

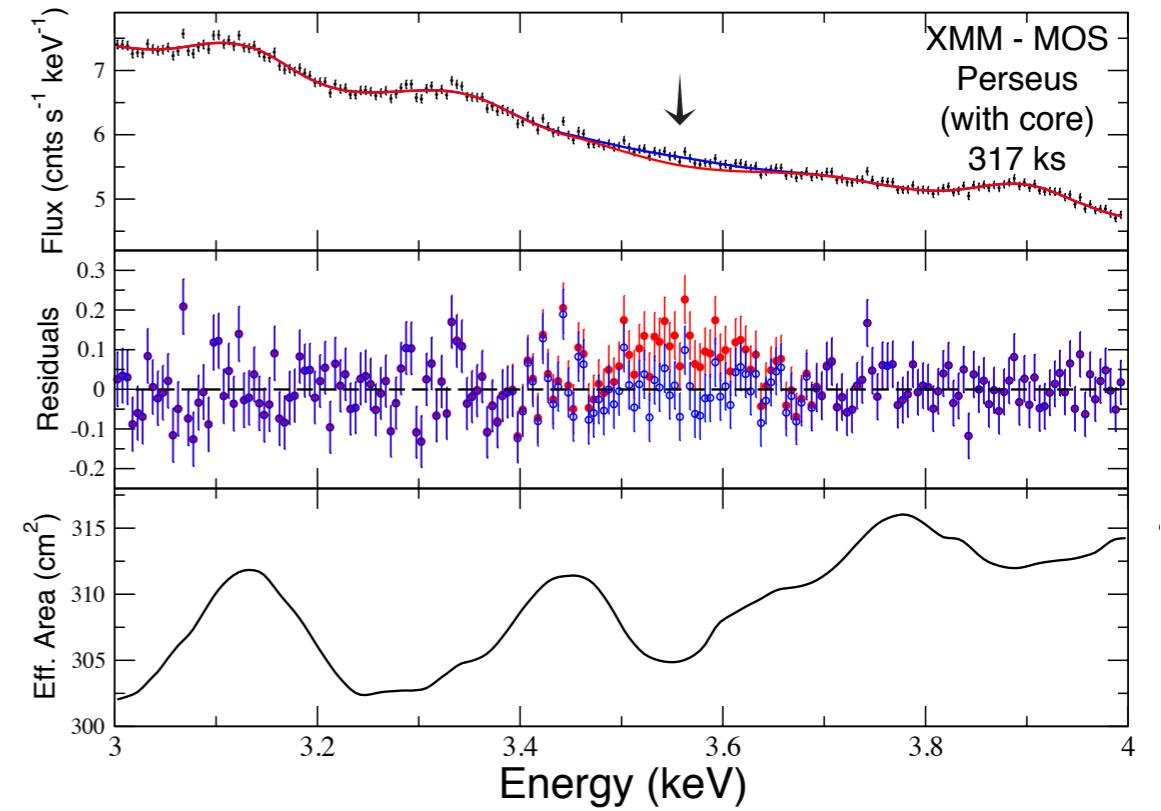
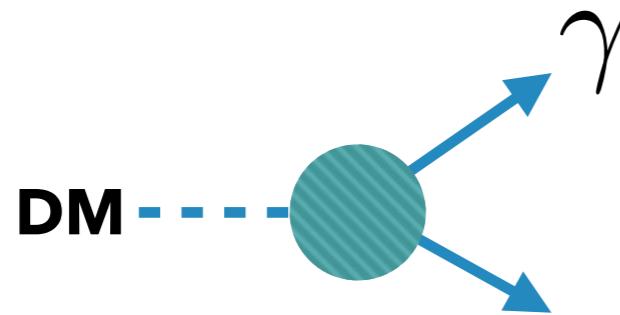
# **Constraints on 7 keV Freeze-in DM: Phase Space Distributions & Ly-a forest constraints**

Keisuke Yanagi (Univ. of Tokyo)

K. J. Bae, A. Kamada, S. P. Liew, K. Y.  
arXiv: 1707.02077, 1707.06418

# 3.5 keV X-ray excess?

- ▶ Rare decay of 7 keV DM?  
From Perseus galaxy cluster, Andromeda galaxy...



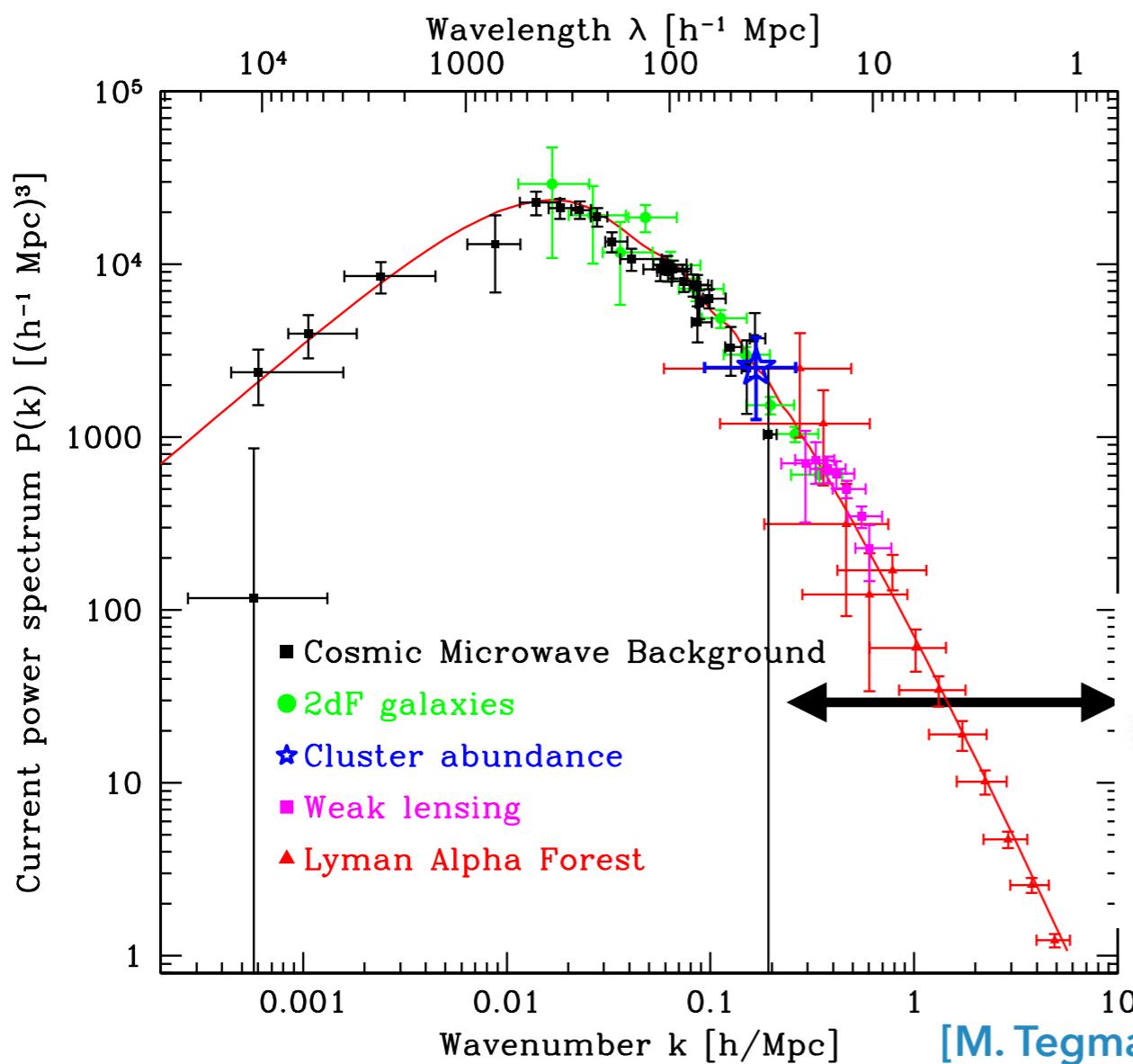
- ▶ Criticism
  - K XVIII explanation ??
  - No signal from dSph, stacked galaxies and groups, M31
- ▶ Even if it is a real signal, lyman-alpha constraint should be studied
- ▶ We take **7 keV axino** as a benchmark

# Matter power spectrum & Ly-alpha forest

- ▶ Matter distribution is calculated by cosmological perturbation theory

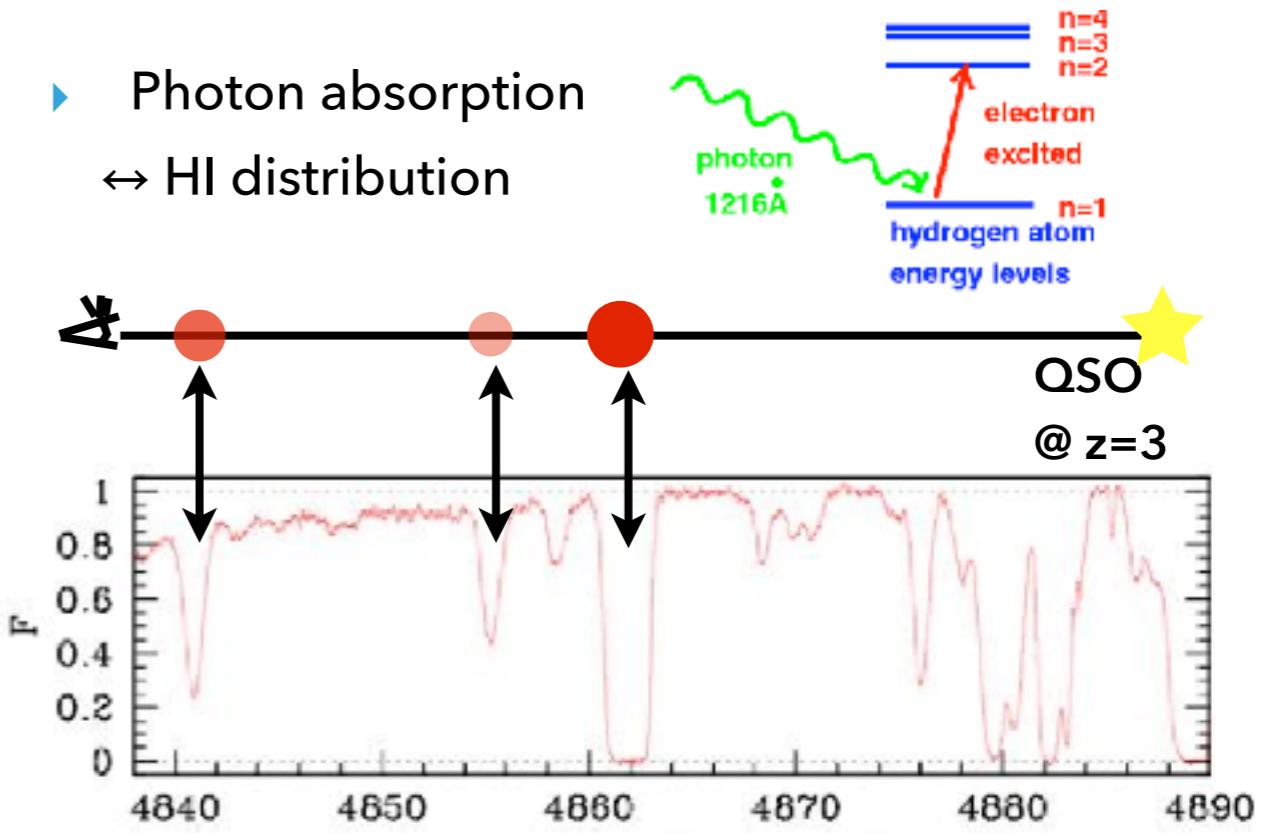
$$\rho_M(\vec{x}, t) = \bar{\rho}_M(t) + \delta\rho_M(\vec{x}, t)$$

