

Recap.

1.  $N \triangleleft G$ .  $gNg^{-1} = N \quad \forall g \in G$ .

c.f. center  $Z(G)$ . :  $gzg^{-1} = z \quad \forall z \in Z(G)$   
 $g \in G$ .

① trivial  $N$ :  $\{1\} \triangleleft G$ .

$G$  has only trivial  $N$  : "simple group"

(Lagrange)  $\mu \nmid \text{simple}$

② any hom  $\mu: G \rightarrow G'$

$K \equiv \ker \mu \triangleleft G$

2. Quotient groups.

cosets  $G/N = \{gN, g \in G\}$

$(g_1N)(g_2N) := (g_1g_2)N$

$\hookrightarrow = g_1g_2 \underbrace{g_2^{-1}N g_2}_{=N} N$   
 $= g_1g_2N$

3. Given  $N$ . construct hom

$\phi: G \rightarrow G/N$  surjective.

$g \mapsto gN$

$\ker \phi = N$

4.  $\mu: G \rightarrow \mu(G) \subset G'$

$\phi \hookrightarrow G/K$  inj.  
surj.

1st iso. thm.  $G/K \cong \mu(G)$  :  $gK \mapsto \mu(g)$

## 6.4 Quotient groups and short exact sequences

$$\dots G_{i-1} \xrightarrow{f_{i-1}} G_i \xrightarrow{f_i} G_{i+1} \xrightarrow{f_{i+1}} \dots$$

$$\text{im } f_{i-1} = \ker f_i$$

short exact sequence.  $1 \xrightarrow{f_0} G_1 \xrightarrow{f_1} G_2 \xrightarrow{f_2} G_3 \xrightarrow{f_3} 1$

$$f_0: \{1\} \rightarrow \{1_{G_1}\} \quad \text{im } f_0 = \{1\}, f_3: G_3 \rightarrow \{1\} \quad \ker f_3 = G_3$$

①  $G_1$ :  $\ker f_1 = \{1_{G_1}\}$  injective

②  $G_2$ :  $\text{im } f_1 = \ker f_2$

③  $G_3$ :  $\text{im } f_2 = G_3$  surjective

Now consider a homomorphism  $\mu: G \rightarrow G'$

$$K = \ker \mu.$$

We have

$$1 \rightarrow K \xrightarrow{i} G \xrightarrow{\mu} \text{im } \mu \rightarrow 1$$

inclusion map

$\cong G/K$

Exactness check:

①  $K$ :  $\ker i = \{1_G\}$  ✓

②  $G$ :  $\ker \mu = \text{im } i = K$  ✓

③  $\text{im } \mu$ :  $\ker(\text{im } \mu \rightarrow 1) = \text{im } \mu$  ✓

1st isomorphism theorem  $\Rightarrow$

$$\boxed{1 \rightarrow K \rightarrow G \rightarrow G/K \rightarrow 1}$$

Remarks.

1. If we have SES.

$$1 \rightarrow N \xrightarrow{f_1} G \xrightarrow{f_2} Q \rightarrow 1$$

then  $N \cong \ker f_2$  (it is iso. to the kernel  
of homomorphism  $G \rightarrow Q$ )

We sometimes write  $Q$  as  $G/f_1(N)$

where  $f_1: N \rightarrow G$  is an injective  
homomorphism.

" $G$  is an extension of  $Q$  by  $N$ "

Example.

$$1. \quad 1 \rightarrow G_1 \xrightarrow{(\quad)} G_1 \times G_2 \xrightarrow{(\quad)} G_2 \rightarrow 1$$

$(G_2) \qquad \qquad \qquad (G_1)$

$$\mu: G_1 \times G_2 \rightarrow G_2 \quad \begin{pmatrix} g_1 \in G_1 \\ g_2 \in G_2 \end{pmatrix}$$

$$(g_1, g_2) \mapsto g_2$$

$$2. \quad \varphi: \mu_4 \rightarrow \mu_2 \quad (\mathbb{Z}_4 \rightarrow \mathbb{Z}_2)$$

$$w \mapsto w^2 \quad w = e^{i\frac{2\pi}{4}}$$

$$\ker \varphi = \{\pm 1\} \cong \mathbb{Z}_2$$

$$1 \rightarrow \mathbb{Z}_2 \rightarrow \mathbb{Z}_4 \rightarrow \mathbb{Z}_2 \rightarrow 1$$

in general  $1 \rightarrow \mathbb{Z}_n \rightarrow \mathbb{Z}_{n^2} \rightarrow \mathbb{Z}_n \rightarrow 1$

$$(\varphi: \mu_{n^2} \rightarrow \mu_n)$$

$$z \mapsto z^n$$

$$3. \quad \det: O(n) \rightarrow \mathbb{Z}_2 \quad AA^T = 1 \Rightarrow \det = \pm 1$$

$$M \mapsto \det(M)$$

$$\ker(\det) = SO(n)$$

$$1 \rightarrow SO(n) \rightarrow O(n) \rightarrow \mathbb{Z}_2 \rightarrow 1$$

Does it work?

