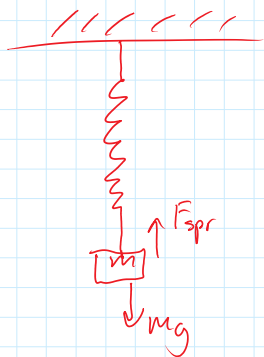


Effect of Constant External Force on a Harmonic Oscillator

Consider a mass on a spring in the vertical position



Force on mass: $F = -k(X - X_e) + mg$

where X_e is the equil. pos'n of the unstretched spring, but gravity will make a new equil. position.

$0 = -K(X_e' - X_e) + mg$ where X_e' is the new equil. position

$\rightarrow X_e' = X_e + \frac{mg}{K}$

Now define displacement from this new equil. pos'n as $x = X - X_e' = X - X_e - \frac{mg}{K}$

$\rightarrow x + \frac{mg}{K} = (X - X_e)$

$\therefore F = -K\left(x + \frac{mg}{K}\right) + mg = -Kx - mg + mg = -Kx$

$\therefore \underline{m\ddot{x} = -Kx}$ will still be e.o.m.

In general, any constant external force applied to a harmonic oscillator just shifts the equilibrium position

Ex: A light spring supports a block of mass m in a vertical position. When in equilibrium, the spring is stretched by an amount D_1 over its unstretched length. If the block is stretched a distance D_2 from the equilibrium position and released at $t=0$, find (a) the resulting motion $x(t)$, (b) the velocity of the block when it passes upwards through the equil. pos'n (c) the acc'n of the block at the top of its oscillatory motion.

First, consider the equil. pos'n: $F = 0 = -KD_1 + mg$ ($x > 0$ down)
 $\therefore K = \frac{mg}{D_1}$

\therefore Ang. freq. of oscillation $\omega_0 = \sqrt{\frac{K}{m}} = \sqrt{\frac{g}{D_1}}$

a) Motion is $x = a \cos(\omega_0 t - \theta)$

@ $t = 0$, $x = D_2 = a \cos(-\theta)$, $\dot{x} = 0 = -a \sin(-\theta)$
 $\rightarrow \theta = 0$

$\therefore a = D_2$

so, $x(t) = D_2 \cos \omega_0 t = D_2 \cos\left(\sqrt{\frac{g}{D_1}} t\right)$ indep. of mass

b) The velocity is $\dot{x}(t) = -D_2 \sqrt{\frac{g}{D_1}} \sin\left(\sqrt{\frac{g}{D_1}} t\right)$

When block is back in equil. position, $x = 0 = D_2 \cos\sqrt{\frac{g}{D_1}} t$

so $\sqrt{\frac{g}{D_1}} t = \frac{\pi}{2}$, so $\dot{x}(\text{equil. pos'n}) = -D_2 \sqrt{\frac{g}{D_1}}$

c) The acceleration is $\ddot{x}(t) = -D_2 \left(\frac{g}{D_1}\right) \cos\left(\sqrt{\frac{g}{D_1}} t\right)$

At the top of the motion, $x = -D_2 = D_2 \cos\left(\sqrt{\frac{g}{D_1}} t\right)$

so $\sqrt{\frac{g}{D_1}} t = \pi$

$\therefore \ddot{x}(\text{top of oscillation}) = \frac{D_2 g}{D_1}$

If $D_1 = D_2$, $\ddot{x} = g$; block is momentarily in free-fall w/ the spring exerting no force.

Ex: Simple Pendulum Revisited

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We've seen the restoring force $F = -mg \sin \theta$

$\therefore m\ddot{s} = -mg \sin \theta$



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For small θ , $s \approx l\theta$ & $\sin \theta \approx \theta$

so e.o.m. is $ml\ddot{\theta} = -mg\theta$

or $\ddot{\theta} + \frac{g}{l}\theta = 0$, same sho ock w/ $\omega_0 = \sqrt{\frac{g}{l}}$

as long as θ is relatively small, the period of oscillation is $T = \frac{2\pi}{\omega_0} = 2\pi\sqrt{\frac{l}{g}}$

Calculate The Time Averaged Kinetic, Potential & Total Energies of a Harmonic Oscillator

$$\langle T \rangle = \frac{1}{T} \int_0^T T(t) dt = \frac{1}{T} \int_0^T \frac{1}{2} m \dot{x}^2 dt$$

$$x = A \cos(\omega_0 t - \theta)$$

$$\dot{x} = -A\omega_0 \sin(\omega_0 t - \theta)$$

$$\text{set } \theta = 0, \quad \langle T \rangle = \frac{1}{T} \left[A^2 \frac{1}{2} m \omega_0^2 \int_0^T \sin^2 \omega_0 t dt \right]$$

$$\text{let } u = \omega_0 t = \left(\frac{2\pi}{T} \right) t \rightarrow du = \left(\frac{2\pi}{T} \right) dt$$

$$\rightarrow \langle T \rangle = \frac{1}{2\pi} \left[\frac{1}{2} m \omega_0^2 A^2 \int_0^{2\pi} \sin^2 u du \right]$$

$$\text{Note that } \frac{1}{2\pi} \int_0^{2\pi} (\sin^2 u + \cos^2 u) du = \frac{1}{2\pi} \int_0^{2\pi} du = 1$$

so $\frac{1}{2\pi} \int \sin^2 u du$ & $\frac{1}{2\pi} \int \cos^2 u du$ are both $\frac{1}{2}$

$$\therefore \boxed{\langle T \rangle = \frac{1}{4} m \omega_0^2 A^2}$$

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$V = \frac{1}{2} K x^2$ so please show that $\langle V \rangle = \frac{1}{4} K A^2 = \frac{1}{4} m \omega_0^2 A^2 = \langle T \rangle$

$$\therefore \langle E \rangle = \langle T \rangle + \langle V \rangle = \frac{1}{2} m \omega_0^2 A^2 = \frac{1}{2} K A^2 = E$$

Time avg. T & V are equal. Avg. E is equal to instantaneous E