

## 2 RTPs on a ring

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### I. MODEL

#### A. Fokker-Planck equations

We consider 2 run-and-tumble particles (RTPs), with swim speed  $v_0$  and tumble rate  $\tau^{-1}$  on a linear ring of length  $L$ . With  $r_1, r_2 \in [0, L]$  the positions and  $\alpha_1, \alpha_2 = \pm$  the states of each particles, and  $V(r_1, r_2) = V(|r_1 - r_2|)$  the potential of interactions between these particles, we have the equations of motion

$$\dot{r}_i = \alpha_i v_0 - \partial_{r_i} V(|r_1 - r_2|) \quad (1)$$

and the Fokker-Planck equations for the joint distribution of positions

$$\begin{aligned} \dot{P}_{\alpha_1, \alpha_2}(r_1, r_2) = & -v_0(\alpha_1 \partial_{r_1} + \alpha_2 \partial_{r_2}) P_{\alpha_1, \alpha_2}(r_1, r_2) && \rightarrow \text{Propulsion} \\ & + \partial_{r_1} (P_{\alpha_1, \alpha_2}(r_1, r_2) \partial_{r_1} V(|r_1 - r_2|)) && \rightarrow \text{Interaction} \\ & + \partial_{r_2} (P_{\alpha_1, \alpha_2}(r_1, r_2) \partial_{r_2} V(|r_1 - r_2|)) \\ & + \tau^{-1} (P_{\bar{\alpha}_1 \alpha_2}(r_1, r_2) + P_{\alpha_1 \bar{\alpha}_2}(r_1, r_2) - 2P_{\alpha_1 \alpha_2}(r_1, r_2)) && \rightarrow \text{Tumble} \\ = & \mathcal{L}_{\alpha_1 \alpha_2}^\dagger P_{\alpha_1 \alpha_2}(r_1, r_2) \end{aligned} \quad (2)$$

with  $\mathcal{L}_{\alpha_1 \alpha_2}^\dagger$  the Fokker-Planck operator.

We introduce  $r \equiv r_2 - r_1 \geq 0$ , such that  $\partial_{r_1} \equiv -\partial_r$  and  $\partial_{r_2} = \partial_r$ , with which we can define the symmetries

$$P_{++}(r) = P_{--}(r), \quad (3)$$

$$P_{+-}(r) = P_{-+}(L - r), \quad (4)$$

and the interaction potential

$$V(r) = \begin{cases} 0 & \text{if } r \in ]0, L[ \\ \infty & \text{otherwise} \end{cases}$$

such that, with Eq. (1), we get for particles with opposite directions

$$-\partial_r V(0) = \partial_r V(L) = v_0. \quad (5)$$

because of mutual hindrance.

We can solve the steady state problem with the normalisation condition

$$\int_0^L dr \sum_{\alpha_1 \alpha_2} P_{\alpha_1 \alpha_2}(r) = 1 \quad (6)$$

the bulk equations

$$\dot{P}_{\alpha_1 \alpha_2}(r \in ]0, L[) = 0 \quad (7)$$

the left boundary equations

$$\int_{0^-}^{0^+} dr \dot{P}_{\alpha_1 \alpha_2}(r) = 0 \quad (8)$$

and the right boundary equations

$$\int_{L^-}^{L^+} dr \dot{P}_{\alpha_1 \alpha_2}(r) = 0. \quad (9)$$

## B. Biased ensemble

We define our biasing quantity

$$\dot{Z} = \sum_{i=1}^2 -\partial_{r_i} V(r) \alpha_i v_0 = v_0 (\alpha_1 - \alpha_2) \partial_r V(r) \quad (10)$$

which is non-zero only at contact, and the tilted generator

$$\mathcal{W}_{\alpha_1 \alpha_2}^\dagger = \mathcal{L}_{\alpha_1 \alpha_2}^\dagger - s \dot{Z} \quad (11)$$

with  $s$  the biasing parameter, which eigenvalue problem

$$\psi P_{\alpha_1 \alpha_2}(r) = \mathcal{W}_{\alpha_1 \alpha_2}^\dagger P_{\alpha_1 \alpha_2}(r) \quad (12)$$

we have to solve for  $\psi(s)$  its largest eigenvalue, which defines the dynamical free energy.

While the symmetries (Eqs. (3, 4)) and the normalisation condition (Eq. (6)) remains unchanged, we have that the bulk equations (Eq. (7)), the left boundary equations (Eq. 8), and the right boundary equations (Eq. 9) become

$$\psi P_{\alpha_1 \alpha_2}(r \in ]0, L[) = \mathcal{W}_{\alpha_1 \alpha_2}^\dagger P_{\alpha_1 \alpha_2}(r \in ]0, L[) \quad (13)$$

$$\int_{0^-}^{0^+} dr \psi P_{\alpha_1 \alpha_2}(r) = \int_{0^-}^{0^+} dr \mathcal{W}_{\alpha_1 \alpha_2}^\dagger P_{\alpha_1 \alpha_2}(r) \quad (14)$$

$$\int_{L^-}^{L^+} dr \psi P_{\alpha_1 \alpha_2}(r) = \int_{L^-}^{L^+} dr \mathcal{W}_{\alpha_1 \alpha_2}^\dagger P_{\alpha_1 \alpha_2}(r) \quad (15)$$

respectively.

## II. UNBIASED STEADY STATE DISTRIBUTION

We use the following ansatz for the unbiased ( $s = 0$ ) steady state distribution

$$P_{\alpha_1 \alpha_2}(r) = a_{\alpha_1 \alpha_2} + b_{\alpha_1 \alpha_2} \delta(r) + c_{\alpha_1 \alpha_2} \delta(L - r) \quad (16)$$

which satisfies the normalisation condition (Eq. (6))

$$\sum_{\alpha_1 \alpha_2} (L a_{\alpha_1 \alpha_2} + b_{\alpha_1 \alpha_2} + c_{\alpha_1 \alpha_2}) = 1 \quad (17)$$

the bulk equations (Eq. (7))

$$\tau^{-1}(a_{+-} + a_{-+} - 2a_{++}) = 0 \quad (18)$$

$$\tau^{-1}(a_{+-} + a_{-+} - 2a_{--}) = 0 \quad (19)$$

$$\tau^{-1}(a_{++} + a_{--} - 2a_{+-}) = 0 \quad (20)$$

$$\tau^{-1}(a_{++} + a_{--} - 2a_{-+}) = 0 \quad (21)$$

the left boundary equations (Eq. (8))

$$\tau^{-1}(b_{+-} + b_{-+} - 2b_{++}) = 0 \quad (22)$$

$$\tau^{-1}(b_{+-} + b_{-+} - 2b_{--}) = 0 \quad (23)$$

$$2v_0 a + \tau^{-1}(b_{++} + b_{--} - 2b_{+-}) = 0 \quad (24)$$

$$-2v_0 a + \tau^{-1}(b_{++} + b_{--} - 2b_{-+}) = 0 \quad (25)$$

the right boundary equations (Eq. (9))

$$\tau^{-1}(c_{+-} + c_{-+} - 2c_{++}) = 0 \quad (26)$$

$$\tau^{-1}(c_{+-} + c_{-+} - 2c_{--}) = 0 \quad (27)$$

$$-2v_0 a + \tau^{-1}(c_{++} + c_{--} - 2c_{+-}) = 0 \quad (28)$$

$$2v_0 a + \tau^{-1}(c_{++} + c_{--} - 2c_{-+}) = 0 \quad (29)$$

and the symmetries (Eqs. (3, 4)).