$$1 \rightarrow \mathbb{Z}_2 \rightarrow O(n) \rightarrow SO(n) \rightarrow 1$$

$$\mu: O(n) \rightarrow SO(n) ?$$

for some axis  $\hat{n}$ :  $R_0 \in SO(n)$ . then  $\sigma R_0 \in O(n) \setminus SO(n) = PSO(n)$   $\sigma$ : reflection

$$\mu(R_0) = \mu(\sigma R_0) = R_0$$

$$\mu(\underbrace{\sigma R_0 \sigma R_0}_{R_{-\theta}}) = \mu(R_{\phi-\theta}) = R_{\phi-\theta} \neq \mu(\sigma R_0) \mu(\sigma R_0)$$

$$4. \quad \det: U(n) \rightarrow U(1) \\ u \mapsto \det u = \lambda \quad \ker? \quad SU(n)$$

$$1 \rightarrow SU(n) \rightarrow U(n) \rightarrow U(1) \rightarrow 1$$

$$5. \quad \pi: SU(2) \rightarrow SO(3)$$

$$u \vec{x} \cdot \vec{\sigma} u^\dagger := (\pi(u) \vec{x}) \cdot \vec{\sigma}$$

$$u \in \ker \pi \quad u \vec{x} \cdot \vec{\sigma} u^\dagger = \vec{x} \cdot \vec{\sigma} \quad u = \lambda \mathbb{1} \\ \lambda = \pm 1$$

$$\pi(u) = \pi(-u)$$

$$\Rightarrow \ker \pi \cong \mathbb{Z}_2$$

$$SU(2)/\mathbb{Z}_2 \cong SO(3)$$

$$1 \rightarrow \mathbb{Z}_2 \rightarrow SU(2) \xrightarrow{\text{quantum spin states}} SO(3) \xrightarrow{\text{rot. of classical vectors}} 1$$

$$\uparrow \text{abelian, and } \mathbb{Z}_2 \subset Z(SU(2))$$

phase ambiguity /

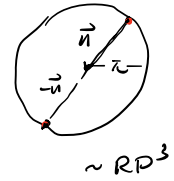
different projective reps

Analogy in 2D?

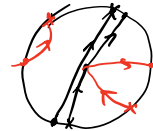
linear  $\leftarrow$  projective

$$1 \rightarrow \pi_1(SO(2)) \rightarrow SO(2) \rightarrow 1$$

$\cong \mathbb{Z}$



$$G/\mathbb{Z} \cong SO(2) \quad G = \mathbb{R}$$



$$1 \rightarrow \mathbb{Z} \rightarrow \mathbb{R} \rightarrow SO(2) \rightarrow 1$$

$$U_3(\theta) = e^{i\theta\sigma} \quad \text{arbitrary phases} \quad \text{anyons}$$

These are examples of:

Definition (central extension)

$$1 \rightarrow A \xrightarrow{i} E \xrightarrow{\pi} G \rightarrow 1$$

1.  $A$  is abelian.

2.  $A \subset Z(E)$   $i(a)b = b i(a) \quad (\forall a \in A, \forall b \in E)$

$$1 \rightarrow N \rightarrow G^{\text{Quantum}} \rightarrow G^{\text{Classical}} \rightarrow 1$$

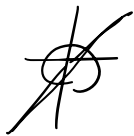
Motivation for such extensions. in QM  $\begin{matrix} 11.1 \\ 15.1 \& 2 \end{matrix}$

In QM. we talk about  $\mathcal{H} \setminus \{0\}$  but distinct states are rep'd by a set of vectors:

$$|\psi_1\rangle \sim |\psi_2\rangle \text{ if } |\psi_1\rangle = \lambda |\psi_2\rangle, \lambda \in \mathbb{C}^*$$

$$\text{"rays"} \quad \underline{\psi} = \{e^{i\alpha} \psi; \alpha \in \mathbb{R}\}$$

This is actually the projective complex plane.  $(z_1, z_2, z_3) \sim \lambda(z_1, z_2, z_3), (\lambda \neq 0) \in \mathbb{CP}^3$   
represented by  $[z_1 : z_2 : z_3]$



|| This is not a linear space:

$$[1:0:0] \text{ ? } [0:1:0] \text{ cannot be defined.}$$

$$\text{no zero vector } [0:0:0]$$

uniquely defined by the density matrix

$$\rho_\psi = \frac{|\psi\rangle\langle\psi|}{\langle\psi|\psi\rangle} \quad (\rho^2 = \rho)$$

projective Hilbert space  $\mathbb{PH} := (\mathcal{H} \setminus \{0\}) / \sim$

Consider symmetry operations on  $\mathbb{PH}$ .

$$\text{The overlap. } \mathcal{O}(\rho_1, \rho_2) = \text{Tr}(\rho_1 \rho_2) = \frac{|\langle\psi_1|\psi_2\rangle|^2}{\|\psi_1\|^2 \|\psi_2\|^2}$$

should be conserved.

