

Optimising NMR Spectroscopy through Method and Software Development

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Chapter 4

NOAH

chpt:noah

This final—and long—chapter describes my work on *NOAH* (NMR by Ordered Acquisition using ^1H detection) *supersequences*, pulse sequences which record multiple 2D datasets ('*modules*') in the time required for one. This is an attractive NMR technique for several reasons: the time savings are clearly a key factor, but the flexibility of being able to combine almost any set of modules also makes NOAH supersequences applicable to a variety of contexts.

I begin by introducing the concepts underlying NOAH supersequences, as well as a general discussion of the time savings (and sensitivity per unit time) benefits thus realised. I then describe the GENESIS (GENEration of Supersequences In Silico) website, which allows users to generate Bruker pulse programmes for almost every imaginable NOAH supersequence. After this, my work on various aspects of the actual sequences themselves is described, with a special focus on newly developed and/or improved modules. Finally, the design of 'parallel' supersequences which use interleaved and/or time-shared modules is discussed.

This work was done in close collaboration with Ēriks Kupče (Bruker UK). However, all results and analysis shown in this chapter are mine. The work in this chapter forms the subject of several publications:

- Yong, J. R. J.; Hansen, A. L.; Kupče, Ē.; Claridge, T. D. W. Increasing sensitivity and versatility in NMR supersequences with new HSQC-based modules. *J. Magn. Reson.* **2021**, 329, 107027, DOI: [10.1016/j.jmr.2021.107027](https://doi.org/10.1016/j.jmr.2021.107027)
- Kupče, Ē.; Yong, J. R. J.; Widmalm, G.; Claridge, T. D. W. Parallel NMR Supersequences: Ten Spectra in a Single Measurement. *JACS Au* **2021**, 1, 1892–1897, DOI: [10.1021/jacsau.1c00423](https://doi.org/10.1021/jacsau.1c00423)
- Yong, J. R. J.; Kupče, Ē.; Claridge, T. D. W. Modular Pulse Program Generation for NMR Supersequences. *Anal. Chem.* **2022**, 94, 2271–2278, DOI: [10.1021/acs.analchem.1c04](https://doi.org/10.1021/acs.analchem.1c04)

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- Yong, J. R. J.; Kupče, Ě.; Claridge, T. D. W. Uniting Low- and High-Sensitivity Experiments through Generalised NMR Supersequences. **2022**, manuscript in preparation

The material in the introductory sections also closely follow two reviews which I have contributed to:

- Kupče, Ě.; Frydman, L.; Webb, A. G.; Yong, J. R. J.; Claridge, T. D. W. Parallel nuclear magnetic resonance spectroscopy. *Nat. Rev. Methods Primers* **2021**, *1*, No. 27, DOI: [10.1038/s43586-021-00024-3](https://doi.org/10.1038/s43586-021-00024-3)
- Yong, J. R. J.; Kupče, Ě.; Claridge, T. D. W. In *Fast 2D solution-state NMR: concepts and applications*, Giraudeau, P., Dumez, J.-N., Eds., forthcoming, 2022

4.1 Introduction

The characterisation of small molecules and biomolecules by NMR spectroscopy relies on a suite of standard 2D NMR experiments, which seek to detect heteronuclear scalar couplings (e.g. HSQC and HMBC), homonuclear scalar couplings (e.g. COSY and TOCSY), or through-space interactions (e.g. NOESY and ROESY). Although 2D experiments provide far superior resolution and information content compared to 1D spectra, they also require substantially longer experiment durations, as the indirect dimension must be constructed through the acquisition of many t_1 increments. This problem is further exacerbated by the fact that structural elucidation or verification often necessitates the acquisition of several different 2D experiments.

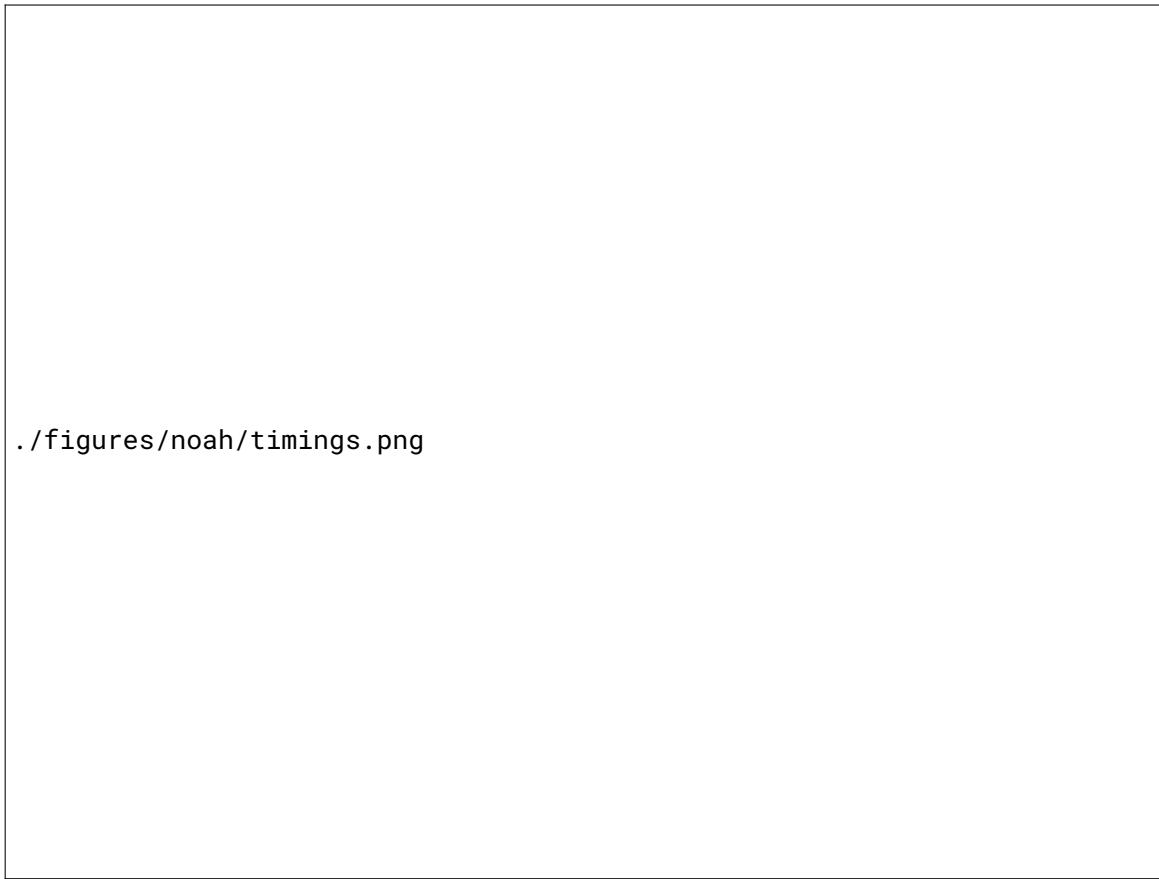
The acceleration of 2D NMR has thus proven to be a popular area of research. We may broadly categorise existing techniques into two classes: firstly, those which seek to directly speed up the acquisition of *individual* 2D spectra, and secondly, *multiple-FID* experiments which aim to collect two or more 2D spectra in the time required for one.* The former category includes methods such as non-uniform sampling (NUS),^{7–10} fast pulsing (i.e. shortening of recovery delays),^{11–14} ultrafast NMR,^{5,15–19} Hadamard encoding,^{20,21} and spectral aliasing;^{22–24} whereas the latter encompasses time-shared NMR,^{25,26} multiple-receiver NMR,^{27–29} and—of course—NOAH supersequences.^{5,30,31}

The scope of this introductory section will be limited to only NOAH supersequences. However, many of these techniques are closely related, and I will introduce concepts from elsewhere as needed. It should be noted that there are several other multiple-FID experiments which, while not explicitly advertised as such, are conceptually identical to NOAH experiments.^{32–35} I do not discuss these here.

4.1.1 Time savings and sensitivity analyses

In a typical 2D NMR experiment, the majority of the experiment duration is taken up by the *recovery delay*—the time required for spins to return to their equilibrium polarisation, such that the next transient or t_1 increment can be recorded. The removal (or shortening) of recovery delays is thus a very effective way of speeding up 2D data acquisition. In NOAH supersequences, 2D experiments ('modules') can be directly concatenated without the addition of extra recovery delays between them: only one overall recovery delay is required for the entire supersequence. This means that, to a first approximation, a supersequence containing N modules ($N \geq 2$) can be acquired in the time needed for just one module. Figure 4.1 shows an example of a NOAH supersequence formed from two modules (HSQC and COSY): the various timings referred to in

*These are by no means mutually exclusive: many of the techniques here can be combined to provide even greater efficiency.



`./figures/noah/timings.png`

Figure 4.1: (a) NOAH-2 SC supersequence, comprising HSQC and COSY modules (see table 4.1 for an explanation of the single-letter module codes used). (b) ‘Conventional’ echo-antiecho HSQC (the same as in fig. 1.7). (c) ‘Conventional’ COSY. The timings referred to in the text are highlighted for all three experiments; I have assumed that d_1 for each experiment is the same. Note that the lengths are not to scale: d_1 is typically far longer than the τ_{ps} and τ_{acq} .

the text which follows are also marked on the diagram.

The duration of an NMR experiment, τ_{exp} , can be expressed as a sum of its parts:

$$\tau_{\text{exp}} = \tau_{\text{ps}} + \tau_{\text{acq}} + d_1, \quad (4.1) \quad \{\text{eq:exp_duration_2d}\}$$

where τ_{ps} is the time required for the pulse sequence itself (typically several milliseconds), τ_{acq} is the acquisition time (several hundred milliseconds), and d_1 is the recovery delay (one or more seconds). The *time-saving factor* ρ_t for a NOAH supersequence, as compared to a series of conventional standalone experiments, is then:

$$\rho_t = \frac{\sum_i \tau_{\text{conv}}^{(i)}}{\tau_{\text{NOAH}}} = \frac{\sum_i (\tau_{\text{ps}}^{(i)} + \tau_{\text{acq}}^{(i)} + d_1^{(i)})}{d_1 + \sum_i (\tau_{\text{ps}}^{(i)} + \tau_{\text{acq}}^{(i)})}, \quad (4.2) \quad \{\text{eq:rho_t}\}$$

where τ_{NOAH} is the duration of the NOAH experiment, τ_{conv} is the duration of a conventional

experiment, and the superscript (i) represents the i -th module or conventional experiment being acquired. The sum runs from $i = 1$ to N , where N is the number of modules. If we assume that $d_1^{(i)} = d_1$ is the same for all N conventional experiments and the supersequence, then in the limit where

$$d_1 \gg \sum_i \tau_{\text{ps}}^{(i)} + \tau_{\text{acq}}^{(i)}, \quad (4.3) \quad \{\text{eq:d1_limit}\}$$

we have that $\rho_t \rightarrow Nd_1/d_1 = N$. This analysis makes plenty of assumptions, and is not entirely valid in practice. For example, *each* τ_{acq} is often around 5–10% of d_1 , so is not entirely negligible, especially in longer supersequences. Furthermore, some modules require longer τ_{ps} : most notable is the NOESY module, which contains a mixing time of several hundred milliseconds. (HMBC, TOCSY, and ROESY spectra are also lesser offenders.) These factors serve to reduce ρ_t from its idealised value of N ; generally, this deviation is larger as N increases, because eq. (4.3) becomes less and less valid. Nevertheless, the general point that time savings are approximately proportional to N stands.

For relatively concentrated samples, where sensitivity is not an issue, we can in fact end the discussion here. In this *sampling-limited regime*, the minimum 2D experiment duration is dictated purely by the number of t_1 increments needed to obtain sufficient resolution in the indirect dimension, as well as the minimum phase cycle required for artefact suppression.* NOAH supersequences are identical to conventional experiments in both aspects, but provide a time-saving factor of $\rho_t \sim N$.

The development of modern NMR instrumentation, including high-field magnets and cryogenic probes, means that the sampling-limited regime continues to be extended to ever lower concentrations. However, it is often not this simple: the opposite *sensitivity-limited regime* is still very commonly encountered, for example with naturally insensitive experiments (e.g. ADEQUATE), low-field benchtop NMR, or most simply, dilute samples.[†]

In such cases, it becomes mandatory to compare the SNRs of the NOAH modules and conventional experiments. To do so, we define for each module an *SNR factor* $A^{(i)}$, which is the SNR of the NOAH module divided by the SNR of a conventional experiment, acquired with the same parameters.[‡] In general, we have that $A \leq 1$, because NOAH modules frequently contain small modifications from conventional experiments (as will be explained in § 4.1.2). The *gain in*

*With modern gradient-enhanced experiments, the minimum phase cycle may well not even be a ‘cycle’; see also fig. 1.8.

[†]If the SNR factor $A^{(i)}$ as discussed below is *very small*, then it is possible that even concentrated samples may be shifted into the sensitivity-limited regime. This is never really the case in practice, though, as will be shown in § 4.1.3.

[‡]The relative SNR will likely vary from peak to peak in the spectrum, and $A^{(i)}$ should in theory be quoted either as an average over all peaks, or as a range. This is what I have done in this thesis. However, comparisons in the literature are not always as thorough.

sensitivity per unit time, $\varepsilon^{(i)}$, is then defined by

$$\varepsilon^{(i)} = A^{(i)} \sqrt{\rho_t}, \quad (4.4) \quad \{eq:varepsilonpsilon_i\}$$

where the square root accounts for the fact that SNR scales only as the square root of the number of scans, or the number of times the experiment can be repeated in a given period. Of course, the exact values calculated for $A^{(i)}$ (and hence $\varepsilon^{(i)}$) will depend on the sample chosen for the comparison. These values should therefore be assumed to be valid only for similar samples.

If $\varepsilon^{(i)} > 1$, as is frequently the case, this means that the NOAH supersequence provides greater sensitivity per unit time in the i -th module compared to a standalone experiment. Equivalently, performing a NOAH experiment allows data of sufficient sensitivity to be obtained in less time. Naturally, this condition is most important for modules which are inherently insensitive, particularly heteronuclear correlation modules. For sensitive (typically homonuclear) modules, it is often perfectly tolerable to have $A < 1$ or even $\varepsilon < 1$, as even with this sensitivity penalty they are still more intense than the heteronuclear modules.

Another issue with NOAH supersequences is that each module is run with the same number of scans (phase cycle). Although this was touted as a benefit in the sampling-limited regime, this may in fact be undesirable in the sensitivity-limited regime, where insensitive experiments need to be run with more scans than sensitive ones. In this case, the *effective* time savings provided by NOAH experiments are smaller:

$$\rho_{t,\text{eff}} = \frac{\sum_i \tau_{\text{conv}}^{(i)}}{\tau_{\text{NOAH}}} = \frac{\sum_i S^{(i)}(\tau_{\text{ps}}^{(i)} + \tau_{\text{acq}}^{(i)} + d_1^{(i)})}{Sd_1 + S \sum_i (\tau_{\text{ps}}^{(i)} + \tau_{\text{acq}}^{(i)})}, \quad (4.5) \quad \{eq:rho_t_eff\}$$

where each standalone experiment is acquired with $S^{(i)}$ scans and the NOAH experiment with S scans. Typically, S is simply the largest of the $S^{(i)}$. If $S^{(i)} = S$ for all i , then eq. (4.5) simply reduces to eq. (4.2); on the other hand, if the $S^{(i)}$'s are different, then we have that $\rho_{t,\text{eff}} < \rho_t$. In such a situation, it is probably more appropriate to describe a NOAH supersequence as ‘measuring the most insensitive module and getting the others for free’. Indeed, if $S = S^{(i)} \gg S^{(j \neq i)}$, then ‘the other’ modules require almost no time to measure (relative to the least sensitive module), and $\rho_{t,\text{eff}}$ tends towards 1, meaning that even the time-saving utility of NOAH vanishes. A corollary of this is that NOAH supersequences are generally best constructed from modules which have similar intrinsic sensitivities and hence similar $S^{(i)}$'s.

As the reader can no doubt appreciate by now, the comparison of NOAH and conventional spectra is fraught with subtleties (which are sometimes glossed over in the literature, but invariably surface in real-life discussions). In fact, it is hardly even difficult to construct yet more edge cases. For example, one may not want to acquire all the individual spectra ‘conventionally’: for example,

NUS may be used for a HSQC experiment but not for others; or d_1 may be varied for different experiments. These will have an impact on both the durations of the experiments, as well as their sensitivities. To make any meaningful quantitative comparisons, it is therefore necessary to restrict the discussion to values of ρ_t , A , and ε , which can be objectively calculated. This is the approach I have taken in this thesis. These should of course be read with the qualitative understanding that depending on the context, these aforementioned factors may lead to *some*—but never a *complete*—decrease in the utility of NOAH experiments.

4.1.2 Magnetisation pools

subsec:noah_magpools Having gotten this relatively dry material out of the way, I now turn to exactly how NOAH supersequences are constructed. Ordinarily, if the recovery delay is removed from an NMR experiment, its sensitivity will be greatly reduced because insufficient magnetisation will have recovered between repetitions; or in other words, $A^{(i)}$ will be very small. Such experiments would only really be useful well in the sampling-limited regime.

The key to avoiding this in NOAH supersequences is to make sure that *each module samples a different source of magnetisation*. For example, a HSQC module can be designed to only sample magnetisation of protons directly bonded to the 1.1%-natural abundance ^{13}C , and leave all other proton magnetisation untouched. Immediately following this, the remainder of the proton magnetisation can then be used to record (say) a COSY module, without needing a separate recovery delay. Using the notation of Orts and Gossert,³⁶ the magnetisation of ^{13}C -bound protons is denoted as $^1\text{H}^{\text{C}}$, and the magnetisation of protons *not* bonded to ^{13}C is denoted as $^1\text{H}^{\text{!C}}$. Protons not directly bonded to NMR-active heteronuclei are labelled $^1\text{H}^{\text{!X}}$, and will often be referred to as ‘bulk’ magnetisation, since (in natural-abundance samples) the majority of protons fall into this category.

Most standard 2D experiments do not preserve unused magnetisation but instead dephase it through CTP gradient selection; thus, NOAH modules often require some modifications compared to standard experiments. For example, compared to the echo–antiecho HSQC (discussed in § 1.4.5), the NOAH HSQC module³⁰ adds an extra CTP gradient so that the bulk magnetisation is refocused and ultimately returned to the $+z$ equilibrium state (fig. 4.2a). (This is largely identical to the ‘symmetrised’ ASAP-HSQC experiment.³⁷)

Sometimes, the modifications required are more extensive, as in the HMBC module. If this module is followed by a HSQC module (or any other module which draws on $^1\text{H}^{\text{C}}$ magnetisation), the initial 90° excitation pulse must be replaced with a zz -filter (fig. 4.2b). This performs an *isotope-selective rotation* in that $^1\text{H}^{\text{C}}$ magnetisation is stored along the z -axis, but $^1\text{H}^{\text{!C}}$ magnetisation is excited (and subsequently detected). In general, sequences which are thus modified have lower



`./figures/noah/hsqc_hmhc_prodops.png`

`fig1goabbbbspopb`

Figure 4.2: (a) NOAH HSQC module. (b) NOAH HMBC module. The 90° pulse highlighted in red is described in § 4.3.7. Delays are set as: $\Delta = 1/(4 \cdot {}^1J_{\text{CH}})$; $\Delta_{\text{LR}} = 1/(2 \cdot {}^nJ_{\text{CH}})$; $\tau_1 = 1/(2 \cdot {}^1J_{\text{CH,max}})$; $\tau_2 = 1/(2 \cdot {}^1J_{\text{CH,min}})$ (see also § 3.4.7 for the LPJF). Phase cycling is performed with $\phi_1 = (x, -x)$, $\phi_2 = (x, x, -x, -x)$, and $\phi_{\text{rec}} = (x, -x, -x, x)$. Gradient amplitudes are $(g_1, g_2, g_3, g_4, g_5) = (80\%, \pm 40.2\%, 15\%, -10\%, -5\%)$. Product operator analysis is provided above both modules for both the ${}^1\text{H}^{\text{C}}$ and ${}^1\text{H}^{\text{!C}}$ magnetisation pools; the notation for this is explained in the *Preface*.

sensitivities (i.e. $A < 1$) than the ‘original’ sequences from which they were derived. This is partly because of imperfect manipulation of magnetisation by the extra pulse sequence elements, and also increased losses due to relaxation during these extended sequences.

In contrast, modules placed towards the *end* of a supersequence do not need to be modified, as they do not need to preserve any magnetisation. This includes virtually all homonuclear modules, which are allowed to simply consume any remaining magnetisation. Although this makes their implementation very straightforward, in general these modules will *also* suffer some losses in sensitivity, because the preceding modules do not perfectly retain all magnetisation.

Thus, in general, it is not possible for any module in a NOAH supersequence to have $A = 1$,

unless it is placed first in the supersequence *and* has not undergone any modifications.* Such cases are very rare, and it is thus necessary to accept some decreases in A , which are often fairly small (on the order of 10–20%). In the sampling-limited regime, sensitivity is not at a premium and this is often perfectly tolerable. In the sensitivity-limited regime, the full time savings ρ_t cannot be realised, but since ε is still typically larger than 1, there is still an overall boost in sensitivity per unit time.

4.1.3 Case studies

Using all that has been described in the previous sections, we now look at a few ‘typical’ supersequences to understand their construction. A quick note about the nomenclature of NOAH supersequences is warranted here. Supersequences are labelled by the number of modules N , plus a series of single-letter codes corresponding to the identity and ordering of the modules involved (table 4.1). Occasionally, superscripts or subscripts are used to qualify the modules involved.[†] Thus, a NOAH supersequence containing three modules—say ^{15}N HMQC, ^{13}C HSQC, and CLIP-COSY—would be referred to as a NOAH-3 M_NSC^c. Table 4.2 provides values of ρ_t and A for each module of several typical supersequences, which will be rationalised in the text which follows.

NOAH-2 SC: HSQC + COSY

We begin with perhaps the simplest example of a NOAH supersequence, one containing the HSQC and COSY modules: this is labelled as a NOAH-2 SC experiment (entry 1, table 4.2). As shown in fig. 4.2a, the HSQC module only samples $^1\text{H}^{\text{C}}$ magnetisation, and leaves $^1\text{H}^{\text{!C}}$ magnetisation along the $+z$ -axis. Although the HSQC experiment has to be modified to preserve this $^1\text{H}^{\text{!C}}$ magnetisation, its sensitivity is practically unaffected as compared to a ‘standard’ HSQC ($A = 0.97$). Furthermore, the COSY module retains *most* of its sensitivity ($A = 0.90$). The small loss here is because the HSQC module does not *perfectly* preserve the $^1\text{H}^{\text{!C}}$ magnetisation: for example, evolution of J-couplings as well as relaxation occur during the HSQC pulse sequence, which are ignored in the product operator analysis in fig. 4.2a.

The value of the time-saving factor, $\rho_t = 1.87$, is very close to the theoretical limit of $N = 2$. This reflects the fact that the pulse sequence itself, τ_{ps} , is fairly short for both the HSQC and COSY

*Of course, this also depends on exactly *what* standalone experiment the NOAH supersequence is being compared against. Sometimes, in the literature, the NOAH experiment has been compared against its constituent modules acquired in a standalone fashion; in this case, the first module will always have $A = 1$. This tells us how much we gain through the act of concatenating modules, but is less meaningful in the ‘real world’ where one is interested in how useful NOAH is relative to ‘typical’ optimised 2D experiments. I therefore prefer to make comparisons against standard-library sequences.