We get

$$P_{++}(r) = a + la\delta(r) + la\delta(L - r) \quad (30)$$

$$P_{--}(r) = a + la\delta(r) + la\delta(L - r) \quad (31)$$

$$P_{+-}(r) = a + 2la\delta(r) \quad (32)$$

$$P_{-+}(r) = a + 2la\delta(L - r) \quad (33)$$

with

$$a = \frac{1}{4(L + 2l)} \quad (34)$$

and the persistence length  $l = v_0\tau$ .

### III. BIASED STEADY STATE DISTRIBUTION

#### A. Exact solution

We use the following ansätze for the steady state distribution

$$P_{++}(r) = P_{--}(r) = \beta(r) + \gamma_-\delta(r) + \gamma_+\delta(L - r) \quad (35)$$

$$P_{+-}(r) = \varepsilon(r) + \zeta\delta(r) \quad (36)$$

$$P_{-+}(r) = \theta(r) + \zeta\delta(L - r) \quad (37)$$

$$(38)$$

with  $\varepsilon(r) = \theta(L - r)$  according to Eq. (4).

We write the eigenvalue equations (Eq. (12))

$$\psi P_{\alpha\alpha}(r) = 2\partial_r(P_{\alpha\alpha}(r)\partial_r V(r)) + \tau^{-1}(P_{+-}(r) + P_{-+}(r) - 2P_{\alpha\alpha}(r)) \quad (39)$$

$$\psi P_{+-}(r) = 2v_0\partial_r P_{+-}(r) + 2\partial_r(P_{+-}(r)\partial_r V(r)) + \tau^{-1}(2P_{\alpha\alpha}(r) - 2P_{+-}(r)) - 2sv_0\partial_r V(r)P_{+-}(r) \quad (40)$$

$$\psi P_{-+}(r) = -2v_0\partial_r P_{-+}(r) + 2\partial_r(P_{-+}(r)\partial_r V(r)) + \tau^{-1}(2P_{\alpha\alpha}(r) - 2P_{-+}(r)) + 2sv_0\partial_r V(r)P_{-+}(r) \quad (41)$$

and integrate from Eqs. (39, 40, 41) the sum of  $\psi(2P_{\alpha\alpha} + P_{+-} + P_{-+})$  between  $0^-$  and  $L^+$  to get

$$\psi = 4sv_0\partial_r V(L)\zeta = 4sv_0^2\zeta \quad (42)$$

linking the dynamical free energy  $\psi$  and the sticking term  $\zeta$ .

We have the bulk equations (Eq. (13))

$$\psi\beta(r) = \tau^{-1}(\varepsilon(r) + \theta(r) - 2\beta(r)) \quad (43)$$

$$\psi\epsilon(r) = 2v_0\epsilon'(r) + \tau^{-1}(2\beta(r) - 2\varepsilon(r)) \quad (44)$$

$$\psi\theta(r) = -2v_0\theta'(r) + \tau^{-1}(2\beta(r) - 2\theta(r)) \quad (45)$$

such that Eq. (43) gives

$$\beta(r) = (2 + \tau\psi)^{-1}(\epsilon(r) + \theta(r)) \quad (46)$$

and the difference and sum of Eqs. (44, 45) give

$$2l(\varepsilon(r) + \theta(r))' = (2 + \tau\psi)(\varepsilon(r) - \theta(r)) \quad (47)$$

$$\begin{aligned} 2l(\varepsilon(r) - \theta(r))' &= -4\beta(r) + (\tau\psi + 2)(\varepsilon(r) + \theta(r)) \\ &= \left( (\tau\psi + 2) - \frac{4}{\tau\psi + 2} \right) (\varepsilon(r) + \theta(r)) \end{aligned} \quad (48)$$

where we set  $A(r) = \varepsilon(r) + \theta(r)$  and  $B(r) = \varepsilon(r) - \theta(r)$ , which on the one hand verify

$$A''(r) - k^2 A(r) = 0 \quad (49)$$

where

$$k^2 l^2 = \tau \psi \left( \frac{\tau \psi}{4} + 1 \right) \quad (50)$$

and which general solution is

$$A(r) = A_+ e^{-kr} + A_- e^{-k(L-r)} \quad (51)$$

and on the other hand

$$B(r) = l(1 + \tau \psi/2)^{-1} A'(r) = kl(1 + \tau \psi/2)^{-1} (A_- e^{-k(L-r)} - A_+ e^{-kr}) \quad (52)$$

from which we infer

$$\begin{aligned} \varepsilon(r) &= \frac{1}{2} (A(r) + B(r)) \\ &= \frac{1}{2} (1 - kl(1 + \tau \psi/2)^{-1}) A_+ e^{-kr} + \frac{1}{2} (1 + kl(1 + \tau \psi/2)^{-1}) A_- e^{-k(L-r)} \end{aligned} \quad (53)$$

$$\begin{aligned} \theta(r) &= \frac{1}{2} (A(r) - B(r)) \\ &= \frac{1}{2} (1 + kl(1 + \tau \psi/2)^{-1}) A_+ e^{-kr} + \frac{1}{2} (1 - kl(1 + \tau \psi/2)^{-1}) A_- e^{-k(L-r)} \end{aligned} \quad (54)$$

where we note that the symmetry condition of Eq. (4) implies that  $A_- = A_+$ , and

$$\begin{aligned} \beta(r) &= (\tau \psi + 2)^{-1} (\varepsilon(r) + \theta(r)) = (\tau \psi + 2)^{-1} A(r) \\ &= \frac{1}{2} (1 + \tau \psi/2)^{-1} (A_+ e^{-kr} + A_- e^{-k(L-r)}) \end{aligned} \quad (55)$$

where we need to determine  $A_+ = A_-$ .

