

# A Quantum Dot Interacting With a Nano-mechanical Resonator

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# Single Atom Laser

- Conventional laser: pump an ensemble of atoms into the excited state to achieve population inversion
- Early experiments on single atom laser:  
e.g. a Rb-85 atom couples to a microwave photon cavity

## One-Atom Maser

D. Meschede, H. Walther, and G. Müller  
Phys. Rev. Lett. **54**, 551 – Published 11 February 1985

- Recent attempts using quantum dots as artificial atoms

## Semiconductor double quantum dot micromaser

Y.-Y. Liu<sup>1</sup>, J. Stehlik<sup>1</sup>, C. Eichler<sup>1</sup>, M. J. Gullans<sup>2</sup>, J. M. Taylor<sup>2,3</sup>, J. R. Petta<sup>1,4,\*</sup>

+ See all authors and affiliations

Science 16 Jan 2015:  
Vol. 347, Issue 6219, pp. 285-287  
DOI: 10.1126/science.aaa2501

- Phonon lasing: quantum dots with mechanical resonators

## Resonant and Inelastic Andreev Tunneling Observed on a Carbon Nanotube Quantum Dot

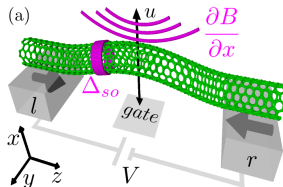
J. Gramich, A. Baumgartner, and C. Schönenberger  
Phys. Rev. Lett. **115**, 216801 – Published 16 November 2015

# Quantum dot with carbon-nanotube

## Control of vibrational states by spin-polarized transport in a carbon nanotube resonator

P. Stadler, W. Belzig, and G. Rastelli

Phys. Rev. B **91**, 085432 – Published 27 February 2015

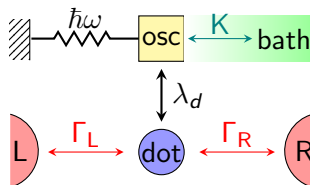


- A carbon nanotube quantum dot suspended between two ferromagnetic leads
- The dot spin couples to the vibration mode of the nanotube because of the nanotube spin-orbit interaction and/or a magnetic field gradient
- Lasing has been shown theoretically with collinearly polarised leads [Mantovani *et al.*, 2019]

Our Goal:

Non-collinearly polarised leads?

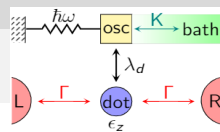
# The System



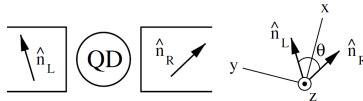
|                            |             |                                   |                                   |                              |
|----------------------------|-------------|-----------------------------------|-----------------------------------|------------------------------|
| Dot State $ \alpha\rangle$ | $ 0\rangle$ | $ \downarrow\rangle$              | $ \uparrow\rangle$                | $ \uparrow\downarrow\rangle$ |
| Dot Energy $E_\alpha$      | 0           | $\epsilon - \frac{\epsilon_z}{2}$ | $\epsilon + \frac{\epsilon_z}{2}$ | $2\epsilon + U$              |

- Basis:  $|\alpha, n\rangle = |\alpha\rangle \otimes |n\rangle$ ;  $n = 0, 1, 2, \dots$  oscillator occupation number
- $E_{\downarrow, n+1} = \epsilon - \frac{\epsilon_z}{2} + (n+1)\hbar\omega$ ;  $E_{\uparrow, n} = \epsilon + \frac{\epsilon_z}{2} + n\hbar\omega$
- When  $\hbar\omega \approx \epsilon_z$ ,  $E_{\downarrow, n+1} \approx E_{\uparrow, n}$ : the dot can exchange energy with the oscillator by flipping its spin.
- Electrons tunnel between the dot and the leads.

