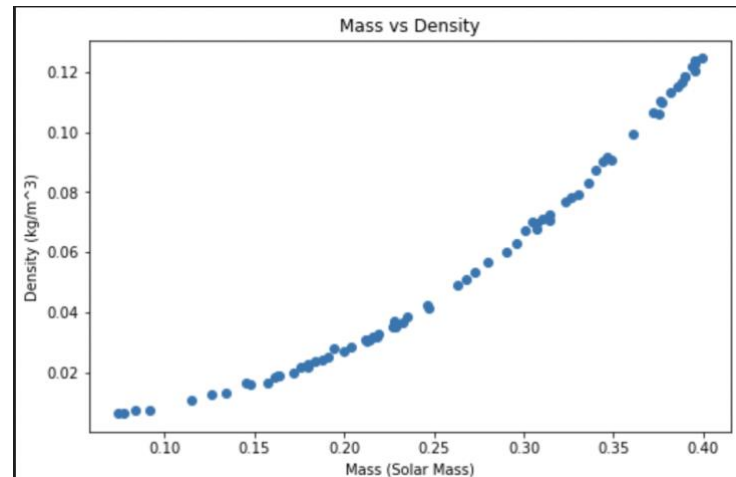
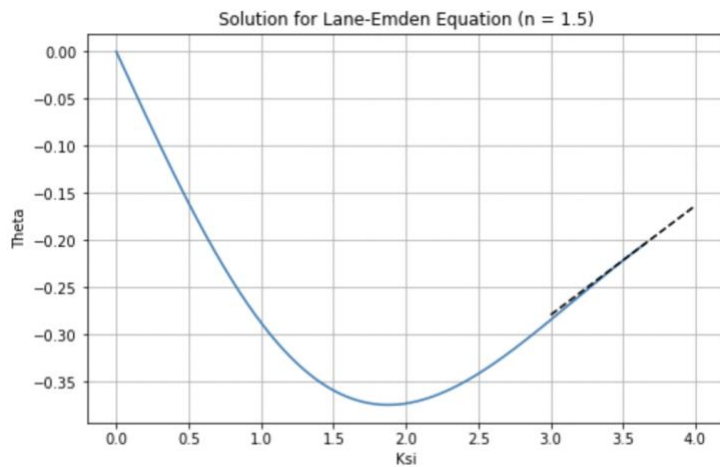
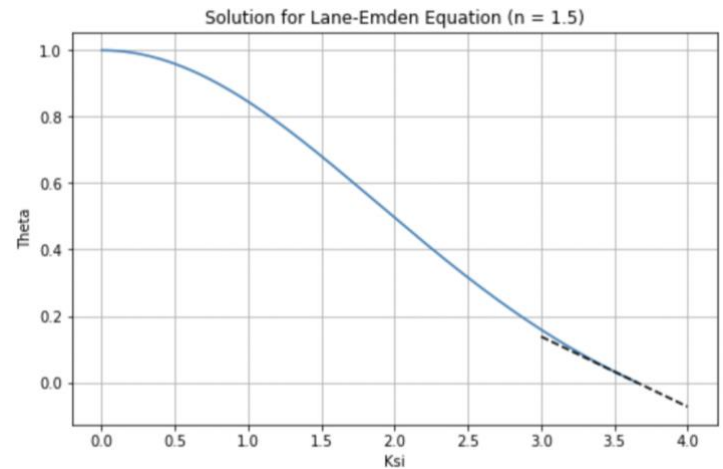
$$\Rightarrow \frac{2}{\xi} \frac{d\theta}{d\xi} + \frac{d^2\theta}{d\xi^2} + \theta^n = 0$$

$$\Rightarrow \frac{d^2\theta}{d\xi^2} = - \underbrace{\left(\frac{2}{\xi} \frac{d\theta}{d\xi} + \theta^n \right)}_{\text{RHS}}$$

From numerics:

$$\xi_n = 3.6581 \quad \theta'(\xi_n) = -0.2033$$

$$K_\star = 5.4058$$



```
ksi_n = 3.653083969542544
dtheta(ksi_n) = -0.20330250655674104
K_star = 5.40582658418237
```

Part(D)

I have implemented this method partially (most functions)