But for all kinds of reasons we want to  
work on linear spaces  $\mathcal{H}$ .  $\text{Aut}(\mathcal{H})$

┌ Wigner's theorem, sym. operations  
are unitary or antiunitary

$$\left\{ \begin{array}{l} \langle U\psi, U\phi \rangle = \langle \psi, \phi \rangle \\ \langle A\psi, A\phi \rangle = \overline{\langle \psi, \phi \rangle} = \langle \phi, \psi \rangle \\ (A^\dagger = -iA, A = UK) \end{array} \right. \quad \begin{array}{l} \\ \\ \rightarrow \text{plx. conj} \end{array} \quad \rfloor$$

Now consider <sup>unitary</sup> symmetry operations on states.  
one only needs

$$U(g_1)U(g_2)\psi = \underline{\alpha(g_1, g_2)} U(g_1 g_2)\psi$$

$$\alpha: G \times G \rightarrow U(1)$$

Not quite a group representation. projective-  
rep.

$$\left( \begin{array}{l} \text{if } \tilde{U}(g) = b(g) U(g) \\ \Rightarrow b(g_1) b(g_2) = f(g_1, g_2) b(g_1 g_2) \quad \forall g_1, g_2 \end{array} \right)$$

reduce back to rep.



rotation of spins:

$$\begin{array}{ccc} \text{classical} & & \text{quantum} \\ SO(3) & \longrightarrow & SU(2) \end{array}$$

Euler angles (xyz)

$$R(\phi, \theta, \psi) \rightarrow e^{i\frac{\phi}{2}\sigma^3} e^{i\frac{\theta}{2}\sigma^1} e^{i\frac{\psi}{2}\sigma^3}$$

$$R(2\pi, 0, 0) = 1 \quad \rightarrow \pm 1 \quad \text{for fermion/boson}$$

a  $\mathbb{Z}_2$  phase

The central extension of  $G$  by  $A$

$$1 \rightarrow A \xrightarrow{\iota} E \xrightarrow{\pi} G \rightarrow 1$$

is classified by the 2-cohomology

group  $H^2(G, A)$

6. finite Heisenberg group.

$$P = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix} \quad Q = \begin{pmatrix} \omega & & & 0 \\ & \omega^2 & & \\ & & \omega^2 & \\ 0 & & & 1 \end{pmatrix}$$

$$\omega = e^{i \frac{2\pi}{t}}$$

$$\underline{QP = \omega PQ}$$

← Weyl relation of  
canonical commutation  
relation.

some background.

$$[f, p] = i\hbar \quad (\hbar=1)$$

$$fP - Pf = i \quad P \cdot f \text{ acts on } f(P)$$

$$\Rightarrow A = e^{i\zeta P} \quad B = e^{i\eta f} \quad (\text{Weyl})$$

$$AB = e^{i\zeta P} \cdot e^{i\eta f}$$

$$e^x e^y = e^z$$

$$z = x + y + \frac{1}{2}[x, y]$$

$$\left\{ \begin{aligned} &= e^{i(\zeta P + \eta f) + \frac{1}{2}[i\zeta P, i\eta f]} \\ &+ \frac{1}{12}([x, [x, y]] - [y, [x, y]] + \dots) \end{aligned} \right.$$

$$BA = e^{i(\zeta P + \eta f) + \frac{1}{2}[i\eta f, i\zeta P]}$$

$$\Rightarrow AB = e^{i\zeta\eta} BA \equiv \omega BA \quad (A, B: n \times n \text{ mats.})$$

$$\det(AB) = \omega^n \det(BA) \Rightarrow \underline{\omega^n = 1}$$

$$\begin{cases} A^k B = \omega^k B A^k \\ A B^l = \omega^l B^l A \end{cases} \Rightarrow A^k B^l = \omega^{kl} B^l A^k$$

$$k=n, d=1 \Rightarrow \underline{A^n B = \omega^n B A^n} \stackrel{?}{\Rightarrow} A^n = 1$$

similarity  $B^n = 1$ .