- ▶ Observable quantity is **matter power spectrum**



**P( $k$ ) at small scale is determined by Ly-alpha forest observation**

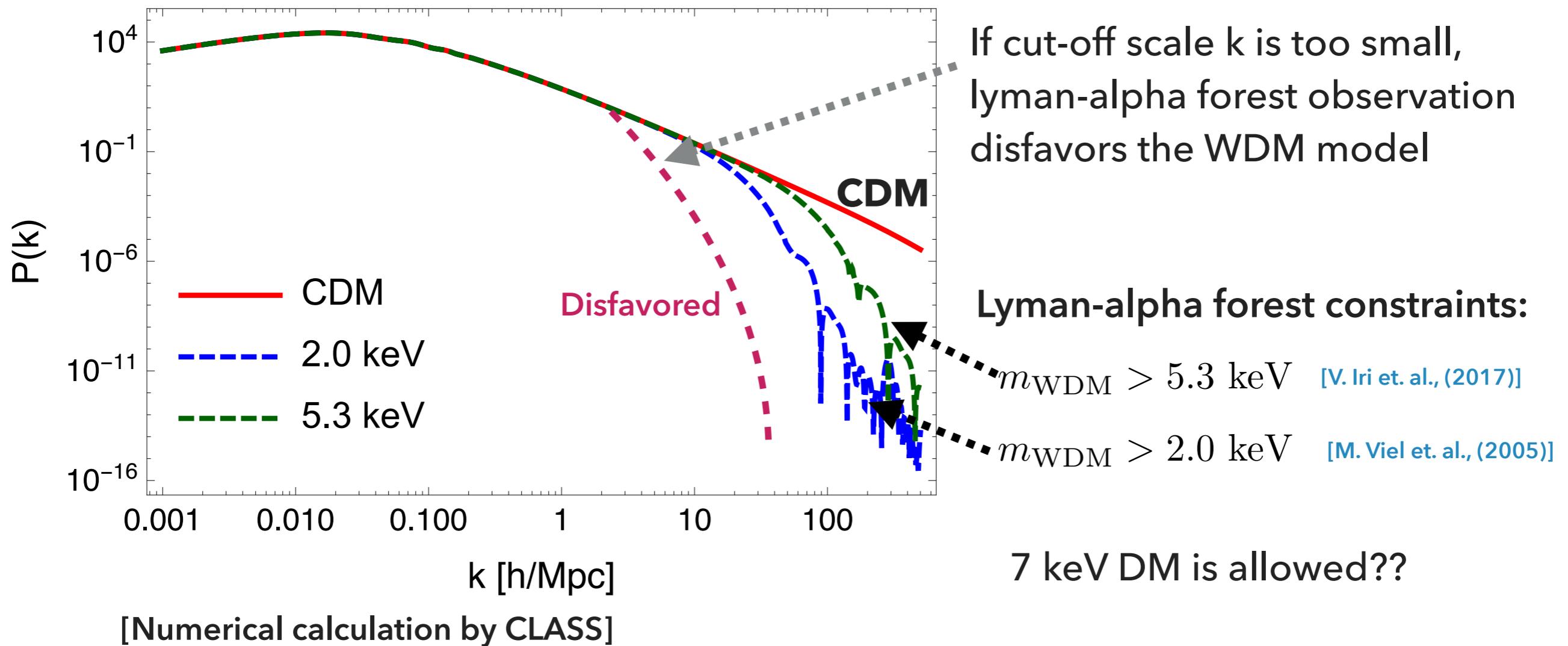
- ▶ Photon absorption  
↔ HI distribution



[M. Tegmark & M. Zaldarriaga (2002)]

# Warm Dark Matter & Ly-alpha constraints

- ▶ DM free-streaming **erases the primordial density contrasts at small scale**
- ▶ Warm Dark Matter (WDM): free-streaming length is kpc - Mpc



# Meaning of "m<sub>WDM</sub> > \*\*\* keV"

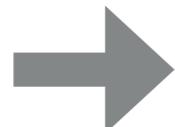
- ▶ Mass bound such as "m<sub>WDM</sub> > 5.3 keV" assumes **conventional WDM**  
i.e., **Fermi-Dirac distribution** w/ spin d.o.f =2

$$n_{\text{WDM}} = 2 \int \frac{d^3 p}{(2\pi)^3} \frac{1}{e^{p/T} + 1}$$

- ▶ DM abundance is determined by 2 parameters: (m<sub>WDM</sub>, T<sub>WDM</sub>)

$$\Omega_{\text{WDM}} h^2 = \left( \frac{m_{\text{WDM}}}{94 \text{eV}} \right) \left( \frac{T_{\text{WDM}}}{T_\nu} \right)^3 = 7.5 \left( \frac{m_{\text{WDM}}}{7 \text{keV}} \right) \left( \frac{106.75}{g_*^{\text{WDM}}} \right)$$

- ▶ For given m<sub>WDM</sub>, T<sub>WDM</sub> is determined to give observed DM abundance.  
Then one compares the theoretical prediction P(k) with the lyman-alpha forest data



m <sub>WDM</sub> > 2.0 keV	[M. Viel et. al., (2005)]
m <sub>WDM</sub> > 3.3 keV	[M. Viel et. al., (2013)]
m <sub>WDM</sub> > 4.09 keV	[J. Baur et. al., (2016)]
m <sub>WDM</sub> > 5.3 keV	[V. Iri et. al., (2017)]

# How can we apply the Ly-alpha constraints to our model?

- Two differences from conventional WDM when considering 7 keV axino

	conventional WDM (Ly-alpha constraints)	7 keV axino (Freeze-in)
$f_{\text{DM}}(p, t)$	Fermi-Dirac	Non-thermal (never thermalized)
$g_*$	~7000	106.75 (SM d.o.f)

$$\Omega_{\text{WDM}} h^2 = \left(\frac{m_{\text{WDM}}}{94 \text{eV}}\right) \left(\frac{T_{\text{WDM}}}{T_\nu}\right)^3 = 7.5 \left(\frac{m_{\text{WDM}}}{7 \text{keV}}\right) \left(\frac{106.75}{g_*^{\text{WDM}}}\right)$$

$$T_{\text{DM}} = \left(\frac{g_*(T)}{g_*(T_{\text{dec}})}\right)^{1/3} T$$

- $P(k)$  is calculated by  $m_{\text{DM}}$ ,  $f_{\text{DM}}(p, t)$ , and  $T_{\text{DM}}$
- In realistic 7 keV DM models, we cannot simply compare 7 keV to  $m_{\text{WDM}} > 5.3 \text{ keV}$
- Axino  $T_{\text{DM}}$  is higher than conventional WDM, the constraint will be more severe

**We need to compare matter power spectrum!**

# Contents

- ▶ 7 keV axino DM in SUSY DFSZ model with R-parity violation
- ▶ Phase space distribution from Boltzmann equation
- ▶ Linear matter power spectrum and Lyman-alpha constraints
- ▶ Entropy production from saxion decay

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# Axino in SUSY DFSZ model

## ► SUSY + PQ

- Motivated by EW naturalness and strong CP problems

- Superpotential

$$W_{\text{DFSZ}} = \frac{y_0}{M_*} X^2 H_u H_d \quad + (\text{PQ SSB sector})$$

**PQ= -1, +1, +1**

- After SSB of PQ symmetry,  $X$  generates axion superfield

$$X = \frac{v_{\text{PQ}}}{\sqrt{2}} e^{A/v_{\text{PQ}}} \quad A = \frac{s + ia}{\sqrt{2}} + \sqrt{2}\theta \tilde{a} + \theta^2 F_A$$

**saxion**      **axion**      **axino (7 keV)**

- 7 keV LSP axino is possible in some models

[e.g., Goto & Yamaguchi (1992); E. J. Chun & A. Lukas (1995)]

## ► Couplings are generated:

- $v_{\text{PQ}} \gtrsim 10^9 \text{ GeV}$      $M_* \sim 10^{16} \text{ GeV}$

- $\mu$  term:  $\mu \sim \frac{y_0 v_{\text{PQ}}^2}{2M_*}$

- Axino-Higgs-Higgsino interaction:  $\frac{2\mu}{v_{\text{PQ}}} \sim 10^{-8}$

**Feebly Interacting Massive Particle (FIMP)**

# R-parity violating decay of axino

- ▶ Bilinear R-parity violation is introduced

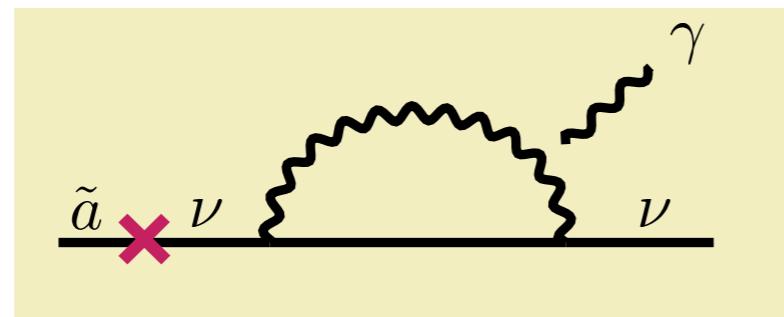
$$W_{\text{bRPV}} = \frac{y'_i}{M_*^2} X^3 L_i H_u \simeq \mu'_i \left(1 + \frac{3A}{v_{\text{PQ}}}\right) L_i H_u, \quad \mu'_i \sim \mu \frac{v_{\text{PQ}}}{M_*}$$

**PQ= -1, +2, +1**

- ▶ After EWSB, bRPV induces **axino-neutrino mixing**

$$|\theta| \simeq \frac{\mu' v_u}{m_{\tilde{a}} v_{\text{PQ}}} \simeq 10^{-5} \left(\frac{\mu'}{4 \text{ MeV}}\right) \left(\frac{7 \text{ keV}}{m_{\tilde{a}}}\right) \left(\frac{10^{10} \text{ GeV}}{v_{\text{PQ}}}\right)$$

- ▶ Axino decays as a sterile neutrino



$\sin^2 2\theta \sim 10^{-10}$  explains 3.5 keV excess!

