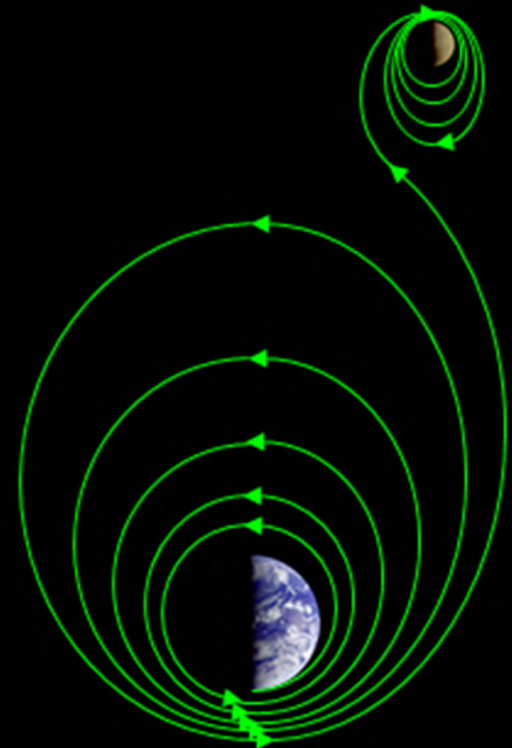


Advanced Topics in Astrodynamics and Trajectory Design

# Continuous Thrust Transfers

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MSc in Astronautics and Space Engineering



# Course Overview

<b>Session 1</b>	Seminar	Lambert Arc Intro
	Workshop	Minimum Energy Transfer Algorithm and F & G Coefficients
<b>Session 2</b>	Seminar	Lambert Arc Session 2
	Workshop	Build your own Lambert-solver
<b>Session 3</b>	Seminar	Pork Chop Plots
	Workshop	Full Earth-Mars Transfer Analysis
<b>Session 4</b>	Seminar	Low Thrust Trajectories
	Workshop	Sub-optimal control law transfers
<b>Session 5</b>	Seminar	CR3BP & LPOs
	Workshop	

# Bibliography

## Some good books ordered by increasing depth and breath :

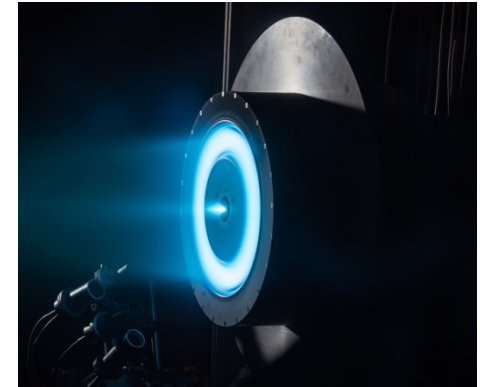
- *Orbital Mechanics*, V.A Chobotov, AIAA Education Series
- *Analytical Mechanics of Space Systems*, H.Shcaub, J.L.Junkins, AIAA Education Series
- *An Introduction to the Mathematics and Methods of Astrodynamics*, R.H.Battin, AIAA Educationan Series
- *Fundamentals of Astrodynamics and Applications*, D.A.Vallado, Space Technology Library
- **Biesbroek, Robin. *Lunar and Interplanetary Trajectories*. Springer International Publishing, 2016.**

# Continuous Thrust Transfers

- Non-Impulsive  $\Delta v$
- Problem Formulation
- Gravity Losses
- Very low thrust hence just a perturbation
- Method of variation of parameters

# Non-Impulsive Manoeuvres

- All space manoeuvres require actually finite burns.
  - ✓ Specially! If we model a electric propulsion system.



- Compute how long would the burn be to achieve escape velocity from a circular 250 km altitude parking orbit if the spacecraft was to use apogee motor such as the one in the picture. Assume a spacecraft with a wet mass of 2000 kg.



