

## Electrodynamics Homework II

### 1. Maths behind magnetic dipole moment

(20 marks)

Consider infinitesimal elements  $d\vec{l}'$  of a closed path are located at position vectors  $\vec{r}'$ . The observation point, where we would like to know the magnetic potential, is at point  $\vec{r}$ . Prove the following identity

$$\oint (\hat{r} \cdot \vec{r}') d\vec{l}' = -\hat{r} \times \int d\vec{a},$$

connecting the integral over the closed loop with the area spanned by the loop.

### 2. Another uniqueness theorem

(30 marks)

Prove that electric field is uniquely defined within a volume if we know charge density in this volume as well as potential  $V$  and its normal derivative  $\frac{\partial V}{\partial n}$  on the boundary.

### 3. Something relaxing

(20 marks)

Two positive point charges  $q_A$  and  $q_B$  (masses  $m_A$  and  $m_B$ ) are at rest, held together by a massless string of length  $a$ . Now the string is cut, and the particles fly off in opposite directions. How fast is each one going when they are far apart?

### 4. Computational problem: electric field inside an atom

(30 marks)

Recall quantum description of the hydrogen atom. Stationary proton is located at the origin and the electron is described by a wave function  $\psi_{nlm}(r, \theta, \phi)$ . Assume that the atom is in the ground state  $n = 1, l = 0, m = 0$  and its charge density has two components: positive point charge  $e$  at the origin (proton) and negative charge distribution  $\rho_e(r, \theta, \phi) = -e|\psi_{100}(r, \theta, \phi)|^2$  describing the electron. Explain qualitatively the electric field of this configuration. Then prepare precise plot of electric field as a function of radial distance.

A few comments:

- Please send your answers in a single file, e.g. photos of calculations or a report.
- Deadline: 17 June 2022.
- Plagiarism: I highly encourage you to discuss these topics, by writing down the answers / calculations should be done individually



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Question 1. Show that.

Griffith 1.108

$$\oint (\hat{r} \cdot \vec{r}') d\vec{l} = -\hat{r} \times \int d\vec{a}$$

Here we first prove a lemma.

lemma 1

$$\oint (\vec{c} \cdot \vec{r}) d\vec{l} = \vec{a} \times \vec{c} \quad (1)$$

Here we first focus on LHS, let  $T = \vec{c} \cdot \vec{r}$  <sup>(2)</sup>, next determine  $\vec{\nabla} T$

$$\vec{\nabla} T = \vec{\nabla} (\vec{c} \cdot \vec{r}) = \vec{c} \times (\vec{\nabla} \times \vec{r}) + (\vec{c} \cdot \vec{\nabla}) \vec{r}$$

Here we know that  $\vec{\nabla} \times \vec{r} = 0$ ,  $(\vec{c} \cdot \vec{\nabla}) \vec{r} = \left( (c_x \hat{x} + c_y \hat{y} + c_z \hat{z}) \cdot (\hat{x} \partial_x + \hat{y} \partial_y + \hat{z} \partial_z) \right) (x \hat{x} + y \hat{y} + z \hat{z})$

$$= (c_x \frac{\partial}{\partial x} + c_y \frac{\partial}{\partial y} + c_z \frac{\partial}{\partial z}) (x \hat{x} + y \hat{y} + z \hat{z})$$

$$= c_x \hat{x} + c_y \hat{y} + c_z \hat{z} = \vec{c}$$

$$\therefore \vec{\nabla} T = \vec{c} \quad (3)$$

We substitute (2)(3) into (1), by theorem of curl.

$$\oint T d\vec{l} = - \int (\vec{\nabla} T) \times d\vec{a}$$

$$\begin{aligned} \oint (\vec{c} \cdot \vec{r}) d\vec{l} &= - \int \vec{c} \times d\vec{a} \\ &= - \vec{c} \times \int d\vec{a} \quad (4) \end{aligned}$$

Now we input  $\hat{r} = \vec{c}$ ,  $\vec{r}' = \vec{r}$   $d\vec{l}' = d\vec{l}$

$$\therefore \oint (\hat{r} \cdot \vec{r}') d\vec{l} = - \hat{r} \times \int d\vec{a}$$



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Question 2

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Show that if  $\rho, Q$  in volume,  $V$  or  $\frac{\partial V}{\partial n}$  on boundary known. Griffith 3.5 Problem.

The Electric field uniquely determined.

Here we apply similar method with 2th uniqueness theorem, we assume there exist 2 set of  $\vec{E}$  field satisfied,  $\vec{E}_1, \vec{E}_2$ .