- ▶ Other options

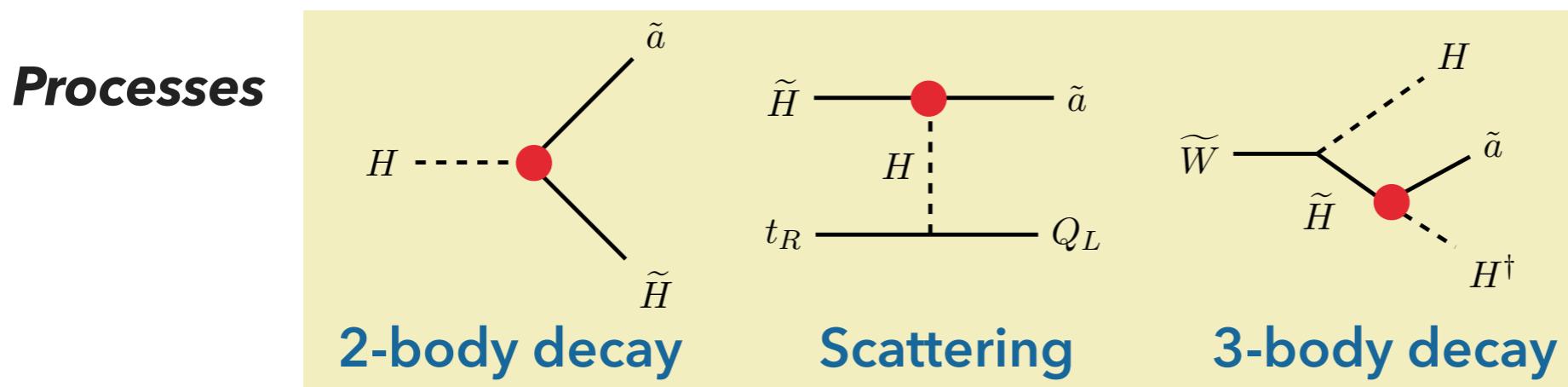
- bRPV in KSVZ model → Light Bino ( $\sim 10$  GeV) or small  $v_{\text{PQ}}$  ( $\sim 10^8$  GeV) required
- sfermion 1-loop in trilinear RPV → Light stau ( $\sim 100$  GeV) required

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# Freeze-in production of axino

- ▶ Axino is gradually produced by the decay or scatterings of MSSM particles
- ▶ Inverse processes are negligible due to small  $n_{\tilde{a}}$



$$W_{\text{DFSZ}} = \frac{y_0}{M_*} X^2 H_u H_d$$

● Feeble coupling  $\sim 10^{-8}$

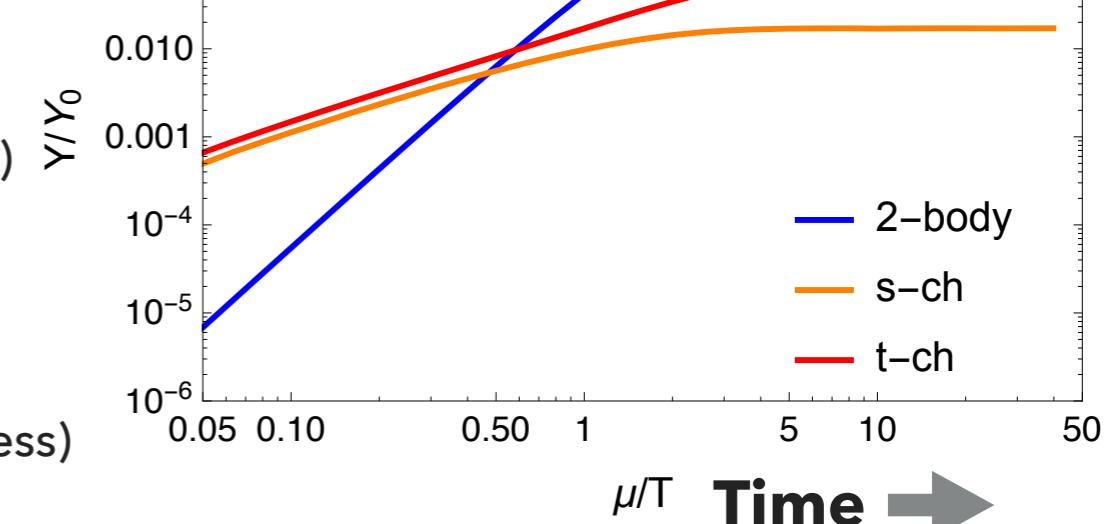
- ▶ For instance, Boltzmann equation for scattering

$$\dot{n}_{\tilde{a}} + 3Hn_{\tilde{a}} = \langle \sigma v \rangle n_1 n_2 \sim T^4 \quad \langle \sigma v \rangle \sim 1/T^2$$

- ▶ Production is **dominated at IR** (mass of other particle)

$$Y_0 = \frac{n_0}{s_0} = \int_0^{t_0} dt \frac{\langle \sigma v \rangle n_1 n_2}{s} \propto \int_m^{T_R} dT M_{pl} \frac{T^4}{T^6} \sim \frac{M_{Pl}}{m}$$

(m=μ for Higgsino process)



→ **NLSP contribution is dominant**

# Boltzmann Equation for phase space distribution

- For matter power spectrum, we need axino **phase space distribution**  $f_{\tilde{a}}(t, p)$
- Axino is never thermalized due to its feeble interaction with MSSM particles
- After production, axino phase space distribution is just **redshifted**

## Boltzmann Eq.

$$\frac{df_{\tilde{a}}(t, p)}{dt} = \frac{\partial f_{\tilde{a}}(t, p)}{\partial t} - \frac{\dot{a}(t)}{a(t)} p \frac{\partial f_{\tilde{a}}(t, p)}{\partial p} = \frac{1}{E_{\tilde{a}}} C(t, p)$$

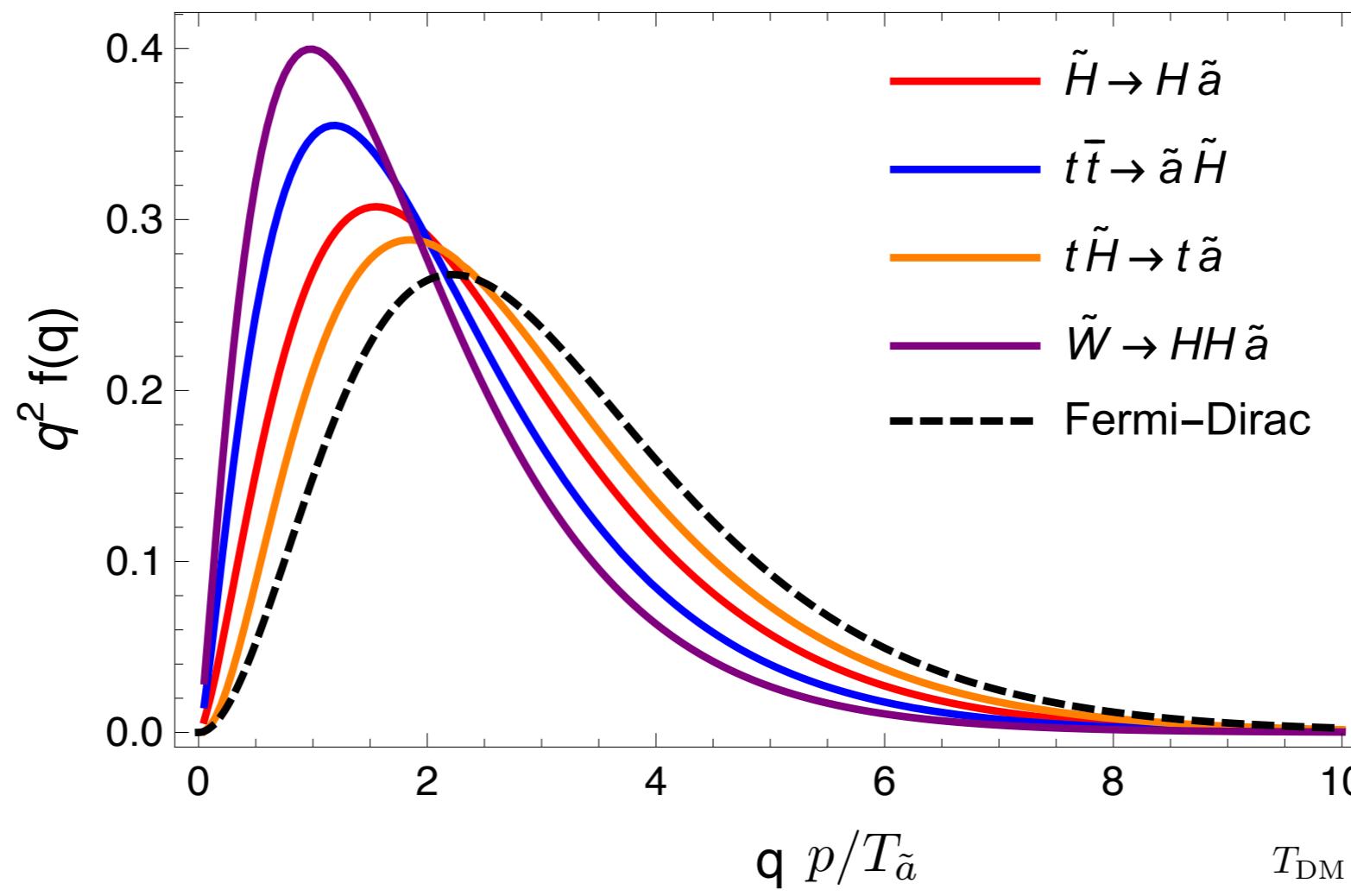
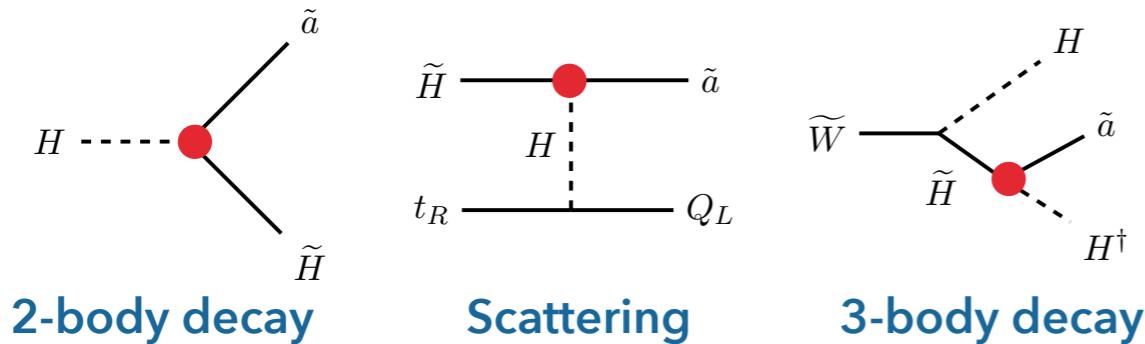
**Collision term** neglect axino density  $f_{\tilde{a}} \simeq 0$

$$\begin{aligned} \frac{g_{\tilde{a}}}{E_{\tilde{a}}} C_{1+2+\dots+\tilde{a}+3+4+\dots}(t, p_{\tilde{a}}) &= \frac{1}{2E_{\tilde{a}}} \int \prod_{I \neq \tilde{a}} \frac{d^3 p_I}{(2\pi)^3 2E_I} (2\pi)^4 \delta^4(\hat{p}_1 + \hat{p}_2 + \dots - \hat{p}_{\tilde{a}} - \hat{p}_3 - \hat{p}_4 - \dots) \\ &\times \sum_{\text{spin}} |\mathcal{M}_{1+2+\dots+\tilde{a}+3+4+\dots}|^2 f_1 f_2 \dots (1 \mp f_3)(1 \mp f_4) \dots \end{aligned}$$

