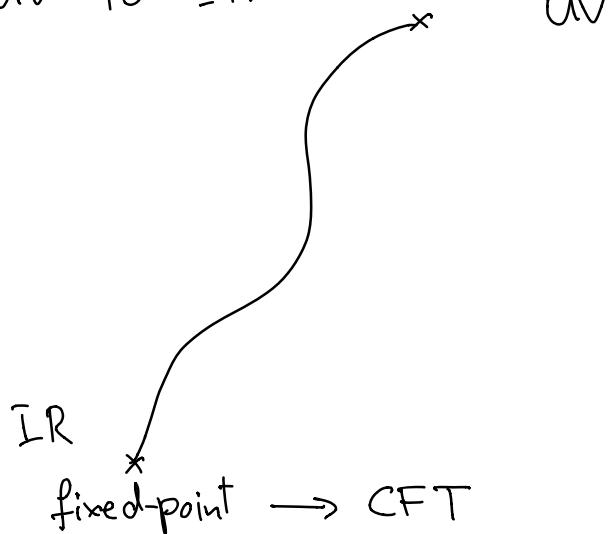
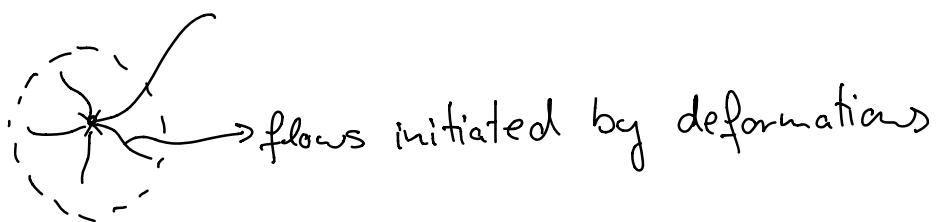


## §2. Deformations of Superconformal Theories

QFT's are renormalization group (RG) flows  
from UV to IR



Given a fixed point, we want to analyze its deformations



3 classes:

1) Adding local operators to the Lagrangian

$$\delta \mathcal{L} = g \mathcal{O}$$

running coupling constant      local operator defined at  $g=0$

2) Gauging a global symmetry:

gauge a continuous flavor symmetry  
with conserved current  $j_\mu$   
 $\rightarrow$  may be obstructed by anomalies

3) Moving onto a moduli space of vacua:

In  $d > 2$  spacetime dimensions CFT may possess non-trivial moduli space of vacua.  
 $\rightarrow$  breaks conformal symmetry spontaneously (restored at origin)

We will focus on deformation 1)

$\mathcal{O}$  must reside in representations of  $SO(d, 2)$

$\rightarrow$  labelled by weights under  $SO(d) \times SO(2)$

Have unique operator  $\mathcal{O}$  of lowest dimension  $\Delta_0$   
Known as conformal primary

$$\begin{array}{ccc} K_\mu \mathcal{O} = 0 & & P_\mu \mathcal{O} \\ \uparrow & & \uparrow \\ \text{of scaling dimension } -1 & & \text{descendants of} \\ & & \text{scaling dimension } \Delta_0 + 1 \\ & & (P_\mu \sim \partial_\mu) \end{array}$$

$\mathcal{O}$  cannot be written as  $D_\mu \mathcal{O}'$

Let us denote  $\Delta_0$  and  $SO(d)$  weights of  $\mathcal{O}$  by  $L_{\mathcal{O}}$ . There is natural inner product  $\langle \mathcal{O}' | \mathcal{O}'' \rangle$

require  $\langle \mathcal{O}' | \mathcal{O}'' \rangle \geq 0$

→ unitarity bound  $\Delta_0 \geq f(L_\mathcal{O})$

→ null-states when bound is saturated

A deformation  $\delta \mathcal{L} \sim g \mathcal{O}$  must be

- a conformal primary

(descendants  $\partial_m \mathcal{O}$  lead to total derivatives)

- $SO(d)$  scalar  $\rightarrow$  preserves Lorentz symmetry

→  $\langle \square \mathcal{O} | \square \mathcal{O} \rangle \geq 0 \Rightarrow \Delta_\mathcal{O} \geq \frac{d-2}{2}$

(recall  $\Delta_\mathcal{O} = h_1 + \frac{d-2}{2}$ )

$\uparrow$   
highest weight of  $SO(d)$   
 $SO(d-2)$ -weights free

bound is saturated for  $\square \mathcal{O} = 0$

Deformation types:

- "relevant deformations" ( $\Delta_\mathcal{O} < d$ ):

CFT at  $g=0$  UV fixed point of RG-flow

→  $g$  grows in the IR (perturbation theory breaks down)

- "Irrelevant deformations" ( $\Delta_\mathcal{O} > d$ ):

CFT is at IR fixed-point of RG-flow

→  $g$  flows to zero

- "Marginal deformations" ( $\Delta_\mathcal{O} = d$ ):

preserve conformal invariance  $\rightarrow$  lead to nearby fixed point

## superconformal theories:

Want to look at deformations preserving  
Poincaré Q-supersymmetries  
(not necessarily S-symmetries)

Recall that we have the following SCA's:

$$d=3 \quad \text{osp}(N|4) \supset SO(3,2) \times SO(N)_R$$

$$d=4 \quad \begin{cases} sl(2,2|N) \supset SO(4,2) \times SU(N)_R \times U(1)_R, N \neq 4 \\ psl(2,2|4) \supset SO(4,2) \times SU(4)_R, N=4 \end{cases}$$

$$d=5 \quad F(4) \supset SO(5,2) \times SU(2)_R, N=1$$

$$d=6 \quad \text{osp}(6,2|N) \supset so(6,2) \times sp(2N)_R$$

$N$  denotes number of d-supercharges

$N_Q$  denotes total number of supercharges

In  $d=3, 4, 5, 6$  minimal  $N=1$  supersymmetry corresponds to  $N_Q = 2, 4, 8, 8$