- $I_{sp}=320$  s
- Thrust 425 N

Socrative: Connect into Room: **GPQK5UNSK**

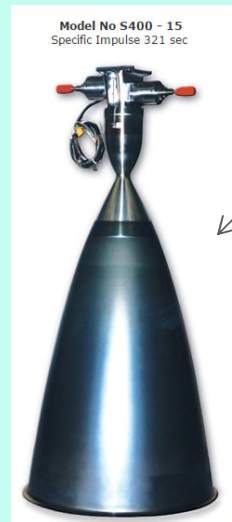
Compute how long would the burn be to achieve escape velocity from a circular 250 km altitude parking orbit if the spacecraft was to use apogee motor such as the one in the picture. Assume a spacecraft with a wet mass of 2000 kg.

Constants:

$R_E = 6378$  km – Radius of the Earth

$\mu_E = 3.986 \times 10^5$  km<sup>3</sup>/s<sup>2</sup> - The gravitational parameter of the Earth

- A) 16 sec
- B) 4.2 h
- C) 2.6 h
- D) <1 sec



- $I_{sp} = 320$  s
- Thrust 425 N

# Low Thrust Transfer: Reminders

- Compute how long would the burn be to achieve escape velocity from a circular 250 km altitude parking orbit if the spacecraft was to use apogee motor such as the one in the picture. Assume a spacecraft with a wet mass of 2000 kg.



- $I_{sp}=320$  s
- Thrust 425 N

$$T = \dot{m}v_e \quad I_{sp} = \frac{T}{\dot{m}g_0} = \frac{v_e}{g_0}$$

$$a_0 = \frac{T}{m_0}$$

$$a(t) = \frac{T}{m(t)}$$

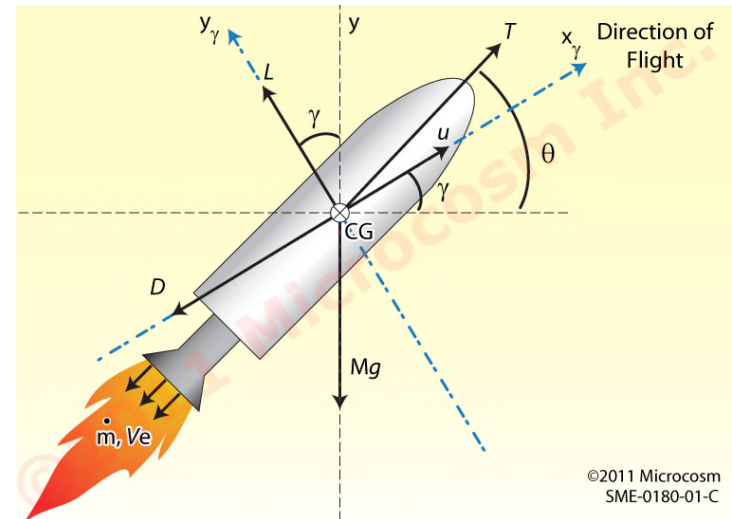


Fig. 18-2. Forces Acting on a Rocket as it Flies.

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SME-0180-01-C

# Low Thrust Transfer: Reminders

- Compute how long would the burn be to achieve escape velocity from a circular 250 km altitude parking orbit if the spacecraft was to use



- Isp=320 s
- Thrust 425 N

apogee motor such as the one in the picture. Assume a spacecraft with a wet mass of 2000 kg.

$$a(t) = \frac{T}{m(t)}$$

$$\Delta v_{acc}(t) = \int_0^t a(t) \cdot dt$$

$$m(t) = m_0 + \int_0^t \dot{m} \cdot dt = m_0 + \dot{m}t$$

$$\Delta v_{acc}(t) = \int_0^t \frac{T}{m_0 + \dot{m}t} \cdot dt$$

$$\Delta v_{acc}(t) = \frac{T}{m_0} \int_0^t \frac{1}{1 + \frac{\dot{m}}{m_0}t} \cdot dt$$



# Low Thrust Transfer: Reminders

- Compute how long would the burn be to achieve escape velocity from a circular 250 km altitude parking orbit if the spacecraft was to use



- $I_{sp}=320$  s
- Thrust 425 N

apogee motor such as the one in the picture. Assume a spacecraft with a wet mass of 2000 kg.