We have the left and right boundary equations (Eqs. (14, 15)) for  $P_{\alpha\alpha}$

$$\psi \gamma_- = \tau^{-1} (\zeta - 2\gamma_-) \quad (56)$$

$$\psi \gamma_+ = \tau^{-1} (\zeta - 2\gamma_+) \quad (57)$$

such that

$$\gamma = \gamma_- = \gamma_+ = (\tau \psi + 2)^{-1} \zeta \quad (58)$$

which with Eq. (42) gives

$$\psi = 4sv_0^2 (\tau \psi + 2) \gamma \quad (59)$$

where we need to determine  $\gamma$ .

We have the left boundary equations (Eq. (14)) for  $P_{\alpha\bar{\alpha}}$

$$\psi \zeta = 2v_0 \varepsilon(0^+) + \tau^{-1} (2\gamma - 2\zeta) - 2sv_0 \partial_r V(0) \zeta \quad (60)$$

$$0 = -2v_0 \varepsilon(L^-) + \tau^{-1} 2\gamma \quad (61)$$

from which we get

$$[(\tau \psi + 2)(\tau \psi + 2 - 2slv_0) - 2] \gamma - l (1 - kl(1 + \tau \psi/2)^{-1}) A_+ - l (1 + kl(1 + \tau \psi/2)^{-1}) e^{-kL} A_- = 0 \quad (62)$$

$$2\gamma - l (1 - kl(1 + \tau \psi/2)^{-1}) e^{-kL} A_+ - l (1 + kl(1 + \tau \psi/2)^{-1}) A_- = 0 \quad (63)$$

and the normalisation condition (Eq. (6))

$$(2\tau\psi + 8)\gamma + \frac{1}{k}(1 - e^{-kL})(1 + (1 + \tau\psi/2)^{-1})(A_+ + A_-) = 1 \quad (64)$$

so we can solve the system of Eqs. (62, 63, 64) with respect to  $\gamma$ ,  $A_+$ ,  $A_-$ , or equivalently with  $A_+ = A_-$

$$[(\tau\psi + 2)(\tau\psi + 2 - 2slv_0) - 2]\gamma - l[(1 - kl(1 + \tau\psi/2)^{-1}) + (1 + kl(1 + \tau\psi/2)^{-1})e^{-kL}]A_- = 0 \quad (65)$$

$$(2\tau\psi + 8)\gamma + \frac{2}{k}(1 - e^{-kL})(1 + (1 + \tau\psi/2)^{-1})A_+ = 1 \quad (66)$$

with respect to  $A_+$  and  $\gamma$ .

We use SAGEMATH to solve Eqs. (65, 66)

```
# variables
psi = var('psi')
tau = var('tau')
k = var('k')
L = var('L')
l = var('l')
v0 = var('v0')
s = var('s')

# system [gamma, A]
system = Matrix([

    [(tau*psi + 2)*(tau*psi + 2 - 2*s*l*v0) - 2,
     -1*((1 - k*l/(1 + (tau*psi)/2)) + (1 + k*l/(1 + (tau*psi)/2))*exp(-k*L))],

    [2*tau*psi + 8,
     (2/k)*(1 - exp(-k*L))*(1 + 1/(tau*psi + 2))]

])

# solution
[[gamma], [A]] = system \ Matrix([[0], [1]])
```

and get

$$\gamma = -\frac{1}{2\left(\frac{(\psi\tau+4)\left(\left(\frac{2kl}{\psi\tau+2}+1\right)e^{(-Lk)}-\frac{2kl}{\psi\tau+2}+1\right)l}{(2lsv_0-\psi\tau-2)(\psi\tau+2)+2} + \frac{\left(\frac{1}{\psi\tau+2}+1\right)(e^{(-Lk)}-1)}{k}\right)} \quad (67)$$

$$A_+ = A_- = \frac{\left(\left(\frac{2kl}{\psi\tau+2}+1\right)e^{(-Lk)} - \frac{2kl}{\psi\tau+2}+1\right)l}{2\left((2lsv_0-\psi\tau-2)(\psi\tau+2)+2\right)\left(\frac{(\psi\tau+4)\left(\left(\frac{2kl}{\psi\tau+2}+1\right)e^{(-Lk)}-\frac{2kl}{\psi\tau+2}+1\right)l}{(2lsv_0-\psi\tau-2)(\psi\tau+2)+2} + \frac{\left(\frac{1}{\psi\tau+2}+1\right)(e^{(-Lk)}-1)}{k}\right)} \quad (68)$$

## B. Scaling regime

We assume there exists a scaling form

$$\psi(s) = L^{-2}\varphi(sL) \quad (69)$$

so that for  $s = \mathcal{O}(L^{-1})$  we have  $\psi(s) = \mathcal{O}(L^{-2})$  and  $k = \mathcal{O}(L^{-1})$  from Eq. (50).

We give an expression of Eqs. (63, 64) at the lowest order of  $L^{-1}$  and using  $A_+ = A_-$

$$2\gamma - l(1 + e^{-kL})A_+ = 0 \quad (70)$$

$$8\gamma + \frac{4}{k}(1 - e^{-kL})A_+ = 1 \quad (71)$$

from which we infer

$$\gamma = \frac{1}{8}kl \frac{1 + e^{-kL}}{1 - e^{-kL}} \quad (72)$$

and an expression of Eq. (59) at the lowest order of  $L^{-1}$

$$\gamma = \frac{\psi}{8sv_0^2} \quad (73)$$

which with Eq.(50) linking  $k$  and  $\psi$ , also at the lowest order of  $L^{-1}$ ,

$$k^2 l^2 = \tau \psi \Rightarrow k = \begin{cases} \frac{1}{l} \sqrt{\tau \psi} & \text{if } s > 0, \\ 0 & \text{if } s = 0, \\ i \frac{1}{l} \sqrt{-\tau \psi} & \text{if } s < 0, \end{cases} \quad (74)$$

give

$$\frac{2}{sLv_0} = \frac{\coth\left(\sqrt{\frac{\psi L^2}{4lv_0}}\right)}{\sqrt{\frac{\psi L^2}{4lv_0}}} \Leftrightarrow \frac{1}{\Lambda} = \frac{\coth(\sqrt{\Psi})}{\sqrt{\Psi}} \quad s > 0, \psi > 0 \quad (75)$$

$$-\frac{2}{sLv_0} = \frac{\cot\left(\sqrt{\frac{|\psi| L^2}{4lv_0}}\right)}{\sqrt{\frac{|\psi| L^2}{4lv_0}}} \Leftrightarrow -\frac{1}{\Lambda} = \frac{\cot(\sqrt{|\Psi|})}{\sqrt{|\Psi|}} \quad s < 0, \psi < 0 \quad (76)$$

where we have set

$$\Lambda = \frac{sLv_0}{2} \quad (77)$$

$$\Psi = \frac{\psi L^2}{4lv_0} = \frac{\varphi}{4lv_0} \quad (78)$$

for convenience, and introduce

$$\Gamma = \frac{\Psi}{\Lambda} = \frac{\psi}{8sv_0^2} \frac{4L}{l} = \gamma \frac{4L}{l} \quad (79)$$

which we plot in Fig. 1.