[†]With the increasing number of modules, and the variety of modern NMR experiments which could be incorporated into NOAH supersequences, keeping these abbreviations short yet meaningful has been a challenge.

$^1\text{H}-^{15}\text{N}$ modules		$^1\text{H}-^{13}\text{C}$ modules		$^1\text{H}-^1\text{H}$ modules	
Module	Code	Module	Code	Module	Code
HMQC	M_N	HSQC	S	COSY	C
HSQC	S_N	seHSQC	S^+	CLIP-COSY	C^c
seHSQC	S_N^+	HSQC-TOCSY	S^T	DQF-COSY	C
HMBC	B_N	HSQC-COSY	S^C	TOCSY	T
		2BOB	O	NOESY	N
		HMBC	B	ROESY	R
		ADEQUATE	A	PSYCHE	P
				TSE-PSYCHE	P^T
				PSYCHE 2DJ	J

tbl:noah_modules

Table 4.1: A (non-exhaustive) list of single-letter module codes for a selection of NOAH modules. Note that, in the literature, the ^{15}N HMQC module has been referred to simply by ‘M’, since the HSQC module is preferred for $^1\text{H}-^{13}\text{C}$ correlations. In this thesis, I include the subscript N throughout to avoid any ambiguity.

Entry	Sequence	τ_{NOAH}	ρ_t	A				
				HMBC	seHSQC	HSQC	COSY	TOCSY
1	SC	15 min 0 s	1.87			0.97	0.90	
2	SCT	16 min 25 s	2.60			1.01	0.99	0.79
3	BS	15 min 40 s	1.82	0.93		0.87		
4	SB	15 min 35 s	1.83	0.99		0.96		
5	BSCT	17 min 48 s	3.22	0.95		0.90	0.36	0.28
6	BS_N^+SCT	18 min 57 s	3.74	0.95	0.71	0.66	0.38	0.30
7	S_N^+BSCT	18 min 56 s	3.75	0.76	0.79	0.74	0.33	0.26

tbl:noah_sensitivities

Table 4.2: Sensitivity and time-saving analyses of several typical NOAH supersequences. All experiments were acquired with 2 scans per increment, 256 t_1 increments, an acquisition time of 67 ms, and a recovery delay of 1.5 s. The HMBC module used here includes the extra ^{13}C 90° pulse described later in § 4.3.7: this has no significant impact on the SNR, and is only mentioned as a technicality. The ^{15}N seHSQC module is that described in § 4.3.3. The CT module here was run with States indirect-dimension quadrature detection, and the individual C module (in entry 1) with echo-antiecho. The following Bruker library sequences were used as the ‘conventional’ experiments: `hmbcetgpl2nd`, `hsqcetf3gpsi2`, `hsqcetgpssp.2`, `cosygpqf`, and `dipsi2gpphzs`. Data code: 7Z-220224.

modules; the deviation therefore chiefly arises from the acquisition time, τ_{acq} . In all respects, this is therefore an example of an ‘ideal’ NOAH supersequence, where the combination of two modules provides time savings without compromising on sensitivity.

It is worth pointing out that the order of the modules cannot be reversed: the COSY module cannot be (easily) modified to preserve $^1\text{H}^C$ magnetisation. In a hypothetical NOAH-2 CS supersequence, the later HSQC module would only be able to use magnetisation recovered

during the COSY FID, leading to substantial sensitivity drops.



`./figures/noah/sc_noah_vs_conv.png`

sc_noah_vs_conv.png

Figure 4.3: (a) HSQC from NOAH-2 SC supersequence. (b) COSY from NOAH-2 SC supersequence. (c) Standalone HSQC. (d) Standalone COSY; off-diagonal artefacts are highlighted in the red box. Data code: 7Z-220224.

A final point to consider would be whether the NOAH data has comparable spectral quality in terms of (for example) artefacts. In this case, the answer is yes: the NOAH HSQC spectrum is virtually identical to the standalone (figs. 4.3a and 4.3c; both spectra have low-level artefacts of different kinds, which do not seriously impede the interpretation and are not shown). On the other hand, the NOAH COSY spectrum seems to actually *improve* on the standalone COSY, in that it better suppresses off-diagonal artefacts (figs. 4.3b and 4.3d). These artefacts likely arise in the standalone COSY because of accidental refocusing of magnetisation which has not completely relaxed between t_1 increments.³⁸ In contrast, the NOAH COSY module has an extra set of HSQC gradients between every repetition of the COSY, so accidental refocusing is less likely. (Similar artefacts have been noted before in the DQF-COSY experiment,^{39,40} and have also shown to

be attenuated in the corresponding NOAH module.⁴¹) That said, such improvements are not always guaranteed: there are sometimes artefacts which arise uniquely in NOAH experiments, some of which are discussed in the following sections.

NOAH-3 SCT: HSQC + COSY + TOCSY

Evidently, the fact that the HSQC preserves almost all $^1\text{H}^{1\text{C}}$ magnetisation means that *any* homonuclear module—or a combination thereof—can be placed after it. In general, since homonuclear modules tend to consume any remaining bulk magnetisation, it is very difficult to create combinations of homonuclear modules which do not lead to significant reductions in sensitivity. The only real exceptions are COSY/X combinations, where X can be NOESY, ROESY, or TOCSY: instead of concatenating the COSY and X modules, the COSY pulse sequence can instead be nested *within* the X module, as was first demonstrated with X = NOESY.^{42,43} Here, we use the COSY/TOCSY combination as an example.⁴⁴ The COSY, TOCSY, and combined COSY/TOCSY modules are shown in fig. 4.4.



ffggttstateset

Figure 4.4: (a) COSY module. (b) TOCSY module; zero-quantum suppression is employed before and after the isotropic mixing period. (c) Combined COSY/TOCSY module, where the COSY FID is acquired immediately before the TOCSY mixing.

As shown in table 4.2 (entry 2), this nesting of the COSY module does not materially affect the TOCSY sensitivity. A small loss of approximately 20% is observed, which is partly due to the imperfect magnetisation preservation by the HSQC, and perhaps also due to relaxation during the COSY acquisition period. As for the time-saving factor, a slightly larger deviation ($\rho_t = 2.60$) is observed from the ideal value of 3. This reflects the addition of a TOCSY mixing period, which contributes to τ_{ps} .

NOAH-2 BS: HMBC + HSQC

As mentioned previously, the HMBC module shown in fig. 4.2b is designed to retain $^1\text{H}^{1\text{C}}$ magnetisation through the addition of the zz -filter. This can be used in a subsequent HSQC module in a NOAH-2 BS supersequence. Entry 3 of table 4.2 shows that the addition of the

zz -filter to the HMBC causes a relatively small 7% decrease in sensitivity; on the other hand, the HSQC loses 13% of its sensitivity because of incomplete magnetisation preservation.

Generally, it has been recommended that less sensitive modules are placed earlier in the supersequence so that they can access a larger proportion of the equilibrium magnetisation. Since the HMBC is less sensitive of the two modules, this rule of thumb suggests that the BS supersequence would be better than the alternative SB supersequence. In fact, the opposite is true, as entry 4 of table 4.2 shows. The HSQC module has a boost in sensitivity because it is placed first in the supersequence, and no longer needs to rely on the $^1\text{H}^{\text{C}}$ magnetisation preserved by the HMBC; and the HMBC also benefits because the zz -filter modification is no longer needed.* Arguably, the ordering of modules in a supersequence should be considered on a case-by-case basis.

NOAH-4 BSCT: HMBC + HSQC + COSY + TOCSY

We now move on to a longer supersequence containing four modules, with a correspondingly larger ρ_t value of 3.22. The sensitivity of the HSQC module is practically the same as in the NOAH-2 BS supersequence just described: however, the COSY and TOCSY modules expose one weakness of the HMBC module which has so far been overlooked. In principle, the HMBC module should only excite magnetisation of protons which are long-range coupled to ^{13}C (which we could, for example, denote as $^1\text{H}^{\text{C(LR)}}$). This magnetisation pool should be separate from both the directly coupled protons ($^1\text{H}^{\text{C}}$), as well as protons which are not coupled to any ^{13}C at all ($^1\text{H}^{\text{!C}}$). Unfortunately, this is not the case: it is not actually possible to separate the $^1\text{H}^{\text{C(LR)}}$ and $^1\text{H}^{\text{!C}}$ magnetisation pools. The HMBC excites both of these magnetisation sources, dephases the latter using CTP gradients, and detects the signal arising from the former.

What this means, of course, is that the COSY/TOCSY module which rely on $^1\text{H}^{\text{!C}}$ magnetisation will have substantially lower sensitivities. The signal detected in these two modules derives only from whatever has recovered during the previous two acquisition periods, as shown in entry 5 of table 4.2: A for COSY and TOCSY is 0.36 and 0.28 respectively. That said, this is in fact not likely to be an issue *even* for sensitivity-limited samples. Because the intrinsic sensitivity of the HMBC is orders of magnitude lower than the COSY and TOCSY, even with these large losses in sensitivity, the COSY and TOCSY spectra still have greater intensities than the HMBC. Thus, as long as the entire supersequence is acquired with enough scans to make the HMBC SNR sufficient, the SNR in the COSY and TOCSY will *also* be acceptable. This is illustrated in fig. 4.5.

A rather more insidious problem is that different signals relax at different rates: thus, the COSY and TOCSY spectra (or indeed, any homonuclear module) will have uneven intensities and are

*In fact, the final 180° pulse in the HMBC module could also be removed: this is likely to give a further boost in SNR, as discussed in § 4.3.7. However, this was not done here.

frequently asymmetric. This can be seen in the COSY spectrum, where a pair of asymmetric crosspeaks are highlighted. Adding a period of isotropic mixing before the COSY module⁴⁵ can help to ameliorate this to some extent (this was not performed when acquiring the data in fig. 4.5).

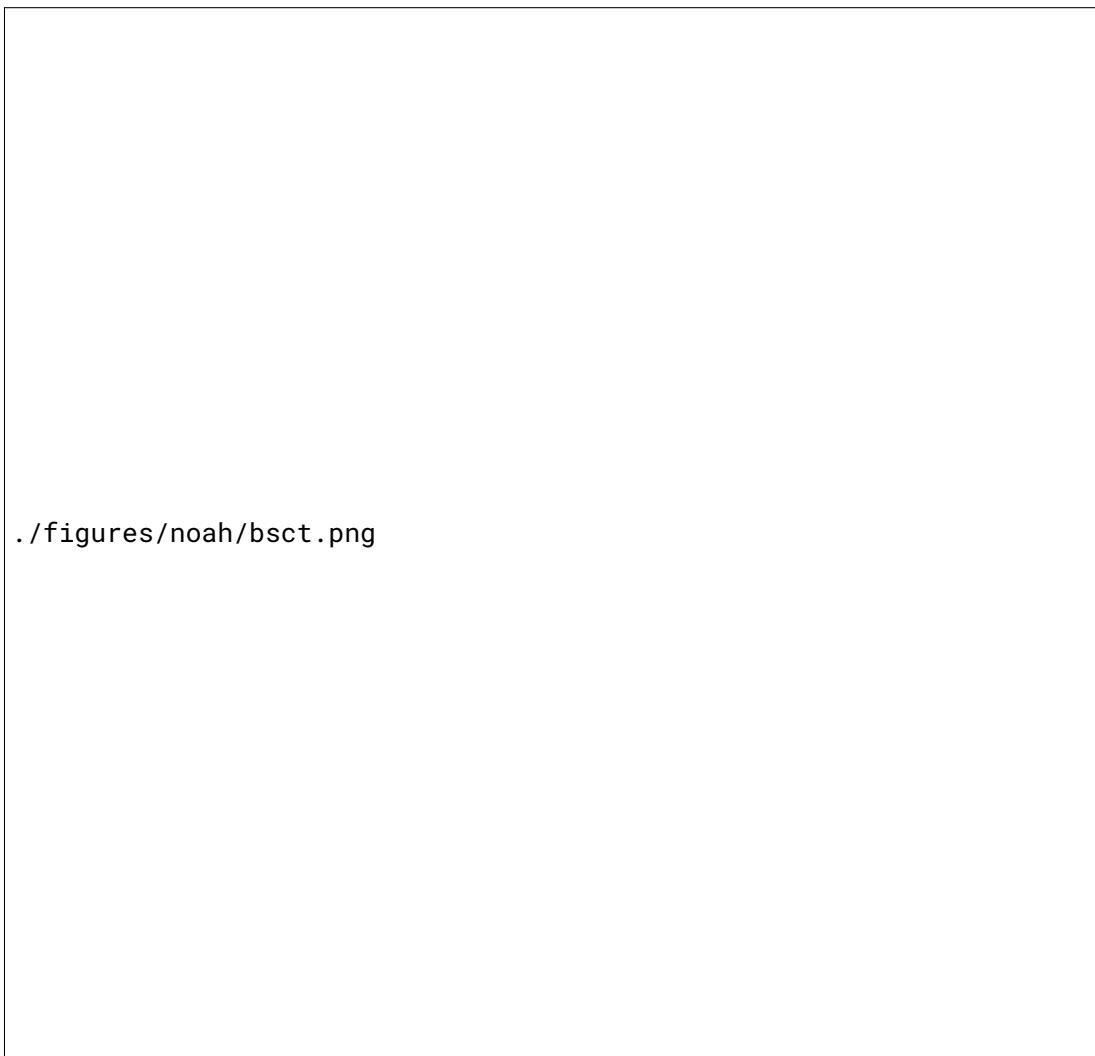


fig4bsctb

Figure 4.5: Spectra obtained from a NOAH-4 BSCT supersequence. (a) HMBC. (b) HSQC. (c) COSY; a pair of asymmetric crosspeaks are highlighted with red arrows. (d) TOCSY (60 ms DIPSI-2 mixing). Despite the COSY and TOCSY having only $\sim 30\%$ sensitivity compared to standalone experiments, the intensity of the spectra obtained is still perfectly acceptable (the contour levels chosen are 1–2 orders of magnitude larger than for the HMBC). *Data code:* 7Z-220224.

NOAH-5 BS_N⁺SCT: HMBC + ¹⁵N seHSQC + HSQC + COSY + TOCSY

As the final example, we add a further magnetisation pool to the mix, namely protons directly coupled to ¹⁵N (i.e. ¹H^N). As of the time of writing, the implementation of multiple-FID experiments on Bruker spectrometers limits N to a maximum of 5, so a supersequence such as

the present NOAH-5 BS_N^+ SCT is the current limit. (However, there is no *scientific* argument forbidding $N > 5$, and it is likely that in future versions of TopSpin this restriction will be lifted.)

The values of A for each module are given in entry 6 of table 4.2. If the HMBC module is placed at the beginning of the supersequence, then in order to preserve *both* ${}^1\text{H}^N$ and ${}^1\text{H}^C$ magnetisation, the zz -filter must be extended to include ${}^{15}\text{N}$ pulses.⁴⁶ As before, the ${}^{15}\text{N}$ seHSQC and ${}^{13}\text{C}$ HSQC modules both suffer drops in sensitivity. For the ${}^{15}\text{N}$ seHSQC, this is partly because of imperfect preservation of ${}^1\text{H}^N$ magnetisation by the HMBC, but also stems from the addition of the zz isotope-selective pulse (ZIP) element to the seHSQC pulse sequence; this is described further in § 4.3.3. On the other hand, for the ${}^{13}\text{C}$ HSQC, the sensitivity loss stems purely from imperfect retention of ${}^1\text{H}^C$ magnetisation. Finally, because the HMBC dephases ${}^1\text{H}^{\text{IX}}$ magnetisation, the COSY and TOCSY at the end have lower sensitivities: however, as discussed above, this is not an issue in practice.

It is also possible to move the ${}^{15}\text{N}$ seHSQC module to the front: this gives it a slightly greater sensitivity, at the cost of the HMBC (entry 7, table 4.2). In general, these two modules tend to have comparable sensitivity, and which of these two arrangements is better depends on which module the sensitivity needs to be optimised for.

Lastly, the value of ρ_t given here of 3.74 represents an effective upper limit on the time-saving factor. Although ρ_t increases with N , the extent to which it deviates from the ideal value of N also increases: it is very difficult to obtain $\rho_t > 4$, even with five modules in the supersequence. Of course, it is possible to increase ρ_t further by lengthening the recovery delay d_1 used for the experiments: for example, if d_1 is increased to 2 s from its present value of 1.5 s, ρ_t increases to 3.94. Obviously, this can only be pushed so far before it becomes meaningless.

4.2 GENESIS: automated pulse programme creation

sec:noah_genesis

In this section, I discuss the development of the GENESIS (GENEration of Supersequences In Silico) website (fig. 4.6): as the name suggests, it automatically generates pulse programmes for arbitrary NOAH supersequences. The website also provides extensive instructions on acquiring and processing NOAH data. Although this may at first glance appear slightly out of chronological order, in that the paper³ was published later than much of the other work in this chapter, an early version of the GENESIS tool was in fact created much earlier (by July 2020).

The present version of GENESIS is available at <https://nmr-genesis.co.uk>; the source code can also be obtained from <https://github.com/yongrenjie/genesis>.



`./figures/noah/genesis_frontpage.png`

fig:genesis_frontpage

Figure 4.6: The front page of the GENESIS website (<https://nmr-genesis.co.uk>), as of 10 September 2022.

4.2.1 Motivation

From the preceding discussion in § 4.1, it is clear that modules which use different magnetisation pools can be combined almost at will. It does not matter *what* modules they are, merely what magnetisation pools they consume (and preserve). Thus, for example, the NOAH-2 SC supersequence can in fact be generalised to any $^1\text{H}^C$ module plus any $^1\text{H}^{1\text{C}}$ module. Very broadly speaking, we may define a generic supersequence as having any or all of the following:

- a HMBC module, which actually uses $^1\text{H}^{1\text{X}}$ magnetisation but can be placed at the front as discussed in the NOAH-3 BSC example above;
- a $^1\text{H}^N$ module;
- one or more $^1\text{H}^C$ modules (it is possible to partition the $^1\text{H}^C$ magnetisation pool between two modules, as will be discussed in § 4.3.4);
- one or more $^1\text{H}^{1\text{X}}$ modules which consume bulk magnetisation.

In the first NOAH paper in 2017,³⁰ a total of 285 ‘viable’ supersequences were already listed. If we further take into account some of the new modules which were developed over the course of my DPhil (§ 4.3), the generic formula above provides for over 4000 viable supersequences.*

*This is also ignoring the ‘parallel’ supersequences, which are discussed in § 4.5. The support for parallel supersequences in GENESIS is not complete: integrating these fully would require substantial changes to the user interface, which I have not had time to do.

(‘Non-viable’ sequences would be those which have undesirable drawbacks: for example, wrongly ordered modules like in a NOAH-2 CS supersequence.)

In spite of this, *only around 45 pulse programmes* had been made available (either via the supplementary information of NOAH papers, or the Bruker User Library). Traditionally, pulse programmes must be written by hand, which is a laborious and fairly error-prone process made worse by the sheer length of NOAH experiments. Doing this for thousands of supersequence is clearly impractical; furthermore, each time a new module is developed, or an old module is improved, updating every relevant supersequence would itself already be a mammoth task. This explains why—although the NOAH concept provides a clear blueprint for how supersequences may be constructed—there is still a huge gap between theory and practice.

To bridge this gap, I turned towards the *programmatic* generation of pulse programmes.* This not only allows for existing supersequences to be generated at will, but also provides an easy way for updates to be rapidly disseminated to the NMR community. Furthermore, a website can serve as a ‘one-stop’ shop where—after downloading pulse programmes—users may download the associated NOAH processing scripts and also access instructions on how to run NOAH experiments. This information did already exist, but was scattered across several different websites and/or journal supplementary information documents, and would have been needlessly confusing to a new user (not to mention the different versions of scripts available in different publications).

4.2.2 Implementation details

I will now describe a few features of GENESIS pulse programmes, as well as how these are implemented. The GENESIS code is written in TypeScript; during deployment, this is compiled to JavaScript, which can then be directly executed in the client’s web browser. No server-side code is required, meaning that the GENESIS web page is actually a static site (it is currently hosted using GitHub Pages).

Overall structure

The algorithm used for pulse programme construction can loosely be separated into three parts, which are shown in listing 4.1:

1. the *preamble*, which consists of everything up until the beginning of the actual pulse sequence (the `ze` command). This includes header comments as well as definitions of

*This is actually a bit of a lie: GENESIS was initially created for my own convenience. Throughout this chapter, I have had to perform many comparisons of different supersequences, and this tool spared me from having to write everything by hand (and—more often than not—subsequently discover mistakes which invalidated the results). Of course, it soon became apparent that it could find much wider use.

```

/* PREAMBLE */
; gns_noah2-SCqf
; 13C HSQC
; 1H magnitude-mode COSY
"d4          = 0.25s/cnst2" ; 13C INEPT
"in0         = inf1/2"      ; 13C HSQC increment
"in11        = 2*dw"       ; COSY increment
; ...
define delay DC_HSQCa
"DC_HSQCa   = d4-p14/2"
"10          = td1/2"       ; TD1/NBL

/* MAIN SECTION */
1 ze
4 d1 st0
; MODULE 1 - HSQC
(p1 ph0):f1
; ...
goscnph30 cpd2:f2
50u do:f2
2m st
; MODULE 2 - COSY
(p1 ph12):f1
; ...
go=2 ph26
1m iu1       ; TD1/NBL counter
1m igrad EA ; echo-antiecho gradients
1m id11      ; COSY t1
30m wr #0 if #0 zd
if "l1 % 2 == 0" {
  1m id0       ; HSQC t1
}
lo to 4 times 10
exit

/* EPILOGUE */
ph0=0
;gpnam4: SMSQ10.100
;gpz4: 70% (13C CTP)
;cpd2:wvm:wudec: cawurst_d-20(220 ppm, 1.4 ms; L2H)
;d4: 1/4J(CH)
;auprog: noah_hsqc:noah_cosy QF
;module identifiers: C_HSQC H_COSY_QF
;pulse programme created by genesis-v2.2.2, https://nmr-genesis.co.uk
;Sun Sep 11 2022 16:04:54 GMT+0800 (Malaysia Time)

```

Listing 4.1: Abridged GENESIS pulse programme for the NOAH-2 SC supersequence shown in fig. 4.1a.

lst:genesis_sc

parameters, such as delays and pulse widths;

2. the *main section*, which contains the actual pulse sequence;
3. the *epilogue*, which contains phase cycle information as well as footer comments describing each parameter. Instructions for generating shaped pulses using Bruker's WaveMaker software are also included here, as well as instructions for automatic processing (§ 4.2.3), and comments indicating how the pulse programme was generated (for reproducibility purposes).

The construction of the preamble and main section is largely accomplished through the collation of module-specific information, the most important of which are:

- information about the module itself, which go into the header comments;
- parameter definitions, which are collated to form the preamble. Duplicates must be removed here to avoid errors; and
- the pulse programmes themselves, which are directly concatenated to form the main section.

These, as well as other smaller bits of information (e.g. relevant citations, appropriate processing scripts), are stored within `NOAHModule` objects (an example is provided in listing 4.2). Each distinct module corresponds to one such object. Therefore, if one wants to add a new module to GENESIS, most of the work can be completed by simply defining a new `NOAHModule` object: no changes to the algorithm itself are needed.