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$$\Delta v_{acc}(t) = \frac{T}{m_0} \int_0^t \frac{1}{1 + \frac{\dot{m}}{m_0}t} \cdot dt$$

$$\int \frac{dx}{ax+b} = \frac{1}{a} \ln(ax+b)$$

$$\Delta v_{acc}(t) = \frac{T}{\dot{m}} \ln \left( 1 + \frac{\dot{m}}{m_0} t \right)$$

## Low Thrust Transfer: Reminders

- Compute how long would the burn be to achieve escape velocity from a circular 250 km altitude parking orbit if the spacecraft was to use



- $I_{sp}=320$  s
- Thrust 425 N

apogee motor such as the one in the picture. Assume a spacecraft with a wet mass of 2000 kg.

$$\Delta v_{acc}(t) = \frac{T}{\dot{m}} \ln \left( 1 + \frac{\dot{m}}{m_0} t \right)$$

$$\Delta t = \frac{m_0}{\dot{m}} \left( e^{\frac{\Delta v \cdot \dot{m}}{T}} - 1 \right) \approx 10,000 \text{ sec}$$

# Low Thrust Engines

- The example above was for a typical “high” thrust engine.
- Low thrust engines have thrust typically  $< 1\text{N}$ . However, they are the most efficient engines available.

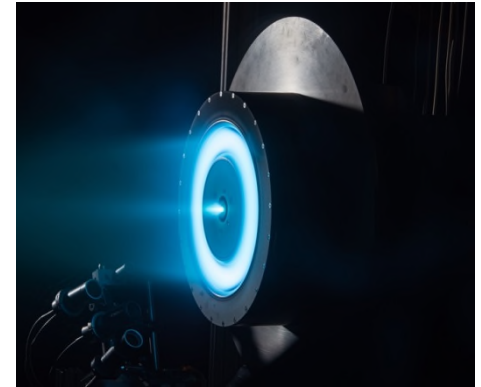
Table extracted from Biesbroek, R. (2016). Lunar and interplanetary trajectories.

**Table 1.3** Overview of engine types and typical values for the specific impulse

Engine type	Typical specific impulse Isp [s]	Applications
Cold-gas	50	Spacecraft attitude control
Mono-propellant	220	Attitude control or spacecraft main engine for all orbit transfers
Solid	250	Launchers and rocket stages
Bi-propellant	320	Larger attitude control or spacecraft main engine for all orbit transfers. More effective than mono-propellant, but more complex too
Cryogenic	450	Launchers (cryogenic engines can only be used until a few days after launch)
Hall-effect (HET)	1650	Deep-space maneuvers/spiraling out to the Moon
Ion	3000	Deep-space maneuvers

# Non-Impulsive Manoeuvres

- All space manoeuvres require actually finite burns.
  - ✓ Specially! If we model a electric propulsion system.



Equation of Motion

$$\ddot{\mathbf{r}} = -\frac{\mu_E}{r^3}\mathbf{r} + \mathbf{a}$$



No closed form solution

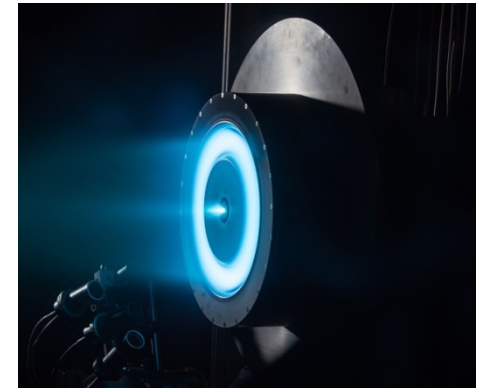
- Low thrust transfers generally require some very complex analysis (e.g., optimal control theories).
- However, some approximations can be done that provide a working solution (e.g. sub-optimal control laws, shape-based methods).