A general element in Heis<sub>N</sub> has the form

$$\omega^a p^b q^c$$

$$(\omega^{a_1} p^{b_1} q^{c_1}) \cdot (\omega^{a_2} p^{b_2} q^{c_2}) = \omega^{a_3} p^{b_3} q^{c_3}$$

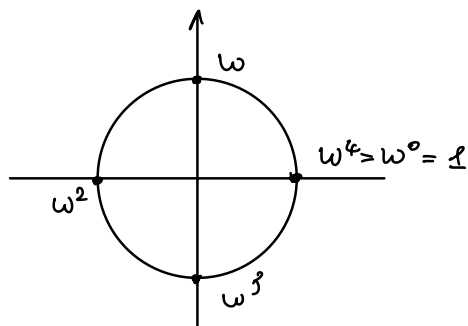
$$\begin{cases} a_3 = a_1 + a_2 + c_1 b_2 \\ b_3 = b_1 + b_2 \\ c_3 = c_1 + c_2 \end{cases}$$

$$\pi: \text{Heis}_N \longrightarrow \mathbb{Z}_N \times \mathbb{Z}_N$$

$$\omega^a p^b q^c \mapsto (b \bmod N, c \bmod N)$$

$$\ker(\pi) = \{ \omega^a \cancel{p^{c_1}} \cancel{q^{c_2}} \stackrel{1}{\omega^{c_1 c_2}} \} \cong \mathbb{Z}_N$$

$$1 \longrightarrow \mathbb{Z}_N \longrightarrow \text{Heis}_N \longrightarrow \mathbb{Z}_N \times \mathbb{Z}_N \longrightarrow 1$$



$$(P \cdot \psi)(\omega^k) := \psi(\omega^{k+1}) \quad \text{translation}$$

$$(Q \psi)(\omega^k) := \omega^k \psi(\omega^k) \quad \text{position operator}$$

$$(Q P) \psi(\omega^k) = \omega^k P \psi(\omega^k) = \omega^k \psi(\omega^{k+1})$$

$$\uparrow (P Q) \psi(\omega^k) = Q \psi(\omega^{k-1}) = \omega^{k-1} \psi(\omega^{k-1})$$

$$\Rightarrow Q P = \omega P Q$$

"Fourier transform" to "plane waves"

$$\text{take } \psi_j(\omega^k) = \omega^{jk}$$

$$P \psi_j(\omega^k) = \psi_j(\omega^{k+1}) = \omega^j \psi_j(\omega^k)$$

$$Q \psi_j(\omega^k) = \omega^k \psi_j(\omega^k) = \psi_j(\omega^{k-1})$$

$P \leftrightarrow Q$  switch places.

$$N \rightarrow \infty : \mathbb{Z}_N \rightarrow U(1)$$

## 7. More on group actions

### 7.1. Some defs and s-o theorem

Recall that the group action of  $G$  on a set  $X$ :

$$\phi : G \times X \rightarrow X$$

① left action:  $\phi(g_1, \phi(g_2, x)) = \phi(g_1 g_2, x)$

$$g_1(g_2 \cdot x) = (g_1 g_2) \cdot x$$

(right action  $(x \cdot g_2) \cdot g_1 = x \cdot (g_2 g_1)$ )

②  $\phi(1_G, x) = x$

mention different forms of L & R actions  
and induced actions on  $\mathbb{F}[X \rightarrow Y]$

A  $G$ -action is: see Moore's note

① effective:  $\forall g \neq 1, \exists x, \text{ s.t. } gx \neq x$

(ineffective  $\exists g \neq 1, \forall x, \text{ s.t. } gx = x$ )

② transitive:  $\forall x, y \in X, \exists g, \text{ s.t. } \underline{y = g \cdot x}$

there is only one orbit

③ free:  $\underline{\forall g \neq 1}, \underline{\forall x}, \underline{g \cdot x \neq x}$

### Definitions.

1. isotropy group (stabilizer group)

$$\begin{aligned} \text{Stab}_G(x) &:= \{ g \in G : g \cdot x = x \} \subset G \\ (\equiv G^x) &\left( \begin{array}{l} g_1, g_2 \in G^x \\ g_1 \cdot x = x \quad g_2(g_1 \cdot x) = g_2 x = x \\ g_2 g_1 \in G^x \end{array} \right) \end{aligned}$$

If the group action of  $G$  is free

$$\Leftrightarrow G^x = \{1\} \quad \forall x \in X.$$

2. If  $\exists g \in G^x \neq 1$   $g \cdot x = x$ .  $x$  is called a fixed point.

$$(X^g \equiv) \text{Fix}_X(g) = \{ x \in X : g \cdot x = x \} \subset X$$

is the fixed point set of  $g$ .

$$\text{free} \Leftrightarrow X^g = \emptyset \quad (g \neq 1)$$

$$3. O_G(x) = \{ g \cdot x : \forall g \in G \}$$