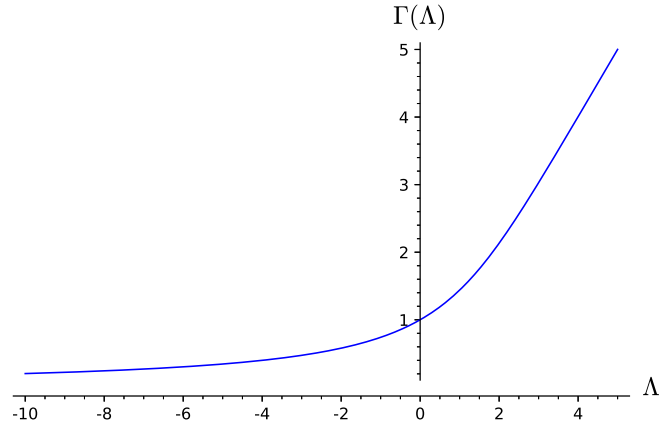


FIG. 1: Rescaled sticking term  $\Gamma = (4L/l)\gamma$  as a function of rescaled biasing parameter  $\Lambda = sLv_0/2$ .

At the lowest order in  $L^{-1}$  we have from Eqs. (53, 54, 55)

$$\varepsilon(r) = \theta(r) = \beta(r) = \frac{1}{2}A_+(e^{-kr} + e^{-k(L-r)}) \quad (80)$$

and from Eqs. (70, 71, 72)

$$A_+ = \frac{2\gamma}{l} \frac{1}{1 + e^{-kL}} = \frac{k}{4(1 - e^{-kL})} \quad (81)$$

so we can write with Eqs. (74, 77, 78, 79)

$$L\varepsilon(r) = \frac{1}{4}\Gamma \begin{cases} \frac{\cosh(\sqrt{\Psi}(1-2r/L))}{\cosh(\sqrt{\Psi})} & \text{if } s > 0, \\ 1 & \text{if } s = 0, \\ \frac{\cos(\sqrt{|\Psi|}(1-2r/L))}{\cos(\sqrt{|\Psi|})} & \text{if } s < 0, \end{cases} \quad (82)$$

which we plot in Fig. 2.

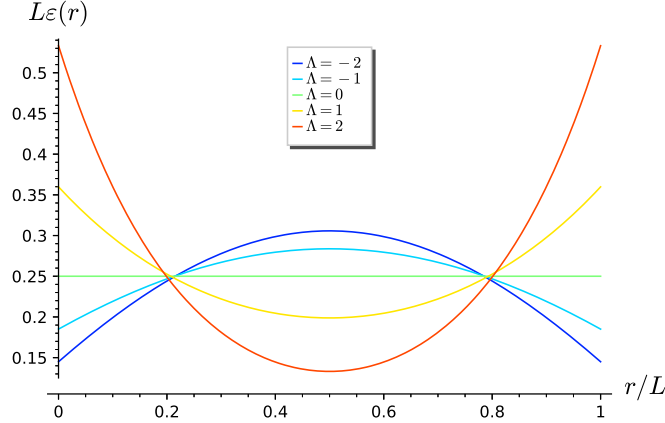


FIG. 2: Rescaled regular part  $L\varepsilon(r)$  as a function of rescaled distance  $r/L$  for different values of the rescaled biasing parameter  $\Lambda = sLv_0/2$ .

#### IV. SOLUTION TO THE ADJOINT EIGENPROBLEM

##### A. General solution

We have computed the left eigenvectors  $P_{\alpha_1\alpha_2}$  of Eq. 12 corresponding to the SCGF  $\psi(\lambda)$  in Sec. III, and will now compute the right eigenvectors  $Q_{\alpha_1\alpha_2}$ . We introduce  $Y_{\alpha_1\alpha_2}(r) = P_{\alpha_1\alpha_2}(r)Q_{\alpha_1\alpha_2}(r)$  which has the same symmetries as  $P_{\alpha_1\alpha_2}(r)$ , so that  $Q_{\alpha_1\alpha_2}(r)$  also has these properties.

We assume that the normalisation condition

$$\int_0^L \sum_{\alpha_1, \alpha_2 = \pm 1} Y_{\alpha_1\alpha_2}(r) dr = 1 \quad (83)$$

as well as the general form

$$Y_{\alpha\alpha}(r) = \hat{\varepsilon}_{\alpha\alpha}(r) + \hat{\gamma}_{\alpha\alpha}\delta(r) + \hat{\gamma}_{\alpha\alpha}\delta(L-r) \quad (84)$$

$$Y_{+-}(r) = \hat{\varepsilon}_{+-}(r) + \hat{\gamma}_{\alpha\bar{\alpha}}\delta(r) \quad (85)$$

$$Y_{-+}(r) = \hat{\varepsilon}_{-+}(r) + \hat{\gamma}_{\alpha\bar{\alpha}}\delta(L-r) \quad (86)$$

also hold for  $Y_{\alpha_1\alpha_2}$ .

We write the hermitian adjoints of Eqs. 39, 40, 41 for the right eigenvectors

$$\psi Q_{\alpha\alpha}(r) = -2\partial_r V(r)\partial_r Q_{\alpha\alpha}(r) + \tau^{-1}(Q_{\alpha\bar{\alpha}}(r) + Q_{\bar{\alpha}\alpha}(r) - 2Q_{\alpha\alpha}(r)) \quad (87)$$

$$\psi Q_{\alpha\bar{\alpha}}(r) = -\alpha 2v_0\partial_r Q_{\alpha\bar{\alpha}}(r) - 2\partial_r V(r)\partial_r Q_{\alpha\bar{\alpha}}(r) + \tau^{-1}(2Q_{\alpha\alpha}(r) - 2Q_{\alpha\bar{\alpha}}(r)) - \alpha 2sv_0\partial_r V(r)Q_{\alpha\bar{\alpha}}(r) \quad (88)$$

following the symmetries we have introduced.