Since we are clear on charge density, therefore, we have Gauss law in  $\vec{\nabla}$  form set

$$\vec{\nabla} \cdot \vec{E}_1 = \vec{\nabla} \cdot \vec{E}_2 = \frac{\rho}{\epsilon_0} \quad (1)$$

And also Gauss law in integral form also set, both inner boundary and outer boundary.

$$\text{inner } \oint \vec{E}_1 \cdot d\vec{a} = \oint \vec{E}_2 \cdot d\vec{a} = \frac{Q_i}{\epsilon_0} \quad (2) \quad \text{outer (total) } \oint \vec{E}_1 \cdot d\vec{a} = \oint \vec{E}_2 \cdot d\vec{a} = \frac{Q_{tot}}{\epsilon_0} \quad (3)$$

Next we let  $\vec{E}_3 = \vec{E}_1 - \vec{E}_2$  as difference of 2 set of  $E$  field

$$\text{From (1)} \quad \vec{\nabla} \cdot \vec{E}_3 = \vec{\nabla} \cdot (\vec{E}_1 - \vec{E}_2) = \frac{\rho}{\epsilon_0} - \frac{\rho}{\epsilon_0} = 0 \quad \Leftrightarrow \quad \vec{\nabla} \cdot \vec{E}_3 = 0$$

$$\text{From (2)(3)} \quad \oint \vec{E}_3 \cdot d\vec{a} = \oint (\vec{E}_1 - \vec{E}_2) \cdot d\vec{a} = \frac{Q - Q}{\epsilon_0} = 0 \quad \Leftrightarrow \quad \oint \vec{E}_3 \cdot d\vec{a} = 0$$

Next we apply the same trick as we do in 2th uniqueness theorem, that make connection  $V_3$  and  $\vec{E}_3$  (Page 123).

$$\vec{\nabla} \cdot (V_3 \vec{E}_3) = V_3 (\vec{\nabla} \cdot \vec{E}_3) + \vec{E}_3 \cdot (\vec{\nabla} V_3) = V_3 (0) - (\vec{E}_3) \cdot (\vec{E}_3) = -|\vec{E}_3|^2 \quad (4)$$





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Next we Integrate in whole volume with (4)

$$\int_V \vec{\nabla} \cdot (V_3 \vec{E}_3) d\tau = \oint_S V_3 \vec{E}_3 \cdot d\vec{a} = - \int_V (E_3)^2 d\tau. \quad (5) \quad \star$$

divergence theorem

Next, we discuss the 2 condition separately

1°  $V$  set, that  $V_1 = V_2$ , that  $V_3 = V_1 - V_2 = 0$

$$\therefore \oint_S V_3 \vec{E}_3 \cdot d\vec{a} = 0 = - \int_V (E_3)^2 d\tau \Leftrightarrow E_3 = 0, \text{ uniqueness exist}$$

2°  $\frac{\partial V}{\partial n}$  set, that  $\frac{\partial V_1}{\partial n} = \frac{\partial V_2}{\partial n}$  that  $\frac{\partial V_3}{\partial n} = \frac{\partial V_1 - V_2}{\partial n} = 0$

that.  $\vec{E}_\perp = 0$  that  $\vec{E} \cdot \vec{a} = 0$

$$\therefore \oint_S V_3 \vec{E}_3 \cdot d\vec{a} = 0 = - \int_V (E_3)^2 d\tau \Leftrightarrow E_3 = 0 \text{ uniqueness exist}$$



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3. 2 point charge.  $q_A, q_B$ , mass  $m_A, m_B$ , length  $a$ , rod cut final speed.

The initial potential  $V_0 = \frac{1}{4\pi\epsilon_0} \frac{q_A q_B}{a}$ , the final potential  $V_\infty \rightarrow 0$ .

$$\begin{cases} \Delta V = V_0 = \frac{1}{4\pi\epsilon_0} \frac{q_A q_B}{a} = \frac{1}{2} m_A V_A^2 + \frac{1}{2} m_B V_B^2 \\ 0 = m_A V_A + m_B V_B \end{cases}$$

Here conservation of momentum  $m$ , A, B move on the same line, we let  $V_A$  as positive.

$$\begin{aligned} V_A &= \frac{1}{m_A} \sqrt{\frac{2 m_A m_B}{m_A + m_B} \frac{1}{4\pi\epsilon_0} \frac{q_A q_B}{a}} \\ V_B &= -\frac{1}{m_B} \sqrt{\frac{2 m_A m_B}{m_A + m_B} \frac{1}{4\pi\epsilon_0} \frac{q_A q_B}{a}} \end{aligned}$$