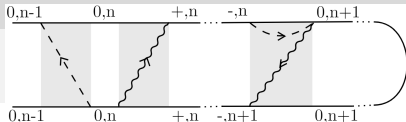
# Assumptions



- Rotation Wave Approximation
  - ▶ the interaction strength between the dot and the oscillator is weak:  $|\lambda_d| < \hbar\omega \approx \epsilon_z$
  - ▶ four eigenstates  $|\chi, n\rangle$  of the dot and the oscillator:
    - $|0, n\rangle, |\uparrow\downarrow, n\rangle$
    - $|+, n\rangle = \sin \theta_n |\uparrow, n\rangle - \cos \theta_n |\downarrow, n+1\rangle$
    - $|-, n\rangle = \cos \theta_n |\uparrow, n\rangle + \sin \theta_n |\downarrow, n+1\rangle$
- Unidirectional Transport (from left to right)
  - ▶ infinitely large voltage on the leads
  - ▶ Fermi functions:  $f_L = 1$  and  $f_R = 0$
- 0 Temperature
  - ▶ the bath cools down the oscillator
  - ▶  $n=0$  when the oscillator is in equilibrium
- Spin Flips in x Axis
  - ▶  $\hat{e}_x = \hat{n}_L + \hat{n}_R$
  - ▶  $\hat{e}_y = \hat{n}_L - \hat{n}_R$
  - ▶  $\hat{e}_z = \hat{n}_L \times \hat{n}_R$



# Analysis



- Reduced Density Matrix of the Dot and the Oscillator

- $$\hat{\rho} = \sum_{\chi_1, n_1; \chi_2, n_2} P_{\chi_2, n_2}^{\chi_1, n_1} |\chi_2, n_2\rangle \langle \chi_1, n_1|$$

- diagonal terms:  $P_{\chi, n} = P_{\chi, n}^{\chi, n}$  and  $\sum_{\chi, n} P_{\chi, n} = 1$

- Master Equation at Steady State

- $$0 = \dot{P}_{\chi_2, n_2}^{\chi_1, n_1} = -\frac{i}{\hbar} (E_{\chi_1, n_1} - E_{\chi_2, n_2}) P_{\chi_2, n_2}^{\chi_1, n_1} + \sum_{\chi'_1, n'_1; \chi'_2, n'_2} W_{\chi_2, n_2; \chi'_2, n'_2}^{\chi_1, n_1; \chi'_1, n'_1} P_{\chi'_2, n'_2}^{\chi'_1, n'_1}$$

- general transition rate:  $W_{\chi_2, n_2; \chi'_2, n'_2}^{\chi_1, n_1; \chi'_1, n'_1} \propto \Gamma$  or  $K$

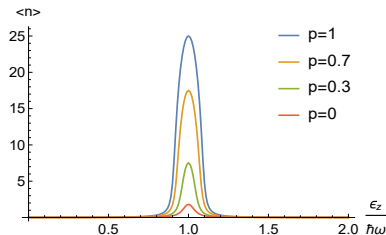
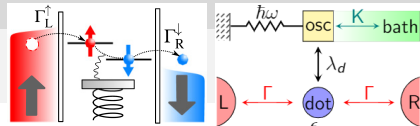
- Keldysh Formalism with Diagrammatic Rule

- $|\chi, n\rangle$  propagates on the Keldysh contour
  - leads electron tunnelling line & bath phonon tunnelling line

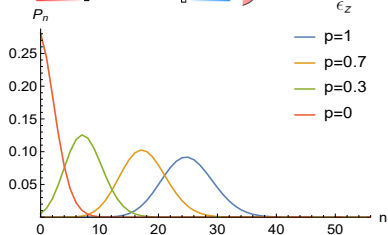
## Numerical Simulation in First Order Expansion ( $n_{\max} = 100$ )

- off-diagonal elements reduce to zero
- diagrammatic results should agree with Fermi Golden Rule ones