To put the epilogue together, the pulse programme constructed so far is scanned for pulse phases, shaped pulses, and all other parameters. Using predefined lookup tables, GENESIS then outputs pulse phase definitions, WaveMaker directives (where appropriate), and comments containing textual descriptions of each parameter. These comments are mostly cosmetic, but are very useful to the user as they are displayed in the `ased` screen when setting up an experiment. Finally, instructions for automatic processing of the NOAH data (explained in § 4.2.1) are added to the bottom, together with the version number and a timestamp (for reproducibility purposes).

Phase/delay incrementation and looping

Since NOAH experiments are 2D experiments, there is one additional complication: the pulse programme must contain appropriate looping statements, together with pulse phase and delay incrementation, in order to correctly generate the indirect dimension. In many existing NOAH pulse programmes, looping in 2D experiments was written using the equivalent of `for` loops (listing 4.3, *left*). Although this suffices for the vast majority of supersequences, whenever anything must be incremented in a different way (e.g. for parallel supersequences, or the PSYCHE module

```

let shortDescription = `; 13C HSQC`

let preamble =
  "d4      = 0.25s/cnst2"           ; 13C INEPT
;
"D[ID]a  = d4-p14/2"    // [ID] is later replaced with the module code
`


let pulprog =
  `; 13C-1H HSQC

  ; INEPT
  (p1 ph0):f1
  ;
  goscnp ph30 cpd2:f2
  50u do:f2
`


const mod = new NOAHModule(
  "C_HSQC",           // internal module code
  "c13",              // module category
  "S",                // single-letter code, Table 4.1
  [],                 // relevant citations (if any)
  "noah_hsqc",        // AU programme for processing
  shortDescription,   // short description
  [AF_EDIT],          // available acquisition flags
  preamble,           // preamble text
  pulprog,            // pulse programme text
  1,                  // number of FIDs
  false               // flag for 'parallel' modules
);
export default mod;

```

Listing 4.2: An excerpt of the NOAHModule object for the HSQC module (internal code C_HSQC).

which uses a different number of t_1 increments), the nested loop structure must be modified. I therefore opted to change the structure to use only one loop, and to control the phase and delay incrementation using modular arithmetic (*listing 4.3, right*). The outcome is entirely equivalent, but edge cases can be implemented simply by adding another check on the loop counter L1.

Parameter standardisation

Each NOAH module contains a number of parameters, including pulse widths, delays, gradient amplitudes, shaped pulse waveforms. Since different modules are stored separately (as differ-

```

"10 = td1/4"

1 ze
3 1m
4 d1 st0

; ... (pulse sequence goes here)

; in inner loop
1m igrad EA ; HSQC gradients
1m id11 ; COSY t1
30m wr #0 if #0 zd
lo to 4 times 2

; in outer loop
1m ip5*2 ; HSQC 13C 90
1m ip30*2 ; HSQC receiver
1m id0 ; HSQC t1
lo to 3 times 10

end

```



```

"10 = td1/2"
"11 = 0"

1 ze
4 d1 st0

; ... (pulse sequence goes here)

; on every pass
1m iu1 ; loop counter
1m igrad EA ; HSQC gradients
1m id11 ; COSY t1
30m wr #0 if #0 zd

; on every second pass
if "11 % 2 == 0" {
    1m ip5*2 ; HSQC 13C 90
    1m ip30*2 ; HSQC receiver
    1m id0 ; HSQC t1
}
lo to 4 times 10

end

```

Listing 4.3: Implementation of phase/delay incrementation and looping in previous NOAH sequences (*left*, using nested loops) and in GENESIS (*right*, using modular arithmetic).

ent objects), directly concatenating their pulse programmes may lead to conflicting parameter definitions. GENESIS avoids this by maintaining a global table of parameter definitions which are applied to all modules: when new modules are added, they must be checked against this to ensure that there are no inconsistencies.

In general, where possible, these parameters are chosen to be consistent with pulse programmes in the Bruker standard library: thus, for example, P1 is the ${}^1\text{H}$ 90° pulse width, and CNST2 is the ${}^1J_{\text{CH}}$ value used for calculating INEPT delays. This makes it easy to read in parameters either from the *prosol* relation tables in TopSpin, or from other existing parameter sets. Some delays are module-specific and do not need to be reused, and in standard library sequences, are often labelled as DELTA1, DELTA2, and so on. To avoid conflicting definitions and also to improve readability, I rename these such that they include the name of the module: thus, in a ${}^{13}\text{C}$ HSQC these may be labelled DC_HSQC_1. Here, C_HSQC is the name associated with the NOAHModule object.

If combined with some caution when adding new modules, these measures ensure that *within* a supersequence there are no parameter clashes: we may view this as a *local uniqueness* of parameters. However, the impact of this design choice is even more far-reaching: since parameters are stored globally, they will always have the same value in *all* possible supersequences (or in other words, the parameters are *globally unique*). This makes it exceptionally easy to set up multiple different supersequences in TopSpin: virtually all of the parameter values may simply be copied from a previous NOAH dataset.

One potential issue with this strategy is that TopSpin provides a finite number of named pulse widths (for example). Thus, there are only so many different parameters which can be stored in a global table before running into inevitable conflicts. A workaround would be to sacrifice the global uniqueness of each parameter, and only have it be unique within a given supersequence. Fortunately, this situation has not (yet) surfaced.*

Parameter descriptions

At the epilogue of the pulse programme, extra comments are added for every parameter present in the pulse programme. Most of these are purely textual, and appear in the TopSpin aed parameter setup screen; this naturally helps to make the pulse programmes as easy to use as possible. However, some of these comments have especial meaning: gradient amplitudes and shapes, for example, are specified in a way which allows them to be automatically populated using the gppp Python script in TopSpin. Furthermore, WaveMaker directives are also specified for some shaped pulses, allowing them to be created in an on-the-fly manner using the wvm command. This means that the user (generally) need not separately download and install a set of shaped pulses.

Module choice

One final issue which must be overcome is the fact that some NOAH modules may be implemented differently depending on the supersequence which it is being used in. The HMBC module described in § 4.1.3 is one such example: the form of the *zz*-filter depends on whether the HMBC module must retain $^1\text{H}^{\text{C}}$ and/or $^1\text{H}^{\text{N}}$ magnetisation for subsequent modules. Thus, in the NOAH-2 BS supersequence the *zz*-HMBC must be used, but in the NOAH-2 SB supersequence the ‘original’ HMBC without a *zz*-filter is preferable. Since they have different pulse programmes, each of these ‘versions’ of the HMBC are described by *separate* NOAHModule objects.

*The number of *pulse phases* in particular, though, is dangerously close to the maximum number of 32. In fact, the global uniqueness criterion is not really important for pulse phases, because—unlike, say, delays—pulse phases are hardcoded in the pulse programme, and cannot be copied from one dataset to another. So, if necessary, I could dispense with the global uniqueness for pulse phases, at the cost of some increased code complexity. I did briefly contemplate this possibility, but since I am at the end of my DPhil and am unlikely to add any new modules soon, this will remain a hypothetical—for now.

However, it is unlikely that the majority of users would want to manually configure the supersequence in such detail. Thus, the GENESIS web interface actually hides these different versions from the user, only showing one button labelled ‘HMBC’. Under the hood, the correct version of the HMBC is automatically chosen based on what other modules the user has chosen (fig. 4.7).



`./figures/noah/hmhc_flowchart.png`

fig:hmhc_flowchart

Figure 4.7: Flowchart showing how the correct ‘version’ of the HMBC module is determined when constructing a supersequence using GENESIS. When developer mode is on, the user directly chooses module codes corresponding to each version of the HMBC. When it is off, the appropriate version is chosen based on which other modules are present in the supersequence.

Should the user indeed desire to control the exact module used, the website also offers a ‘developer mode’ switch: turning this on allows the user to directly choose the desired module code. (Enabling developer mode also reveals a handful of extra modules which were used exclusively for my DPhil and are not generally of interest.)

Tests

One risk inherent in any NMR experiment is the possibility of causing spectrometer damage when executing malformed pulse programmes. To minimise the chances of this, each version of GENESIS is checked against a series of automatic tests before being released. The most im-

portant tests ensure that each module containing decoupling statements also turn off decoupling immediately after acquisition. On top of this, there are also a series of regression tests where (Some) supersequences are checked against those generated by previous versions: any differences in these are automatically flagged for review.

None of this can actually *stop* wrong pulse programmes from being written. For example, if a module are incorrectly programmed (which is still very easy), then any supersequence containing that module will not be fully functional. However, the inclusion of tests minimises the chance of creating faulty pulse programmes.

It should be mentioned that this is not a *new* problem which GENESIS introduces: errors in pulse programmes can arise equally easily (if not *even more easily*) when they are handwritten. There is no way to really circumvent this, except through careful inspection of the pulse programme prior to execution. The real difference with GENESIS is that users may be more likely to blindly use the pulse programmes it creates, without applying an equal amount of caution. The tests are therefore not meant to ensure perfection, but rather to avoid some of the worst consequences which may surface. Of course, it is also not possible for GENESIS to prevent users from setting up parameters incorrectly (e.g. overly long acquisition times).

Reproducibility

The final section of a GENESIS pulse programme consists of comments intended solely for reproducibility purposes. These comments indicate exactly which NOAHModule objects were used to create the pulse programme (these can be manually entered via developer mode), as well as a GENESIS *version number*. Since GENESIS has been updated with some regularity during my DPhil, each release is assigned a version number, which broadly follows the principles of *semantic versioning*. Old versions of GENESIS may be accessed by navigating to the specific URL <https://nmr-genesis.co.uk/X/Y/Z> for version X.Y.Z. Together, these two pieces of information allow all pulse programmes to be regenerated whenever necessary: this is important for ensuring that the spectra thus acquired are reproducible (as far as possible).

How smart is GENESIS?

Having written several pages of text about the *features* present in GENESIS, it is tempting to think that it is ‘intelligent’ in its design of pulse programmes. At various points in time, extensions to the concept have even been proposed: in GENESIS, the building blocks used for pulse programme construction are NOAH modules, but one could envision breaking up these into the smallest possible units of pulses, delays, gradients, and so on, and using something like GENESIS to construct *arbitrary* pulse programmes. This was even speculatively mentioned in the paper.

However, in truth, GENESIS is a long way from being able to do such things. In particular, it does not have any actual understanding of the Bruker pulse programming syntax: *almost* all of the pulse programme instructions are hardcoded as strings.* Thus, the creation of pulse programmes is not *truly* being done from the ground up: instead of ‘combining pulse sequence elements’, a more accurate description is that the *text* corresponding to different pulse sequence elements is being concatenated. There is actually a substantial amount of brute force involved. In a similar way, it is not possible for the automated tests to really check for ‘correctness’ of the pulse programmes.

Of course, it is hardly a trivial task to write something more sophisticated: I would need to construct abstract representations of each pulse sequence element, and functions which could transform these into actual pulse programme commands. I feel that I would need rather more knowledge in computer science, especially a course on compilers and/or programming language theory, to do this. However, this was something I would have really liked to do had I had more time!

4.2.3 Processing improvements

Finally, I touch on a few small improvements in processing NOAH data: these are not strictly made possible by the GENESIS website, but having a unified source for pulse programmes and processing scripts does make it substantially easier to propagate these developments.

AU processing scripts

When processing NOAH experiments, the first step is to split up the FIDs of different modules into different files: this is done by the `splitx_au` AU programme. However, after that is done, `splitx_au` *also* calls on auxiliary AU programmes to process each of the resulting datasets. Each module is associated with a specific auxiliary AU programme: for example, HSQC data is processed with `noah_hsqc`, COSY data with `noah_cosy`, and so on.

Previously, these auxiliary AU programmes had to be specified manually using the `USERP1` through `USERP5` processing parameters, which respectively apply to the first through fifth modules in a supersequence. Although this was a perfectly serviceable setup, I realised that it was a source of annoyance when running multiple different supersequences, as modules would often be processed incorrectly if the processing parameters were inadvertently copied. Even when just running a single supersequence, this was still an extra cognitive load, especially for users who may not have been familiar with NOAH experiments.

*There are several exceptions where some parameters (for example, loop counters in DIPSI mixing) are dynamically generated, mainly to avoid clashes between different parts of the supersequence.

Since the pulse programme specifies which modules are run, and the ‘correct’ processing parameters also depend on the modules, it stands to reason that the processing parameters can be encoded in the pulse programme itself. Thus, all GENESIS pulse programmes contain a line near the bottom specifying which AU programmes are to be used for each of the modules present. I also modified the `splitx_au` AU programme to parse the pulse programme for this extra information. Thus, the `USERP1` through `USERP5` parameters need no longer be specified, as long as a GENESIS pulse programme and a recent version of `splitx_au` are used. These parameters still serve a purpose, though; if they are non-empty, they will override the ‘default’ processing instructions found in the pulse programme. This maintains backwards compatibility with, for example, parameter sets which have been previously saved; it also allows advanced users to customise the processing, if so desired.

Non-uniform sampling

Another small improvement relates to the implementation of non-uniform sampling (NUS) in NOAH supersequences. Because the indirect dimension in NOAH pulse programmes is generated through explicit looping (listing 4.3), instead of the `mc` macro used in most 2D experiments, the standard NUS implementation in TopSpin cannot be used. Instead, one must explicitly define a list of t_1 increments to sample (stored as the `VCLIST` parameter). On top of this, the original NUS implementation⁴¹ used a Python script (`noah_nus.py`) to modify the NOAH pulse programme such that it only sampled t_1 values from this list. This, however, necessitated storing two copies of the same pulse programme, one with NUS and one without.

There is no way to circumvent the requirement for the `VCLIST` parameter, but in GENESIS pulse programmes, I opted to implement NUS as an acquisition flag which could be toggled. This means that there is no need to store two versions of the same pulse programme, and also makes it easier to seamlessly switch between non-uniform and uniform sampling as desired.

Unfortunately, this change is slightly more problematic than the `splitx_au` change: this new NUS implementation is not backwards-compatible, so requires a new NUS setup script which I have called `noah_nus2.py`. This script cannot be used with old pulse programmes, and the old script cannot be used with GENESIS pulse programmes. Nevertheless, it does represent a real improvement over the previous implementation: the GENESIS website also helps to make the changeover as smooth as possible.*

*That said, there were still some bugs in the new NUS implementation right up until August 2022, so this change was not as smooth-sailing as I had hoped for.

4.3 Discussion of individual modules

sec:noah__modules

The GENESIS website makes it possible to easily implement and distribute entirely new modules, as well as updates to old modules: simply changing the underlying `NOAHModule` objects is sufficient to propagate changes to all supersequences generated using those modules. In this section, I discuss a number of new pulse sequence developments made in the course of my DPhil; all of these have been successfully implemented in GENESIS and are available to download.

4.3.1 ^{13}C sensitivity-enhanced HSQC

subsec:noah__sehsqc_c

The first of these is the sensitivity-enhanced HSQC (seHSQC) experiment, which provides up to $2\times$ increased SNR over a standard (echo–antiecho) HSQC.* In the original version of the seHSQC (fig. 4.8a), developed by Cavanagh, Rance, and Kay ('CRK'),^{48,49} this is accomplished through the so-called *preservation of equivalent pathways* (PEP) technique.⁵⁰

Just like in a standard HSQC, the experiment begins with an INEPT block and t_1 evolution. At the end of t_1 , there are two terms which are cosine- and sine-modulated with respect to Ω_S . In the standard HSQC, only the $2I_zS_y$ term is returned into observable ^1H magnetisation by the reverse INEPT block (see § 1.4.4). In contrast, the PEP block transfers *both* terms back to ^1H and subsequently detects both. Specifically, in the echo experiment it accomplishes the transfer $2I_zS_y \rightarrow I_y$ and $2I_zS_x \rightarrow I_x$, meaning that the density operator at the end of t_1

$$-2I_zS_y \cos(\Omega_S t_1) - 2I_zS_x \sin(\Omega_S t_1) \quad (4.6) \quad \{eq:sehsqc_t1_modulation\}$$

is transformed into

$$-I_y \cos(\Omega_S t_1) - I_x \sin(\Omega_S t_1) \quad (4.7) \quad \{eq:sehsqc_before_detection\}$$

just prior to acquisition. During the FID, only the -1 -coherence component is detected:

$$\frac{1}{2i}I_- \cos(\Omega_S t_1) - \frac{1}{2}I_- \sin(\Omega_S t_1) = \frac{1}{2i}I_- \exp(-i\Omega_S t_1), \quad (4.8) \quad \{eq:sehsqc_detection_m\}$$

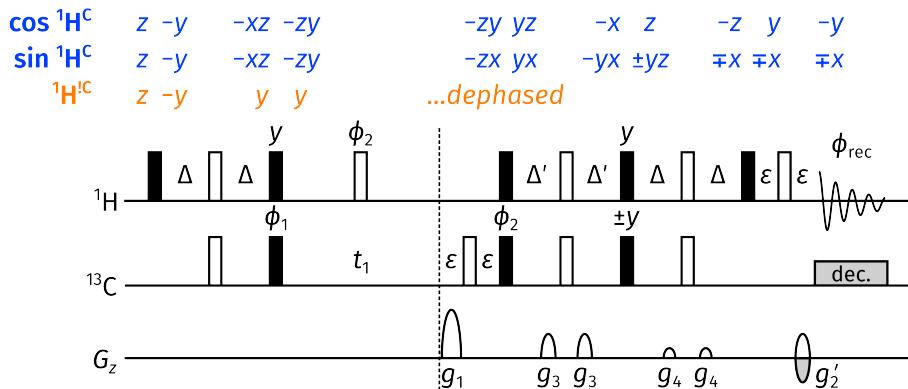
yielding a signal of

$$\frac{1}{2i} \exp(-i\Omega_S t_1) \exp(i\Omega_I t_2), \quad (4.9) \quad \{eq:s_echo_sehsqc\}$$

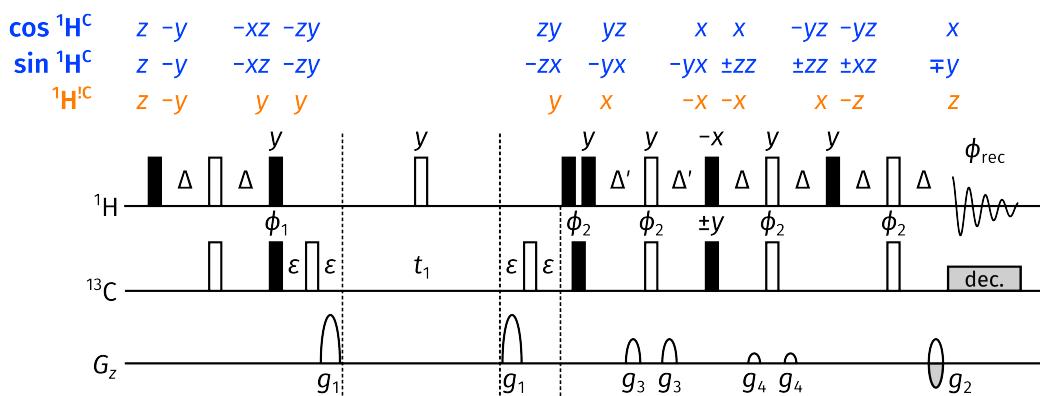
which has a $2\times$ larger amplitude than the original EA HSQC (eq. (1.96)). The antiecho experiment can be similarly analysed. Since the seHSQC has the same noise level as in an EA HSQC, this also corresponds to a $2\times$ increase in SNR.

*Since a States HSQC has $\sqrt{2}$ times the SNR of an EA HSQC (as shown in fig. 1.8), this also means that the seHSQC has a $\sqrt{2}$ SNR improvement over a States HSQC. The literature can be somewhat confusing on this point: sometimes the gain in signal is even conflated with the gain in SNR. The clearest exposition I have found is that provided by Kontaxis et al.⁴⁷

(a) CRK seHSQC



(b) NOAH seHSQC1



(c) NOAH seHSQC2

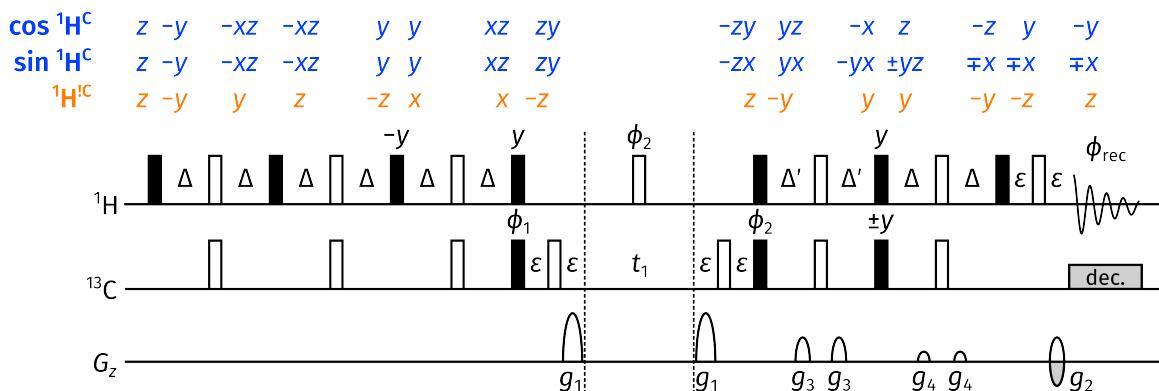


Figure 4.8: Sensitivity-enhanced HSQC sequences discussed in this section, along with product operator analysis. This analysis is provided only for the first step of the phase cycle, and assumes an IS spin pair with $\Delta' = 1/(4 \cdot ^1J_{\text{CH}})$. ‘ $\cos^1\text{H}^C$ ’ refers to the component of the $^1\text{H}^C$ magnetisation which is cosine-modulated during t_1 . (a) Cavanagh–Rance–Kay seHSQC. (b) NOAH seHSQC1 module. (c) NOAH seHSQC2 module. Phase cycling is performed with $\phi_1 = (x, -x)$, $\phi_2 = (x, x, -x, -x)$, and $\phi_{\text{rec}} = (x, -x, -x, x)$. The delay Δ is set to $1/(4 \cdot ^1J_{\text{CH}})$; see the text for a discussion of Δ' . The pulses marked with a phase of $\pm y$ are applied with a phase of y in the echo experiment and $-y$ in the antiecho. Gradient amplitudes are $(g_1, g_2, g'_2, g_3, g_4) = (80\%, \pm 40.2\%, \pm 20.1\%, 11\%, -5\%)$.

It should, however, be noted that the PEP transfer is only fully attained for *IS* spin pairs if the delay Δ' is set to $1/(4 \cdot {}^1J_{\text{CH}})$. For I_2S or I_3S spin systems, no gain in sensitivity is accomplished with this setting of $\Delta' = 1/(4 \cdot {}^1J_{\text{CH}})$. It is more common, therefore, to shorten Δ' to $1/(8 \cdot {}^1J_{\text{CH}})$: this sacrifices some transfer efficiency for *IS* systems, but allows for some sensitivity enhancement for I_2S and I_3S systems (table 4.3). In this work, all experiments were run assuming a value of 145 Hz for ${}^1J_{\text{CH}}$.

