

Notes

Something

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Caprice

1 Fall 2018

This is my note for some non-trivial but not systematic problems which involves some interesting physics or maths.

| Section 1. Fall 2018

1.1 Walkway equilibrium

Suppose the mass of the objects attached to each end of the rope are m_1 and m_2 , The angles between each segment of the rope, bended by the central object which has mass M , with the horizontal plane are θ and ϕ . The distance between two pulleys is L , and what we want to know is the vertical displacement d of the central object. Thus we can obtained the equations for d when the system is at equilibrium.

$$L = d(\cot \theta + \cot \phi), \quad (1.1)$$

$$m_1 g \cos \theta = m_2 g \cos \phi, \quad (1.2)$$

$$m_1 g \sin \theta + m_2 g \sin \phi = Mg, \quad (1.3)$$

From (1.2), we have $\cos \phi = \frac{m_1}{m_2} \cos \theta$, thus (1.3) can be written as

$$m_1 \sin \theta + m_2 \sqrt{1 - \frac{m_1^2}{m_2^2} (1 - \sin^2 \theta)} = M, \quad (1.4)$$

such that we can solve for $\sin \theta$ and $\cos \theta$

$$\sin \theta = \frac{M^2 + m_1^2 - m_2^2}{2Mm_1}, \quad (1.5)$$

$$\cos \theta = \sqrt{1 - \sin^2 \theta} = \frac{1}{2Mm_1} \sqrt{[(m_1 + m_2)^2 - M^2][M^2 - (m_1 - m_2)^2]}, \quad (1.6)$$

$$\cot \theta = \frac{\sqrt{[(m_1 + m_2)^2 - M^2][M^2 - (m_1 - m_2)^2]}}{M^2 + m_1^2 - m_2^2}, \quad (1.7)$$

together with $\sin \phi$ and $\cos \phi$

$$\cos \phi = \frac{m_1}{m_2} \cos \theta = \frac{1}{2Mm_2} \sqrt{[(m_1 + m_2)^2 - M^2][M^2 - (m_1 - m_2)^2]}, \quad (1.8)$$

$$\sin \phi = \sqrt{1 - \cos^2 \phi} = \frac{M^2 - m_1^2 + m_2^2}{2Mm_2}, \quad (1.9)$$

$$\cot \phi = \frac{\sqrt{[(m_1 + m_2)^2 - M^2][M^2 - (m_1 - m_2)^2]}}{M^2 - m_1^2 + m_2^2}. \quad (1.10)$$

Therefore we can plug into (1.1) and obtain the expression of d as follows

$$d = \frac{L[M^4 - (m_1^2 - m_2^2)^2]}{2M^2 \sqrt{[(m_1 + m_2)^2 - M^2][M^2 - (m_1 - m_2)^2]}}. \quad (1.11)$$

The equilibrium condition in this case is

$$|m_1 - m_2| < M < (m_1 + m_2). \quad (1.12)$$

such that the argument under the square root is positive. Also we can easily check that if $m_1 = m_2 = m$ then this result reduces to our former result

$$d = \frac{LM}{2\sqrt{4m^2 - M^2}}. \quad (1.13)$$

1.2 A derivation of Gamma function from Fourier transform

It is well-known that $\Gamma(n+1) = n!$ for any natural number $n \in \mathbb{N}$. It is natural to ask what is $\Gamma(x)$ for any real number $x \geq 1$. Our purpose is to show that Gamma function can be express as an integral

$$\Gamma(x) = \int_0^\infty dx t^{x-1} e^{-t}, \quad (1.14)$$

given that

$$\Gamma(x+1) = x\Gamma(x), \quad (1.15)$$

which is the most essential property and motivation to define the Gamma function.

Using Talor expansion we can easily show that

$$f(x + \Delta x) = \sum_{n=0}^{\infty} \frac{f^{(n)}(x)}{n!} (\Delta x)^n = \exp\left(\Delta x \frac{d}{dx}\right) f(x) \quad (1.16)$$

where Δx is a constant of translation. Therefore $\Gamma(x+1) = e^{\frac{d}{dx}} \Gamma(x)$, and we can rewrite (1.17) as

$$e^{\frac{d}{dx}} \Gamma(x) = x\Gamma(x). \quad (1.17)$$

Consider doing Fourier transform of (1.17), such that $\frac{d}{dx} \rightarrow i\omega$, $x \rightarrow i\frac{d}{d\omega}$, $\Gamma(x) \rightarrow \tilde{\Gamma}(\omega)$, and

$$\tilde{\Gamma}(\omega) = \mathcal{F}[\Gamma(x)] = \int_{-\infty}^{\infty} dx \Gamma(x) e^{-i\omega x}, \quad (1.18)$$

$$\Gamma(x) = \mathcal{F}^{-1}[\tilde{\Gamma}(\omega)] = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \tilde{\Gamma}(\omega) e^{i\omega x}, \quad (1.19)$$

$$e^{i\omega} \tilde{\Gamma}(\omega) = i \frac{d}{d\omega} \tilde{\Gamma}(\omega). \quad (1.20)$$

Solve the above differential equation of $\tilde{\Gamma}(\omega)$ we find

$$\tilde{\Gamma}(\omega) = C \exp(-e^{i\omega}), \quad (1.21)$$

$$\Gamma(x) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} C \exp(-e^{i\omega}) e^{i\omega x}. \quad (1.22)$$

However, (1.22) does not converge, since (1.21) is a nonzero periodic function.