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Question 4

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$$\psi_{100} = \frac{1}{\sqrt{4\pi}} \cdot \frac{1}{a_0^{3/2}} \cdot 2 e^{-r/a_0}$$

$$\rho = -e|\psi|^2 = -e \frac{1}{4\pi} \frac{1}{a_0^3} \cdot 4 e^{-2r/a_0}$$

$$\oint \vec{E} \cdot d\vec{\alpha} = (e + \int \rho (r^2 \sin\theta dr d\theta d\phi)) / \epsilon_0$$

$$= \frac{e}{\epsilon_0} \left( 1 - \frac{1}{4\pi} \frac{1}{a_0^3} \cdot 4 \int e^{-2r/a_0} r^2 dr \cdot \cancel{2} \cdot 2\pi \right)$$

$$= \frac{e}{\epsilon_0} \left( 1 - \frac{4}{a_0^3} \int_0^r e^{-2r/a_0} r^2 dr \right)$$

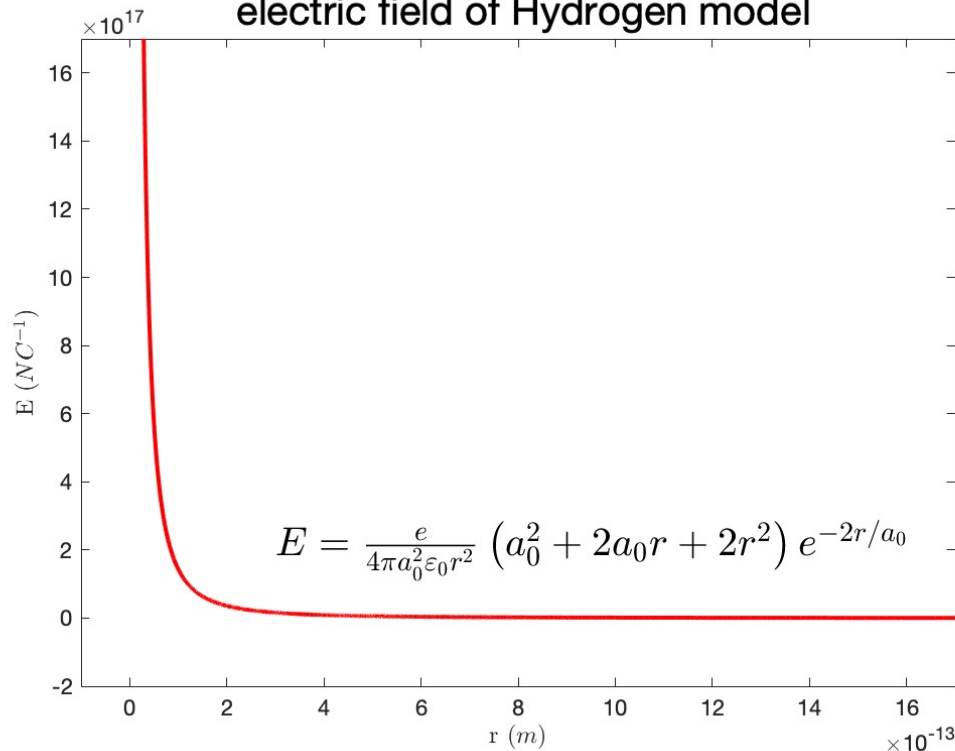
$$= \frac{e}{\epsilon_0} \left( 1 - \frac{4}{a_0^3} \cdot \frac{1}{4} (a_0^3 - a_0^2 e^{-2r/a_0} (a_0^2 + 2a_0 r + 2r^2)) \right)$$

$$= \frac{e}{\epsilon_0} \left( 1 - 1 + \frac{1}{a_0^2} (a_0^2 + 2a_0 r + 2r^2) e^{-2r/a_0} \right)$$

$$E \cdot 4\pi r^2 = \frac{e}{a_0^2 \epsilon_0} (a_0^2 + 2a_0 r + 2r^2) e^{-2r/a_0}$$

$$E = \frac{e}{4\pi a_0^2 \epsilon_0 r^2} (a_0^2 + 2a_0 r + 2r^2) e^{-2r/a_0}$$

The electric field image with distance from center of electric field of Hydrogen model



The electric field goes to positive infinity as the distance from the proton point decreases. And it rapidly decreases to zero as the radius increases, because the hydrogen atom is neutral, therefore when the distance goes to infinity, the total electric field tends to zero.

The code and figure are attached in github. [https://github.com/yimingio/Lecture\\_Electrodynamics/tree/main/Homework\\_2](https://github.com/yimingio/Lecture_Electrodynamics/tree/main/Homework_2)