Spin system	Theoretical sensitivity enhancement		
	$\Delta' = 1/(4 \cdot {}^1J_{\text{CH}})$	$\Delta' = 1/(8 \cdot {}^1J_{\text{CH}})$	$\Delta' = 1/(12 \cdot {}^1J_{\text{CH}})$
<i>IS</i>	2	1.71	1.5
I_2S	1	1.41	1.37
I_3S	1	1.21	1.25

tbl:sehsqc_theory

Table 4.3: Theoretical sensitivity enhancements for *IS*, I_2S , and I_3S spin systems in the seHSQC, as a function of the delay Δ' . The values are taken from Schleucher et al.⁵¹

NOAH seHSQC versions

Just like the original EA HSQC (fig. 1.7), the CRK seHSQC must be modified in order to be compatible with NOAH supersequences. In particular, the CTP gradients in the CRK seHSQC dephase the bulk ${}^1\text{H}^{1\text{C}}$ magnetisation pool: we would very much like it to return that to $+z$ instead, so that it can be sampled in later homonuclear modules (or, indeed, a HMBC module). This experiment is more tricky to adapt than the HSQC, because there are *three* different magnetisation components to juggle: the cosine-modulated ${}^1\text{H}^{\text{C}}$, sine-modulated ${}^1\text{H}^{\text{C}}$, and ${}^1\text{H}^{1\text{C}}$. Nevertheless, there are at least two ways of doing so; these two modified modules are labelled seHSQC1 and seHSQC2 respectively.

The seHSQC1 module (fig. 4.8b) was developed by me.* It retains the same general structure of the CRK seHSQC up until t_1 , but immediately after t_1 a composite ${}^1\text{H}$ pulse is used to effect the transformations

$$I_z \rightarrow I_y; \quad I_y \rightarrow I_x, \tag{4.10} \quad \text{(eq:sehsqc1_dp)}$$

where the former is required for the ${}^1\text{H}^{\text{C}}$ pool and the latter for ${}^1\text{H}^{1\text{C}}$. This has the effect of storing the sine-modulated term as $2I_zS_z$ magnetisation for one spin echo (as opposed to the CRK seHSQC, which stores cosine-modulated term as I_z). At the beginning of the final spin echo (containing the rephasing gradient), this is transformed into antiphase magnetisation of the form $2I_xS_z$; thus, this spin echo must be lengthened to a total duration of 2Δ in order to fully refocus ${}^1J_{\text{CH}}$. The final modification involves the addition of an extra gradient immediately before to t_1 :

*Through a great deal of trial and error, and certainly *not* intelligent design.

this ensures that the bulk ${}^1\text{H}^{1\text{C}}$ magnetisation is not dephased, just like in the NOAH HSQC module (fig. 4.2a).

The seHSQC2 module (fig. 4.8c), on the other hand, was initially reported by Hansen et al.⁵² In this pulse sequence, the initial ${}^1\text{H}$ 90° excitation pulse is replaced with a zz isotope-selective pulse element (ZIP element), which is very similar to the zz -filter but has different pulse phases. This acts as a 90°_x pulse on the ${}^1\text{H}^{1\text{C}}$ magnetisation pool and thus ultimately leads to the same signals being detected (save for a trivial 180° phase shift). However, on ${}^1\text{H}^{1\text{C}}$ magnetisation, the ZIP element acts as a 90°_y pulse; it turns out that this modification alone is sufficient to return the bulk ${}^1\text{H}^{1\text{C}}$ magnetisation to $+z$ at the end of the sequence. The ZIP element therefore represents an isotope-specific rotation on ${}^1\text{H}$ spins, where the rotation axis depends on whether it is coupled to ${}^{13}\text{C}$ or not.

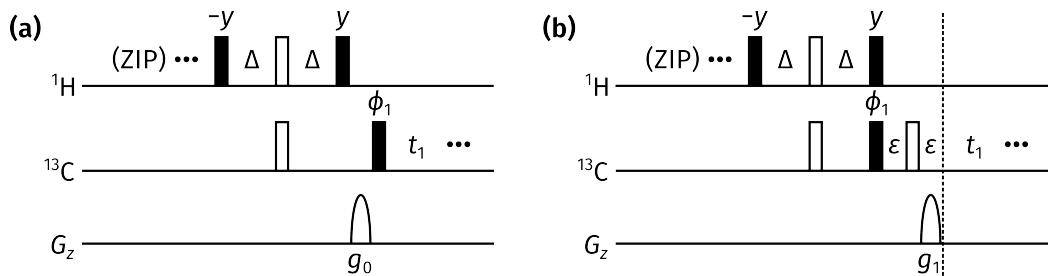


Figure 4.9: Comparison of CTP gradient schemes in seHSQC2 module. (a) As reported in Hansen et al.⁵² g_0 is a purge gradient with arbitrary amplitude; the amplitude of the final CTP gradient must be halved. (b) The version used in this work (corresponding to fig. 4.8c).

In my work, I made one change to the experiment, namely the addition of an extra gradient echo prior to t_1 (thus leading to a symmetric gradient scheme similar to that in seHSQC1). Instead of this, the original paper had in fact inserted a purge gradient between the ${}^1\text{H}$ and ${}^{13}\text{C}$ 90° pulses just after the INEPT spin echo (fig. 4.9). The scheme I used leads to similar results, but has one advantage in that it allows the amplitude of the final CTP gradient (g_2 in fig. 4.8) to be twice as large: this is particularly relevant to ${}^{15}\text{N}$ experiments, as will be discussed in § 4.3.2.

The primary consideration when evaluating the different seHSQC versions in fig. 4.8 is naturally sensitivity. Figure 4.10 compares the sensitivities of the various possible NOAH-2 S^+C^c super-sequence against a NOAH-2 SC^c experiment. As can be seen, all three seHSQC implementations yield improvements in the sensitivities of CH peaks: the original CRK seHSQC has the largest effect here, followed by seHSQC2 and seHSQC1. For CH_2 and CH_3 peaks, the CRK seHSQC and seHSQC2 perform better than the standard HSQC, but the average sensitivities in seHSQC1 have dipped slightly below. This is not entirely surprising in light of the discussion in § 4.1.1: the modified seHSQC1 and seHSQC2 sequences are expected to have lower sensitivities than the original experiment from which they are derived. In this case, it appears that the modifications

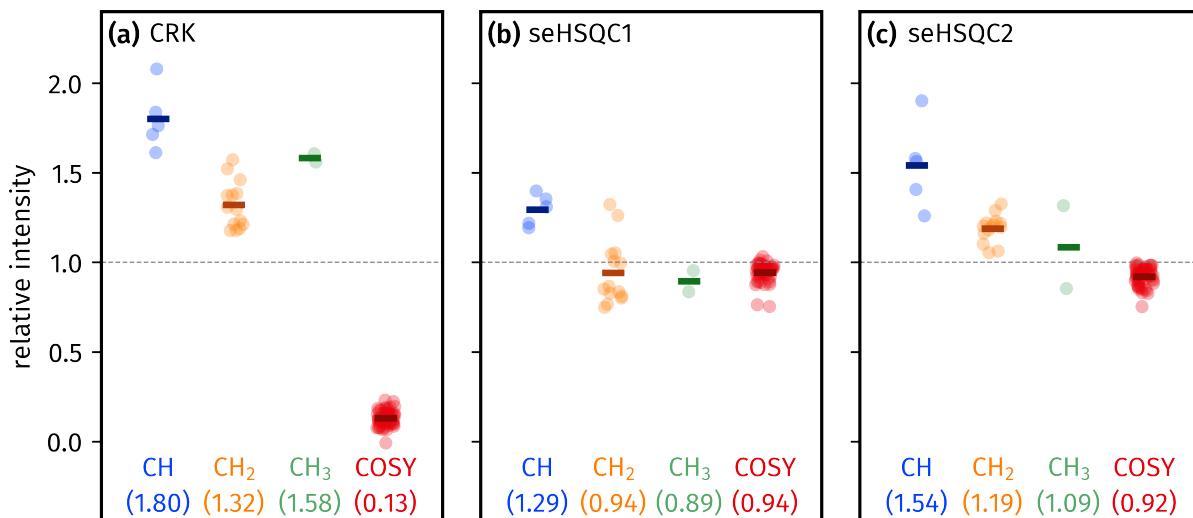


Figure 4.10: Sensitivities of seHSQC and CLIP-COSY modules in the NOAH-2 S^+C^c supersequences (run using $\Delta' = 1/(8 \cdot {}^1J_{CH})$). Intensities are reported relative to the corresponding peaks in a NOAH-2 SC^c supersequence; HSQC peaks are further broken down by their multiplicity. Each dot represents one peak; the solid lines as well as the numbers at the bottom of each plot represent averages over all peaks. (a) Using the CRK seHSQC. (b) Using the seHSQC1 module. (c) Using the seHSQC2 module. Data code: 7A-201115.

made to the seHSQC1 sequence have caused sensitivity losses which outweigh the benefits of using a sensitivity-enhanced experiment.

In the context of a NOAH supersequence, it is also important to consider how well each module preserves ${}^1H\text{--}{}^{13}C$ magnetisation for subsequent homonuclear module(s). In fig. 4.10, the CLIP-COSY module is used as the ‘reporter’ module, but the conclusions drawn here are applicable to *any* homonuclear module (or HMBC). The CRK seHSQC naturally does poorly in this respect, as it dephases the bulk magnetisation: this leads to almost an order-of-magnitude sensitivity loss in the CLIP-COSY. Both the other NOAH seHSQC modules, however, perform almost identically to the original HSQC module, as expected from the product operator analysis in fig. 4.8.

From fig. 4.10, we can draw some very simple conclusions: if the seHSQC is to be used at (or near) the *beginning* of a NOAH supersequence, where it needs to preserve ${}^1H\text{--}{}^{13}C$ magnetisation, then the seHSQC2 module is the best choice. On the other hand, if the seHSQC module is placed at the *end* (say in a NOAH-2 BS⁺ supersequence), then the CRK seHSQC is optimal. This logic is encoded in GENESIS.

COSY-type artefacts

One (minor) way in which the seHSQC1 module outperforms seHSQC2 is in the suppression of COSY-type artefacts in the seHSQC spectrum.⁵³ These arise due to evolution of ${}^nJ_{HH}$ during the

second-last spin echo of the CRK seHSQC. Specifically, if we assume the presence of a separate ^1H spin K which is coupled to spin I , the sine-modulated $2I_yS_z$ term can evolve under both $^1J_{IS}$ and $^nJ_{IK}$ into $2I_yK_z$, and the last ^1H 90° pulse transfers this coherence to the spin K . This leads to peaks correlating spin S with K . This is also true of the seHSQC2 module, as the product operator analysis for the CRK seHSQC is entirely applicable to it.

However, for the seHSQC1 module, the corresponding term giving rise to these artefacts would be the cosine-modulated I_x term (at the beginning of the second-last spin echo). During this spin echo, this can again evolve under both J-couplings into $4I_yK_zS_z$, and the final 90° pulse would transform this to $-4I_zK_yS_z$. However, crucially, this term is antiphase with respect to the heteronucleus S : therefore, when decoupling is applied during acquisition, this term should not be observed. This analysis can be verified experimentally by inspection of the seHSQC spectra thus obtained. The COSY-type artefacts, labelled with red boxes in the seHSQC2 spectrum (fig. 4.11b), are largely absent in the seHSQC1 spectrum (fig. 4.11a). Some artefacts still remain, which perhaps arise due to pulse imperfections (or perhaps more complicated spin systems than the three-spin system considered here; I did not analyse this in any further detail).

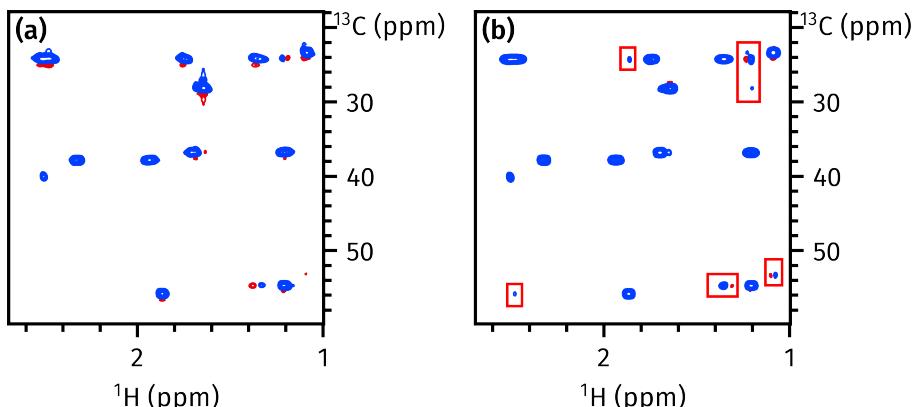


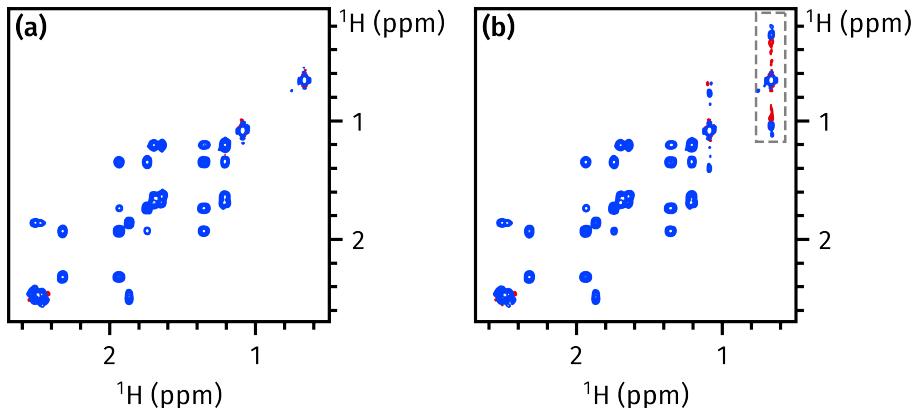
Figure 4.11: Comparison of COSY-type artefacts in NOAH seHSQC modules. (a) seHSQC1. (b) seHSQC2. The artefacts are highlighted in red boxes. Data code: 7A-201115.

Wing artefacts

One point worth considering is whether the extra gradient in the seHSQC2 module is truly needed. In the HSQC and seHSQC1 modules, the bulk magnetisation is placed along the transverse axis during t_1 ; therefore, if a pair of symmetric gradients are not used, this magnetisation will be dephased (as in the CRK seHSQC). However, in the seHSQC2 module, the bulk magnetisation is longitudinal, which suggests that this gradient is not necessary.

This gradient can indeed be removed in the seHSQC2 module, and doing so does not lead to any poorer preservation of $^1\text{H}^{13}\text{C}$ magnetisation. However, the omission of this gradient leads to artefacts in the CLIP-COSY spectrum, which I term ‘wing artefacts’ because of the way they flank

peaks in the CLIP-COSY (fig. 4.12). These artefacts can most clearly be seen in the diagonal peaks of methyl groups, but are present across the entire spectrum.



sehsqcswbgcavtsg2grad

Figure 4.12: Wing artefacts in CLIP-COSY spectra, taken from NOAH-2 S^+C^c supersequences. (a) Using the seHSQC2 module as shown in fig. 4.8c. (b) Using a modified seHSQC2 module where the gradient echo just before t_1 is removed. The wing artefacts flanking the two rightmost diagonal peaks can be clearly seen; the boxed artefacts are studied further in fig. 4.13. Data code: 7A-201115.

A clue as to their origin is provided by the characteristic frequencies of the artefacts:

$$\left(\Omega_1, \Omega_2 = \Omega_I \pm \frac{\Omega_I \cdot [t_{1,\text{HSQC}}/2]}{t_{1,\text{COSY}}}, \Omega_I \right) \quad (4.11)$$

where $t_{1,X}$ refers to the value of t_1 in experiment X (it does not matter which increment is used, because the ratio of t_1 in the fraction above is constant). In accordance with this, when the seHSQC indirect-dimension spectral width is reduced, the artefacts are displaced outwards (fig. 4.13e). This suggests that these artefacts arise from $^1\text{H}^{13}\text{C}$ magnetisation which is frequency-labelled in two different t_1 periods: specifically, it evolves once during *half* of the seHSQC t_1 period, and then again during the CLIP-COSY t_1 period. In fact, it is necessary to apply CTP gradients on both sides of t_1 in order to suppress evolution of this stray magnetisation during both halves: if either or both of the gradients are removed, the artefacts are present (figs. 4.13a to 4.13d).

These wing artefacts are unique to fast-pulsing experiments such as NOAH supersequences: in a typical 2D experiment, these are less likely to arise because of the long(er) recovery delays between data acquisition periods. In the sections which follow, we will see these surface repeatedly. Fortunately, it proves to be relatively easy to suppress them through the judicious use of gradients.

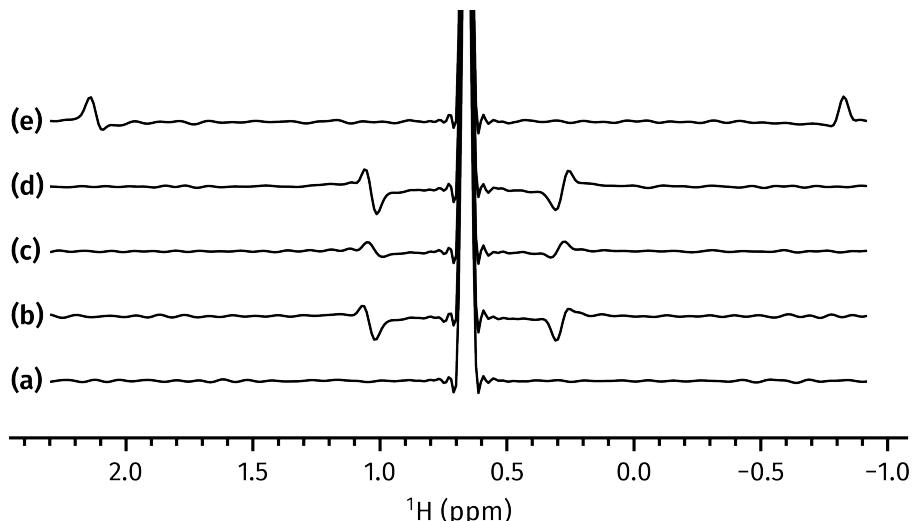


Figure 4.13: A closer look at wing artefacts in CLIP-COSY spectra, taken from NOAH-2 S⁺C^c supersequences with various modifications made to the seHSQC2 module. The spectra being plotted are 1D slices through $F_2 = 0.67$ ppm of the 2D CLIP-COSY spectrum; this corresponds to the boxed region in fig. 4.12b. (a) Using the default seHSQC2 module, with one gradient before and one after t_1 . (b) The same as (a), but with the gradient before t_1 turned off (the spin echo is still present, but the gradient intensity is set to 0). (c) The same as (a), but with the gradient after t_1 turned off. (d) The same as (a), but with both gradients turned off. (e) The same as (b), but with the seHSQC spectral width reduced by a factor of 4. *Data code: 7A-201115.*

Multiplicity editing

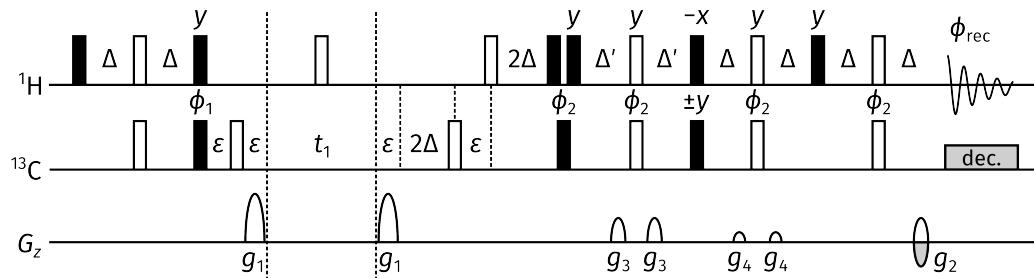
Multiplicity editing is reasonably easy to include in all of the seHSQC sequences above. It suffices to add a spin echo of total duration $4\Delta = 1/\gamma J_{\text{CH}}$ just after t_1 , while also making sure to change some pulse phases to account for the extra ^1H 180° pulse (fig. 4.14).

The inclusion of editing does not make a substantial difference in the sensitivity comparisons (fig. 4.15). However, it is interesting to note that the edited seHSQC2 in fact performs *better* than the edited HSQC in terms of preserving bulk magnetisation, as evidenced by the larger COSY intensities in fig. 4.15c. This can be explained by the fact that in the editing period of the HSQC experiment (and seHSQC1), the bulk magnetisation is placed in the transverse plane. The evolution of homonuclear couplings will thus lead to a small loss in the amount of $^1\text{H}^{13}\text{C}$ magnetisation preserved. In the seHSQC2 sequence, the bulk magnetisation is longitudinal during the editing period, so does not evolve under J_{HH} .

Choice of Δ'

As discussed previously, there are several possible values for the delay Δ' . I thus also investigated the possibility of setting $\Delta' = 1/(4 \cdot \gamma J_{\text{CH}})$. The corresponding sensitivity comparisons are shown in fig. 4.16.

(a) NOAH seHSQC1, edited



(b) NOAH seHSQC2, edited

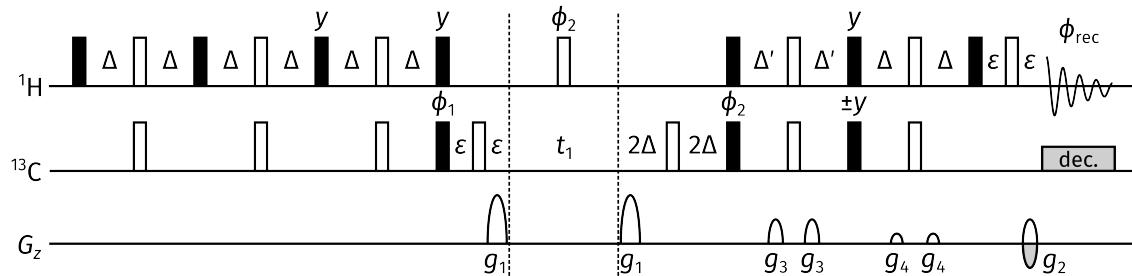
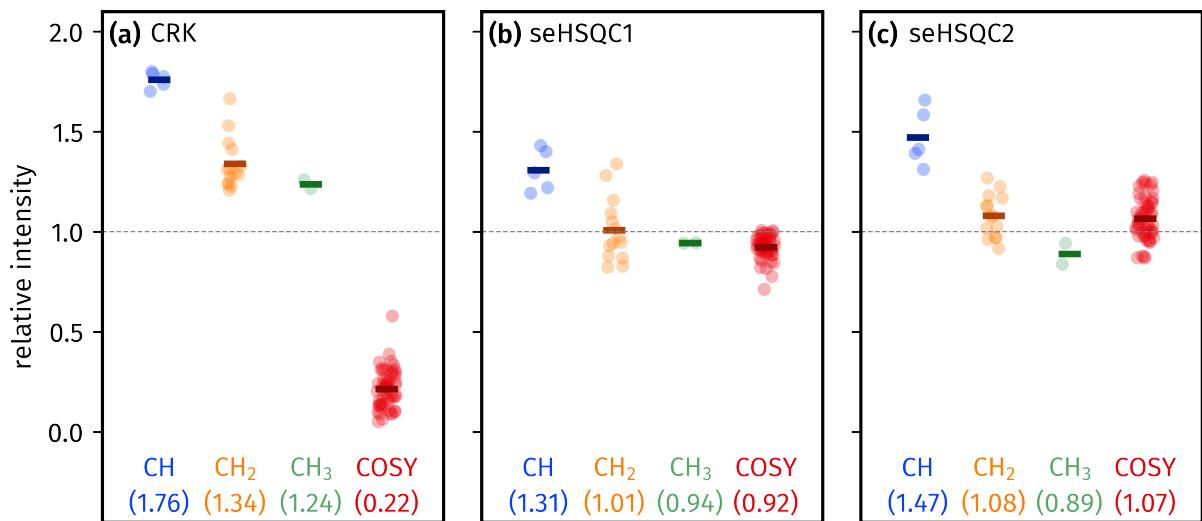


Figure 4.14: Multiplicity-edited NOAH seHSQC modules. (a) Edited seHSQC1. (b) Edited seHSQC2.