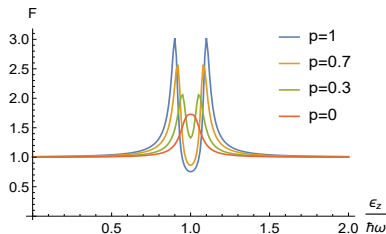
# Anti-parallel Polarisation



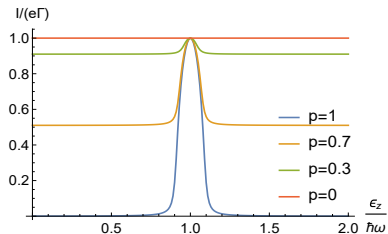
(a) average occupation number



(b) probability distribution at resonance

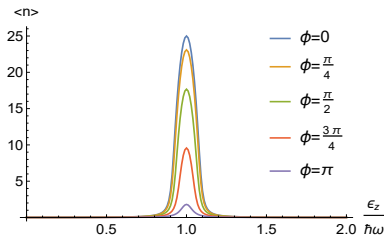
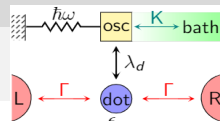


(c) Fano factor ( $F$ ) =  $\frac{\langle (\Delta n)^2 \rangle}{\langle n \rangle}$

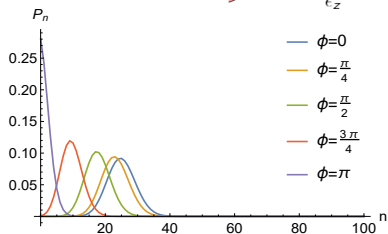


(d) the current in  $e\Gamma$  units

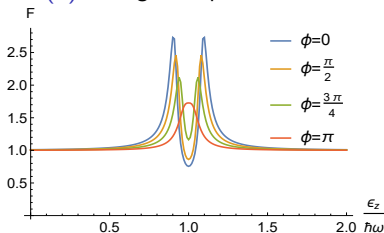
# Non-collinear Polarisation (1st order, $p=1$ )



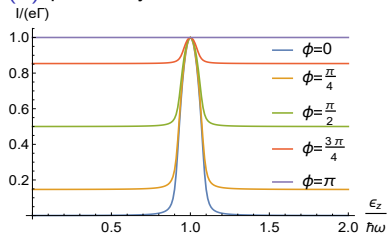
(a) average occupation number



(b) probability distribution at resonance



(c) Fano factor ( $F$ ) =  $\frac{\langle(\Delta n)^2\rangle}{\langle n \rangle}$



(d) the current in  $e\Gamma$  units



# Conclusion

## Carbon-nanotube quantum dot with ferromagnetic leads

- Reproduced collinear polarisation results
  - ▶ Demonstrated the single atom lasing when the polarisation is sufficiently large
- Extended to non-collinear case
  - ▶ Developed diagrammatic rules
  - ▶ Demonstrated the single atom lasing when the angle between two polarisation axes is sufficiently small with full polarisation strength

• **Thank You** 😊 ☕

# The Hamiltonian

$$\hat{H}_{\text{leads}} = \sum_{\eta, \sigma, k} \epsilon_{k, \sigma} \hat{c}_{\eta, \sigma, k}^{\dagger} \hat{c}_{\eta, \sigma, k},$$

$$\hat{H}_{\text{tun}} = \sum_{\eta, \sigma, k} \left( V_{\eta, \sigma} \hat{c}_{\eta, \sigma, k}^{\dagger} \hat{d}_{\sigma} + V_{\eta, \sigma}^{*} \hat{d}_{\sigma}^{\dagger} \hat{c}_{\eta, \sigma, k} \right)$$

$$\hat{H}_{\text{dot}} = \sum_{\sigma} \epsilon_{\sigma} \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{\uparrow} \hat{n}_{\downarrow}$$