- Integrating Boltzmann Eq., we obtain phase space distribution

$$f_{\tilde{a}}(t_f, p) = \int_{t_i}^{t_f} dt \frac{1}{E_{\tilde{a}}} C \left( t, \frac{a(t_f)}{a(t)} p \right)$$

# Axino momentum distribution: channel dependence

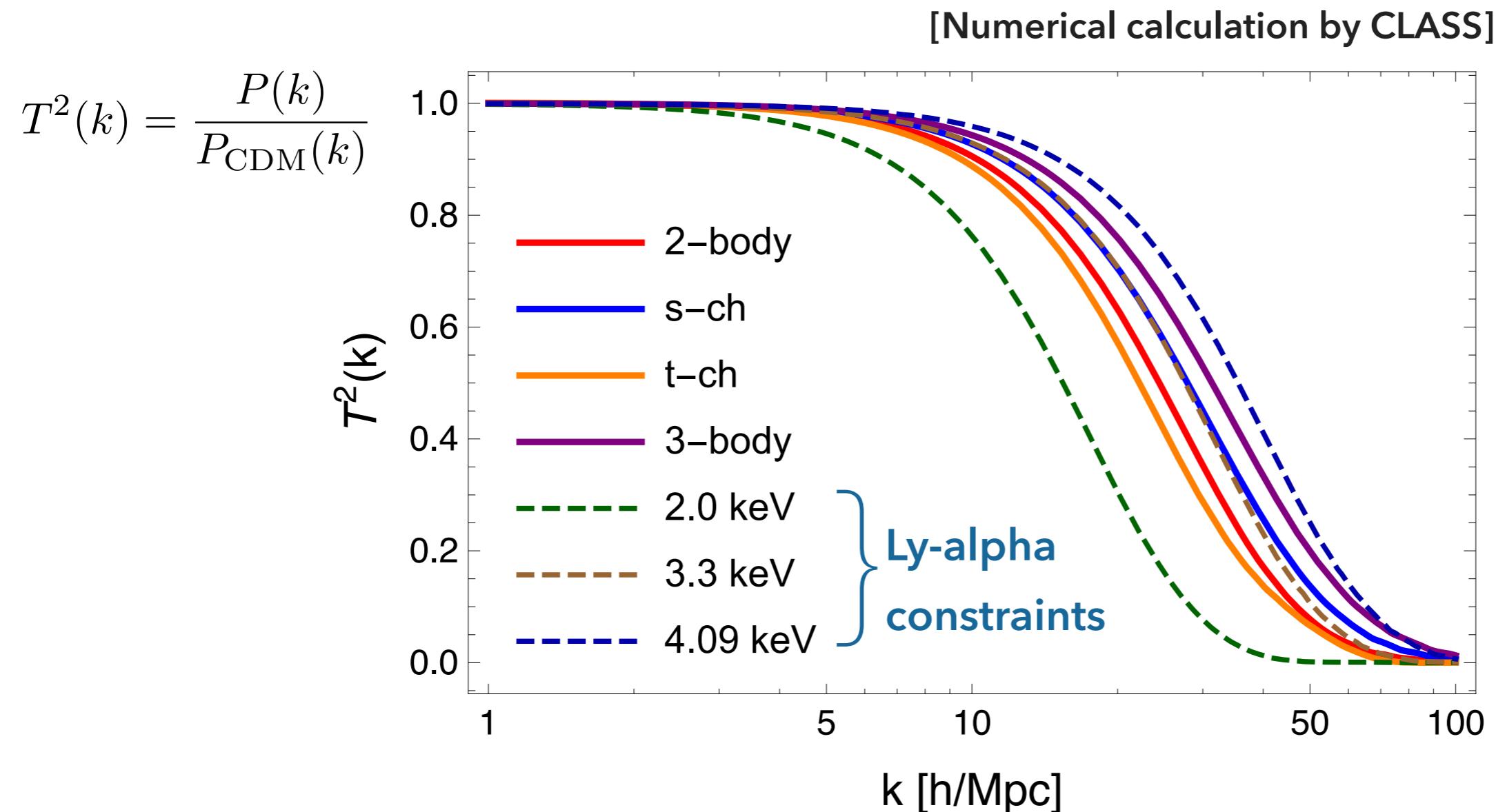


$$T_{\text{DM}} = \left( \frac{g_*(T)}{g_*(T_{\text{dec}})} \right)^{1/3} T$$

- ▶ Distribution is different for different production process
- ▶ **Colder** than Fermi-Dirac distribution

# Matter power spectrum: channel dependence

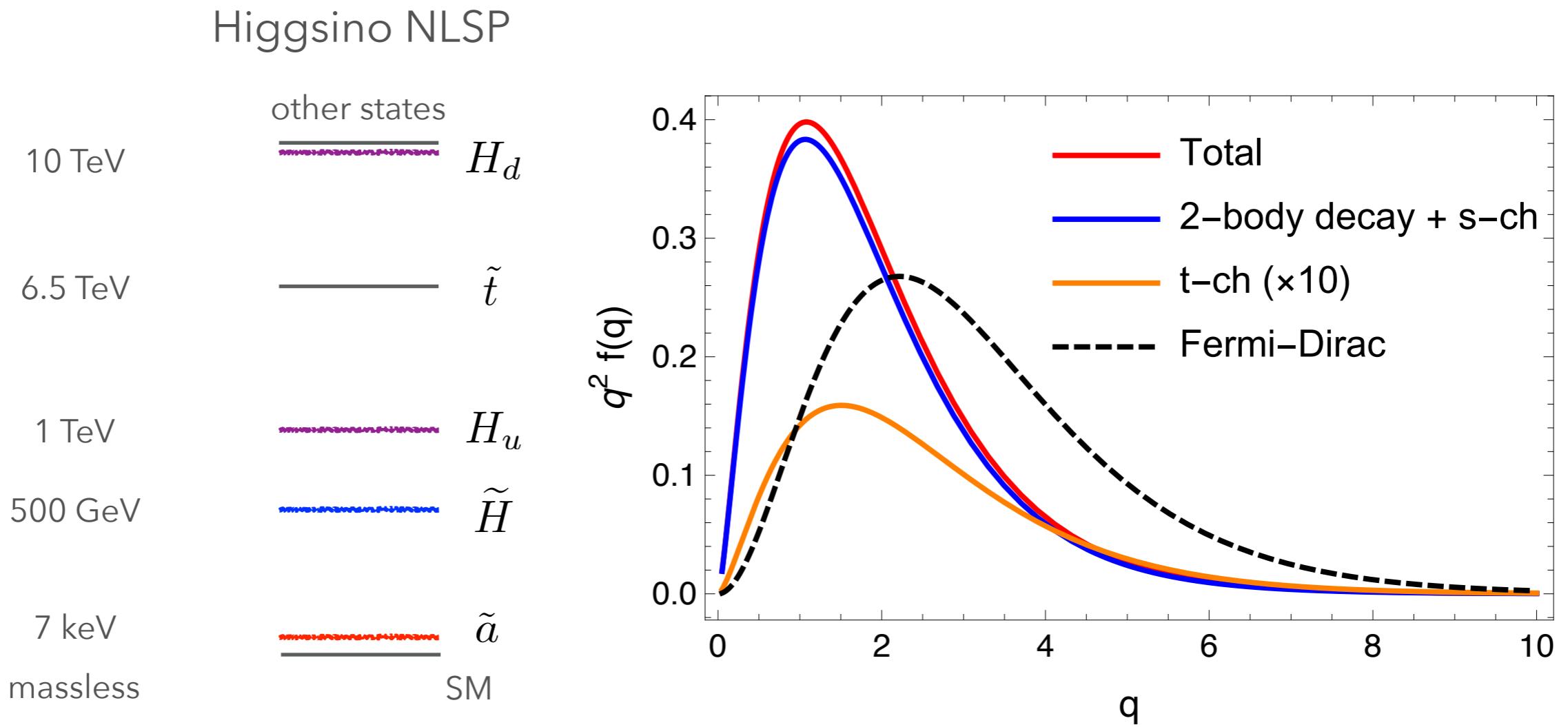
- ▶ We solve cosmological perturbation equations taking axino distributions as inputs



- ▶ We may expect that  $m_{\text{WDM}} > 4.09 \text{ keV}$  excludes 7 keV axino DM  
**even for the coldest distribution from 3-body decay**

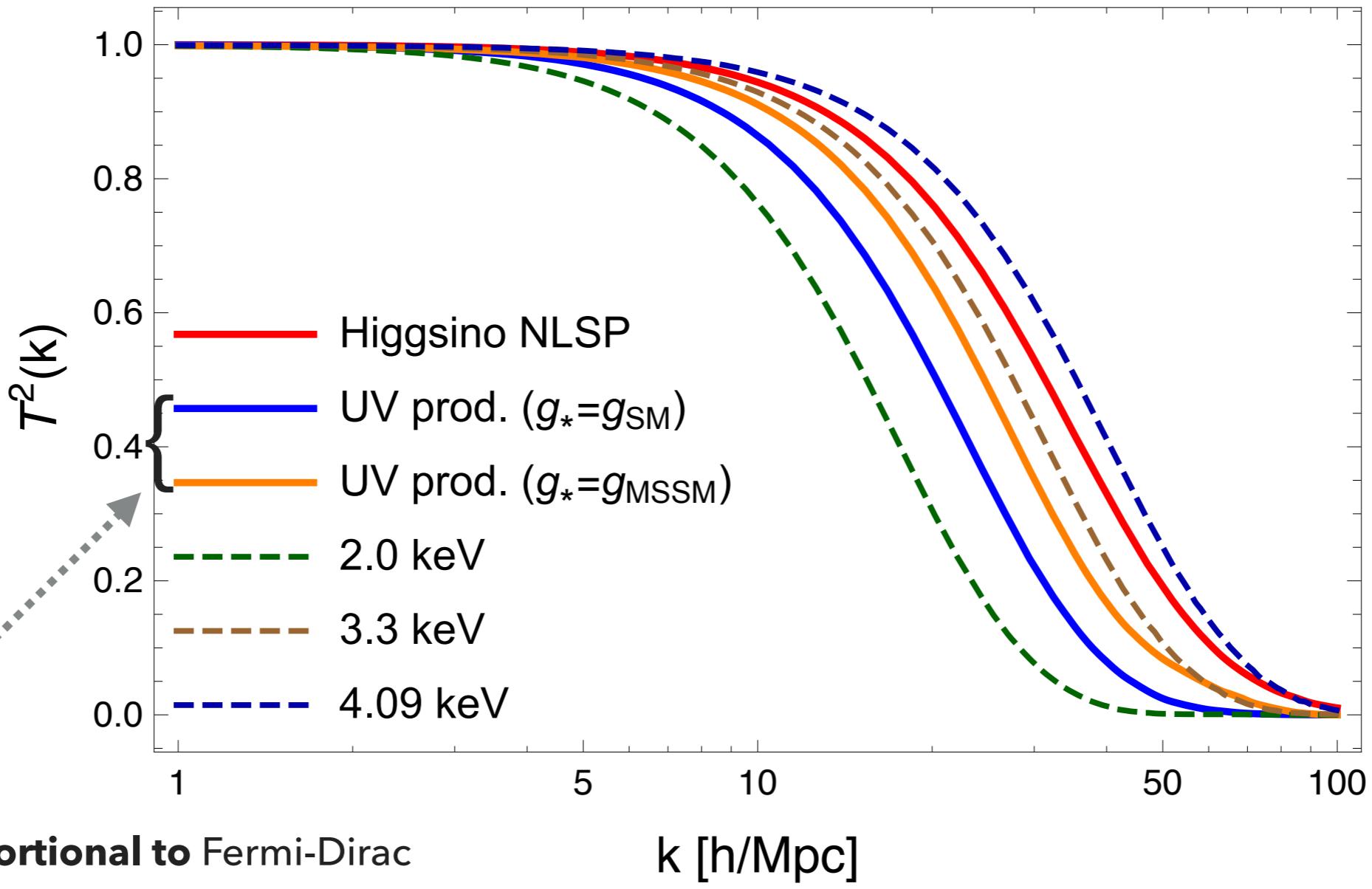
# Axino momentum distribution: a realistic model