# Electric Propulsion Fundamentals

- The propellant is converted into charged particles or plasma, and these are accelerated using electrostatic or electromagnetic forces.
- The propellant is ejected at exhaust velocities about 10 times greater than those of conventional chemical rockets.
- However, the propellant mass-flow is extremely small, and thus the thrust is also very small.

$$T = \dot{m}v_e$$

- Consequently, the Electric Propulsion system must operate nearly continuously so that the low acceleration produces noticeable changes in the trajectory.



# Electric Propulsion Fundamentals

$$T = \dot{m}v_e$$

$$v_e = g_0 \cdot I_{sp}$$

$$P_{out} = \frac{1}{2} \dot{m}v_e^2$$

Thruster efficiency  $\eta = \frac{P_{out}}{P_{in}}$

$$\eta = \frac{P_{out}}{P_{in}} = \frac{\frac{1}{2} \dot{m}v_e^2}{P_{in}} = \frac{Tv_e}{2P_{in}}$$



$$T = \frac{2\eta P_{in}}{g_0 I_{sp}}$$

# Low Thrust Transfer: Problem Formulation

- Let  $\mathbf{X}_0$  be an initial state vector such as  $(x_0, y_0, z_0, v_{x0}, v_{y0}, v_{z0})$ , and  $\mathbf{u}(t)$  the force vector provided by our propulsion system.
- $\mathbf{u}(t)$  is generally referred as *control inputs*.
- Our spacecraft is subject to 2-body motion, thus:

Equation of Motion

$$\ddot{\mathbf{r}} = -\frac{\mu_E}{r^3} \mathbf{r} + \mathbf{a}$$

$$\mathbf{a} = \frac{\mathbf{u}}{m}$$

# Low Thrust Transfer: Problem Formulation

- Let  $\mathbf{X}_0$  be an initial state vector such as  $(x_0, y_0, z_0, v_{x0}, v_{y0}, v_{z0})$ , and  $\mathbf{u}(t)$  the force vector provided by our propulsion system.
- $\mathbf{u}(t)$  is generally referred as *control inputs*.
- Our spacecraft is subject to 2-body motion, thus:

$$\ddot{x} = -\frac{\mu_E}{r^3} x + a_x$$

$$\ddot{y} = -\frac{\mu_E}{r^3} y + a_y$$

$$\ddot{z} = -\frac{\mu_E}{r^3} z + a_z$$



# Low Thrust Transfer: Problem Formulation

- Let  $\mathbf{X}_0$  be an initial state vector such as  $(x_0, y_0, z_0, v_{x0}, v_{y0}, v_{z0})$ , and  $\mathbf{u}(t)$  the force vector provided by our propulsion system.
- $\mathbf{u}(t)$  is generally referred as *control inputs*.
- Our spacecraft is subject to 2-body motion, thus:

$$\frac{dx}{dt} = v_x$$

$$\frac{dy}{dt} = v_y$$

$$\frac{dz}{dt} = v_z$$

$$\frac{dv_x}{dt} = -\frac{\mu_E}{r^3} x + a_x$$

$$\frac{dv_y}{dt} = -\frac{\mu_E}{r^3} y + a_y$$

$$\frac{dv_z}{dt} = -\frac{\mu_E}{r^3} z + a_z$$

# Low Thrust Transfer: Problem Formulation

- Let  $\mathbf{X}_0$  be an initial state vector such as  $(x_0, y_0, z_0, v_{x0}, v_{y0}, v_{z0})$ , and  $\mathbf{u}(t)$  the force vector provided by our propulsion system.
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$$\frac{dv_x}{dt} = -\frac{\mu_E}{r^3} x + a_x$$

$$\frac{dv_y}{dt} = -\frac{\mu_E}{r^3} y + a_y$$

$$\frac{dv_z}{dt} = -\frac{\mu_E}{r^3} z + a_z$$

$$\mathbf{a} = \frac{\mathbf{u}}{m} \quad \dot{m} = \frac{T}{v_e}$$

$$\frac{dm}{dt} = \frac{T}{v_e} \quad \longrightarrow \quad (x_0, y_0, z_0, v_{x0}, v_{y0}, v_{z0}, m_0)$$