In order to solve ivp we need the relation between $\frac{dP}{dr}$ and $\frac{dp}{dr}$
for R & D

mass density relation

Use the EOS:

$$P = C [x(2x^2 - 3) \sqrt{x^2 + 1} - 3 \sinh^{-1} x]$$

$$x = \left(\frac{\rho}{D} \right)^{1/9}$$

$$\Rightarrow \frac{dP}{dr} = \frac{dP}{dx} \cdot \frac{dx}{d\rho} \cdot \frac{d\rho}{dr}$$

using Mathematica:

$$\frac{dP}{dx} = \frac{8Cx^4}{\sqrt{1+x^2}}$$

$$\frac{dx}{d\rho} = \frac{1}{D^{1/9}} \cdot \frac{1}{9} \rho^{\frac{1}{9}-1}$$

the original equation:

$$\frac{dP}{dr} = - \frac{Gm(r)\rho(r)}{r^2}$$

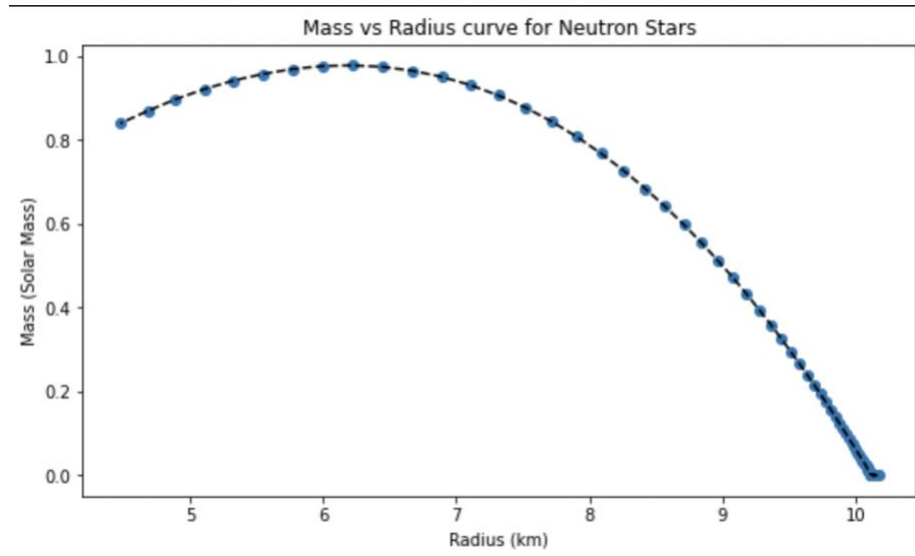
We obtain:

$$\left| \frac{d\rho}{dr} = \frac{\sqrt{1+x^2}}{8Cx^4} \cdot 9 \cdot D^{1/9} \cdot \rho^{1-\frac{1}{9}} \cdot \left(- \frac{Gm\rho}{r^2} \right) \right|$$

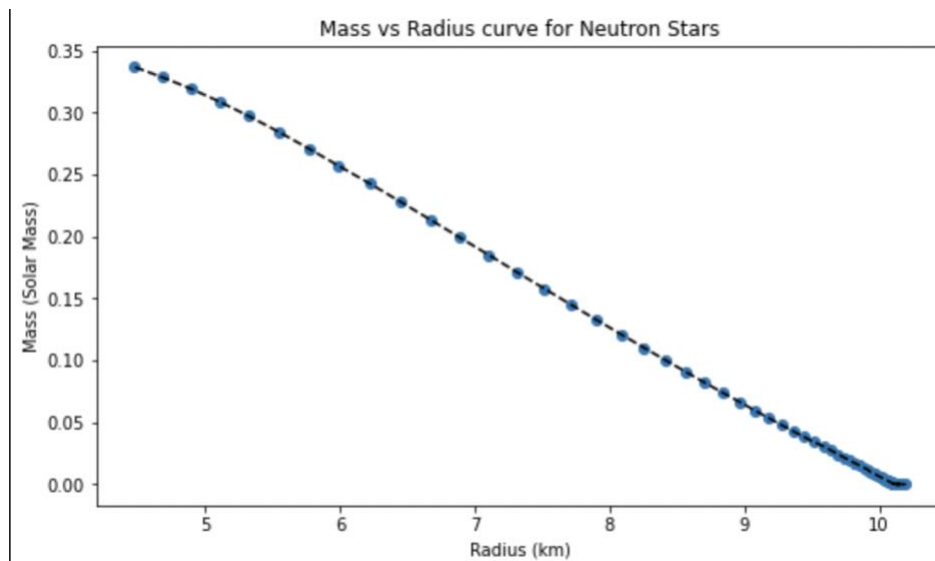
Einstein Part

I think I have scaling issue (around 2) for this part which I realized at the last part when finding valid K values. I couldn't address the reason.