We notice that the bulk equations for Eqs. 87, 88 are the bulk equations (43), (44), (45) under the transformation  $\alpha \rightarrow \bar{\alpha}$  so we get on  $]0; L[$

$$Q_{\alpha\alpha}(r) = A'(\tau\psi + 2)^{-1}(e^{-kr} + e^{-k(L-r)}) \quad (89)$$

$$2Q_{\alpha\bar{\alpha}}(r) = \left(1 + \frac{\alpha 2kl}{\tau\psi + 2}\right) A' e^{-kr} + \left(1 - \frac{\alpha 2kl}{\tau\psi + 2}\right) A' e^{-k(L-r)} \quad (90)$$

where  $k$  is defined by Eq. 50 and  $A'$  is a new constant.

### 1. Complete determination method

We write for  $Y_{\alpha_1\alpha_2}$  the equations

$$\begin{aligned} \psi Y_{\alpha\bar{\alpha}}(r) \frac{P_{\alpha\alpha}(r)}{P_{\alpha\bar{\alpha}}(r)} = & [-\alpha 2v_0 \partial_r Q_{\alpha\bar{\alpha}}(r) - 2\partial_r V(r) \partial_r Q_{\alpha\bar{\alpha}}(r)] P_{\alpha\alpha}(r) \\ & + \tau^{-1} \left( 2Y_{\alpha\alpha}(r) - 2Y_{\alpha\bar{\alpha}}(r) \frac{P_{\alpha\alpha}(r)}{P_{\alpha\bar{\alpha}}(r)} \right) - \alpha 2sv_0 \partial_r V(r) Y_{\alpha\bar{\alpha}}(r) \frac{P_{\alpha\alpha}(r)}{P_{\alpha\bar{\alpha}}(r)} \end{aligned} \quad (91)$$

$$\begin{aligned} \psi Y_{\alpha\bar{\alpha}}(r) = & -\alpha 2v_0 P_{\alpha\bar{\alpha}}(r) \partial_r Q_{\alpha\bar{\alpha}}(r) - 2\partial_r V(r) P_{\alpha\bar{\alpha}}(r) \partial_r Q_{\alpha\bar{\alpha}}(r) + \tau^{-1} \left( 2Y_{\alpha\alpha}(r) \frac{P_{\alpha\bar{\alpha}}(r)}{P_{\alpha\alpha}(r)} - 2Y_{\alpha\bar{\alpha}}(r) \right) \\ & - \alpha 2sv_0 \partial_r V(r) Y_{\alpha\bar{\alpha}}(r) \end{aligned} \quad (92)$$

$$\begin{aligned} \psi Y_{\alpha\bar{\alpha}}(r) = & \alpha 2v_0 Q_{\alpha\bar{\alpha}}(r) \partial_r P_{\alpha\bar{\alpha}}(r) + 2Q_{\alpha\bar{\alpha}}(r) \partial_r (Q_{\alpha\bar{\alpha}}(r) \partial_r V(r)) + \tau^{-1} \left( 2Y_{\alpha\bar{\alpha}}(r) \frac{P_{\alpha\alpha}(r)}{P_{\alpha\bar{\alpha}}(r)} - 2Y_{\alpha\bar{\alpha}}(r) \right) \\ & - \alpha 2sv_0 \partial_r V(r) Y_{\alpha\bar{\alpha}}(r) \end{aligned} \quad (93)$$

where (91) is derived from the multiplication of (88) by  $P_{\alpha\alpha}(r)$ , (92) is derived from the multiplication of (88) by  $P_{\alpha\bar{\alpha}}(r)$ , (93) is derived from the multiplication of (40, 41) by  $Q_{\alpha\bar{\alpha}}(r)$ , and finally subtract Eqs. 92, 93

$$0 = -\alpha 2v_0 \partial_r Y_{\alpha\bar{\alpha}}(r) - 2\partial_r (Y_{\alpha\bar{\alpha}}(r) \partial_r V(r)) + \tau^{-1} \left( 2Y_{\alpha\alpha}(r) \frac{P_{\alpha\bar{\alpha}}(r)}{P_{\alpha\alpha}(r)} - 2Y_{\alpha\bar{\alpha}}(r) \frac{P_{\alpha\alpha}(r)}{P_{\alpha\bar{\alpha}}(r)} \right) \quad (94)$$

to get equations where derivatives should have nice integrals.

We integrate Eqs. 91, 93 between  $0^-$  and  $0^+$

$$\psi \hat{\gamma}_{\alpha\bar{\alpha}} \frac{\gamma_{\alpha\alpha}}{\gamma_{\alpha\bar{\alpha}}} = \tau^{-1} \left( 2\hat{\gamma}_{\alpha\alpha} - 2\hat{\gamma}_{\alpha\bar{\alpha}} \frac{\gamma_{\alpha\alpha}}{\gamma_{\alpha\bar{\alpha}}} \right) + 2sv_0^2 \hat{\gamma}_{\alpha\bar{\alpha}} \frac{\gamma_{\alpha\alpha}}{\gamma_{\alpha\bar{\alpha}}} \quad (95)$$

$$0 = -v_0 \hat{\varepsilon}_{+-}(0^+) + \tau^{-1} \left( \hat{\gamma}_{\alpha\alpha} \frac{\gamma_{\alpha\bar{\alpha}}}{\gamma_{\alpha\alpha}} - \hat{\gamma}_{\alpha\bar{\alpha}} \frac{\gamma_{\alpha\alpha}}{\gamma_{\alpha\bar{\alpha}}} \right) \quad (96)$$

where we have discarded in (95) the integral of the first term of the right hand side of (91) since  $\partial_r V(0) = -v_0$ , and



where

$$\begin{aligned}\hat{\varepsilon}_{+-}(r) &= Q_{+-}(r)P_{+-}(r) \quad (r \in ]0; L[) \\ &= \frac{AA'}{4} \left[ 1 - \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right] (e^{-2kr} + e^{-2k(L-r)}) \\ &\quad + \frac{AA'}{4} \left[ \left( 1 + \frac{2kl}{\tau\psi + 2} \right)^2 + \left( 1 - \frac{2kl}{\tau\psi + 2} \right)^2 \right] e^{-kL}\end{aligned}\tag{97}$$

$$\begin{aligned}\hat{\varepsilon}_{+-}(0^+) &= \lim_{r \rightarrow 0^+} \hat{\varepsilon}_{+-}(r) = \frac{AA'}{4} \left[ 1 - \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right] (1 + e^{-2kL}) + \frac{AA'}{2} \left[ 1 + \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right] e^{-kL} \\ &= \frac{AA'}{2} e^{-kL} \left[ \left( 1 - \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) \cosh(kL) + \left( 1 + \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) \right]\end{aligned}\tag{98}$$