## Low Thrust Transfer: Problem Formulation

- Let  $\mathbf{X}_0$  be an initial state vector such as  $(x_0, y_0, z_0, v_{x0}, v_{y0}, v_{z0})$ , and  $\mathbf{u}(t)$  the force vector provided by our propulsion system.
- $\mathbf{u}(t)$  is generally referred as *control inputs*.
- Our spacecraft is subject to 2-body motion, thus:

$$\ddot{\mathbf{r}} = -\frac{\mu_E}{r^3} \mathbf{r} + \frac{\mathbf{u}(t)}{m}$$

- Solving the optimal control problem requires us to find the *control history*  $\mathbf{u}(t)$  generally subject to a series of *constraints* and *boundary conditions*.
- This is out of the scope of this course (but some notions are provided).

# Trajectory Optimization

- The standard form of an optimization problem is:

$$\begin{aligned} & \underset{x}{\text{minimize}} && f(x) \\ & \text{subject to} && g_i(x) \leq 0, \quad i = 1, \dots, m \\ & && h_j(x) = 0, \quad j = 1, \dots, p \end{aligned}$$

- ✓ The **objective function**  $f(x): \mathbb{R}^n \rightarrow \mathbb{R}$
- ✓ Subject to **inequality** and **equality constraints**.
- **Nonlinear programming** (NLP) is generally referred to the process of solving an **optimization problem**.

# Optimal Control Problem for a continuous thrust transfer

- NLP provides a finite set of variables  $\mathbf{x}$  that minimize a given objective function.
- Instead, an optimal control problem involves finding the continuous function  $\mathbf{u}(t)$  that minimizes the objective function.
- Since solving an infinite set of variables becomes rather cumbersome, the general approach is to **transcribe** or convert the *infinite-dimensional optimal control* problem into a *finite-dimensional NLP*.

# Optimal Control Problem for a continuous thrust transfer

- NLP provides a finite set of variables  $\mathbf{x}$  that minimize an objective func.
- An optimal control problem involves finding the continuous function  $\mathbf{u}(t)$ .
- **Transcribe** the *infinite-dimensional optimal control* problem into a *finite-dimensional NLP*.

Gradient-based methods

$$f = c(x) \quad c(x^*) = 0$$

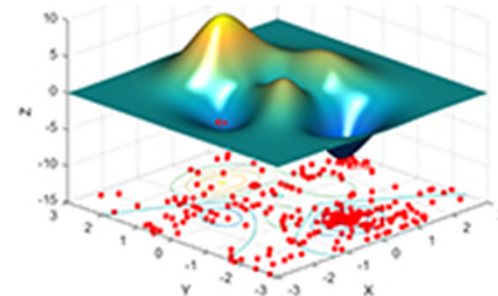
$$c(\bar{x}) = c(x) + \frac{dc}{dx}(\bar{x} - x)$$

$$c(\bar{x}) = c(x^*) = 0$$

$$0 = c(x) + \frac{dc}{dx}(\bar{x} - x)$$

$$\bar{x} = x - \left[ \frac{dc}{dx} \right]^{-1} c(x)$$

Heuristic methods



A **heuristic** technique is any **approach** to problem solving that employs a practical method not guaranteed to be optimal (e.g. genetic algorithms, Particle Swarm Optimization, etc)