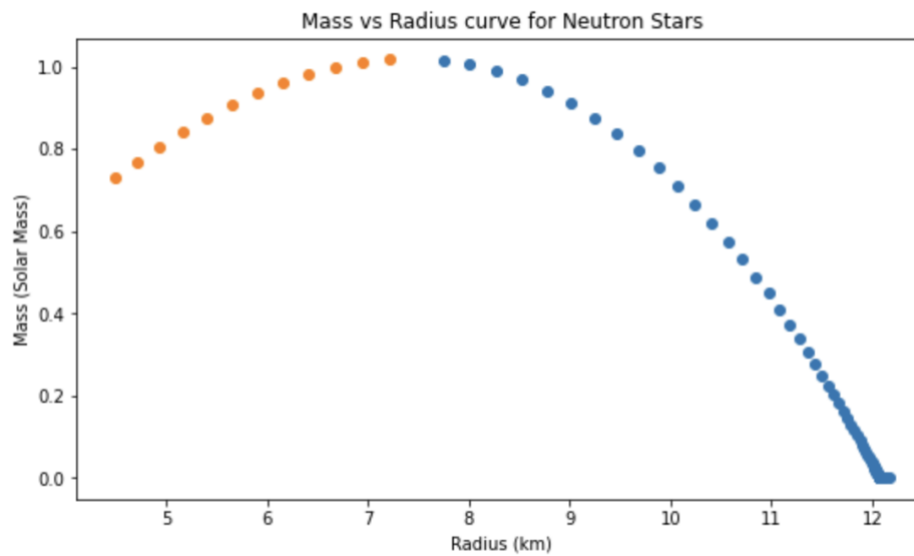
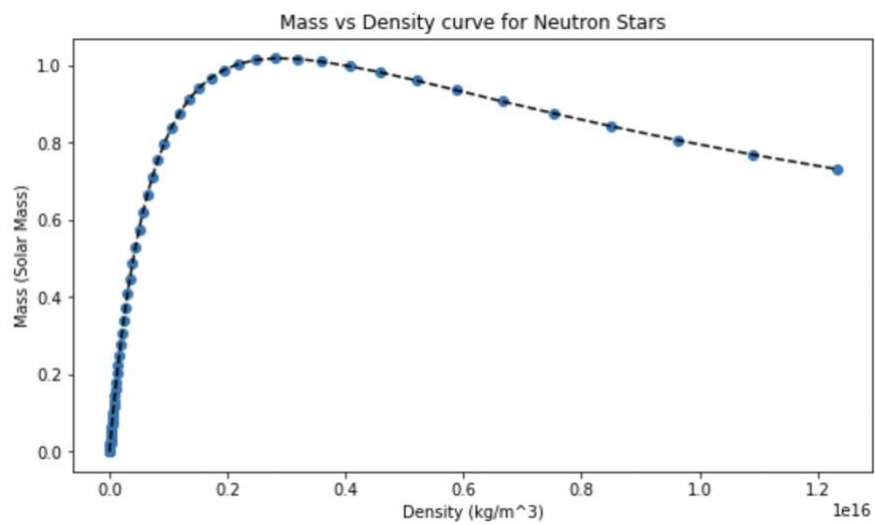
Part(A)



Part(B)



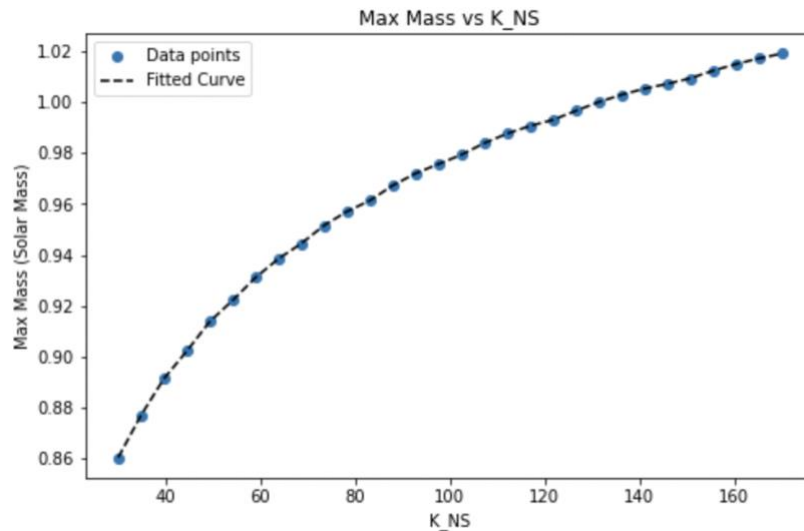
Part(C)



The maximal NS mass: 1.018979660888874 Solar Masses

Part(D)

I scaled with my hand the observed mass value to match with my computation



Max allowed K value: 124.87487487487488

Part(E)

```
In[6]:= FullSimplify[DSolve[v'[r] == (2*M)/(r*(r-2*M)), v[r], r]]
```

```
Out[6]= {{v[r] -> c1 - Log[r] + Log[-2 M + r]}}
```

The obtained solution by mathematica:

$$v(r) = C - \ln(r) + \ln(-2M+r)$$

$$\Rightarrow v(r) = \ln\left(-\frac{2M}{r} + 1\right) + C$$

At $r=R$ we have:

$$v(R) = \ln\left(-\frac{2M}{R} + 1\right) + C$$

$$\Rightarrow C = v(R) - \ln\left(-\frac{2M}{R} + 1\right)$$

Then by substituting this for the generic equation with $r > R$ condition we obtain:

$$v(r > R) = \ln\left(1 - \frac{2M}{r}\right) - \ln\left(1 - \frac{2M}{R}\right) + v(R)$$