Figure 4.15: Sensitivity comparisons for multiplicity-edited seHSQC in NOAH-2 S⁺C^c supersequences. The delay Δ' was set to $1/(8 \cdot {}^1J_{\text{CH}})$. (a) Using the edited CRK seHSQC. (b) Using the edited seHSQC1 module. (c) Using the edited seHSQC2 module. Data code: 7A-201115.

As can be appreciated, the sensitivity boosts obtained for CH groups is higher than in the corresponding spectra with $\Delta' = 1/(8 \cdot {}^1J_{\text{CH}})$ (fig. 4.10). However, in this case, there is a small sensitivity loss for CH₂ and CH₃ groups as compared to the original HSQC (which is consistent with previous studies⁵¹). For CH₃ groups this is unlikely to be of any consequence, but particularly for diastereotopic CH₂ groups this may not be desirable.

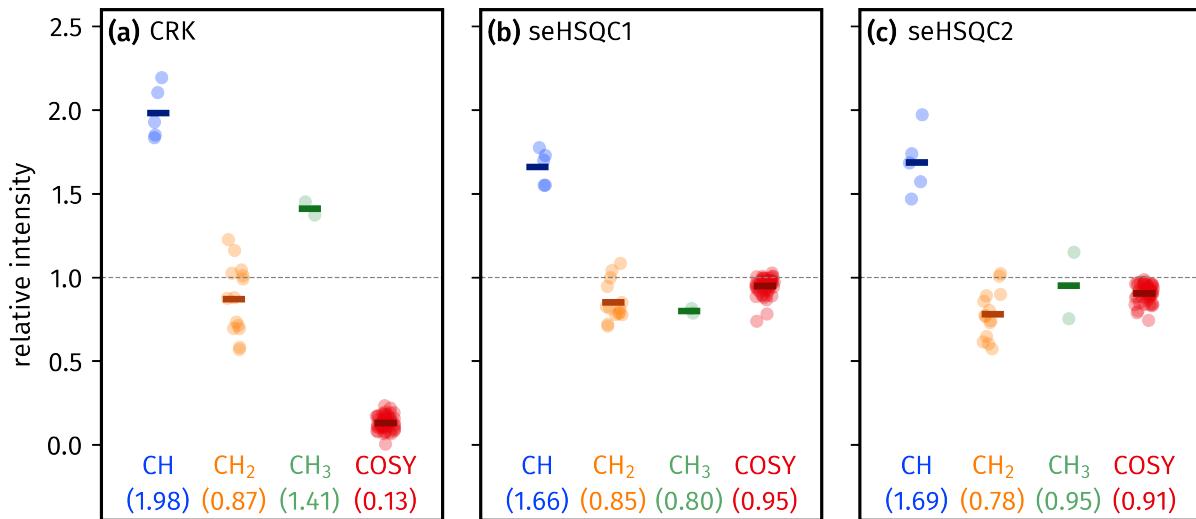


Figure 4.16: Sensitivity comparisons with Δ' set to $1/(4 \cdot {}^1J_{CH})$. Multiplicity editing was not used. (a) Using the CRK seHSQC. (b) Using the seHSQC1 module. (c) Using the seHSQC2 module. Data code: 7A-201115.

It is interesting to note that the performance of the seHSQC1 module is much closer to that of the seHSQC2 here. The fact that seHSQC1 performs worse with the reduced value of $\Delta' = 1/(8 \cdot {}^1J_{CH})$ suggests that there are some inefficiencies in this section of the pulse sequence; however, there is no immediate explanation for this in the product operators (the relevant terms in the seHSQC1 are entirely similar to that in the CRK, save for a minus sign).

The conclusions drawn are entirely similar when multiplicity editing is enabled (fig. 4.17), so will not be further discussed.

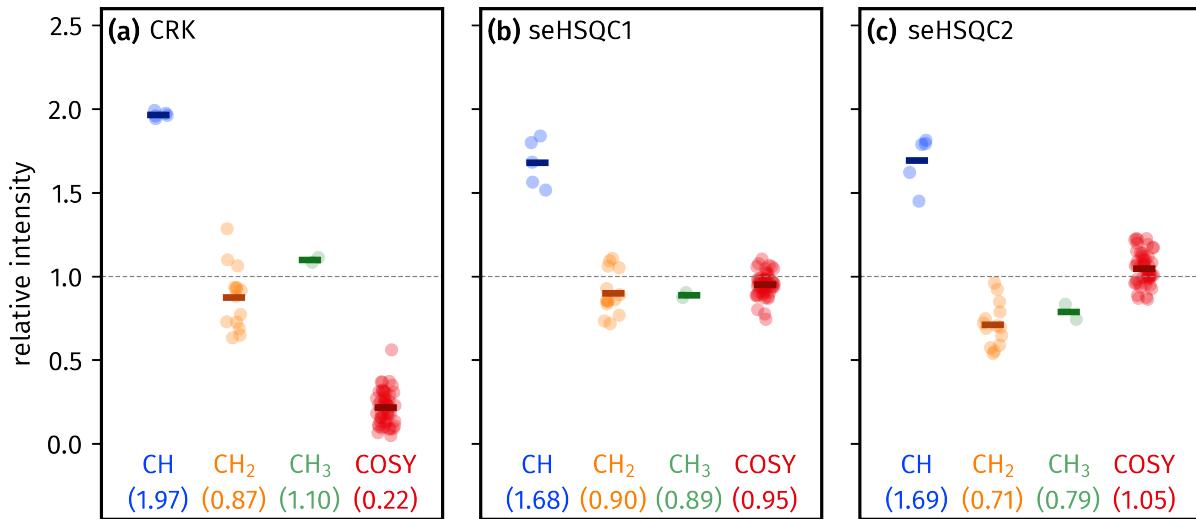


Figure 4.17: Sensitivity comparisons with Δ' set to $1/(4 \cdot {}^1J_{CH})$ and with multiplicity editing. (a) Using the edited CRK seHSQC. (b) Using the edited seHSQC1 module. (c) Using the edited seHSQC2 module. Data code: 7A-201115.

Optimal control for seHSQC1

One unresolved question is why the seHSQC1 module has a lower sensitivity than the seHSQC2, despite being shorter and containing fewer 180° pulses. There is also nothing in the product operator analysis to explain why this should be the case. One remaining possibility is the presence of some non-ideality in the pulses themselves, and the composite ^1H pulse is an obvious candidate for investigation.*

In order to investigate the extent to which this was responsible, I sought to use optimal control theory to develop a *universal rotation pulse* (URP) which could replace this composite pulse. Such a pulse must accomplish the transformations in eq. (4.10), namely $z \rightarrow y$ and $y \rightarrow x$.† This optimisation was performed using an interior-point algorithm (the default in Matlab `fmincon`) and the ESCALADE method⁵⁵ for the calculation of analytic derivatives. Pulse fidelity was calculated as an average over 51 ^1H spins, evenly distributed across a frequency range of 11 kHz (corresponding to 15.7 ppm at 700 MHz); the duration of the pulse was fixed at 200 μs , and 200 pulse points were used.

The ^{13}C 90° pulse which is simultaneously applied during the pulse sequence was not optimised together with this: because of the longer duration of the ^1H URP, the sequence must be modified slightly (fig. 4.18). In particular, the delays α and β must be chosen in order to satisfy the relations $2\varepsilon = \alpha + \beta$ (to ensure $^1\text{H}^{13}\text{C}$ magnetisation is refocused), and $\alpha = \beta + \tau$ (to ensure that for $^1\text{H}^{13}\text{C}$ magnetisation, the ^{13}C chemical shift evolves for a total duration of t_1).

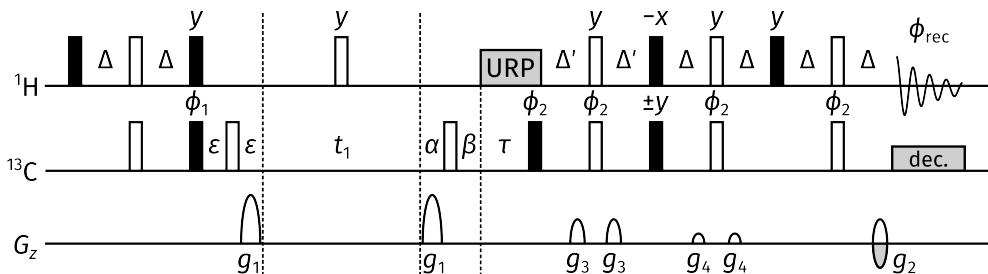


fig:sehsqc1_urp

Figure 4.18: seHSQC1 pulse sequence using a ^1H URP in place of the double- 90° composite pulse. τ is the difference in duration between the URP and the ^{13}C 90° hard pulse. Delays are set as: $\alpha = \varepsilon + \tau/2$; $\beta = \varepsilon - \tau/2$.

This optimisation process yielded a pulse with 99.99% fidelity: the (theoretical) performance of this pulse on z - and y -magnetisation is shown in fig. 4.19. However, when tested in the actual seHSQC1 experiment, this failed to yield any substantial difference compared to the original

*The concept of a *composite pulse*⁵⁴ is usually associated with greater efficiency / uniformity, but here I have used the term to loosely refer to a consecutive series of pulses.

†The effect of this pulse on x -magnetisation is not important for the seHSQC1, but is in fact fully determined by these two constraints: $UI_xU^\dagger = -iU(I_yI_z - I_zI_y)U^\dagger = -i(UI_yU^\dagger UI_zU^\dagger - UI_zU^\dagger UI_yU^\dagger) = -i(I_xI_y - I_yI_x) = I_z$. This reveals a geometric interpretation of this pulse element: it is actually a 120° rotation about the vector $(1, 1, 1)$, which is closely related to the C_3 symmetry operation in the octahedral point group.

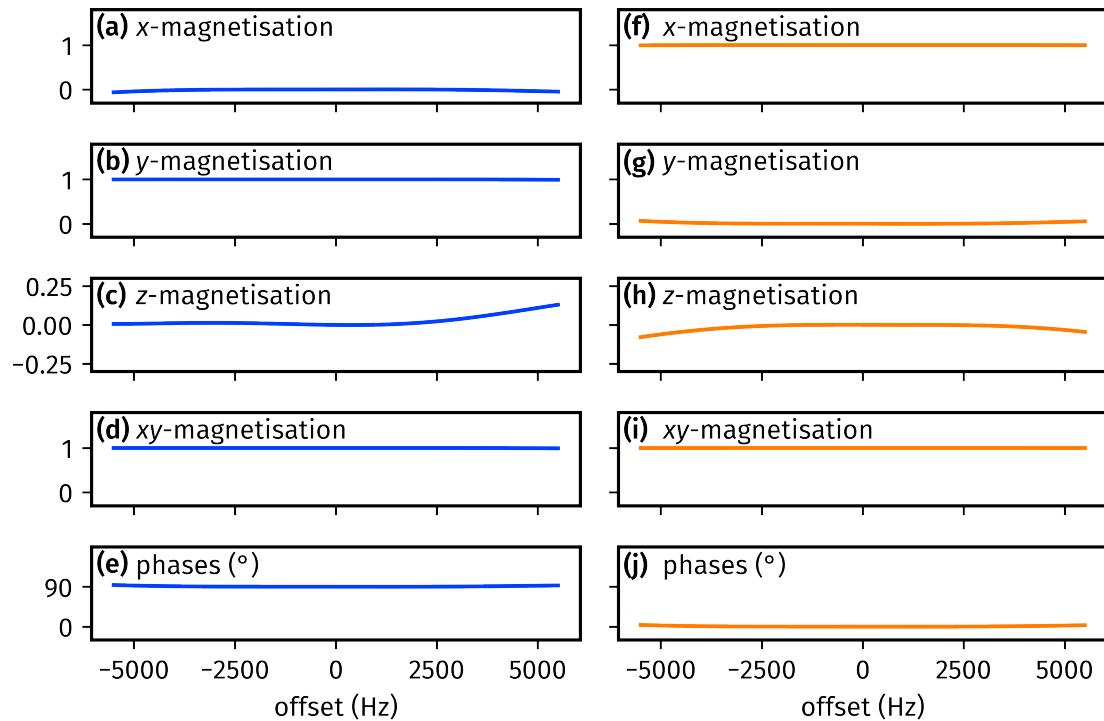


Figure 4.19: Simulated performance of URP used for seHSQC1 module on z - and y -magnetisation. The pulse fidelity was 99.99%. (a)–(e) Using z -magnetisation as input: the plots respectively show the amount of x -magnetisation, y -magnetisation, z -magnetisation, transverse magnetisation ($M_{xy} = \sqrt{M_x^2 + M_y^2}$), and the phase of the transverse magnetisation generated, as a function of offset frequency. (f)–(j) The same, but using y -magnetisation as input.

double 90° pulse (figs. 4.20a and 4.20b). Importantly, the performance still falls below that of the seHSQC2 module (fig. 4.20c). The reason for the poorer sensitivity therefore likely lies elsewhere.

BIG-BIRD versus ZIP

The final point in this section pertains to the implementation of the seHSQC2 module. As it stands, the isotope-specific rotation element placed at the start of the module is the ZIP element: its role is to effect 90° rotations with different phases on the ${}^1\text{H}^{\text{C}}$ and ${}^1\text{H}^{\text{!C}}$ magnetisation pools. However, such pulse elements have been known for a long time: these include TANGO,⁵⁶ BANGO,⁵⁷ BIRD,^{58–60} BIG-BIRD,⁶¹ and TIG-BIRD.⁶² In particular, the BIG-BIRD element can be designed to accomplish the same overall effect as the ZIP element. For the seHSQC2 without multiplicity editing, we can therefore replace the ZIP element with

$$45^\circ_{45^\circ}({}^1\text{H}) - 2\Delta - 180^\circ({}^1\text{H}, {}^{13}\text{C}) - 2\Delta - 45^\circ_{225^\circ}({}^1\text{H}), \quad (4.12) \quad \text{(eq:big_bird_unedited)}$$

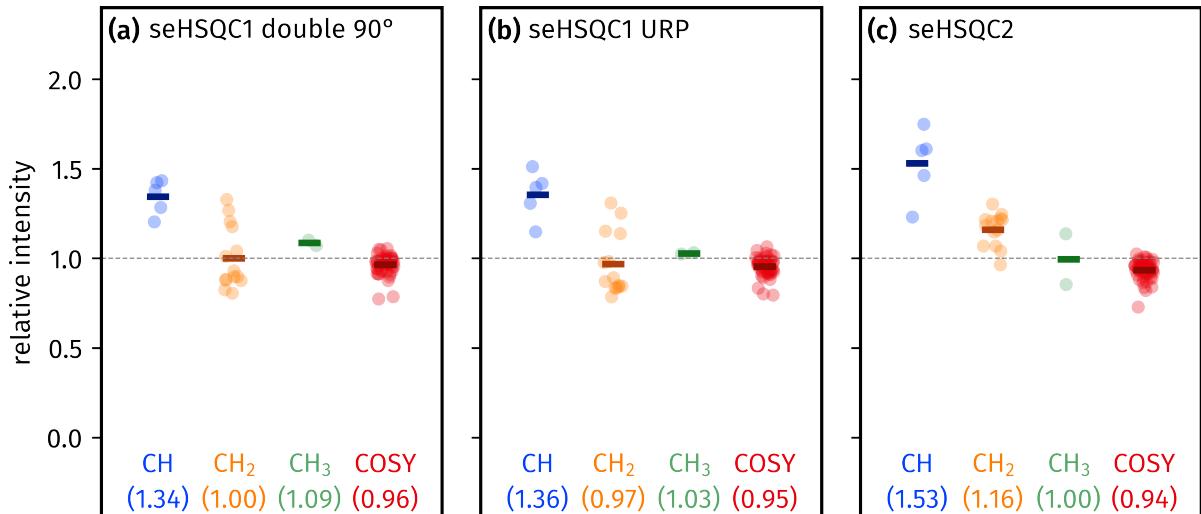


Figure 4.20: Sensitivity comparisons using the seHSQC1 URP. Δ' was set to $1/(8 \cdot ^1J_{\text{CH}})$; no multiplicity editing was used. Note that a different dataset was used for this figure and fig. 4.10, so the numbers are very slightly different. (a) Original seHSQC1 module with double ^1H 90° pulse. (b) seHSQC1 using the optimised URP shown in fig. 4.19. (c) seHSQC2 module for comparison. *Data code:* 7A-220110.

where β_ϕ represents a hard pulse with flip angle β and phase ϕ . For the edited seHSQC, the phases must be altered slightly:

$$45^\circ_{315^\circ}({}^1\text{H}) - 2\Delta - 180^\circ({}^1\text{H}, {}^{13}\text{C}) - 2\Delta - 45^\circ_{135^\circ}({}^1\text{H}). \quad (4.13) \quad \{\text{eq:big_bird_edited}\}$$

This BIG-BIRD version of the seHSQC2 was also evaluated. However, its performance in all respects was not as good as the ZIP version: both the seHSQC sensitivity itself, as well as the sensitivity of the later CLIP-COSY, were lower when the BIG-BIRD element was used.

4.3.2 ${}^{15}\text{N}$ HMQC

In § 4.3.3, I will discuss how the sensitivity-enhanced HSQC modules developed above may be adapted into ${}^1\text{H}-{}^{15}\text{N}$ experiments. However, before that, I make a slight detour to cover the ${}^{15}\text{N}$ HMQC experiment, which (up until my DPhil) was the experiment of choice for detecting one-bond ${}^1\text{H}-{}^{15}\text{N}$ correlations.

CTP gradient scheme

The HMQC module is based on the ASAP-HMQC reported by Kupče and Freeman,¹³ which used a symmetric gradient scheme similar to that in the seHSQC modules previously described (fig. 4.22a). However, in the NOAH module,³⁰ bipolar gradient pulse pairs were placed before and after t_1 (fig. 4.22b). This was likely implemented in order to allow the final gradient, g_2 , to have as large an amplitude as possible. In heteronuclear experiments, this final gradient is particularly

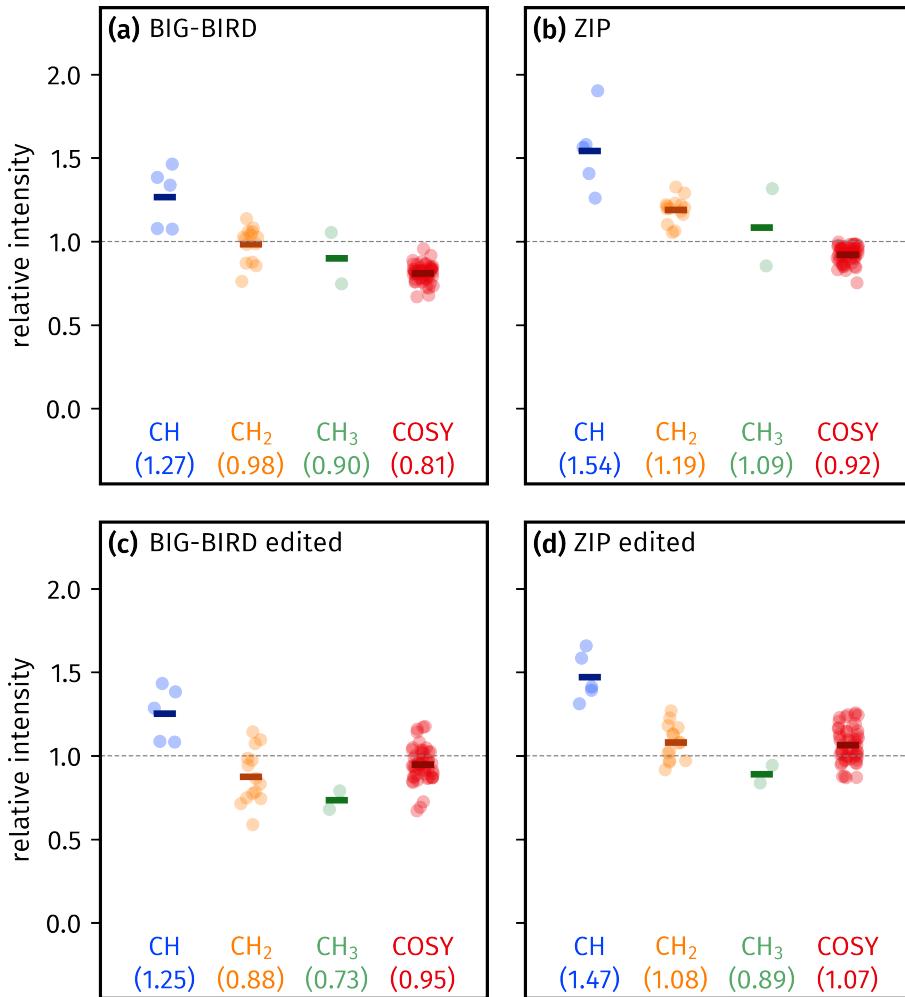


Figure 4.21: Sensitivity comparisons of BIG-BIRD and ZIP pulse elements in seHSQC2 module. The reference dataset being compared against is still the NOAH-2 SC^c (but with editing in (c) and (d)). The delay Δ' was set to $1/(8 \cdot ^1J_{\text{CH}})$. (a) Unedited seHSQC2 using BIG-BIRD. (b) Unedited seHSQC2 using ZIP. (c) Edited seHSQC2 using BIG-BIRD. (d) Edited seHSQC2 using ZIP. Data code: 7A-201115.

important for dephasing bulk magnetisation which is transverse just prior to detection (due to pulse imperfections or relaxation). If this gradient is too weak, this unwanted magnetisation will be incompletely dephased, leading to artefacts in the resulting spectrum.

Strategies to maximise this gradient amplitude are particularly crucial in ^1H - ^{15}N experiments (as compared to ^1H - ^{13}C experiments) for two reasons. Firstly, the natural abundance of ^{15}N (0.36%) is even smaller than ^{13}C (1.1%), meaning that better suppression must be achieved in order for the artefacts to not obscure the signal. Secondly, the gyromagnetic ratio of ^{15}N is also smaller: thus, since $g_2/g_1 \propto \gamma_{\text{N}}/\gamma_{\text{H}}$, an unmodified pulse sequence will naturally have a smaller g_2 .