Notation: Will consider conformal primaries

labelled by  $L_\beta, R_\beta, \Delta_\beta$

↑  
Lorentz-charge

↑  
R-sym charge

←  
scaling dim

$$\mathcal{O} \in [L_\beta]_{\Delta_\beta}^{(R_\beta)}$$

## Superconformal primaries:

a conformal primary  $V$  with

- lowest scaling dimension  $\Delta_V$
- irreducible under R-sym

→ annihilated by  $S$  (scaling dim  $-\frac{1}{2}$ )  
and  $K_n$  (scaling dim.  $-1$ )

other conformal primaries are superconformal descendants of  $V$  (by action of  $Q$ )

Unitarity bound:

$$\Delta_V \geq f(L_V, R_V)$$

saturation leads to null-states

→ descendants form sub-representation

notation: "short" representation

## Deformations!

conformal primary  $O$  with

$$Q O = \partial_m (\dots)$$

→ must transform to descendant  
under action of  $Q$

→ cannot be superconformal primary

→ must be a "top" component

$Q^l V$  are "level"  $l$  conformal primaries  
( $l$  nested anti-commutators)

$$\{Q_i, Q_j\} \sim 0, \quad i, j = 1, \dots, N_Q$$

(we are dropping  $\partial_\mu$ -terms, correspond to conformal descendants)

We have:

$$0 \leq l \leq l_{\max} \text{ with } l_{\max} \leq N_Q$$

→ saturated for "long" multiplets (no "null"-states)

$Q^{N_Q} V$  is top component

Lorentz- and R-sym singlet

"D-term" deformation  $\mathcal{L}_D = Q^{N_Q} V$

$$\Delta(\mathcal{L}_D) = \frac{1}{2} N_Q + \Delta_V$$

example: Kähler potential  $\int d^4 \Theta K(\bar{\Phi}, \bar{\Phi})$  in  $d=4$   
( $V=1$ )

→ F-terms correspond to "short" multiplets (BPS)

$V_{BPS}$  is annihilated by half of supercharges

$$\mathcal{L}_F = Q^{\frac{1}{2} N_Q} V_{BPS}$$

example: chiral superpotential  $\int d^2 \Theta W$

→ sporadic very short multiplets

example: consider the stress-tensor multiplet in

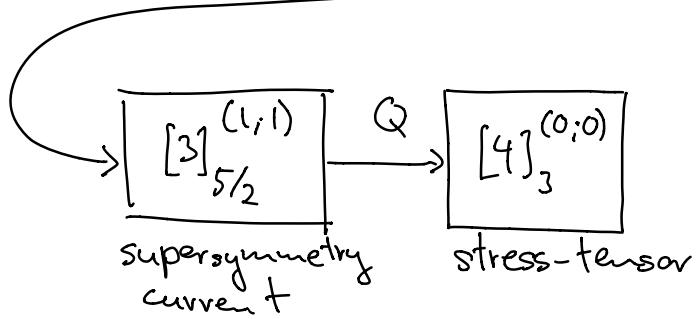
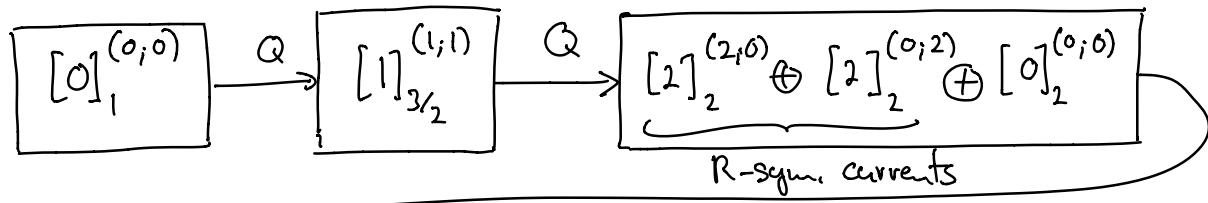
$N=4$   $d=3$  SCFT's

Lorentz-sym.:  $SU(2)$

R-sym.:  $SU(2)_R \times SU(2)'_R$

$$\rightarrow [ij]_{\Delta}^{(R; R')}$$

$Q_\alpha^{i,i'}$  transform as  $[1]_{1/2}^{(1;1)}$   $\rightarrow$  conformal prim. decomp.:



$[0]_2^{(0,0)}$  at level 2 is Lorentz-scalar

$$Q [0]_2^{(0,0)} \sim 0$$

$\rightarrow [0]_2^{(0,0)}$  gives relevant deformation with scaling dimension  $\Delta = 2$ .

Table of deformations in various dimensions:

see next page

$d$	$N$	Relevant	Marginal	Irrelevant $\Delta_{\min}$
3	1	D-term	D-term	$\Delta_{\min} > 3$
	2	Flavor current, F-t.	F-term	$\Delta_{\min} > 3$
	3	Flavor current	—	4
	4	stress t., fl. curr.	—	4
	5, 6	stress tensor	—	5
	8	stress tensor	—	6
4	1	F-term	F-term	$\Delta_{\min} > 4$
	2	fl. curr., F-term	F-term	$\Delta_{\min} > 4$
	3	—	—	$\Delta_{\min} > 4$
	4	—	stress tensor	8
5	1	flavor curr.	—	8
6	$(1,0)$ $(2,0)$	—	—	10 12