- ▶ Axino distribution in a realistic model is *superposition* of several processes
- ▶ We study a benchmark point to confirm the tension with lyman-alpha forest



Dominant production channel is  $H \rightarrow \tilde{H} + \tilde{a}$

# Matter power spectrum: a realistic model



distribution **proportional to** Fermi-Dirac  
(for comparison)

- As expected,  $m_{\text{WDM}} > 4.09 \text{ keV}$  disfavors the benchmark point
- 7 keV axino DM has tension with the lyman-alpha observation

# Contents

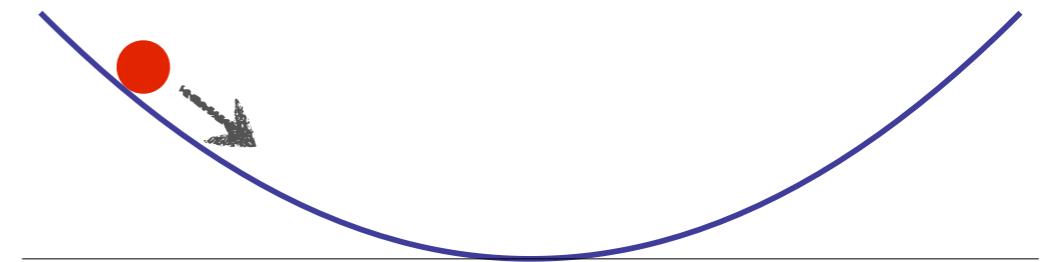
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# How to evade Ly-alpha constraint?

## - Entropy production

- ▶ **Saxion is produced by coherent oscillation**

$$Y_s^{\text{CO}} \simeq 1.9 \times 10^{-6} \left( \frac{\text{GeV}}{m_s} \right) \left( \frac{\min[T_R, T_s]}{10^7 \text{ GeV}} \right) \left( \frac{s_0}{10^{12} \text{ GeV}} \right)^2$$

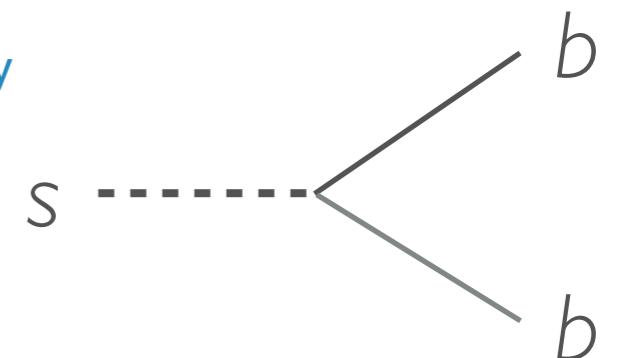


- ▶ **At some later time, saxion dominates the universe**

$$T_e^s = \frac{4}{3} m_s Y_s^{\text{CO}} \simeq 2.5 \times 10^2 \text{ GeV} \left( \frac{\min[T_R, T_s]}{10^7 \text{ GeV}} \right) \left( \frac{s_0}{10^{16} \text{ GeV}} \right)^2$$

- ▶ **Then saxion decays **after axino production** and heats up only SM plasma**

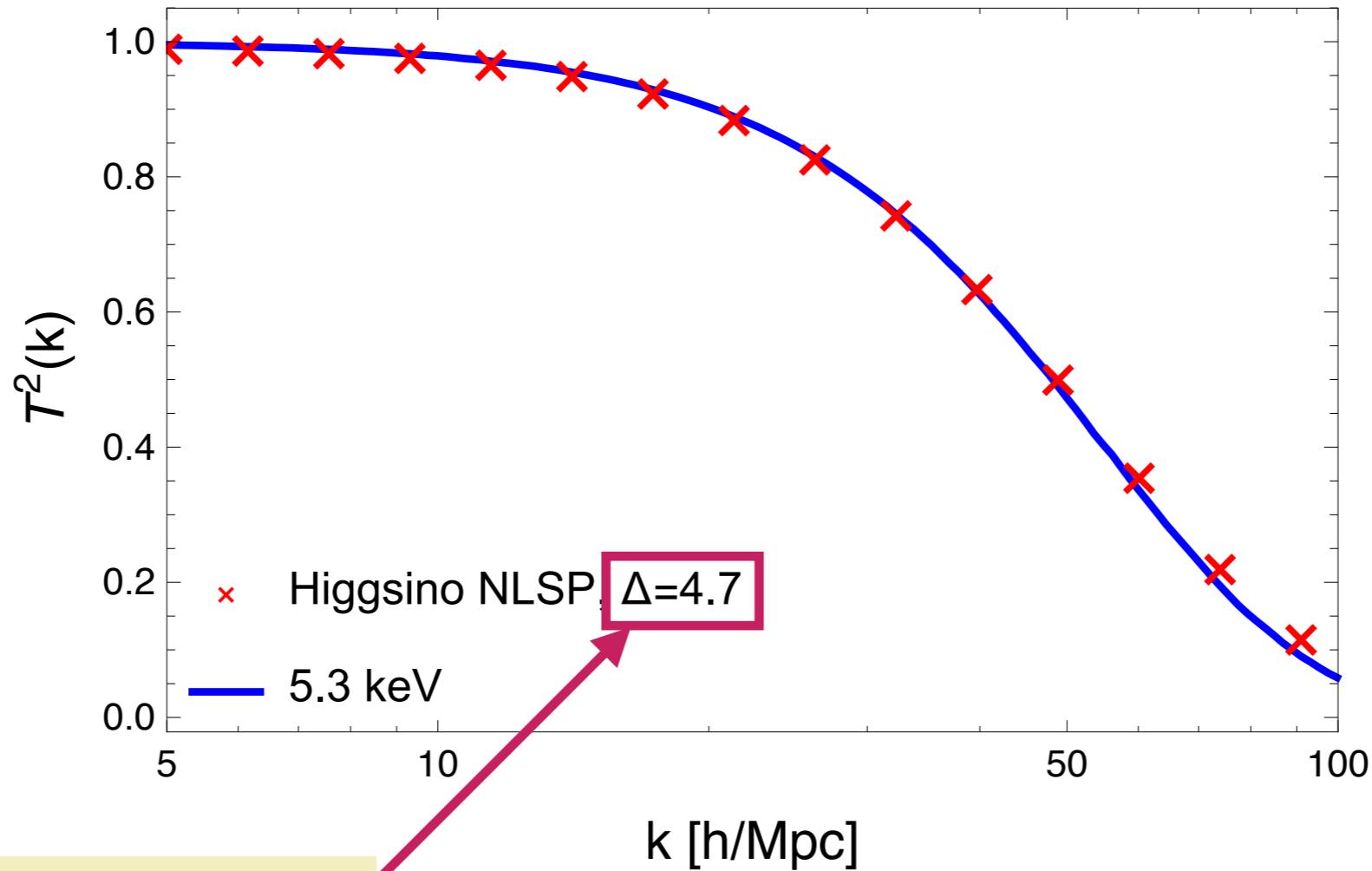
- Entropy production  $\Delta = \frac{s_f}{s_i} \simeq \frac{T_e^s}{T_D^s}$  Temperature after saxion decay
  - Axino temperature decreases  $T_{\tilde{a}} \rightarrow T_{\tilde{a}} \Delta^{-1/3}$



# How to evade Ly-alpha constraint?

## - Entropy production

- With  $\Delta=4.7$ , the strongest Ly-alpha constraint  $m_{\text{WDM}} > 5.3 \text{ keV}$  is evaded!



$v_{\text{PQ}} = 2.5 \times 10^{10} \text{ GeV}$   
 $m_s \simeq 110 \text{ GeV}$   
 $T_D^s \simeq 53 \text{ GeV}$

$k$  [h/Mpc]

DM abundance is also explained!

$$\Omega_{\tilde{a}} h^2 \simeq 0.1 \left( \frac{4.7}{\Delta} \right) \left( \frac{2.5 \times 10^{10} \text{ GeV}}{v_{\text{PQ}}} \right) \left( \frac{m_{\tilde{a}}}{7 \text{ keV}} \right)$$