In this respect, the four-gradient scheme in fig. 4.22b is superior to the two-gradient scheme,

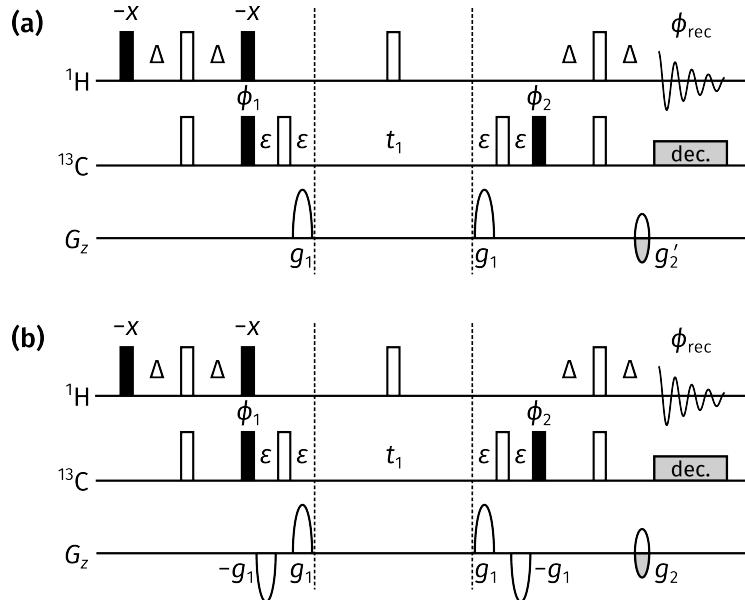


fig:noah_hmqc_hgmd

Figure 4.22: (a) With two encoding gradients around t_1 . (b) With four encoding gradients around t_1 . Phase cycling is performed with $\phi_1 = (x, -x)$, $\phi_2 = (x, x, -x, -x)$, and $\phi_{rec} = (x, -x, -x, x)$. The delay Δ is set to $1/(4 \cdot {}^1J_{NH})$. Gradient amplitudes are: $g_1 = 80\%$; $g_2 = \pm 32.4\%$; $g'_2 = g_2/2$.

because the gradient g_2 will have an amplitude of $4\gamma_N g_1/\gamma_H$. However, when used in a NOAH supersequence, this leads to wing artefacts in downstream modules, since bulk ${}^1H^{15N}$ magnetisation effectively does not experience any coherence order selection during t_1 . This motivates a return to the two-gradient scheme of fig. 4.22a. To compensate for the fact that the decoding gradient g'_2 has half of the amplitude of g_2 , all CTP gradients were instead *lengthened* from their usual duration of 1 ms to 2.5 ms. This ensures that any stray transverse bulk magnetisation at the end of the HMQC module is effectively dephased.

The HMQC spectra thus obtained are shown in fig. 4.23. These, and all other ${}^1H-{}^{15}N$ spectra in this chapter, were acquired assuming a ${}^1J_{NH}$ value of 90 Hz. The first column shows the spectra obtained with the original four-gradient scheme: although the artefacts at 2.2 ppm are reasonably well-suppressed in the HMQC module (fig. 4.23b), the CLIP-COSY spectrum clearly has a set of wing artefacts (fig. 4.23c). (Note that the wing artefacts occur at different frequencies compared to the ${}^{13}C$ seHSQC case, because the t_1 increment in the ${}^{15}N$ HMQC module is different.)

The second column shows what happens when the two-gradient scheme is adopted without changing the gradient duration. Although the CLIP-COSY wing artefacts disappear (fig. 4.23f), the HMQC artefacts are over twice as intense (fig. 4.23e), and (in this case) have comparable intensity to the desired peaks.

By using the two-gradient scheme and increasing the gradient duration to 2.5 ms, we obtain the best of both worlds: the HMQC artefacts are well-suppressed (fig. 4.23h, in fact even better than

in the original spectrum), and the CLIP-COSY is free of wing artefacts (fig. 4.23*i*). The only drawback is a slight loss in signal intensity in the HMQC, which arises due to diffusion and relaxation during the longer pulse sequence. However, this decrease is only on the order of 5% (for this particular case), which is a totally acceptable price to pay in return for the improved spectral quality.

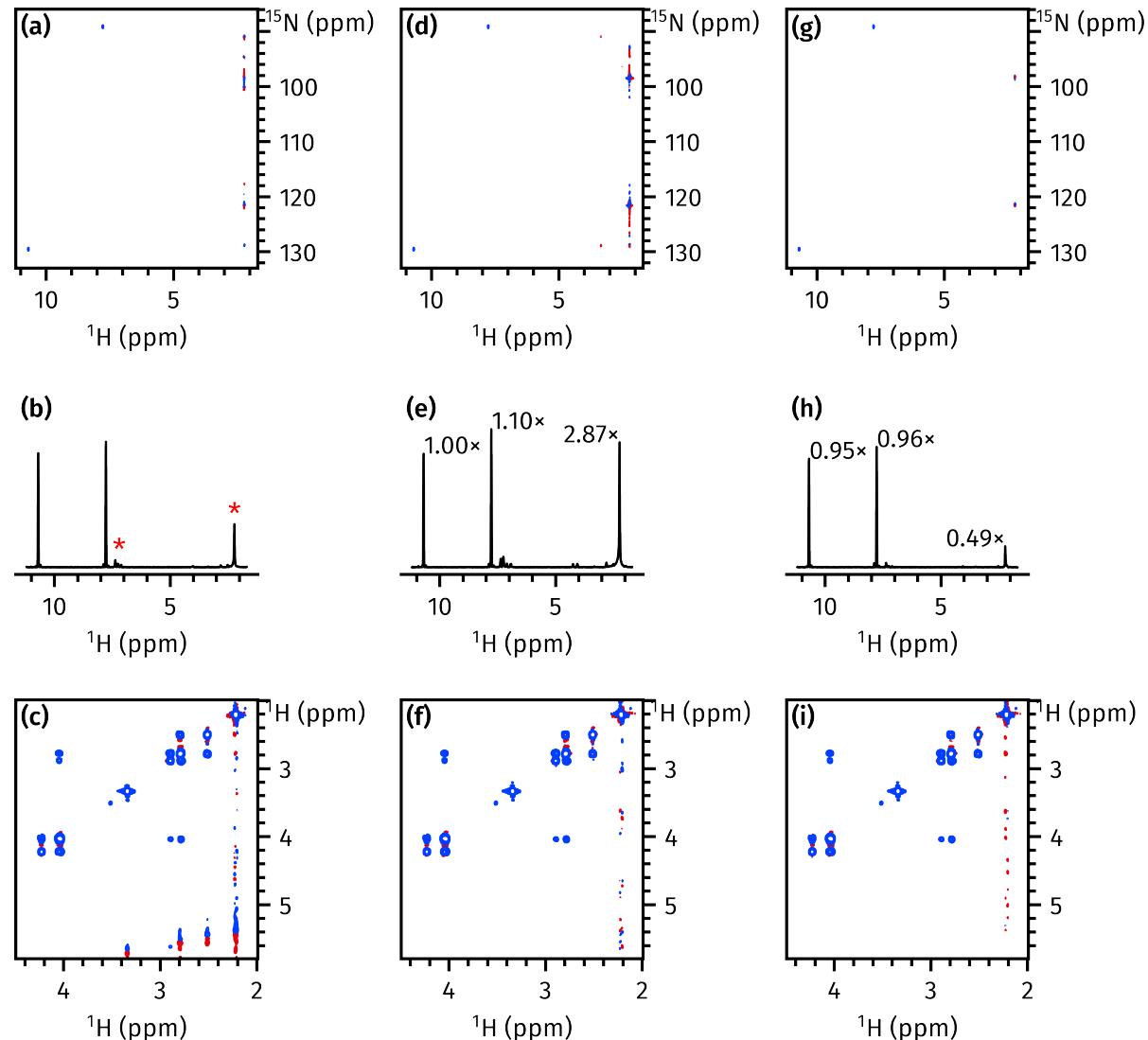


Figure 4.23: Comparison of HMQC and CLIP-COSY spectra obtained from NOAH-3 $M_N S^+ C^c$ supersequences, acquired using different HMQC gradient schemes. In the first row, the HMQC spectrum itself is shown. In the second row, the positive projection of the HMQC spectrum onto the F_2 axis is shown; the numbers indicate peak intensities with respect to the reference dataset (the left column). The asterisks indicate artefacts arising from bulk magnetisation which is not sufficiently dephased by the final gradient. In the third row, (an inset of) the CLIP-COSY spectrum is shown. (a)–(c) Using the four-gradient scheme of fig. 4.22*b*, with 1 ms gradients. (d)–(f) Using the two-gradient scheme of fig. 4.22*a*, with 1 ms gradients. (g)–(i) Using the two-gradient scheme of fig. 4.22*a*, with 2.5 ms gradients. Data code: 7Z-200926.

It is of some interest to check whether such a long gradient is truly required. I therefore ran the two-gradient HMQC experiment with a series of gradient durations from 1 ms to 2.5 ms; the signal and artefact intensities in each of these experiments are plotted in fig. 4.24. Generally, there is little variation in the intensities of the two desired peaks. The artefact intensity is more erratic, which possibly reflects the fact that gradient dephasing varies sinusoidally with the gradient duration τ due to the spatial integral

$$\frac{1}{L} \int_{-L/2}^{L/2} \exp(i\gamma G z \tau) dz = \frac{\sin(\gamma G L \tau / 2)}{\gamma G L \tau / 2}, \quad (4.14) \quad \text{eq:gradient_dephasing}$$

where L is the sample length and G the gradient amplitude (see also eq. (1.93)). Nevertheless, there is a clear decrease in the signal intensity as the gradient duration increases, which (at least in these datasets) is greatest with 2.5 ms gradients. Since the signal intensity is not affected much, I deemed this to be a perfectly suitable value.

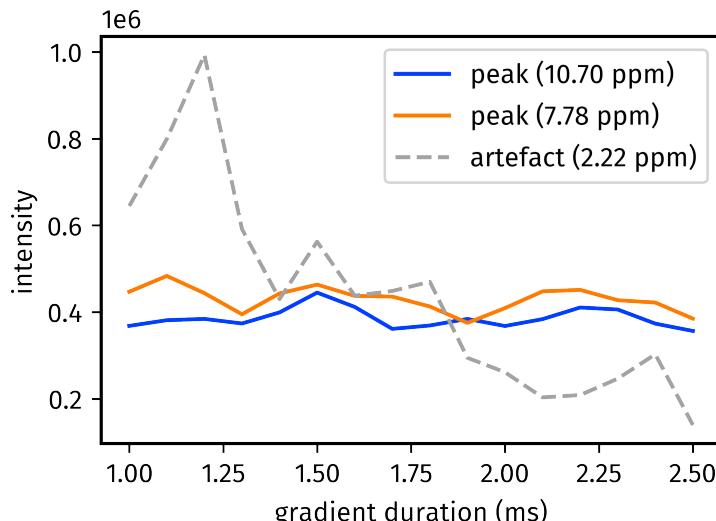


fig:hmqc_cnst16

Figure 4.24: Variation of HMQC signal and artefact intensity with CTP gradient duration.
Data code: 7Z-200926.

k- and SW-scaling

$^1\text{H}-^{15}\text{N}$ spectra of small molecules are often relatively sparse in the indirect dimension, and do not typically need to be acquired with the same resolution as a $^1\text{H}-^{13}\text{C}$ experiment. However, if these experiments were to be combined together within a supersequence, it is not ordinarily possible to toggle the resolution of each module individually. One way of circumventing this issue is to decrease the frequency with which the $^{15}\text{N} t_1$ duration is incremented in the pulse programme: this was first proposed by Parella et al. in the context of time-shared NMR,^{26,63} and in this thesis is called ‘k-scaling’. This means that each t_1 increment is acquired several times before t_1 is incremented; effectively leading to k times fewer t_1 increments, but with each

increment having k times the number of scans (after the data have been combined). An alternative is to simply increase the spectral width of the ^{15}N module (which is encoded as the CNST40 parameter in GENESIS).

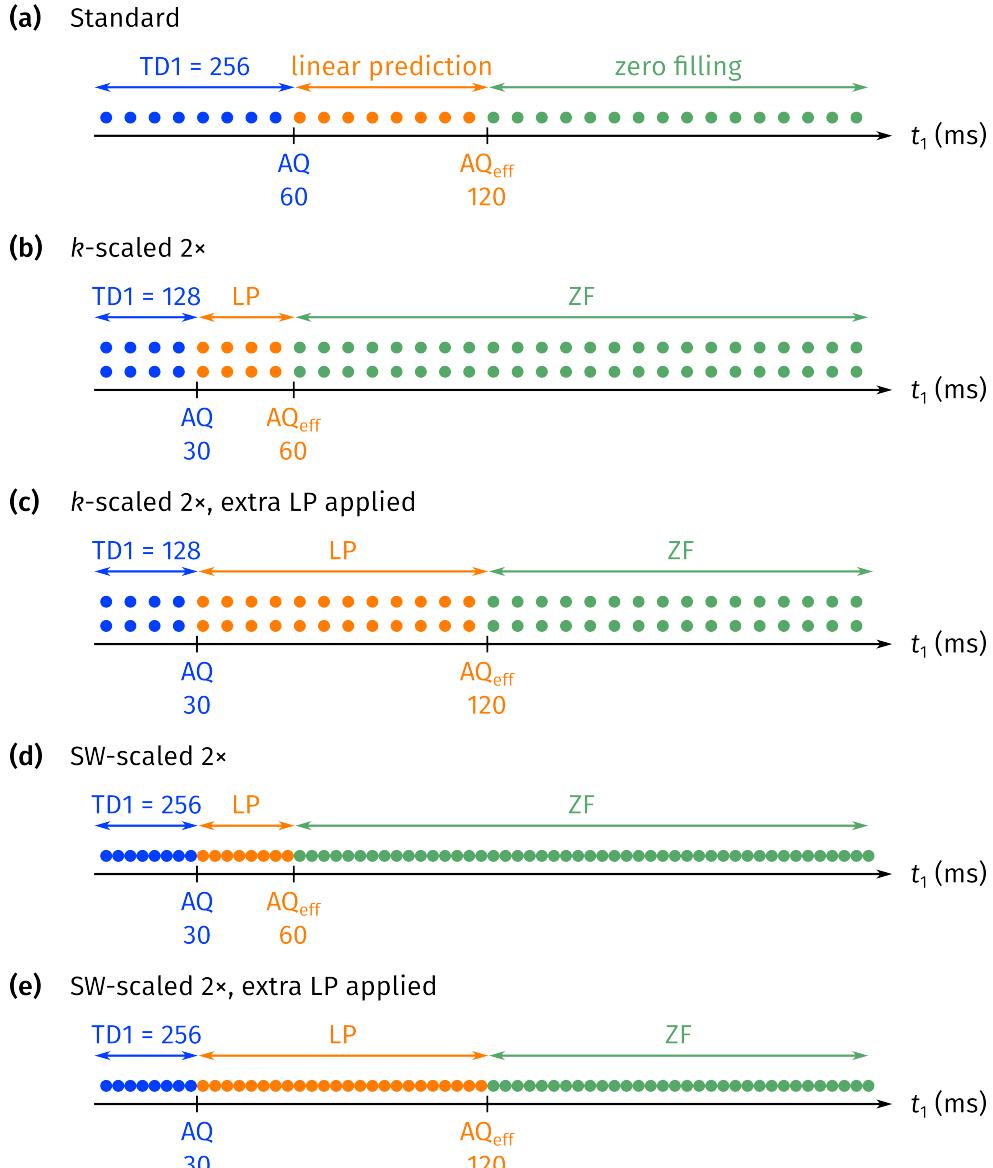


Figure 4.25: Pictorial representation of k - and SW-scaling in NOAH ^{15}N modules. Each blue dot represents 32 t_1 increments physically acquired as part of the experiment; the experimental time is proportional to the number of blue dots, and is constant across all of these experiments. Orange dots represent 32 t_1 increments obtained through forward linear prediction, and green dots represent zeroes. (a) The default experiment. (b)–(c) k -scaled experiment, without and with extra linear processing to restore the original value of AQ_{eff} . (d)–(e) SW-scaled experiment, without and with extra linear processing to restore the original value of AQ_{eff} .

Figure 4.25 illustrates all of these possibilities in detail. In a ‘typical’ NOAH supersequence (fig. 4.25a), 256 t_1 increments would be recorded (corresponding to the eight blue dots); for

the gramicidin sample used here, the largest value of t_1 (denoted as AQ) is around 60 ms. The default TopSpin processing routine for 2D NMR applies forward linear prediction (LP)^{64–67} in the indirect dimension, such as to double the number of points: the data points obtained through LP are indicated by orange dots. This leads to an *effective* AQ which is twice as large, denoted as AQ_{eff} . Finally, zero filling (ZF) is applied to increase the digital resolution to a given value, as represented by the green dots.

In the k -scaled spectrum, both AQ (and thus AQ_{eff}) are halved, in return for a doubling of the number of scans (fig. 4.25b). The resolution in the indirect dimension is proportional to AQ_{eff} ; thus, this also corresponds to a 2 \times decrease in resolution. However, this can be counteracted through the use of extra linear prediction (fig. 4.25c) to extend AQ_{eff} back to its original value. The SW-scaled spectra (figs. 4.25d and 4.25e) are almost entirely equivalent to the k -scaled spectra in terms of their effect on AQ and AQ_{eff} , except that the sampling schedule is very slightly modified.

Figure 4.26 shows the effects of performing k - or SW-scaling on the ^{15}N HMQC module (without any extra linear prediction beyond the default). In this case, these scaling procedures can in fact lead to significant gains in SNR (as measured by peak heights), with up to 2.6 \times improvements being observed at the extreme of 8-fold k - or SW-scaling. This likely arises because the J_{HH} modulation in the indirect dimension (visible in the standard spectrum, fig. 4.26b) is no longer being resolved: the largest gains are attained for the peak at 123 ppm, where this modulation is especially prominent.

The effect of adding extra linear prediction is shown in fig. 4.27. In general, the combination of scaling plus LP leads to improvements in spectral SNR of up to 6 \times . However, this is accompanied by distortions in the F_1 multiplet structure, especially for the k -scaled spectra (figs. 4.27c to 4.27h): it is possible that the corresponding SW-scaled spectra are not so heavily distorted because there are more data points to extrapolate from.

However, caution should be exercised when interpreting these results. Although this improvement in SNR is genuine, it is not necessarily the case that this represents a true improvement in *detection sensitivity*: in other words, processing techniques such as LP (and to some extent, NUS) do not always allow for better discrimination between signal and noise.^{68,69*} Further tests would need to be done on more dilute samples[†] to ascertain the benefits of the scaling-plus-LP routine on spectra with less intrinsic SNR.

*In recent years, this issue has been investigated more thoroughly in the context of NUS.⁷⁰

[†]Or, perhaps, less sensitive instruments—although the 700 MHz cryoprobe used here was the only option for these supersequences.

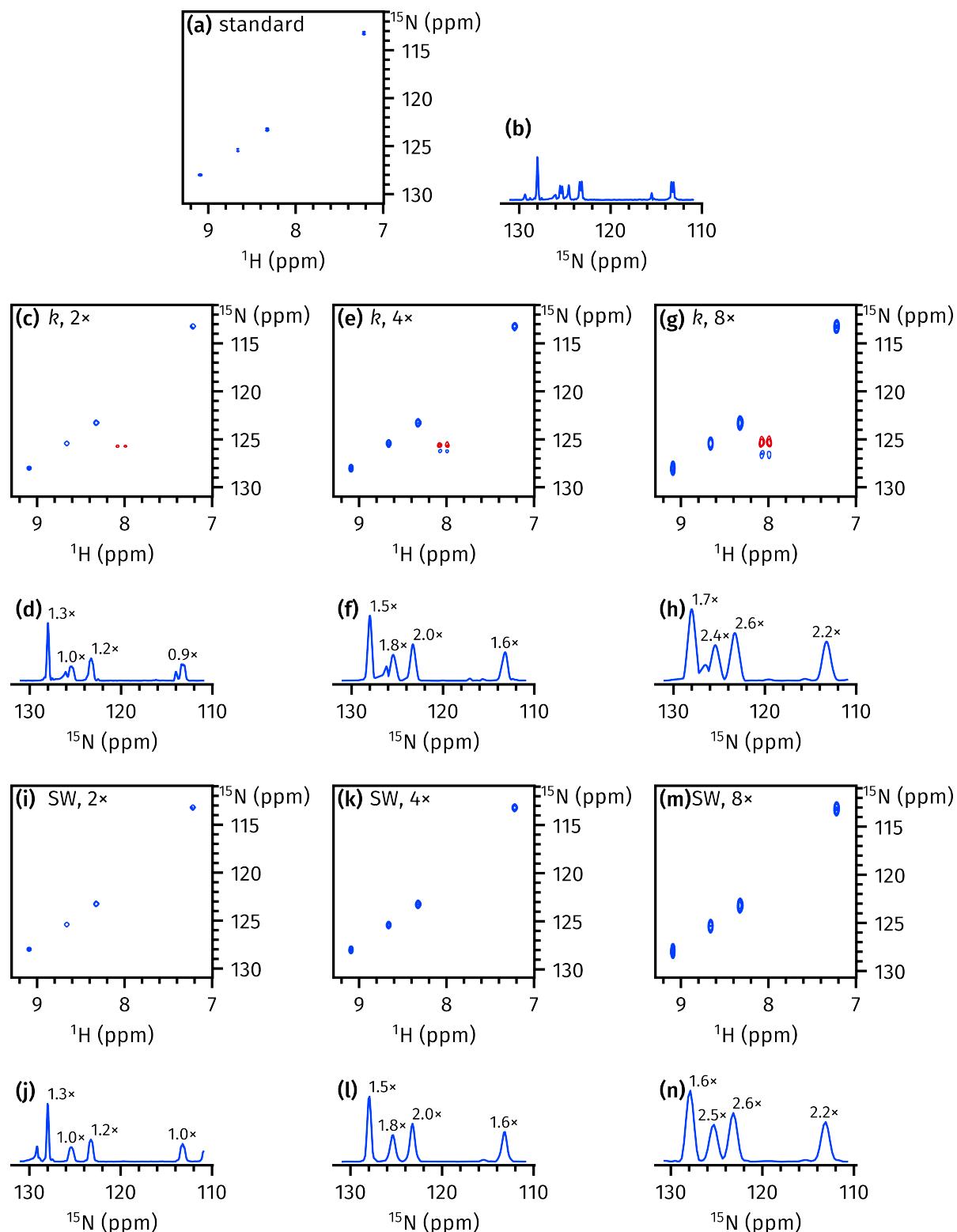


Figure 4.26: Effects of k - and SW-scaling on NOAH HMQC spectrum (taken from NOAH-3 $\text{M}_\text{N}\text{S}^+\text{C}^\text{c}$ supersequences). Each HMQC spectrum is shown together with a positive projection onto the F_1 axis. The relative SNR of each peak, with respect to the standard spectrum, is indicated on each of the other projections. (a)–(b) Standard spectrum. (c)–(h) k -scaled spectra. (i)–(n) SW-scaled spectra. Data code: 7G-210310.