# Low Thrust Transfer: Sub-optimal Control Laws.

## Lagrange's Planetary Equations

*Gaussian form*

$$\frac{da}{dt} = \frac{2a^2 v}{\mu} a_t$$

$$\frac{de}{dt} = \frac{1}{v} \left[ 2(e + \cos f) a_t - \frac{r}{a} \sin f a_n \right]$$

$$\frac{di}{dt} = \frac{r \cos \theta}{h} a_h$$

$$\frac{d\Omega}{dt} = \frac{r \sin \theta}{h \sin i} a_h$$

$$\frac{d\omega}{dt} = \frac{1}{ev} \left[ 2 \sin f a_t + \left( 2e + \frac{r}{a} \cos f \right) a_n \right] - \frac{r \sin \theta \cos i}{h \sin i} a_h$$

$$\frac{dM}{dt} = n - \frac{b}{eav} \left[ 2 \left( 1 + \frac{e^2 r}{p} \right) \sin f a_t + \frac{r}{a} \cos f a_n \right]$$

$f \equiv$  true anomaly

$\vartheta \equiv$  the argument of latitude (i.e.,  $\omega + f$ )

The acceleration vector  $(a_t, a_n, a_h)$  is expressed in a Cartesian rotating frame  $(i_t, i_n, i_h)$ , where  $i_t$  is the direction along the orbit velocity vector,  $i_h$  is parallel to the orbital momentum vector, thus the out-of-plane direction, and  $i_n$  completes the right-hand coordinates system

## Low Thrust Transfer: Sub-optimal Control Laws.

What then would be the *best* control law to achieve a escape trajectory?

$$\frac{da}{dt} = \frac{2a^2 v}{\mu} a_t$$

$$\frac{de}{dt} = \frac{1}{v} \left[ 2(e + \cos f) a_t - \frac{r}{a} \sin f a_n \right]$$

$$\frac{di}{dt} = \frac{r \cos \theta}{h} a_h$$

$$\frac{d\Omega}{dt} = \frac{r \sin \theta}{h \sin i} a_h$$

$$\frac{d\omega}{dt} = \frac{1}{ev} \left[ 2 \sin f a_t + \left( 2e + \frac{r}{a} \cos f \right) a_n \right] - \frac{r \sin \theta \cos i}{h \sin i} a_h$$

$$\frac{dM}{dt} = n - \frac{b}{eav} \left[ 2 \left( 1 + \frac{e^2 r}{p} \right) \sin f a_t + \frac{r}{a} \cos f a_n \right]$$

$f \equiv$  true anomaly

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The acceleration vector  $(a_t, a_n, a_h)$  is expressed in a Cartesian rotating frame  $(i_t, i_n, i_h)$ , where  $i_t$  is the direction along the orbit velocity vector,  $i_h$  is parallel to the orbital momentum vector, thus the out-of-plane direction, and  $i_n$  completes the right-hand coordinates system



## Low Thrust Transfer: Simplified Control Scenarios.

- Let  $\mathbf{X}_0$  be an initial state vector such as  $(x_0, y_0, z_0, v_{x0}, v_{y0}, v_{z0}, m_0)$ .

$$\frac{dx}{dt} = v_x$$

$$\frac{dy}{dt} = v_y$$

$$\frac{dz}{dt} = v_z$$

$$\frac{dv_x}{dt} = -\frac{\mu_E}{r^3} x + \frac{T_{\max}}{m} \frac{v_x}{|v|}$$

$$\frac{dv_y}{dt} = -\frac{\mu_E}{r^3} y + \frac{T_{\max}}{m} \frac{v_y}{|v|}$$

$$\frac{dv_z}{dt} = -\frac{\mu_E}{r^3} z + \frac{T_{\max}}{m} \frac{v_z}{|v|}$$

$$\frac{dm}{dt} = -\frac{T_{\max}}{v_e}$$

- Propagate, using Matlab, the escape trajectory described in the previous exercise using a constant thrust in tangential direction.

## Low Thrust Trajectory - Exercise 2

Propagate, using Matlab ODE solvers, the escape trajectory described in the previous exercise using a constant thrust in tangential direction.