from Eqs. 53, 90, with  $A = A_{\pm}$  which appear in Eqs. 53, 54, 55. We solve Eqs. 95, 96 for  $\hat{\gamma}_{\alpha\alpha}$  and  $\hat{\gamma}_{\alpha\bar{\alpha}}$

$$\hat{\gamma}_{\alpha\alpha} = \frac{\left[ \frac{1}{2}(\tau\psi + 2) - slv_0 \right] \frac{\gamma_{\alpha\alpha}}{\gamma_{\alpha\bar{\alpha}}} l \hat{\varepsilon}_{+-}(0^+)}{\frac{1}{2}(\tau\psi + 2) - slv_0 - \frac{\gamma_{\alpha\alpha}}{\gamma_{\alpha\bar{\alpha}}}}\tag{99}$$

$$\hat{\gamma}_{\alpha\bar{\alpha}} = \frac{l \hat{\varepsilon}_{+-}(0^+)}{\frac{1}{2}(\tau\psi + 2) - slv_0 - \frac{\gamma_{\alpha\alpha}}{\gamma_{\alpha\bar{\alpha}}}}\tag{100}$$

as a function of  $\hat{\varepsilon}_{+-}(0^+)$ .

We can finally write the normalisation condition (83)

$$1 = 4\hat{\gamma}_{\alpha\alpha} + 2\hat{\gamma}_{\alpha\bar{\alpha}} + 2 \int_0^L \hat{\varepsilon}_{\alpha\alpha}(r) dr + 2 \int_0^L \hat{\varepsilon}_{\alpha\bar{\alpha}}(r) dr\tag{101}$$

where we have

$$\hat{\varepsilon}_{\alpha\alpha}(r) = P_{\alpha\alpha}(r \in ]0; L[) Q_{\alpha\alpha}(r \in ]0; L[) = \frac{AA'}{(\tau\psi + 2)^2} \left( e^{-2kr} + e^{-2k(L-r)} + 2e^{-kL} \right)\tag{102}$$

according to Eqs. 55, 89. We compute from (102)

$$\int_0^L \hat{\varepsilon}_{\alpha\alpha}(r) dr = \frac{2AA'}{(\tau\psi + 2)^2} e^{-kL} \left[ \frac{1}{k} \sinh(kL) + L \right]\tag{103}$$

and from (97)

$$\int_0^L \hat{\varepsilon}_{\alpha\bar{\alpha}}(r) dr = \frac{AA'}{2} e^{-kL} \left[ \left( 1 - \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) \frac{1}{k} \sinh(kL) + \left( 1 + \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) L \right]\tag{104}$$

on the hand, and use on the other hand (98) to re-write (101)

$$\begin{aligned}
1 &= \left[ 4 \frac{\left(\frac{1}{2}(\tau\psi + 2) - slv_0\right) \frac{\gamma_{\alpha\alpha}}{\gamma_{\alpha\bar{\alpha}}}}{\frac{1}{2}(\tau\psi + 2) - slv_0 - \frac{\gamma_{\alpha\alpha}}{\gamma_{\alpha\bar{\alpha}}}} + 2 \frac{l}{\frac{1}{2}(\tau\psi + 2) - slv_0 - \frac{\gamma_{\alpha\alpha}}{\gamma_{\alpha\bar{\alpha}}}} \right] \\
&\quad \times \frac{AA'}{2} e^{-kL} \left[ \left( 1 - \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) \cosh(kL) + \left( 1 + \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) \right] \\
&\quad + \frac{4AA'}{(\tau\psi + 2)^2} e^{-kL} \left[ \frac{1}{k} \sinh(kL) + L \right] \\
&\quad + AA' e^{-kL} \left[ \left( 1 - \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) \frac{1}{k} \sinh(kL) + \left( 1 + \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) L \right] \\
&= AA' e^{-kL} \\
&\quad \times \left[ \frac{l}{\frac{1}{2}(\tau\psi + 2) - slv_0 - \frac{\gamma_{\alpha\alpha}}{\gamma_{\alpha\bar{\alpha}}}} \left( 2 \left( \frac{1}{2}(\tau\psi + 2) - slv_0 \right) \frac{\gamma_{\alpha\alpha}}{\gamma_{\alpha\bar{\alpha}}} + 1 \right) \right. \\
&\quad \times \left( \left( 1 - \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) \cosh(kL) + \left( 1 + \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) \right) \\
&\quad \left. + \frac{4}{(\tau\psi + 2)^2} \left( \frac{1}{k} \sinh(kL) + L \right) + \left( 1 - \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) \frac{1}{k} \sinh(kL) + \left( 1 + \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) L \right]
\end{aligned} \tag{105}$$

and conclude

$$\begin{aligned}
A' &= \frac{e^{kL}}{A} \\
&\quad \times \left[ \frac{l}{\frac{1}{2}(\tau\psi + 2)^2 - slv_0(\tau\psi + 2) - 1} (2(\tau\psi + 2) - 2slv_0) \right. \\
&\quad \times \left( \left( 1 - \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) \cosh(kL) + \left( 1 + \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) \right) \\
&\quad \left. + \frac{4}{(\tau\psi + 2)^2} \left( \frac{1}{k} \sinh(kL) + L \right) + \left( 1 - \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) \frac{1}{k} \sinh(kL) + \left( 1 + \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) L \right]^{-1}
\end{aligned} \tag{106}$$

where we have used  $\gamma_{\alpha\bar{\alpha}}/\gamma_{\alpha\alpha} = \tau\psi + 2$  (58) so we finally know  $Y_{\alpha_1\alpha_2}$  completely.

## 2. Polarisation-oriented method

We have that the form of Eqs. 87, 88 indicates that there cannot be  $\delta$ -function terms in the right eigenvectors  $Q_{\alpha_1\alpha_2}$ , therefore these are completely described by Eqs. 89, 90. We can thus write

$$\hat{\varepsilon}_{\alpha\alpha}(r) = \frac{A'}{A} \varepsilon_{\alpha\alpha}(r)^2 \tag{107}$$

$$\hat{\varepsilon}_{\alpha\bar{\alpha}}(r) = \frac{A'}{A} \varepsilon_{\alpha\bar{\alpha}}(r) \varepsilon_{\bar{\alpha}\alpha}(r) \tag{108}$$

$$\hat{\gamma}_{\alpha\alpha} = \gamma_{\alpha\alpha} \varepsilon_{\alpha\alpha}(0^+) \tag{109}$$

$$\hat{\gamma}_{\alpha\bar{\alpha}} = \gamma_{\alpha\bar{\alpha}} \varepsilon_{+-}(L^-) \tag{110}$$

with the notations of Eqs. 84, 85, 86.