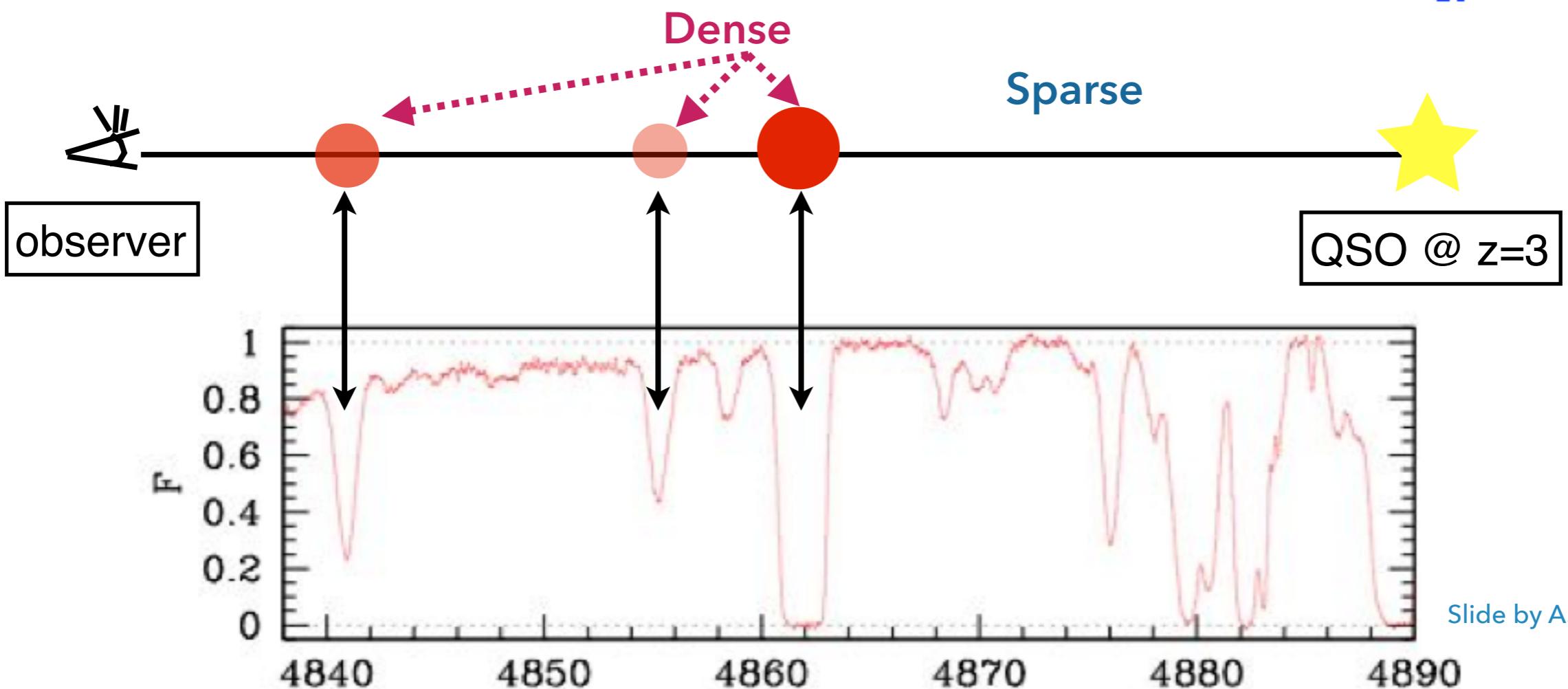
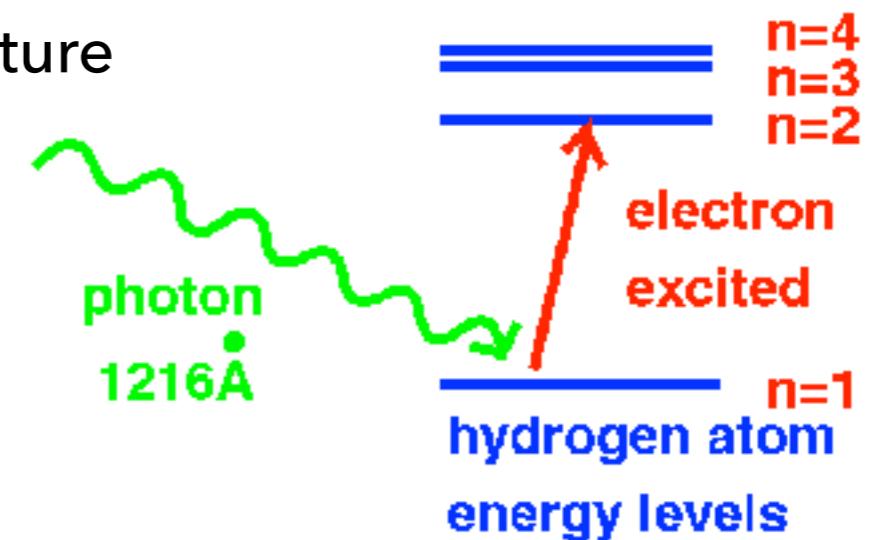
# Summary

- ▶ 7 keV DM (motivated by 3.5 keV excess) models should be tested by the lyman-alpha forest observation
- ▶ To do so, we need to calculate DM phase space distribution and resultant matter power spectrum
- ▶ 7 keV axino WDM (and possibly other proposed 7 keV DMs) has tensions with the recent strong lyman-alpha forest constraint  $m_{\text{WDM}} > 5.3 \text{ keV}$
- ▶ In SUSY axion model, inherent entropy production from saxion decay alleviates the tension

# **Backup**

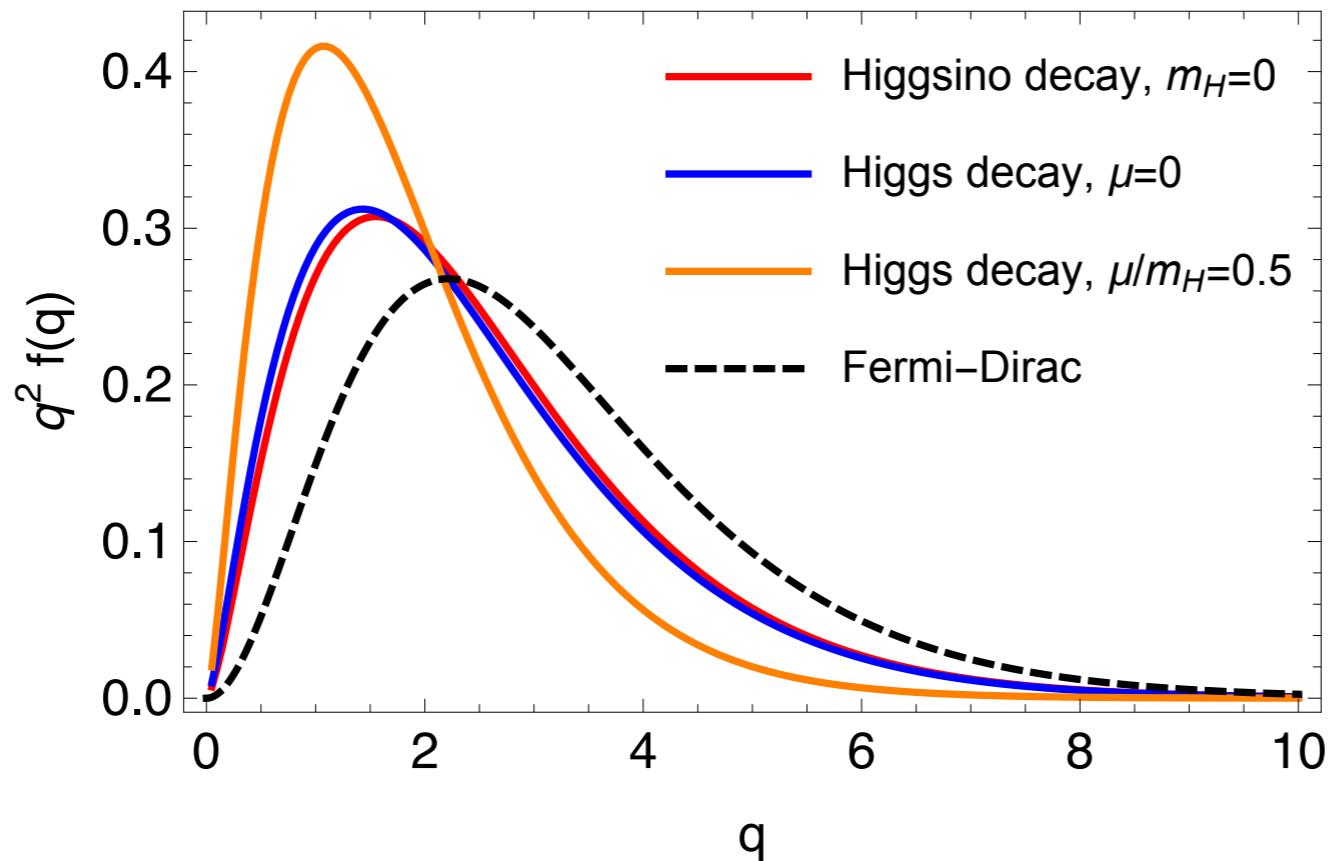
# Ly-a forest observation

- ▶ Important constraints on light DM from cosmic structure
- ▶ absorption intensity/frequency  
↔ HI distribution along the line-of-sight
- ▶ **Matter distribution** of the universe

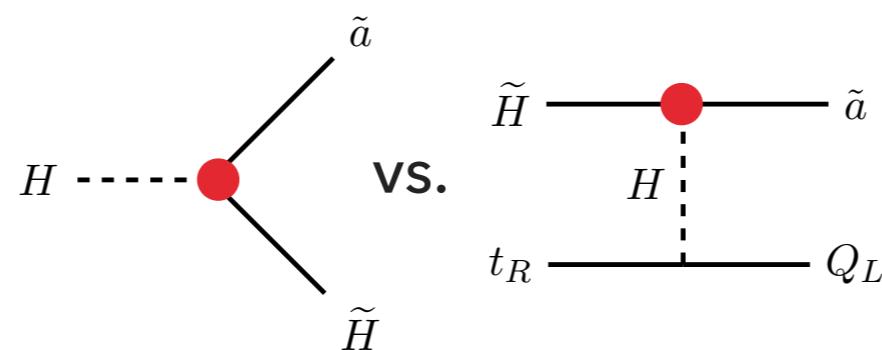


# How to evade Ly-alpha constraint?

## - Mass degeneracy



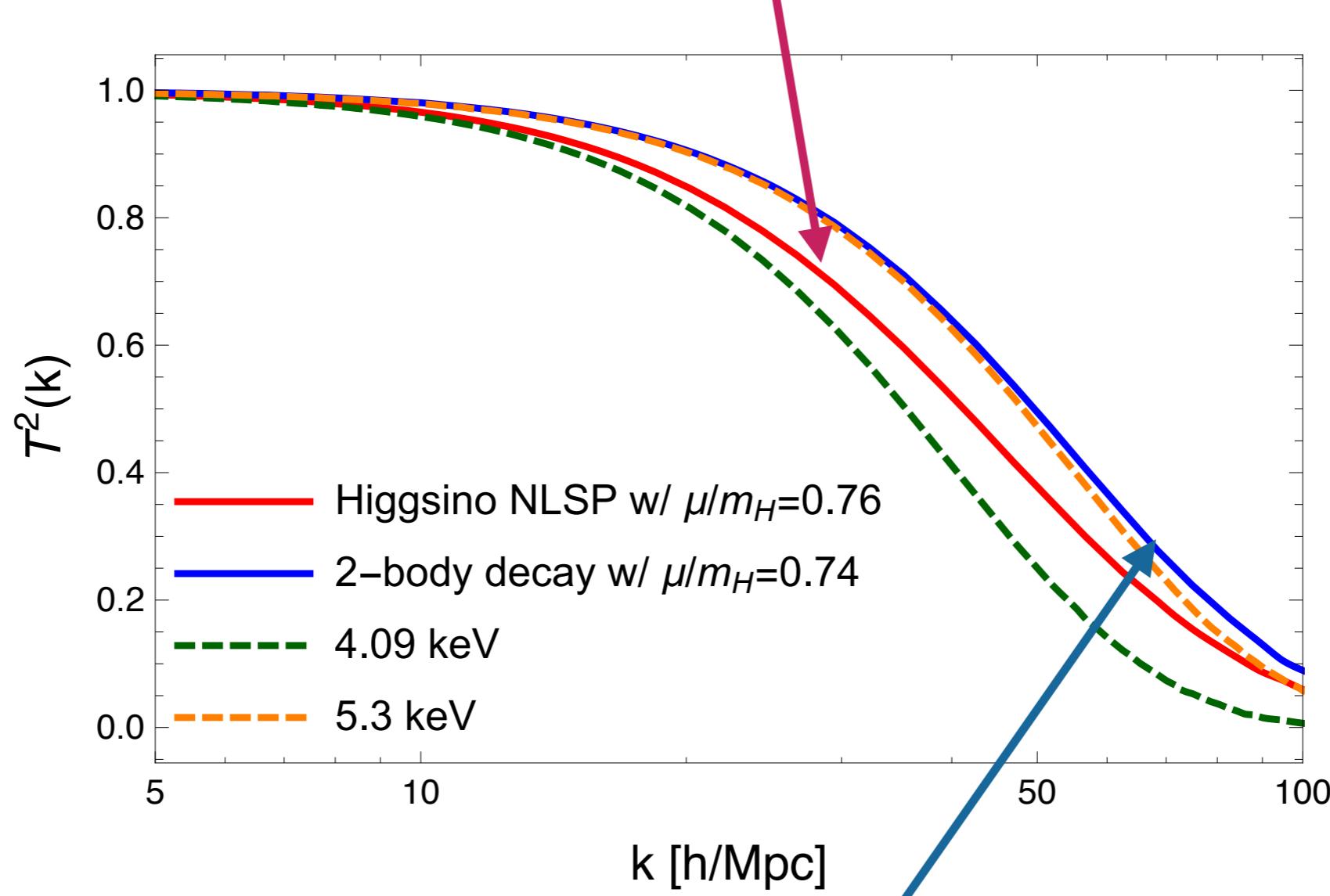
- ▶ Axino distribution is colder for compresses mass spectrum
- ▶ However, **scattering > 2-body decay** for very degenerate case