Initial Orbit: Circular 250 km altitude parking orbit.

Apogee motor specs: Thrust 425 N, Isp 320s.

Spacecraft wet mass : 2000 kg.

# Low Thrust Transfer: Simplified Control Scenarios.

$$\frac{dx}{dt} = v_x$$

$$\frac{dy}{dt} = v_y$$

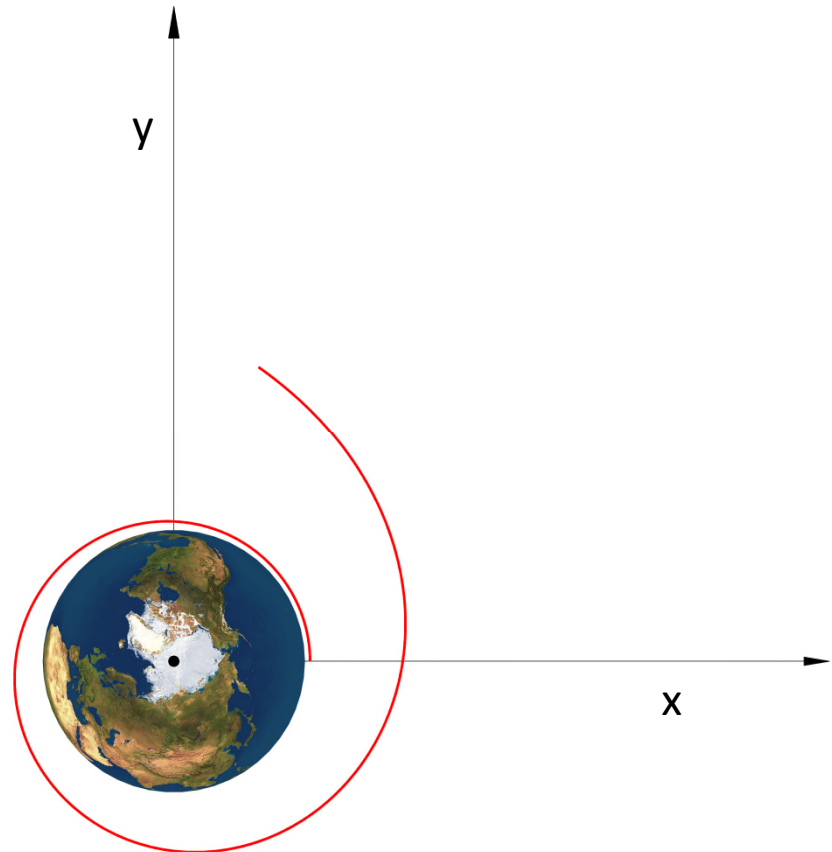
$$\frac{dz}{dt} = v_z$$

$$\frac{dv_x}{dt} = -\frac{\mu_E}{r^3} x + \frac{T_{\max}}{m} \frac{v_x}{|v|}$$

$$\frac{dv_y}{dt} = -\frac{\mu_E}{r^3} y + \frac{T_{\max}}{m} \frac{v_y}{|v|}$$

$$\frac{dv_z}{dt} = -\frac{\mu_E}{r^3} z + \frac{T_{\max}}{m} \frac{v_z}{|v|}$$

$$\frac{dm}{dt} = \frac{T}{v_e}$$



# Low Thrust Transfer: Simplified Control Scenarios.

$$\frac{dx}{dt} = v_x$$

$$\frac{dy}{dt} = v_y$$

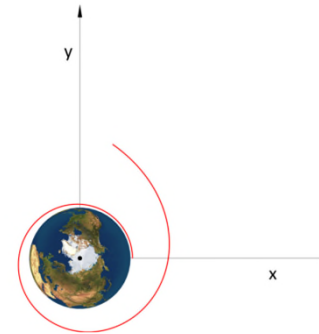
$$\frac{dz}{dt} = v_z$$

$$\frac{dv_x}{dt} = -\frac{\mu_E}{r^3}x + \frac{T_{\max}}{m} \frac{v_x}{|v|}$$

$$\frac{dv_y}{dt} = -\frac{\mu_E}{r^3}y + \frac{T_{\max}}{m} \frac{v_y}{|v|}$$

$$\frac{dv_z}{dt} = -\frac{\mu_E}{r^3}z + \frac{T_{\max}}{m} \frac{v_z}{|v|}$$

$$\frac{dm}{dt} = \frac{T}{v_e}$$



- Is the final state vector in a escape trajectory?

```
%-----
% Is the final state vector in a escape trajectory?
%[T,X,TE,YE,IE]=ode45(F,[0,tf],SV0,options);
SV_final=X(end,:); % final state vector

% Final distance to Earth
rf=sqrt(SV_final(1)^2+SV_final(2)^2+SV_final(3)^2);
% Escape velocity at the final distance
vesc=sqrt(2*muEarth/rf);
% Our final velocity
vf=sqrt(SV_final(4)^2+SV_final(5)^2+SV_final(6)^2);
% is vf>vesc
vf>vesc
% no
%-----
```

## Low Thrust Trajectory - Exercise 2B

Ensure the spacecraft thrust until reaching escape condition.

Initial Orbit: Circular 250 km altitude parking orbit.

Apogee motor specs: Thrust 425 N, Isp 320s.