We want to compute the average polarisation

$$\nu = 4\hat{\gamma}_{\alpha\alpha} - 2\hat{\gamma}_{\alpha\bar{\alpha}} + 2 \int_0^L \hat{\varepsilon}_{\alpha\alpha}(r) dr - 2 \int_0^L \hat{\varepsilon}_{\alpha\bar{\alpha}}(r) dr \tag{111}$$

corresponding to the propensity of particles to be aligned, and use the normalisation condition (101) to re-write it

$$\begin{aligned} \nu &= \frac{4\hat{\gamma}_{\alpha\alpha} - 2\hat{\gamma}_{\alpha\bar{\alpha}} + 2\int_0^L \hat{\varepsilon}_{\alpha\alpha}(r) dr - 2\int_0^L \hat{\varepsilon}_{\alpha\bar{\alpha}}(r) dr}{4\hat{\gamma}_{\alpha\alpha} + 2\hat{\gamma}_{\alpha\bar{\alpha}} + 2\int_0^L \hat{\varepsilon}_{\alpha\alpha}(r) dr + 2\int_0^L \hat{\varepsilon}_{\alpha\bar{\alpha}}(r) dr} \\ &= \frac{\frac{\gamma_{\alpha\alpha}}{A} \left[ 4\frac{\varepsilon_{\alpha\alpha}(0^+)}{A} - 2(\tau\psi + 2)\frac{\varepsilon_{+-}(L^-)}{A} \right] + 2\int_0^L \frac{\varepsilon_{\alpha\alpha}(r)^2}{A^2} dr - 2\int_0^L \frac{\varepsilon_{\alpha\bar{\alpha}}(r)\varepsilon_{\bar{\alpha}\alpha}(r)}{A^2} dr}{\frac{\gamma_{\alpha\alpha}}{A} \left[ 4\frac{\varepsilon_{\alpha\alpha}(0^+)}{A} + 2(\tau\psi + 2)\frac{\varepsilon_{+-}(L^-)}{A} \right] + 2\int_0^L \frac{\varepsilon_{\alpha\alpha}(r)^2}{A^2} dr + 2\int_0^L \frac{\varepsilon_{\alpha\bar{\alpha}}(r)\varepsilon_{\bar{\alpha}\alpha}(r)}{A^2} dr} \end{aligned} \quad (112)$$

where we have used  $\gamma_{\alpha\bar{\alpha}} = (\tau\psi + 2)\gamma_{\alpha\alpha}$  (58) from the first to the second line.

We can compute from Eqs. 53, 61

$$\frac{\gamma_{\alpha\alpha}}{A} = \frac{1}{2}l(1 + e^{-kL}) + l\frac{kl}{\tau\psi + 2}(1 - e^{-kL}) \quad (113)$$

from Eqs. 53, 54

$$\frac{\varepsilon_{\alpha\alpha}(0^+)}{A} = \frac{1}{\tau\psi + 2}(1 + e^{-kL}) \quad (114)$$

$$\begin{aligned} \frac{\varepsilon_{+-}(L^-)}{A} &= \frac{1}{2} \left( 1 - \frac{2kl}{\tau\psi + 2} \right) e^{-kL} + \frac{1}{2} \left( 1 + \frac{2kl}{\tau\psi + 2} \right) \\ &= \frac{1}{2}(1 + e^{-kL}) + \frac{kl}{\tau\psi + 2}(1 - e^{-kL}) \end{aligned} \quad (115)$$

and from Eqs. 102, 97

$$\int_0^L \frac{\varepsilon_{\alpha\alpha}(r)^2}{A^2} dr = \frac{2}{(\tau\psi + 2)^2} \left[ \frac{1}{2k}(1 - e^{-2kL}) + Le^{-kL} \right] \quad (116)$$

$$\int_0^L \frac{\varepsilon_{\alpha\bar{\alpha}}(r)\varepsilon_{\bar{\alpha}\alpha}(r)}{A^2} dr = \left( 1 - \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) \frac{1}{4k}(1 - e^{-2kL}) + \frac{1}{2} \left( 1 + \left( \frac{2kl}{\tau\psi + 2} \right)^2 \right) Le^{-kL} \quad (117)$$

so that we can write

$$\frac{1}{(1 + e^{kL})^2} \frac{\gamma_{\alpha\alpha}}{A} \frac{\varepsilon_{\alpha\alpha}(0^+)}{A} = \frac{1}{4k}\Omega + \frac{1}{4k}\Omega^2 \frac{1 - e^{-kL}}{1 + e^{-kL}} \quad (118)$$

$$\frac{1}{(1 + e^{-kL})^2} \frac{\gamma_{\alpha\alpha}}{A} \frac{\varepsilon_{+-}(L^-)}{A} = \frac{1}{4}l + \frac{1}{2}l\Omega \frac{1 - e^{-kL}}{1 + e^{-kL}} + \frac{1}{4}l\Omega^2 \left( \frac{1 - e^{-kL}}{1 + e^{-kL}} \right)^2 \quad (119)$$

$$\frac{1}{(1 + e^{-kL})^2} \int_0^L \frac{\varepsilon_{\alpha\alpha}(r)^2}{A^2} dr = \frac{1}{4k^3l^2}\Omega^2 \frac{\frac{1}{2}(e^{kL} - e^{-kL})}{1 + \frac{1}{2}(e^{kL} + e^{-kL})} + \frac{L}{2k^2l^2}\Omega^2 \frac{1}{2(1 + \frac{1}{2}(e^{kL} + e^{-kL}))} \quad (120)$$

$$\frac{1}{(1 + e^{-kL})^2} \int_0^L \frac{\varepsilon_{\alpha\bar{\alpha}}(r)\varepsilon_{\bar{\alpha}\alpha}(r)}{A^2} dr = \frac{1}{4k}(1 - \Omega^2) \frac{\frac{1}{2}(e^{kL} - e^{-kL})}{1 + \frac{1}{2}(e^{kL} + e^{-kL})} + \frac{1}{2}L(1 + \Omega^2) \frac{1}{2(1 + \frac{1}{2}(e^{kL} + e^{-kL}))} \quad (121)$$

where we have introduced

$$\Omega = \frac{2kl}{\tau\psi + 2} \quad (122)$$

to ease the notations.