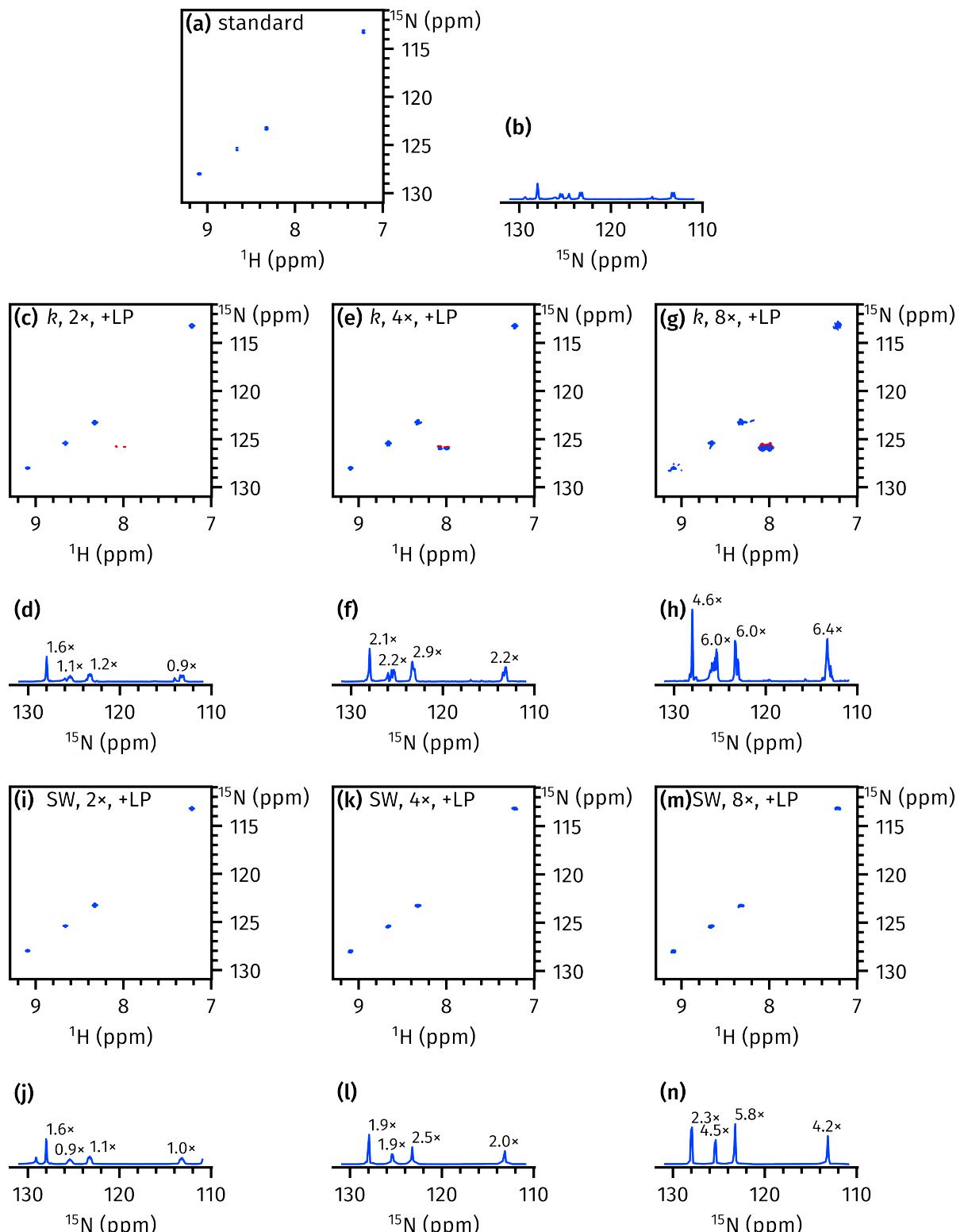


Figure 4.27: The same as in fig. 4.26, but with extra linear prediction applied to all scaled spectra to bring AQ_{eff} up to its original value in the standard spectrum. Linear prediction of k times more points leads to a \sqrt{k} increase in noise; to account for this, all spectra are plotted with the same noise level. (a)–(b) Standard spectrum. (c)–(h) k -scaled spectra. (i)–(n) SW-scaled spectra. Data code: 7G-210310.

4.3.3 ^{15}N sensitivity-enhanced HSQC

I now move on to the ^{15}N seHSQC experiment, which (for small molecules) is an improvement on the HMQC experiment for two reasons: firstly, the PEP transfer grants additional sensitivity (and can be optimised for NH groups); and secondly, splittings due to $^nJ_{\text{HH}}$ are not observed in the indirect dimension.

Unsurprisingly, much of the discussion in the HMQC section is also relevant here. For example, although both seHSQC modules already use two-gradient schemes and therefore do adequately suppress wing artefacts in later modules, the CTP gradient duration must still be lengthened in order to suppress bulk artefacts in the seHSQC itself.

Sensitivity analysis

To quantify these sensitivity gains, I ran several NOAH-3 XS⁺C^c experiments, where the first module was a ^{15}N HSQC, seHSQC1, or seHSQC2 module (fig. 4.28). Including the HSQC module in this analysis allows us to determine how much of the sensitivity gain is due to multiplet collapse, and how much is due to the sensitivity enhancement. The value of both delays Δ and Δ' were set to $1/(4 \cdot ^1J_{\text{NH}})$; no multiplicity editing was used. The peak at 8.0 ppm is folded and is therefore omitted from this analysis.

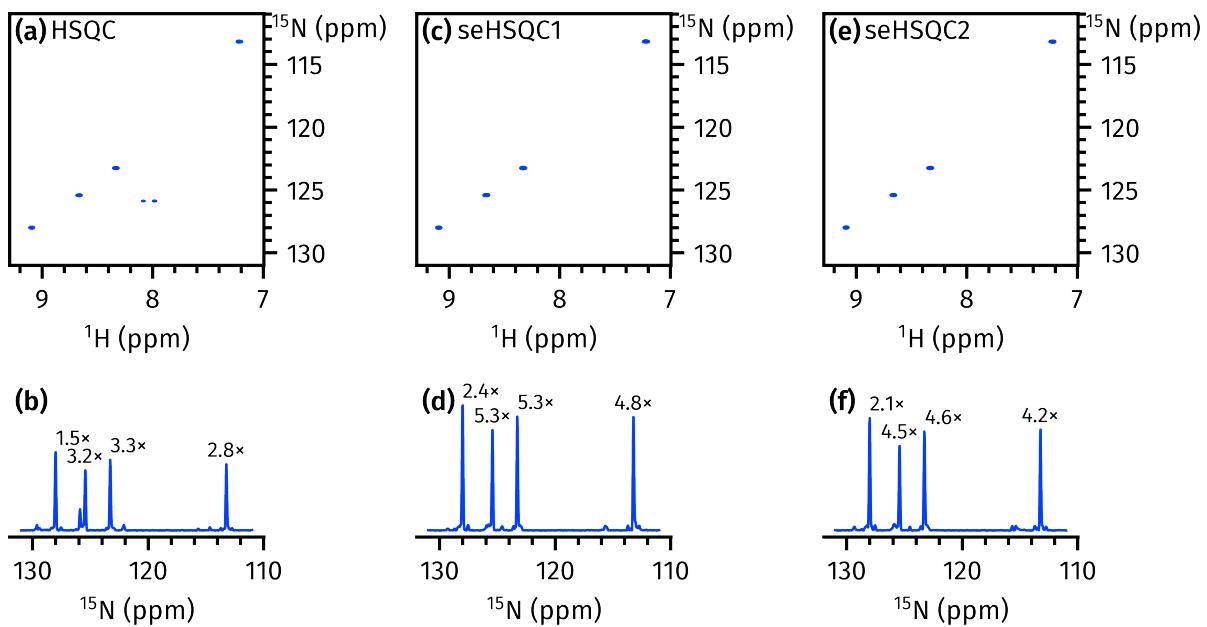


Figure 4.28: Comparison of NOAH ^1H - ^{15}N modules; the 2D spectra and their positive projections onto the F_1 axis are shown. Numbers above each peak in the projections indicate sensitivity increases compared to the HMQC module (see, e.g., figs. 4.26a and 4.26b). (a)-(b) HSQC module. (c)-(d) seHSQC1 module. (e)-(f) seHSQC2 module.

As can be seen, the HSQC alone allows for substantial sensitivity increases over the HMQC with

over 3 \times gains being accomplished simply due to multiplet collapse. The seHSQC experiments push this even further, with an approximate 1.5 \times sensitivity improvement for all peaks. In contrast to the ^{13}C case, the ^{15}N seHSQC1 module outperforms seHSQC2 in terms of raw sensitivity, though (as before) there is no clear explanation for this subtle difference.

[Compare other modules](#)

CTP gradient duration

k -Scaling and linear prediction

4.3.4 HSQC-TOCSY

subsec:noah_hsqctocsy
HSQC + DIPSI + HSQC combos

Extension to HSQC-TOCSY

Cite ASAP work (Luy)

4.3.5 HSQC-COSY

subsec:noah_hsqccosy
Comparison of several versions of HSQC-COSY (JACS Au SI)

4.3.6 2DJ and PSYCHE

subsec:noah_2djpsyche
'Normal' homonuclear modules are almost trivial to include in NOAH supersequences: since these are placed at the end of supersequences, there is rarely any need to modify them, as they do not need to preserve any magnetisation. One exception to this is the family of 2DJ and pseudo-2D pure shift experiments (here typified by PSYCHE), where the spectral width in the indirect dimension is extremely small, on the order of 50 Hz. In such cases, the number of t_1 increments required is far smaller than for a typical 2D experiment. Since—by default—each module in a NOAH supersequence is acquired with the same number of increments, directly adding such modules to the end of a supersequence would therefore prove suboptimal.

However, as described previously in § 4.3.3, it is possible to reduce the number of t_1 increments for a particular module, and in exchange, increase the number of scans recorded for that module. In the context of ^{15}N modules, this was a 'special' procedure referred to as k -scaling; however, for these homonuclear modules, it is natural and necessary. One difference in the implementation is that, instead of specifying a value k by which TD1 is scaled down by and NS scaled up by, the user is allowed to directly specify the number of t_1 increments as an integer (CNST37 in TopSpin). This value must be a divisor of TD1/NBL, i.e., the number of t_1 increments for all other modules. The indirect-dimension spectral width is specified as CNST38.



Figure 4.29: Spectra obtained from a NOAH-3 $S_N^+S^+J$ supersequence. (a) ^{15}N seHSQC (256 t_1 increments and 2 scans per increment). (b) ^{13}C seHSQC. (c) PSYCHE 2DJ (32 t_1 increments and 8 scans per increment). Data code: 7G-201028.

The modules thus implemented include the standard magnitude-mode 2DJ experiment, the PSYCHE 2DJ experiment,⁷¹ the original 1D PSYCHE,⁷² and the 1D TSE-PSYCHE.⁷¹ An example of the data thus obtained (with the PSYCHE 2DJ) is shown in fig. 4.29. One particular advantage of including PSYCHE-type modules in NOAH supersequences is that sensitivity is not likely to be at a premium: this is partly because of the increased number of scans, but also partly because other 2D experiments have comparably low sensitivity (meaning that in the time needed to acquire a HSQC, for example, the PSYCHE experiment will also have sufficient sensitivity). This allows the user to choose a relatively small flip angle (ca. 10°) for the PSYCHE saltire pulses in order to minimise artefacts from imperfect decoupling.

On top of that, for the 1D (TSE-)PSYCHE sequences, the extra transients can be used to perform SAPPHIRE averaging:⁷³ in this procedure, each chunk of the pure shift interferogram is collected multiple times while varying the point in time where the J-couplings are perfectly refocused. Summation of these data leads to the suppression of artefacts which arise due to the periodic J-modulation in the interferogram. This averaging is somewhat analogous to a phase cycle, and performing an 8-step SAPPHIRE averaging procedure (for example) would often require the experiment to be lengthened beyond the duration which is truly necessary: however, in the context of NOAH this is obtained essentially ‘for free’. In the GENESIS pulse programmes,

this feature is enabled by default (it can be turned off by selecting the original modules using developer mode).

```
./figures/noah/psyche_wrongsw1.png
```

Figure 4.30: A series of 1D PSYCHE spectra obtained from the NOAH-3 $S_N^+ S^+ P$ super-sequence (saltire flip angle of 15°). (a) 1D PSYCHE spectrum acquired without SAPPHIRE averaging and an incorrect chunk size which is not an integral number of complex data points. (b) 1D PSYCHE spectrum with 8-step SAPPHIRE averaging and an incorrect chunk size; this manifests as phase errors which cannot be corrected. (c)–(d) The same as (a) and (b), but with the chunk size automatically corrected in the pulse programme. The impact of the SAPPHIRE averaging is less obvious in these datasets (compare (c) and (d)), but is more striking when larger chunk sizes are used, as discussed in the original paper.⁷³ Data code: 7C-211123.

A final and more prosaic implementation detail is that in the 1D pure shift modules, the chunk size is automatically rounded to the nearest even multiple of the dwell time τ_{dw} (DW in TopSpin). This ensures that each chunk consists of an integral number of complex data points. Although this can be set by the user manually, it is very easy to forget, especially for someone not fully acquainted with the experiment; the results can be quite different, as illustrated in fig. 4.30. When an n -step SAPPHIRE averaging is used, the requirement is even stricter: the chunk size must be a multiple of $2n\tau_{dw}$. This is also encoded in the pulse programmes.

4.3.7 HMBC

The HMBC module is one which in fact does not fit perfectly into the NOAH principle of only exciting magnetisation which is needed. As described in § 4.1.3, the HMBC module should only

require magnetisation of protons which have long-range couplings to ^{13}C ; however, it ends up exciting *all* $^1\text{H}^{13}\text{C}$ magnetisation. This leads to sharply reduced, and also unbalanced, intensities of homonuclear modules which come later in the supersequences.

I made some early (and brief) attempts at devising a pulse sequence which sought to discriminate these two magnetisation components using a perfect echo.⁷⁴ However, this was quickly abandoned as it proved very difficult to *also* retain $^1\text{H}^{13}\text{C}$ magnetisation. I do not claim here that it is impossible to come up with a pulse sequence which does this, but it is certainly not easy, and ultimately I turned my focus to improving (rather than replacing) the HMBC module.

Suppression of one-bond artefacts

One of the issues with the NOAH HMBC module was that there were an unusual amount of one-bond artefacts, which arise from $^1\text{H}^{13}\text{C}$ magnetisation which is allowed to evolve during the pulse sequence. Generally, HMBC experiments seek to suppress this through the use of a low-pass J-filter (LPJF, see also § 3.4.7). The NOAH HMBC module *additionally* contains a *zz*-filter, which stores $^1\text{H}^{13}\text{C}$ magnetisation along $+z$ before the LPJF. Thus, in theory, one-bond artefacts should be suppressed in the NOAH HMBC to an even greater extent.

However, this expectation is not borne out: in some cases, the NOAH HMBC in fact has *more intense* one-bond artefacts when compared against a standard HMBC experiment (figs. 4.31a and 4.31c). Even performing an optimisation of the LPJF delays, as described in § 3.4.7, did not lead to any reduction in artefacts. I hypothesised instead that these artefacts arose from imperfect manipulation of $^1\text{H}^{13}\text{C}$ magnetisation by the *zz*-filter.* In particular, any *antiphase* magnetisation (of the form $I_x S_z$ or $I_y S_z$) generated after the *zz*-filter would be reconverted into in-phase I_x or I_y terms, which would not be destroyed by the LPJF. (The LPJF works based on the assumption that there is in-phase magnetisation at the beginning of it, which is true if the excitation element is just a ^1H 90° pulse.)

Such antiphase terms can be easily removed by the addition of a ^{13}C 90° pulse, which transforms them into a mixture of double- and zero-quantum terms: these are either unobservable, or can be efficiently dephased by CTP gradients. This technique is used in the CLIP-HSQC family of experiments,^{75,76} as well as the LPJF itself. In this case, we simply need to add an additional ^{13}C 90° pulse at the end of the *zz*-filter: this pulse is highlighted in fig. 4.2b. This results in a striking reduction of the one-bond artefacts, as shown in fig. 4.31b. The suppression accomplished with the addition of this 90° pulse is superior to that in the standard HMBC (fig. 4.31c), and is comparable to a standard HMBC with a third-order LPJF (fig. 4.31d). Of course, the NOAH module—which by default uses a second-order LPJF—can also be ‘upgraded’ to use a third-order

*It would be nice to back this up with simulations, but I did not have the time to run these.

LPJF. In GENESIS pulse programmes, this can be done using the `-DLP3` acquisition flag (although the results were not evaluated here).

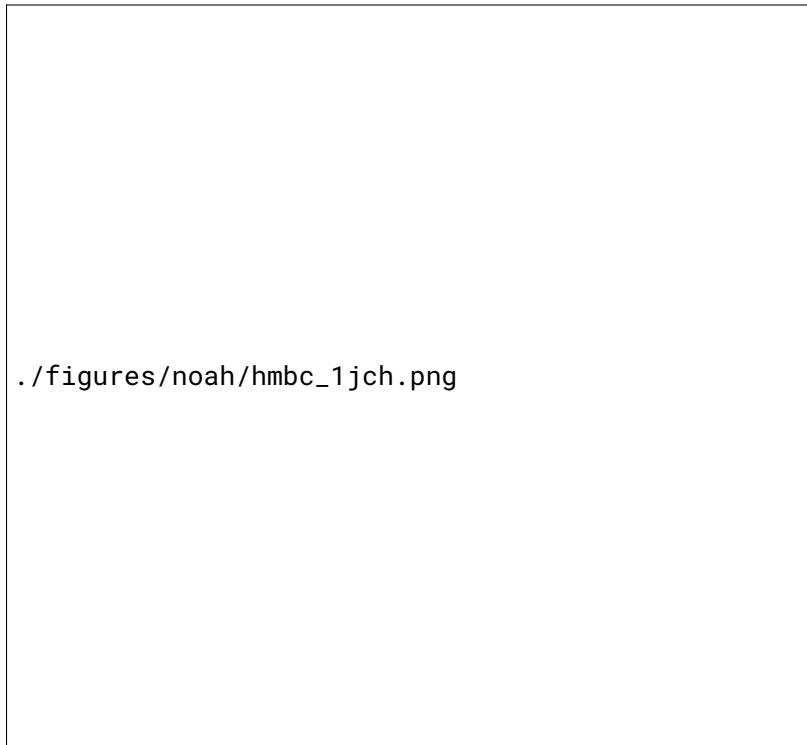


Figure 4.31: (a) NOAH zz-HMBC module without the additional 90° pulse. One-bond artefacts are highlighted in red. (b) NOAH zz-HMBC with the 90° pulse. (c) Standard library HMBC with a second-order LPJF. (d) Standard library HMBC with a third-order LPJF. Data code: 7A-210916.

Gradient selection schemes

Another point which was investigated (but bore less fruit) was the gradient scheme used for CTP selection. The NOAH module, as shown in fig. 4.2b, uses a ‘symmetric’ scheme where two gradients of equal amplitude surround the t_1 period: this encoding is later decoded by a third gradient just prior to acquisition. However, other choices exist: for example, the Bruker standard library HMBC (derived from the work of Cicero et al.⁷⁷) uses only two gradients in total, which have unequal amplitudes.

This gradient scheme cannot be directly used in a NOAH HMBC module, though. This is because the *zz*-filter element places $^1\text{H}^{\text{C}}$ magnetisation along the $+z$ axis just before the HMBC J-evolution delay (see fig. 4.2b). This magnetisation later experiences the ^1H 180° pulse in the middle of t_1 , which means that an extra 180° pulse must be added at the very end of the sequence to finally return it to $+z$. It is necessary to ensure that there is at least one gradient placed after this 180° pulse to ensure proper CTP selection (in the ‘symmetric’ scheme of fig. 4.2b, this is fulfilled by the gradient g_2).

```
./figures/pp/hmhc/noah_grads.png
```

Figure 4.32: Alternative CTP gradient schemes investigated for the NOAH HMBC module. The coherences selected for during each gradient are indicated above each gradient, using the same notation for product operators as described in the *Preface*: the ‘upper’ term (e.g. $+$ in \pm) refers to the echo experiment, and the ‘lower’ term to the antiecho experiment. So, for example, $+\pm$ refers to selection of I_+S_+ during the echo experiment and I_+S_- during the antiecho experiment. Gradient amplitudes are as described in the text. (a) ‘Asymmetric scheme A’, modified from the standard library sequence to include an additional 180° pulse and gradient. (b) ‘Asymmetric scheme B’, where the 180° pulse is shifted forward to the end of t_1 . (c) ‘Asymmetric scheme C’, which modifies the zz -filter instead of using an extra 180° pulse. (d) The original ‘symmetric’ scheme (the same as in fig. 4.2b), placed here for convenience.

Using this knowledge, it is possible to construct several ‘asymmetric’ gradient schemes:

- ‘Scheme A’ (fig. 4.32a) is modified from the Bruker standard library to include a 180° pulse and gradient at the end. The presence of an additional gradient means that there is a free parameter, here denoted as α , which can be used to control the relative amplitudes of these three CTP gradients. The gradient amplitudes are chosen as follows:

$$\text{echo: } g_1 = gc_1 \quad g_2 = g \quad g_3 = gc_2 \quad (4.15) \quad \{\text{eq: noah_hmhc_grads_bg}\}$$

$$\text{antiecho: } g_1 = g \quad g_2 = gc_1 \quad g_3 = gc_2 \quad (4.16) \quad \{\text{eq: noah_hmhc_grads_bg}\}$$

where $c_1 = -\alpha(\gamma_H - \gamma_C)/(\gamma_H + \gamma_C)$ and $c_2 = (1 - \alpha)(\gamma_H - \gamma_C)/\gamma_H$. In principle g is also a free parameter; for maximum suppression of artefacts I chose a relatively large value of 80%.

- In ‘Scheme B’ (fig. 4.32b), the 180° pulse is shifted to immediately after t_1 , before any of the CTP gradients have been applied. This means that there is no need for a third gradient, and the CTP gradient amplitudes can be directly taken from the standard library sequence:

$$\text{echo: } g_1 = g \quad g_2 = gc \quad (4.17) \quad \{\text{eq: noah_hmhc_grads_bg}\}$$

$$\text{antiecho: } g_1 = gc \quad g_2 = g \quad (4.18) \quad \{\text{eq: noah_hmhc_grads_bg}\}$$

where $c = -(\gamma_H - \gamma_C)/(\gamma_H + \gamma_C)$ and $g = 80\%$.

- ‘Scheme C’ (fig. 4.32c) simply does not add a 180° pulse, but instead modifies the phases of the zz -filter in order to place $^1\text{H}^C$ magnetisation along the $-z$ axis during the HMBC J-evolution delay. Here, the gradient amplitudes are the same as those in the standard library sequence as well as in scheme B.

It is of interest to note two limiting cases of scheme A: when $\alpha = (\gamma_H + \gamma_C)/(\gamma_H - \gamma_C) \approx 1.67$, we have that $g_1 : g_2 : g_3 = 1 : -1 : \pm 2\gamma_C/\gamma_H$, which mimics the original ‘symmetric’ scheme (fig. 4.32d); and when $\alpha = 1$, we have that $g_3 = 0$, i.e. a return to the two-gradient tactic of schemes B and C. In the tests which follow, I ran scheme A with $\alpha = 1.67$, $\alpha = 0.6$, and $\alpha = 0.3$.