# How to evade Ly-alpha constraint?

## - Mass degeneracy

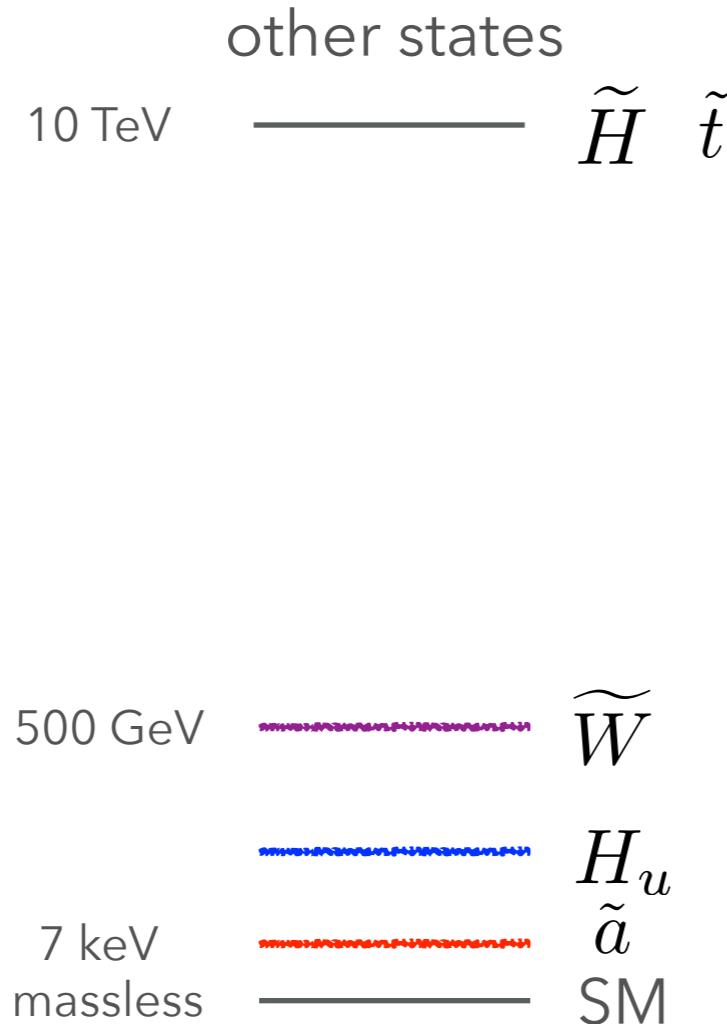
Realistic axino models w/ Higgsino NLSP cannot become colder than this line



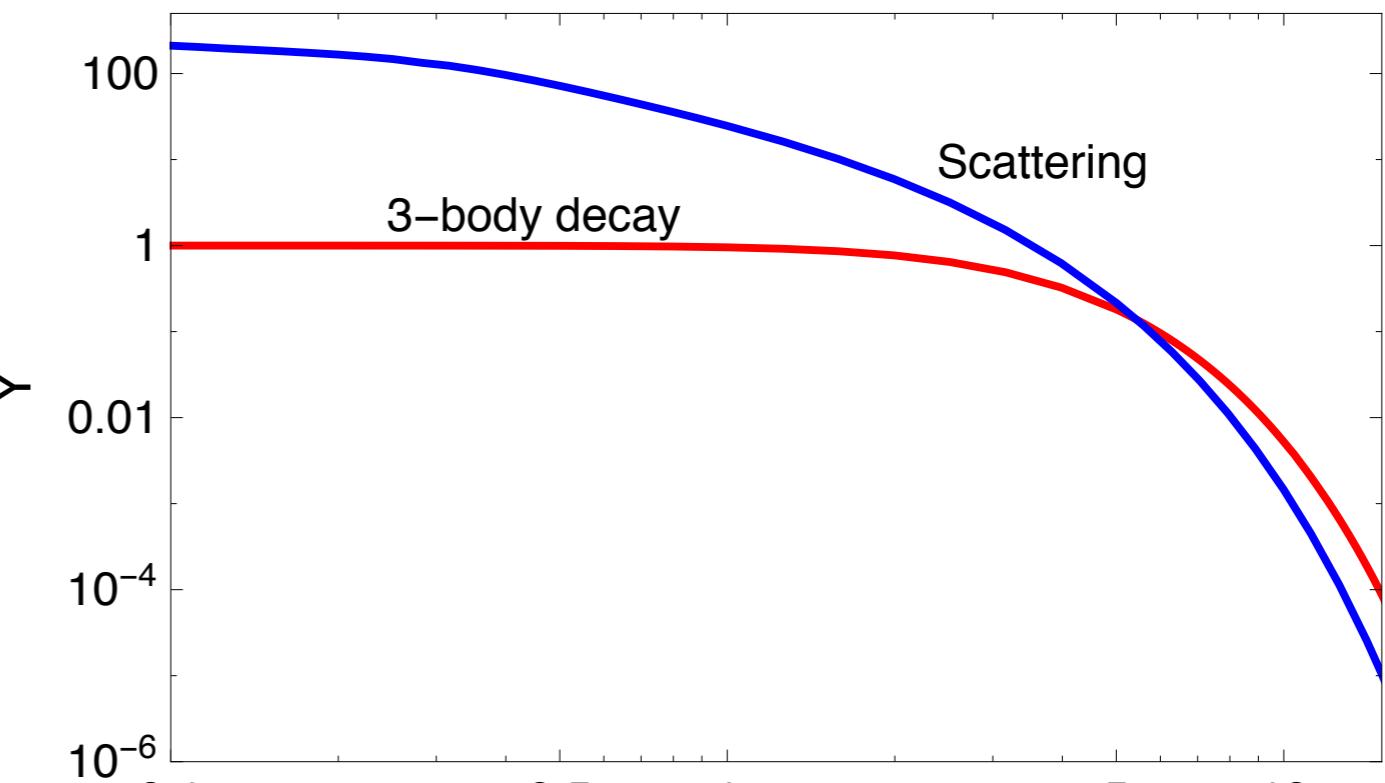
If only 2-body decay exists,  $m_{\text{WDM}} > 5.3 \text{ keV}$  can be evaded

# Another benchmark

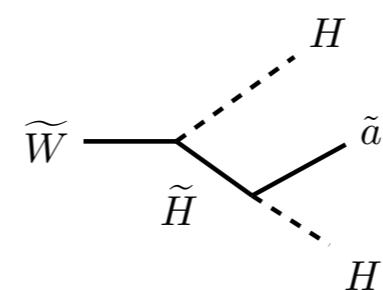
Wino NLSP



3-body decay vs. scattering



$M_2/T_R$



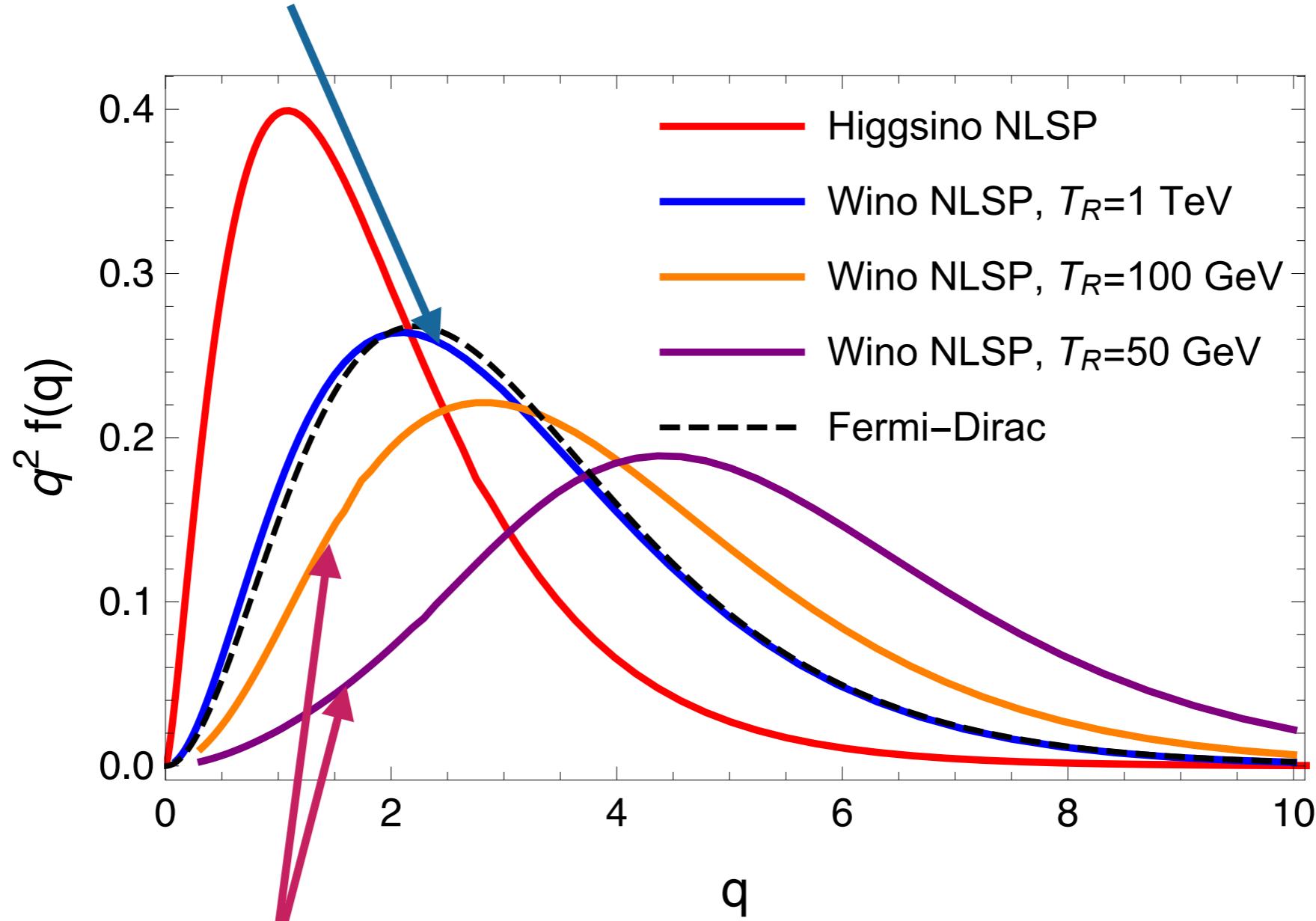
= dim. 5

→ efficient at UV

- ▶  $T_R > M_2$  : 2-to-2 scattering dominates
- ▶  $T_R < M_2/5$  : Wino 3-body decay dominates

