Forces in vertex models

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I. CELL POTENTIAL ENERGY

We introduce for each cell i a reference perimeter P_i^0 and a reference area A_i^0 , and for each of its cell corner $\mu \in \Omega_i$ a force whose effect is to bring the cell's perimeter P_i and area A_i to their reference quantities. A possible force derives from the following potential energy [1–4]

$$E_{\text{VM}} = \sum_{\text{colls } i} \left[\frac{1}{2} K_i (A_i - A_i^0)^2 + \frac{1}{2} \Gamma_i (P_i - P_i^0)^2 \right], \tag{1}$$

where K_i and Γ_i are respectively area and perimeter elastic constants. We denote r_{μ} , r_i the position of vertices μ , i (see Fig. 1). The cell's perimeter can be written

$$P_i = \sum_{\mu \in \Omega_i} |\boldsymbol{r}_{\mu} - \boldsymbol{r}_{\mu-1}|,\tag{2}$$

and the cell's area can be computed with the shoelace formula

$$A_i = \sum_{\mu \in \Omega_i} \frac{1}{2} [(\boldsymbol{r}_{\mu} - \boldsymbol{r}_i) \times (\boldsymbol{r}_{\mu-1} - \boldsymbol{r}_i)] \cdot \hat{\boldsymbol{e}}_z.$$
(3)

Note that, by convention of ordering of cell corners, each term in this sum must be positive. With these notations, we thus write the force acting on vertex μ

$$\mathbf{F}_{\mathrm{VM},\mu} = -\nabla_{\mu} E_{\mathrm{VM}},\tag{4}$$

where $\nabla_{\mu} \equiv \partial/\partial \boldsymbol{r}_{\mu}$. We compute

$$\nabla_{\mu} (|\mathbf{r}_{\mu} - \mathbf{r}_{\mu-1}|) = \frac{\mathbf{r}_{\mu} - \mathbf{r}_{\mu-1}}{|\mathbf{r}_{\mu} - \mathbf{r}_{\mu-1}|},$$
 (5)

as well as

$$\nabla_{\mu} \left(\left[\left(\boldsymbol{r}_{\mu} - \boldsymbol{r}_{i} \right) \times \left(\boldsymbol{r}_{\mu-1} - \boldsymbol{r}_{i} \right) \right] \cdot \hat{\boldsymbol{e}}_{z} \right) = \left(\boldsymbol{r}_{\mu-1} - \boldsymbol{r}_{i} \right) \times \hat{\boldsymbol{e}}_{z} \tag{6}$$

to write (4) in its full form

$$F_{\text{VM},\mu} = -\sum_{\text{cells i, }\mu \in \Omega_{i}} \left[\frac{1}{2} K_{i} (A_{i} - A_{i}^{0}) \left[(\boldsymbol{r}_{\mu+1} - \boldsymbol{r}_{i}) \times \hat{\boldsymbol{e}}_{z} - (\boldsymbol{r}_{\mu-1} - \boldsymbol{r}_{i}) \times \hat{\boldsymbol{e}}_{z} \right] + \Gamma_{i} (P_{i} - P_{i}^{0}) \left[\frac{\boldsymbol{r}_{\mu} - \boldsymbol{r}_{\mu-1}}{|\boldsymbol{r}_{\mu} - \boldsymbol{r}_{\mu-1}|} + \frac{\boldsymbol{r}_{\mu} - \boldsymbol{r}_{\mu-1}}{|\boldsymbol{r}_{\mu} - \boldsymbol{r}_{\mu-1}|} \right] \right]$$

$$= \sum_{\text{cells i, }\mu \in \Omega_{i}} \left[\frac{1}{2} K_{i} (A_{i} - A_{i}^{0}) \underbrace{(\boldsymbol{r}_{\mu-1} - \boldsymbol{r}_{\mu+1}) \times \hat{\boldsymbol{e}}_{z}}_{\text{towards cell interior}} + \Gamma_{i} (P_{i} - P_{i}^{0}) \left[\frac{\boldsymbol{r}_{\mu-1} - \boldsymbol{r}_{\mu}}{|\boldsymbol{r}_{\mu-1} - \boldsymbol{r}_{\mu}|} + \frac{\boldsymbol{r}_{\mu+1} - \boldsymbol{r}_{\mu}}{|\boldsymbol{r}_{\mu+1} - \boldsymbol{r}_{\mu}|} \right] \right],$$

$$\text{towards neighbour}$$

where underbraced vectors are represented in Fig. 1. It is noteworthy that, with potential (1), the force acting on each cell centre is zero

$$\mathbf{F}_{\text{VM},i} = -\nabla_i E_{\text{VM}} = -\frac{1}{2} K_i (A_i - A_i^0) \sum_{\mu \in \Omega_i} (\mathbf{r}_{\mu} - \mathbf{r}_{\mu-1}) \times \hat{\mathbf{e}}_z = 0.$$
 (8)

^[1] Reza Farhadifar, Jens-Christian Röper, Benoit Aigouy, Suzanne Eaton, and Frank Jülicher, "The Influence of Cell Mechanics, Cell-Cell Interactions, and Proliferation on Epithelial Packing," Current Biology 17, 2095–2104 (2007).

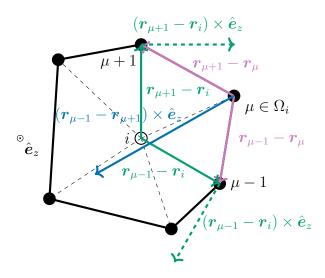


FIG. 1. Vertex representation of a single cell. By convention, cell centres are denoted by latin indices, and cell corners are denotes by greek indices. We denote Ω_i the ensemble of cell corners μ of the cell whose centre is i. By convention, $\mu+1$ (respectively $\mu-1$) denotes the next cell corner of i after μ in anticlockwise (respectively clockwise) order.

- [2] Alexander G. Fletcher, Miriam Osterfield, Ruth E. Baker, and Stanislav Y. Shvartsman, "Vertex Models of Epithelial Morphogenesis," Biophysical Journal 106, 2291–2304 (2014).
- [3] Dapeng Bi, Xingbo Yang, M. Cristina Marchetti, and M. Lisa Manning, "Motility-Driven Glass and Jamming Transitions in Biological Tissues," Physical Review X 6, 021011 (2016).
- [4] Rastko Sknepnek, Ilyas Djafer-Cherif, Manli Chuai, Cornelis Weijer, and Silke Henkes, "Generating active T1 transitions through mechanochemical feedback," eLife 12, e79862 (2023).