We can first check

$$\lim_{s \rightarrow 0} \frac{1}{(1 + e^{kL})^2} \frac{\gamma_{\alpha\alpha}}{A} \frac{\varepsilon_{\alpha\alpha}(0^+)}{A} = \frac{l}{4} \quad (123)$$

$$\lim_{s \rightarrow 0} \frac{1}{(1 + e^{-kL})^2} \frac{\gamma_{\alpha\alpha}}{A} \frac{\varepsilon_{+-}(L^-)}{A} = \frac{l}{4} \quad (124)$$

$$\lim_{s \rightarrow 0} \frac{1}{(1 + e^{-kL})^2} \int_0^L \frac{\varepsilon_{\alpha\alpha}(r)^2}{A^2} dr = \frac{L}{8} \quad (125)$$

$$\lim_{s \rightarrow 0} \frac{1}{(1 + e^{-kL})^2} \int_0^L \frac{\varepsilon_{\alpha\bar{\alpha}}(r)\varepsilon_{\bar{\alpha}\alpha}(r)}{A^2} dr = \frac{L}{8} \quad (126)$$

so that, replacing in (111), we get  $\lim_{s \rightarrow 0} \nu = 0$  as expected.

We finally give the full expressions for  $s < 0$

$$\begin{aligned} \nu(s < 0) = & \frac{\frac{|\Omega|}{|\tilde{k}|} - \frac{1}{2}(\tilde{\psi} + 2) + \left[ \frac{|\Omega|}{|\tilde{k}|} - (\tilde{\psi} + 2) \right] |\Omega| \tan \left( \frac{|\tilde{k}| \tilde{L}}{2} \right) - \frac{1}{2}(\tilde{\psi} + 2) |\Omega|^2 \tan \left( \frac{|\tilde{k}| \tilde{L}}{2} \right)^2}{\frac{|\Omega|}{|\tilde{k}|} + \frac{1}{2}(\tilde{\psi} + 2) + \left[ \frac{|\Omega|}{|\tilde{k}|} + (\tilde{\psi} + 2) \right] |\Omega| \tan \left( \frac{|\tilde{k}| \tilde{L}}{2} \right) + \frac{1}{2}(\tilde{\psi} + 2) |\Omega|^2 \tan \left( \frac{|\tilde{k}| \tilde{L}}{2} \right)^2} \\ & + \frac{\left[ \frac{1}{2|\tilde{k}|^3} |\Omega|^2 - \frac{1}{2|\tilde{k}|} (1 + |\Omega|^2) \right] \frac{\sin(|\tilde{k}| \tilde{L})}{1 + \cos(|\tilde{k}| \tilde{L})} + \left[ \frac{\tilde{L}}{|\tilde{k}|^2} |\Omega|^2 - \tilde{L} (1 - |\Omega|^2) \right] \frac{1}{2(1 + \cos(|\tilde{k}| \tilde{L}))}}{\left[ \frac{1}{2|\tilde{k}|^3} |\Omega|^2 + \frac{1}{2|\tilde{k}|} (1 + |\Omega|^2) \right] \frac{\sin(|\tilde{k}| \tilde{L})}{1 + \cos(|\tilde{k}| \tilde{L})} + \left[ \frac{\tilde{L}}{|\tilde{k}|^2} |\Omega|^2 + \tilde{L} (1 - |\Omega|^2) \right] \frac{1}{2(1 + \cos(|\tilde{k}| \tilde{L}))}} \end{aligned} \quad (127)$$

and for  $s > 0$

$$\begin{aligned} \nu(s > 0) = & \frac{\frac{\Omega}{\tilde{k}} - \frac{1}{2}(\tilde{\psi} + 2) + \left[ \frac{\Omega}{\tilde{k}} - (\tilde{\psi} + 2) \right] \Omega \tanh \left( \frac{\tilde{k} \tilde{L}}{2} \right) - \frac{1}{2}(\tilde{\psi} + 2) \Omega^2 \tanh \left( \frac{\tilde{k} \tilde{L}}{2} \right)^2}{\frac{\Omega}{\tilde{k}} + \frac{1}{2}(\tilde{\psi} + 2) + \left[ \frac{\Omega}{\tilde{k}} + (\tilde{\psi} + 2) \right] \Omega \tanh \left( \frac{\tilde{k} \tilde{L}}{2} \right) + \frac{1}{2}(\tilde{\psi} + 2) \Omega^2 \tanh \left( \frac{\tilde{k} \tilde{L}}{2} \right)^2} \\ & + \frac{\left[ \frac{1}{2\tilde{k}^3} \Omega^2 - \frac{1}{2\tilde{k}} (1 - \Omega^2) \right] \frac{\sinh(\tilde{k} \tilde{L})}{1 + \cosh(\tilde{k} \tilde{L})} + \left[ \frac{\tilde{L}}{\tilde{k}^2} \Omega^2 - \tilde{L} (1 + \Omega^2) \right] \frac{1}{2(1 + \cosh(\tilde{k} \tilde{L}))}}{\left[ \frac{1}{2\tilde{k}^3} \Omega^2 + \frac{1}{2\tilde{k}} (1 - \Omega^2) \right] \frac{\sinh(\tilde{k} \tilde{L})}{1 + \cosh(\tilde{k} \tilde{L})} + \left[ \frac{\tilde{L}}{\tilde{k}^2} \Omega^2 + \tilde{L} (1 + \Omega^2) \right] \frac{1}{2(1 + \cosh(\tilde{k} \tilde{L}))}} \end{aligned} \quad (128)$$

where we have introduced the dimensionless quantities  $\tilde{\psi} = \tau\psi$ ,  $\tilde{k} = kl$  and  $\tilde{L} = L/l$ .

## B. Unbiased ensemble check

We have in the unbiased ensemble  $\psi(s = 0) = 0$ , therefore  $k = 0$  according to (50). It follows that

$$\hat{\varepsilon}_{\alpha\alpha}(r, s = 0) = \hat{\varepsilon}_{\alpha\bar{\alpha}}(r, s = 0) = AA' \quad (129)$$

from Eqs. 97, 102. We know from Eqs. 30, 31, 32, 33 that

$$\left. \frac{\gamma_{\alpha\alpha}}{\gamma_{\alpha\bar{\alpha}}} \right|_{s=0} = \frac{1}{2} \quad (130)$$

therefore we get

$$\hat{\gamma}_{\alpha\alpha}(s = 0) = lAA' \quad (131)$$

$$\hat{\gamma}_{\alpha\bar{\alpha}}(s = 0) = 2lAA' \quad (132)$$

from Eqs. 99, 100. We can then write the normalisation condition (101)

$$1 = 4lAA' + 4lAA' + 2AA'L + 2AA'L \quad (133)$$

such that

$$A' = \frac{1}{A(8l + 4L)} = 1 \quad (134)$$

where the second equality derives from Eq. 34, and which consequence is that

$$Q_{\alpha\alpha}(r, s = 0) = Q_{\alpha\bar{\alpha}}(r, s = 0) = 1 \quad (135)$$

according to Eqs. 89, 90.

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