To resolve this difficulty of convergence, we expand the domain of Gamma function to the complex plane, such that $\Gamma(z+1) = z\Gamma(z)$, where $z \in \mathbb{C}$. Consider a pure imaginary number $z = ix$, where $x \in \mathbb{R}$, we can rewrite the recursion relation (1.15) as

$$e^{-i\frac{d}{dx}} \Gamma(ix) = ix\Gamma(ix) \quad (1.23)$$

Again using Fourier transform we have

$$e^{\omega} \mathcal{F}[\Gamma(ix)] = -\frac{d}{d\omega} \mathcal{F}[\Gamma(ix)], \quad (1.24)$$

where $\mathcal{F}[\Gamma(ix)]$ is the Fourier transform of $\Gamma(ix)$

$$\mathcal{F}[\Gamma(ix)] = \int_{-\infty}^{\infty} dx \Gamma(ix) e^{-i\omega x}. \quad (1.25)$$

Solve (1.24) we have

$$\mathcal{F}[\Gamma(ix)] = C \exp(-e^{\omega}), \quad (1.26)$$

$$\Gamma(ix) = \frac{C}{2\pi} \int_{-\infty}^{\infty} d\omega \exp(-e^{\omega}) e^{i\omega x}. \quad (1.27)$$

Thus

$$\begin{aligned} \Gamma(z) &= \frac{C}{2\pi} \int_{-\infty}^{\infty} d\omega \exp(-e^{\omega}) e^{\omega z} \\ &= \frac{C}{2\pi} \int_{-\infty}^{\infty} de^{\omega} \exp(-e^{\omega}) e^{\omega(z-1)} \\ &= \frac{C}{2\pi} \int_{-\infty}^{\infty} dt t^{z-1} e^{-t}. \end{aligned}$$

To determine the constant C we use the fact that $\Gamma(1) = 0! = 1$, thus $C/2\pi = 1$, and we obtain the final integral expression of Gamma function

$$\Gamma(z) = \int_{-\infty}^{\infty} dt t^{z-1} e^{-t} \quad (1.28)$$

where $z \in \mathbb{C}$.

1.3 Euler's reflection formula

In mathematics, a reflection formula or reflection relation for a function f is a relationship between $f(a-x)$ and $f(x)$. A famous relationship is Euler's reflection formula

Proposition 1.1.

$$\Gamma(z)\Gamma(1-z) = \frac{\pi}{\sin \pi z}, \quad z \notin \mathbb{Z}. \quad (1.29)$$

Proof. To prove this reflection formula, we first notice a relationship between Gamma function and Beta function

$$B(q, p) = \frac{\Gamma(q)\Gamma(p)}{\Gamma(q+p)}, \quad (1.30)$$

where

$$B(q, p) = \int_0^1 dt t^{q-1} (1-t)^{p-1}, \quad q, p \neq 0, -1, -2, \dots \quad (1.31)$$

This can be shown by performing a variable transformation of $\Gamma(q)\Gamma(p)$

$$\begin{aligned}
 \Gamma(q)\Gamma(p) &= \int_0^\infty du e^{-u} u^{q-1} \int_0^\infty dv e^{-v} v^{p-1} \\
 u = zt, v = z(1-t), \quad &= \int dz dt \left| \frac{\partial(u, v)}{\partial(z, t)} \right| e^{-z} (zt)^{q-1} [z(1-t)]^{p-1} \\
 &= \int_0^\infty dz z^{q+p-1} e^{-z} \int_0^1 dt t^{q-1} (1-t)^{p-1} \\
 &= \Gamma(q+p) B(q, p).
 \end{aligned}$$

Therefore

$$\Gamma(z)\Gamma(1-z) = B(z, 1-z) = \int_0^1 dt t^{z-1} (1-t)^{-z}. \quad (1.32)$$

In order to prove Proposition 1.1, we only need to prove

$$\int_0^1 dt t^{z-1} (1-t)^{-z} = \frac{\pi}{\sin \pi z}. \quad (1.33)$$

Perform a variable substitution $t \rightarrow \frac{x}{1+x}$, such that $dt = dx / (1+x)^2$ and

$$\int_0^1 dt t^{z-1} (1-t)^{-z} = \int_0^\infty dx \frac{x^{z-1}}{1+x} \quad (1.34)$$

Consider the following integral

$$\int_0^\infty dx \frac{x^{\alpha-1}}{x + e^{i\phi}}, \quad 0 < \alpha < 1, \quad -\pi < \phi < \pi. \quad (1.35)$$

□