All of the different HMBC versions above, plus the original ‘symmetric’ scheme in fig. 4.2b, were tested in the context of a NOAH-2 BS supersequence using the andrographolide sample (fig. 4.33). Since the HMBC module is the first module in this supersequence, the values here are an accurate reflection of their intrinsic sensitivities. As can be seen, there is not much at all which separates the different versions (outliers with $> 2\times$ ‘sensitivity improvements’ can be attributed to different J-modulation in the multiplet). The most sensitive of these is asymmetric scheme C, which may be explained by the fact that it has one fewer 180° pulse: however, this comes with an

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./figures/noah/hmhc_grad_sens.png
```

`fig:hmhc_grad_sens`

Figure 4.33: Sensitivities of various asymmetric HMBC gradient schemes, as compared to the symmetric scheme in fig. 4.2b. Each dot indicates one crosspeak in the HMBC spectrum; the numbers in parentheses are the average over all peaks. *Data code:* 7A-211226.

immediate drawback. Since scheme C places the $^1\text{H}^{\text{C}}$ magnetisation along $-z$ during the LPJF as well as the J-evolution delay Δ_{LR} (a total of ca. 70 ms), relaxation losses during this period lead to poorer retention of $^1\text{H}^{\text{C}}$ magnetisation for later modules, as shown by the decreased HSQC sensitivities in fig. 4.34. In contrast, all the other gradient schemes retain $^1\text{H}^{\text{C}}$ magnetisation equally well.

```
./figures/noah/hmhc_grad_sens_hsqc.png
```

`fig:hmhc_grad_sens_hsqc`

Figure 4.34: Sensitivities of the HSQC module in a NOAH-2 BS supersequence, where the HMBC module is implemented using the gradient schemes of fig. 4.33. *Data code:* 7A-211226.

The final point worth studying is the quality of the HMBC spectrum itself. To do this, we need to look at the actual spectra (fig. 4.35). For the most part, the spectra are all the same; however,

there is a notable set of artefacts present in fig. 4.35b (scheme A with $\alpha = 1.67$) as well as fig. 4.35e (scheme B), highlighted in red boxes. These artefacts occur at the frequencies

$$(\Omega_1, \Omega_2) = \left(\Omega_S \pm \frac{\Omega_I}{2}, \Omega_I \right), \quad (4.19) \quad \text{eq:hmhc_wing_artefact}$$

and are in fact ‘wing’ artefacts similar to that observed in other modules [insert reference](#). In this case, they arise due to imperfect refocusing of the ^1H chemical shift during t_1 : specifically, whenever $I_z S_{\pm}$ terms are present during the second half of t_1 . Extra evidence for the origin of these artefacts comes from the observation that when the 180° pulse in the middle of t_1 is phase cycled, the artefacts are removed. In a standard HMBC, these terms would not be detected in the final FID; however, in this case, the addition of an extra 180° pulse after t_1 provides an opportunity for these to be converted back into observable spin- I magnetisation (through off-resonance effects or miscalibration).

The poorer performance may therefore be understood as follows: when scheme A is acquired with $\alpha = 1.67$, the gradients g_1 and g_2 have equal and opposite amplitudes, and so do not enforce any coherence selection on spin I during the second half of t_1 . Likewise, scheme B contains no gradients during the second half of t_1 .

The characteristics of these gradient schemes are summarised in table 4.4. As can be seen, the ‘best’ schemes are either the original symmetric scheme, or asymmetric scheme A with $\alpha \neq 1.67$. However, there is not much difference between these: it is not clear whether the improvement in sensitivity is reproducible across a wide range of samples, and in any case, the gains are extremely marginal.

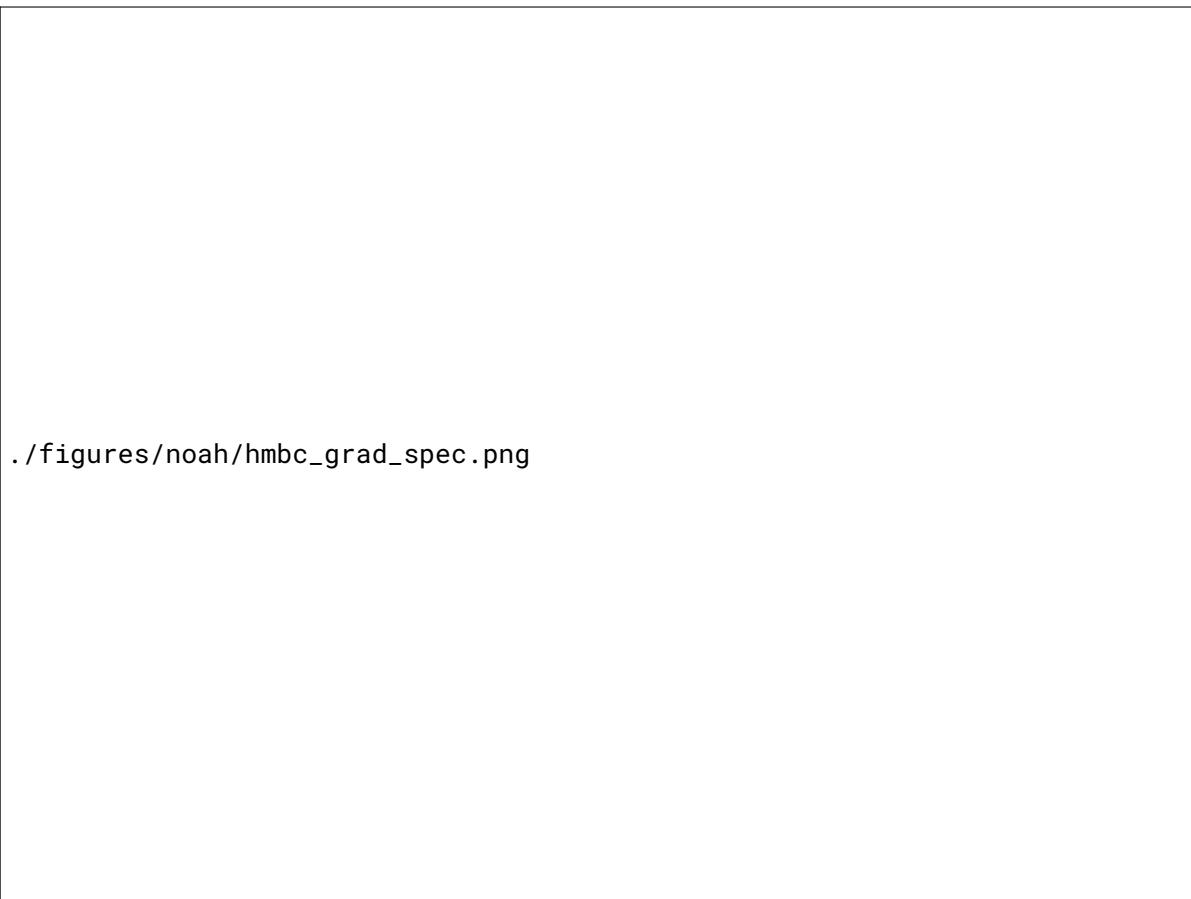
Gradient scheme	HMBC sensitivity	HSQC sensitivity	Wing artefacts
Symmetric	1	1	No
Asymmetric A, $\alpha = 1.67$	1.05	0.99	Yes
Asymmetric A, $\alpha = 0.6$	1.04	0.99	No
Asymmetric A, $\alpha = 0.3$	1.05	0.99	No
Asymmetric B	1.06	1.01	Yes
Asymmetric C	1.09	0.71	No

tbl:hmhc_grads

Table 4.4: Comparison of HMBC gradient schemes discussed in this section: the data are a summary of figs. 4.33 to 4.35. Data code: 7A-211226.

Other artefacts

It has been established that the HMBC module (and supersequences containing it) are not fully ideal in terms of magnetisation preservation. However, there are also some other curious phenomena which have not been fully described in the literature. One of these is the presence



`./figures/noah/hmhc_grad_spec.png`

Figure 4.35: HMBC spectra acquired with the gradient schemes of fig. 4.33. Extra ‘wing’ artefacts present in two of the spectra (asymmetric scheme A with $\alpha = 1.67$, (b), and asymmetric scheme B, (e)) are highlighted in red boxes. *Data code:* 7A-211226.

of *inverted peaks* in the homonuclear X module(s) in a NOAH-3 BSX supersequence: this is illustrated in fig. 4.36a with the CLIP-COSY module ($X = C^c$). It is not clear why this occurs, because the HMBC module (and the gradients which follow) should dephase all $^1H^{1C}$ magnetisation. Although this leads to reduced sensitivity in the homonuclear module, in that the signal derives from polarisation which has recovered during the preceding FIDs, it is not clear why this polarisation should be *negative*. One clue lies in the fact that these peaks are very sensitive to the 1H 90° pulse width: simply changing this by 0.5 μs is sufficient to restore the correct signal sign (fig. 4.36b).

The modules placed between the HMBC and the homonuclear module also play an important role. When *two* HSQC modules are used, i.e. a NOAH-4 BSSC^c supersequence (using $f = 0.7$ as described in § 4.3.4—although this is unlikely to matter), the negative peaks are no longer observed (fig. 4.36c). In fact, having *no modules* between the HMBC and the homonuclear module is also (at least sometimes) acceptable: a separate set of data shows that the inverted peaks in an NOAH-3 BS⁺C^c experiment are not seen in a NOAH-3 S⁺BC^c supersequence

```
./figures/noah/hmhc_invert_1.png
```

Figure 4.36: (a) CLIP-COSY from a NOAH-3 BSC^c supersequence, acquired with a ¹H 90° pulse width of 11.28 μ s (this value was obtained using the POISE calibration described in § 3.4.1). An inverted peak is visible at 2.25 ppm. (b) The same, but acquired using a 90° pulse width of 11.78 μ s. (c) CLIP-COSY from a NOAH-4 BSSC^c supersequence. The 90° pulse width was 11.28 μ s, the same as in (a). *Data code:* 7Z-220214.

(figs. 4.37a and 4.37b). The use of ASAP mixing just before the homonuclear module does not remedy this (figs. 4.37c and 4.37d). Unfortunately, a good explanation for these artefacts has remained elusive.

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./figures/noah/hmhc_invert_2.png
```

Figure 4.37: (a) CLIP-COSY from a NOAH-3 BS⁺C^c supersequence. An inverted diagonal peak can be seen at 6.6 ppm. (b) From a NOAH-3 S⁺BC^c supersequence. (c)–(d) The same as (a) and (b), but with 40 ms ASAP mixing placed just before the CLIP-COSY module. *Data code:* 7A-211227.

^{15}N HMBC module

The entirety of this section has—until now—been devoted to the ^{13}C HMBC module. However, the techniques used in constructing this, including the implementation of the *zz*-filter, are equally applicable to a ^{15}N HMBC. For simplicity, the NOAH ^{15}N HMBC module uses a first-order LPJF (since $^1J_{\text{NH}}$ pairs are less common); this may be omitted if desired. To minimise the number of pulses, a simple magnitude-mode version of the HMBC is used (fig. 4.38). The implementation of this module within supersequences is discussed in greater detail within the context of *generalised supersequences*, in REF.

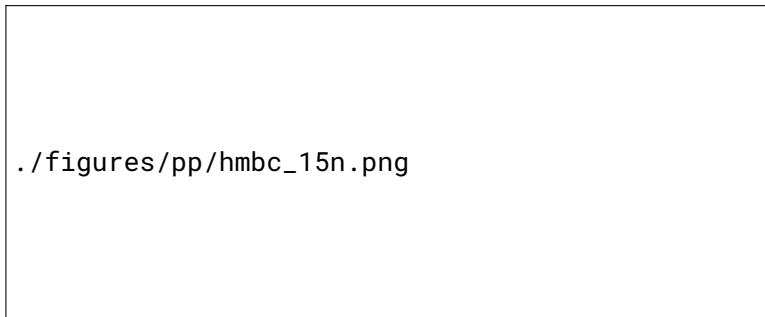


fig:noah_15n_hmhc

Figure 4.38: NOAH ^{15}N HMBC module. The *zz*-filter can be implemented as necessary in the same way as for the ^{13}C HMBC module (a final ^1H 180° pulse may also be required, but is not shown here). Delays are set as follows: $\Delta_N = 1/(2 \cdot ^1J_{\text{NH}})$; $\Delta_{\text{LR},N} = 1/(2 \cdot ^nJ_{\text{NH}})$. Phase cycling is performed using $\phi_1 = \phi_{\text{rec}} = (x, -x)$ and $\phi_2 = (x, x, -x, -x)$. Gradient amplitudes are $(g_1, g_2, g_3, g_4) = (5\%, 70\%, 30\%, 50.1\%)$.

4.3.8 ADEQUATE

The final ‘new’ module to be described in this section is the 1,1-ADEQUATE experiment,⁷⁸ which provides two-bond C–H correlations through a combination of $^1J_{\text{CH}}$ and $^1J_{\text{CC}}$. Because of the requirement for two adjacent ^{13}C nuclei, the sensitivity of the ADEQUATE experiment is considerably lower than almost all other NOAH modules considered so far. As established in § 4.1.1, sensitivity imbalances between different modules lead to less ideal supersequences with smaller effective time savings $\rho_{t,\text{eff}}$.

In this section, I only consider the design of the NOAH ADEQUATE module itself, which turns out to be extraordinarily simple. The ^1H pulses in the original ADEQUATE experiment mirror that of the CRK seHSQC almost perfectly, save for one extra 180° pulse. In particular, the INEPT excitation and the PEP transfer element are the same. Thus, the ADEQUATE module can be modified to preserve $^1\text{H}^{13}\text{C}$ magnetisation using exactly the same strategy in seHSQC2: namely, replacing the first ^1H excitation pulse with the ZIP element (fig. 4.39). In order to compensate for the extra 180° pulse (which arises in the constant-time period), the pulse phases in the ZIP element are modified accordingly: thus, the third ^1H 90° pulse has a phase of y (instead of $-y$ in

seHSQC2).

```
./figures/pp/adequate.png
```

`fig:adequate_noah`

Figure 4.39: NOAH 1,1-ADEQUATE module. The grey filled bar is a 120° pulse, optimised for ^{13}C double-quantum to single-quantum coherence transfer.⁷⁹ Delays are set as follows: $\Delta = 1/(4 \cdot {}^1\text{J}_{\text{CH}})$, $\Delta' = 1/(8 \cdot {}^1\text{J}_{\text{CH}})$, and $\tau = 1/(4 \cdot {}^1\text{J}_{\text{CC}})$. Phase cycling is performed using $\phi_1 = (x, -x)$, $\phi_2 = (x, x, -x, -x)$, $\phi_3 = (y, y, -y, -y)$, $\phi_4 = (x, x, x, x, -x, -x, -x, -x)$, and $\phi_{\text{rec}} = (x, -x, -x, x, -x, x, x, -x)$. Gradient amplitudes are $(g_1, g_2, g_3) = (78.5\%, 77.6\%, \mp 59\%)$. Echo-antiecho selection is achieved by inverting the sign of g_3 as well as the pulse phase ϕ_3 .

An example of a NOAH-2 AB supersequence is shown in fig. 4.40. This is not the best setting to use the ADEQUATE module in, because the ^{13}C HMBC module has roughly 100 times the sensitivity of the ADEQUATE; however, it demonstrates that the ADEQUATE module does indeed work as intended. More constructive uses of the ADEQUATE module are deferred to [Section REF](#): in there, I discuss combinations with the (lower-sensitivity) ^{15}N HMBC module, and show how other modules may be added to form so-called ‘generalised’ supersequences.

```
./figures/noah/ab.png
```

`fig:noah_ab`

Figure 4.40: Spectra from a NOAH-2 AB supersequence. (a) 1,1-ADEQUATE module. (a) HMBC module. *Data code:* 7T-211231.

4.4 Solvent suppression in NOAH

`sec:noah__solvsupp`

Solvent suppression is one of the most important aspects of modern NMR, and has been incorporated into a large number of experiments. It is not a particularly difficult task to implement

some basic solvent suppression techniques in NOAH supersequences, as will be described in this section. All suppression techniques shown here can be enabled via acquisition flags (the TopSpin ZGOPTNS parameter); this means that there is no need to switch pulse sequences.

4.4.1 Presaturation

:noah_solvsupp_presat

A simple first step is to implement simple presaturation of the solvent resonance during the recovery delay of the sequence. This often provides excellent suppression in the first module in a supersequence, and to a lesser extent, later modules (as the solvent magnetisation recovers during acquisition periods).

Presaturation is also included during long delays, in particular mixing times in NOESY experiments. The use of presaturation during the HMBC J-evolution delay was also tested, but was found to be generally unnecessary, especially since the HMBC module is typically placed first in supersequences.

4.4.2 Intrinsic suppression

:noah_solvsupp_intrinsic

A number of NOAH modules in fact come with an *intrinsic* form of solvent suppression, in that they return $^1\text{H}^{1\text{X}}$ magnetisation (including that of solvents) to the $+z$ axis at the end of the module. This is generally true of many of the HSQC-based NOAH modules, which seek to only sample $^1\text{H}^X$ magnetisation pool. Thus, these modules already have far better solvent suppression properties compared to standard library sequences: there is no need for further modification.

For these sequences, it is interesting to go one step further and to see how the solvent suppression varies with transmitter offset. This is less relevant if the solvent is just water (in which case it suffices to put the water peak on resonance), but may be important for samples in protonated solvents or mixtures thereof. I tested this by recording the first increment of the NOAH HSQC and seHSQC modules on an aqueous sample of sucrose, and integrating the water peak. The results (fig. 4.41) indicate that the HSQC module provides the best suppression across a range of offsets; of the two seHSQC modules, seHSQC1 provides more uniform suppression, whereas seHSQC2 has some regions ('spikes') where suppression is poorer than average.

It is not surprising that the more complicated seHSQC sequences have slightly poorer suppression: there are substantially more pulses involved which lead to cumulative errors, especially for peaks which are further away from resonance. However, it is not entirely clear why seHSQC2 has such large spikes. Some further investigation showed that the poorer suppression partly arises from imperfect cancellation of the water signal by phase cycling. This is, however, not a full explanation; it merely replaces one mystery with another. Nevertheless, it should be reiterated



Figure 4.41: Residual water peaks observed in the first increment of NOAH HSQC and seHSQC modules, as a function of transmitter offset. The y -axes of the three subplots are the same. (a) HSQC module. (b) seHSQC1 module. (c) seHSQC2 module. Data code: 7S-201014.

that all of these sequences still provide excellent water suppression. The standard library seHSQC sequences, for example, do not even come close to this level of suppression (although these do dephase water magnetisation prior to acquisition, radiation damping still leads to a very large signal being detected).

4.4.3 Excitation sculpting

Since solvent suppression in HMBC and HSQC-type modules can be adequately accomplished through presaturation and intrinsic suppression, it remains to consider the suppression in homonuclear modules. These modules invariably occur at the end of supersequences, and thus almost any form of solvent suppression can be used: there is no need to consider how magnetisation needs to be preserved for other modules.

In practice, I chose to implement excitation sculpting⁸⁰ (ES) just prior to acquisition: the refocusing element chosen was a combination of a selective 180° sinc pulse and a hard 180° pulse.

This worked perfectly for almost all of the homonuclear modules used in NOAH supersequences, with only a few adjustments needed, e.g. for the PSYCHE modules to ensure that chemical shifts and J-couplings evolved for the correct amount of time.

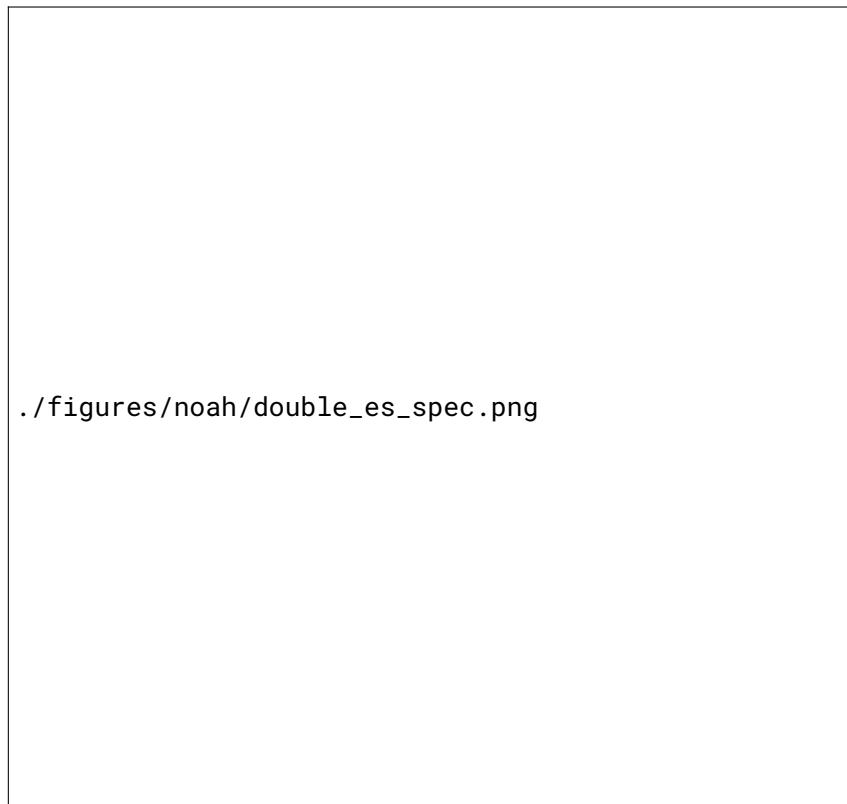
The case of the ‘double’ COSY + X homonuclear modules, however, proved to be more intricate. This is illustrated here using X = TOCSY, but the considerations below are equally applicable to X = NOESY or ROESY. It is tempting to simply place ES blocks prior to both FIDs: in the COSY module, this dephases transverse solvent magnetisation as desired, and (in principle) should leave longitudinal magnetisation untouched, meaning that the TOCSY module which later consumes this should be unaffected.

However, this is not totally true. *Inside* the bandwidth of the selective 180° pulse, any longitudinal magnetisation experiences an 720° rotation; and *outside* of the bandwidth, it experiences a 360° rotation. However, *between* these two extremes, there is a crossover point where longitudinal magnetisation is sent into the transverse plane and subsequently dephased by gradients. These lead to nulls in fig. 4.42b, and peaks which fall within this range will be lost in the TOCSY module. This is visible in the spectra of fig. 4.43: although the water suppression obtained using this ‘double ES’ scheme is better, several peaks in the TOCSY spectrum have disappeared, because they fall precisely into these nulls. It proves better to omit the ES in the COSY module: an adequate degree of water suppression in the COSY can still be attained thanks to the use of presaturation.

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./figures/noah/double_es_sim.png
```

Figure 4.42: Simulations showing the proportion of magnetisation retained by the excitation sculpting block (using a 2 ms sinc pulse). (a) Retention of transverse magnetisation: this was obtained by simulating the spectrum of a 90°–ES–detect pulse sequence. (b) Retention of longitudinal magnetisation: this was obtained using an ES–90°–detect pulse sequence.

Of course, it may well be that *no* peaks fall within this null, and thus the better solvent suppression may be obtained at no cost. This was the case when the experiments in fig. 4.43 were reacquired on a 700 MHz spectrometer; or when the sinc pulse was lengthened to 4 ms. However, such fortuitousness cannot always be relied on.



./figures/noah/double_es_spec.png

g : double_deshapes2spec

Figure 4.43: (a)–(b) COSY and TOCSY spectra obtained from a NOAH-3 SCT supersequence where excitation sculpting was placed before both the COSY and TOCSY FIDs. Although the water suppression is better, some peaks in the TOCSY module are lost. (c)–(d) The same, except that excitation sculpting was applied only in the TOCSY module. For all spectra shown here, presaturation of the water resonance was applied during the recovery delay. Data code: 4S-211105.

4.5 Parallel and generalised NOAH supersequences

sec:noah_parallel

Blah.

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