# Another benchmark

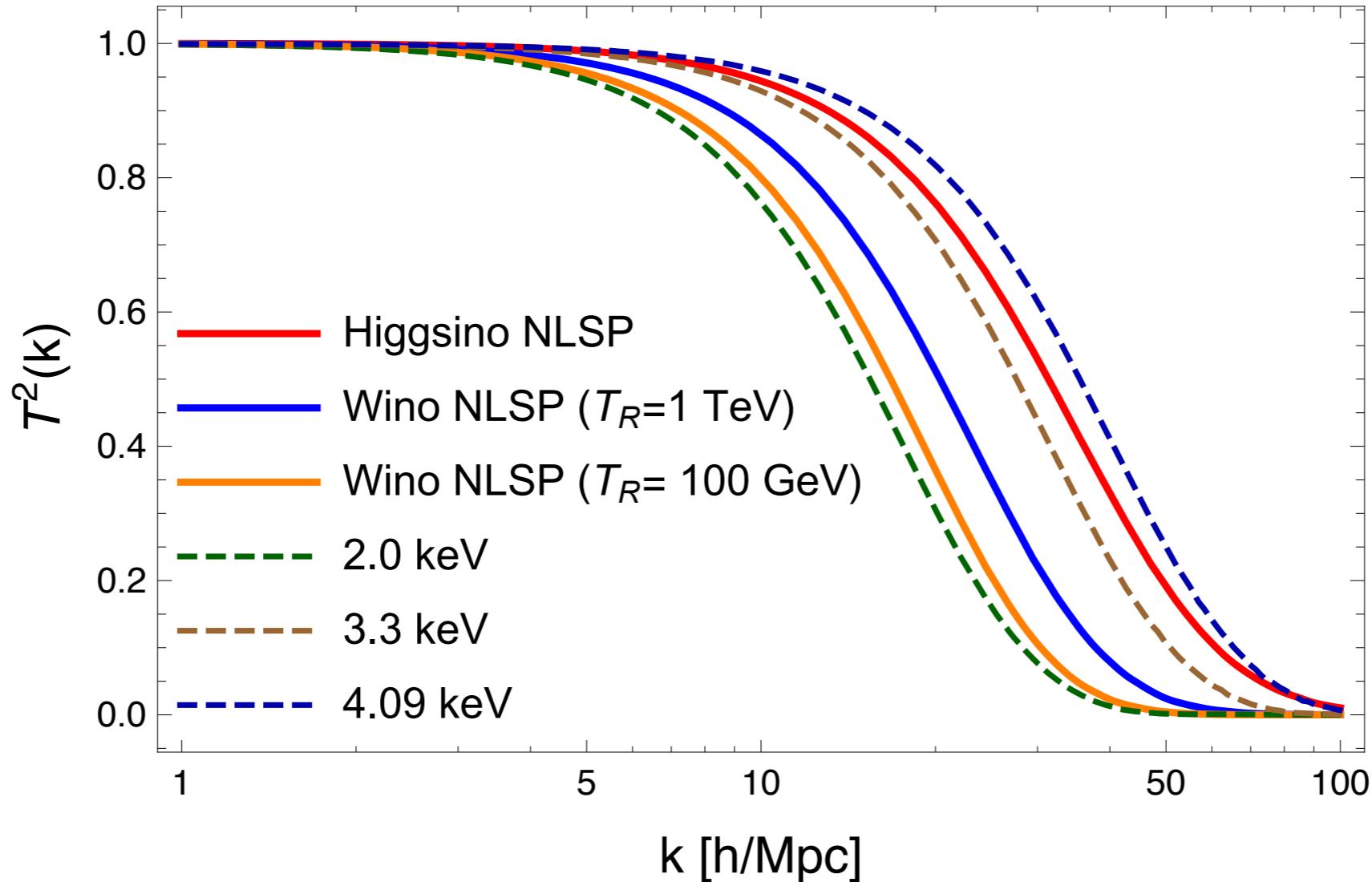
UV production produces thermal-like distribution



3-body decay dominant, but **hotter** than thermal one

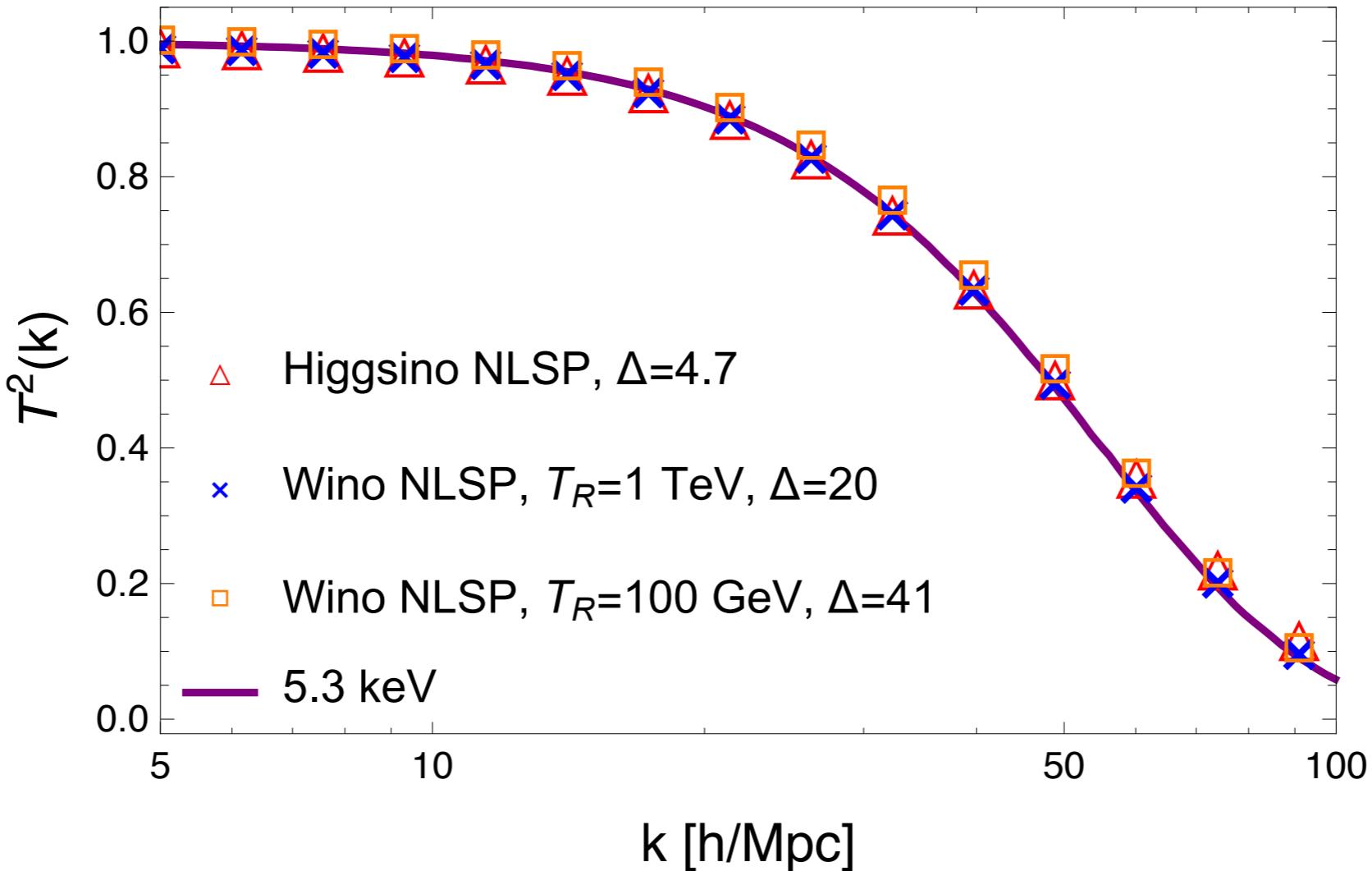
$T_R < M_2/5$  implies axino energy ( $\sim M_2/3$ ) is larger than temperature

# Another benchmark



- ▶ Wino NLSP cases are disfavored by  $m_{\text{WDM}} > 3.3$  keV

# Another benchmark



- ▶ Saxion decay alleviates the tension
- ▶ But  $\Delta$  is large and difficult to realize without spoiling BBN

$T_e^s \sim 2.5 - 25$  MeV (saxion domination)

- ▶ 110 GeV saxion decay into SM fermion pair, and dominantly into  $bb$
- ▶ Because it comes from mixing with Higgs, decay rate is proportional to yukawa
- ▶ Anomaly term for gluon comes from MSSM quark (squark) loop, so it is not proportional to reheating temperature unless it is below squark or gluino mass

# How can we apply the Ly-alpha constraints to our model?

- ▶ Two differences from conventional WDM when considering 7 keV axino

## 1. Phase space distribution is not always Fermi-Dirac (thermal)

Distribution is very **model-dependent** (e.g., freeze-in DM is never thermalized )

Free-streaming length is different for different distributions

## 2. DM “Temperature” is determined by $g_*$ at the DM decoupling

Typically,  $g_{\text{SM}} = 106.75$  or  $g_{\text{MSSM}} = 226.75$  is adopted

But bounds such as  $m_{\text{WDM}} > 5.3 \text{ keV}$  assume **very large  $g_*$**  (large entropy production)

$$\Omega_{\text{WDM}} h^2 = \left( \frac{m_{\text{WDM}}}{94 \text{ eV}} \right) \left( \frac{T_{\text{WDM}}}{T_\nu} \right)^3 = 7.5 \left( \frac{m_{\text{WDM}}}{7 \text{ keV}} \right) \left( \frac{106.75}{g_*^{\text{WDM}}} \right)$$

→  $g_{*, \text{WDM}} \sim 7000$

- ▶ In realistic 7 keV DM models, we cannot simply compare 7 keV to  $m_{\text{WDM}} > 5.3 \text{ keV}$

**We need to compare matter power spectrum!**