Spacecraft wet mass : 2000 kg.

Notes: Use MATLAB's ODE event feature to stop the propagation when it reaches escape velocity. For information on how to do this go to:

<https://uk.mathworks.com/help/matlab/math/ode-event-location.html>

- This is the suggested event function:

```
function [value,isterminal,direction]=event_Escape(t,x)
%-----
muEarth = getAstroConstants('Earth','mu');
%-----
% Current velocity
v=sqrt(x(4)^2+x(5)^2+x(6)^2);
% Escape velocity at altitude
vesc=sqrt(2*muEarth/sqrt(x(1)^2+x(2)^2+x(3)^2));

value = v-vesc; % when s/c reaches scape velocity value will be zero
isterminal = 1;
direction = 0;
end
```

# Low Thrust Transfer: Simplified Control Scenarios.

$$\frac{dx}{dt} = v_x$$

$$\frac{dy}{dt} = v_y$$

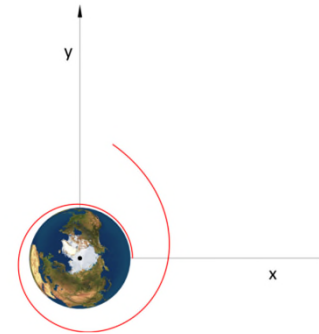
$$\frac{dz}{dt} = v_z$$

$$\frac{dv_x}{dt} = -\frac{\mu_E}{r^3}x + \frac{T_{\max}}{m} \frac{v_x}{|v|}$$

$$\frac{dv_y}{dt} = -\frac{\mu_E}{r^3}y + \frac{T_{\max}}{m} \frac{v_y}{|v|}$$

$$\frac{dv_z}{dt} = -\frac{\mu_E}{r^3}z + \frac{T_{\max}}{m} \frac{v_z}{|v|}$$

$$\frac{dm}{dt} = -\frac{T}{v_e}$$



- How much is the gravity loss generated if we plan to thrust until escape condition without apogee motor?

$$\Delta v_{\text{GravityLoss}} \approx 1.5 \text{ km/s}$$

## Earth-Moon Smart-1 Trajectory

Smart-1 spacecraft (367 kg) is launched into a 250 km circular altitude orbit. Smart-1 propulsion system uses a Hall-effect Plasma Thruster with  $I_{sp}$  of 3,000s and thruster efficiency of  $\eta=0.7$ . The input power is of 1190 W. What would it be the total DV to reach Lunar altitude? And what the propellant mass used?



[www.cranfield.ac.uk](http://www.cranfield.ac.uk)