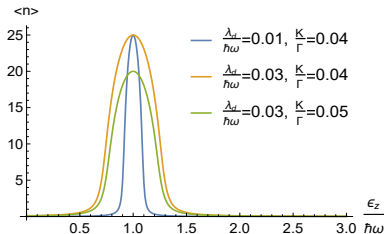
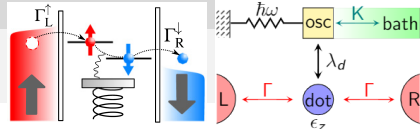
$$\hat{H}_{\text{intdot}} = -\lambda_d \left( \hat{d}_{\downarrow}^{\dagger} \hat{d}_{\uparrow} + \hat{d}_{\uparrow}^{\dagger} \hat{d}_{\downarrow} \right) \left( \hat{b}^{\dagger} + \hat{b} \right)$$

$$\hat{H}_{\text{osc}} = \hbar \omega \hat{b}^{\dagger} \hat{b},$$

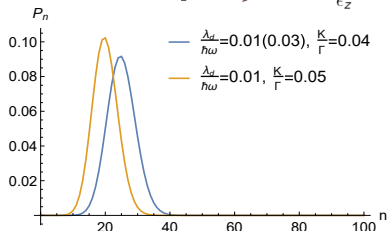
$$\hat{H}_{\text{intbath}} = \lambda_b \sum_q \left( a_q^{\dagger} b + b^{\dagger} a_q \right)$$

$$\hat{H}_{\text{bath}} = \sum_q \hbar \omega_q a_q^{\dagger} a_q$$

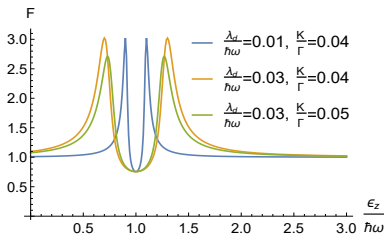
# Anti-parallel Polarisation ( $p=1$ )



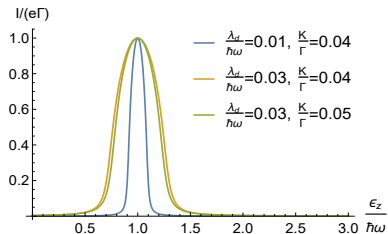
(a) average occupation number



(b) probability distribution at resonance

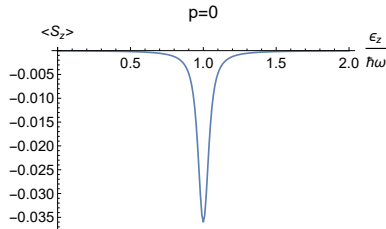
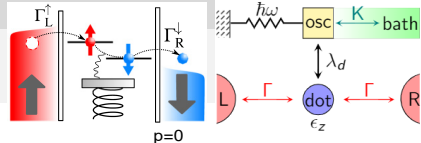


(c) Fano factor ( $F$ ) =  $\frac{\langle (\Delta n)^2 \rangle}{\langle n \rangle}$

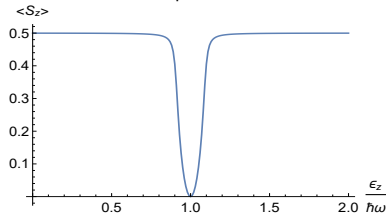


(d) the current in  $e\Gamma$  units

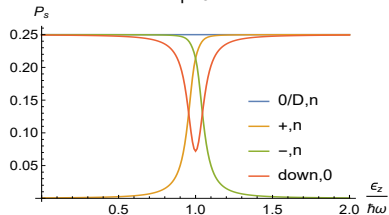
# $p=0$ vs $p=1$



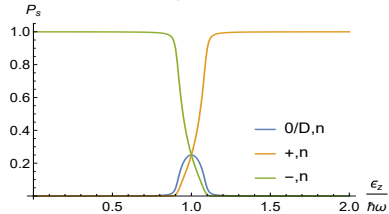
(a) average spin in the dot  
 $p=1$



(c) average spin in the dot



(b) probability of each state  
 $p=1$



(d) probability of each state