

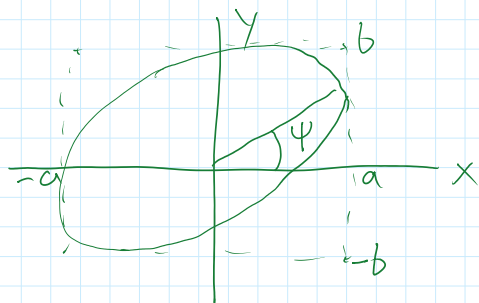
$$\therefore \frac{x^2}{a^2} + \frac{y^2}{b^2} - \frac{2xy \cos \Delta}{ab} = \sin^2 \Delta$$

Recall the quadratic equation $ax^2 + bxy + cy^2 + dx + ey + f = 0$ represents a conic section. See Appendix B for a review of conics.

If the discriminant ($b^2 - 4ac$) is negative, then the shape is an ellipse.

$$\text{In this case } b^2 - 4ac = \frac{4 \cos^2 \Delta}{a^2 b^2} - 4 = -\frac{4}{a^2 b^2} (1 - \cos^2 \Delta) = -\frac{4 \sin^2 \Delta}{a^2 b^2}$$

which is always negative, so the path is an ellipse



$$\text{If } \Delta = \frac{\pi}{2}, \text{ then } \frac{x^2}{a^2} + \frac{y^2}{b^2} = 1$$

which is an ellipse coincident w/ the axes
semi-major axis a ; semi-minor axis b

$$\text{If } \Delta = 0 \text{ or } \pi, \text{ then } \frac{x^2}{a^2} \pm \frac{2xy}{ab} + \frac{y^2}{b^2} = 0$$

$$\left(\frac{x}{a} \pm \frac{y}{b}\right)^2 = 0 \Rightarrow y = \pm \frac{b}{a} x \quad \begin{matrix} (+ \text{ sign if } \Delta = 0) \\ (- \text{ sign if } \Delta = \pi) \end{matrix}$$

The path is a straight line

$$\text{See §4.1 in text to see that } \tan 2\psi = \frac{2ab \cos \Delta}{a^2 - b^2}$$

In 3D, the path is also an ellipse, but in a plane off the xy plane.

... is an ellipse

Nonisotropic Oscillator

If the magnitudes of the components of the restoring force depend on the direction of the displacement, this is a nonisotropic oscillator.

For a suitable choice of axes in 2D

$$m\ddot{x} = -K_1 x$$

$$m\ddot{y} = -K_2 y$$

$$\text{So, } \omega_1 = \sqrt{\frac{K_1}{m}} \text{ ; } \omega_2 = \sqrt{\frac{K_2}{m}} \text{ and}$$

$$x = a \cos(\omega_1 t + \alpha) \text{ , } y = b \cos(\omega_2 t + \beta)$$

The oscillation lies entirely within a rectangle w/ sides $2a, 2b$ centered on the origin. If ω_1 & ω_2 are commensurate

ie if $\frac{\omega_1}{n_1} = \frac{\omega_2}{n_2}$ where n_1 & n_2 are integers then the path

is called a Lissajous figure, and will be closed. After a

time $\frac{2\pi n_1}{\omega_1} = \frac{2\pi n_2}{\omega_2}$ the particle will return to the initial pos'n.

However, if ω 's are not commensurate, the path is not closed & the particle will eventually 'fill the box'

Example: A particle of mass m moves in 2D under the following PE function $V(\vec{r}) = \frac{1}{2} K(x^2 + 4y^2)$. Find the resulting motion, given the initial conditions at $t=0$: $x=a, y=0, \dot{x}=0, \dot{y}=v_0$

$$\text{The force is } \vec{F} = -\vec{\nabla} V = -Kx\hat{i} - 4Ky\hat{j} = m\vec{\ddot{r}}$$

Separable, but not isotropic

$$\text{e.o.m's : } m\ddot{x} + Kx = 0 \text{ ; } m\ddot{y} + 4Ky = 0$$

The x-motion has angular freq. $\omega_x = \sqrt{\frac{k}{m}}$ while the y-motion has freq. $\omega_y = \sqrt{\frac{4k}{m}} = 2\omega_x$. The general solution is then

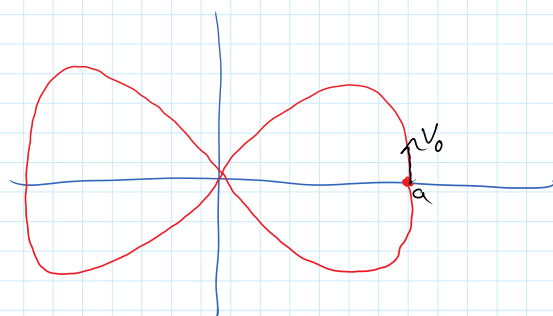
$$\begin{aligned} x &= a_1 \cos \omega_x t + b_1 \sin \omega_x t \\ y &= a_2 \cos 2\omega_x t + b_2 \sin 2\omega_x t \end{aligned} \quad @t=0 \quad \begin{aligned} x &= a = a_1 \\ y &= 0 = a_2 \end{aligned}$$

So, the velocities are

$$\begin{aligned} \dot{x} &= -a_1 \omega_x \sin \omega_x t + b_1 \omega_x \cos \omega_x t \\ \dot{y} &= -2a_2 \omega_x \sin 2\omega_x t + 2b_2 \omega_x \cos 2\omega_x t \end{aligned} \quad @t=0 \quad \begin{aligned} \dot{x} &= 0 = b_1 \omega_x \rightarrow b_1 = 0 \\ \dot{y} &= V_0 = 2b_2 \omega_x \rightarrow b_2 = \frac{V_0}{2\omega_x} \end{aligned}$$

\therefore Final eqns are $x = a \cos \omega_x t$; $y = \frac{V_0}{2\omega_x} \sin 2\omega_x t$

Path is a Lissajous figures



Central Force ; Gravitation

The harmonic oscillator is one example of a central, conservative force. More generally, in 3D a conservative force has

Components $F_r = -\frac{\partial V}{\partial r}$, $F_\theta = -\frac{1}{r} \frac{\partial V}{\partial \theta}$, $F_\phi = -\frac{1}{r \sin \theta} \frac{\partial V}{\partial \phi}$

\therefore The force will be central iff V is independent of θ & ϕ

Then $\vec{F} = -\hat{r} \frac{\partial V}{\partial r}$

Thus, the electric & gravitational forces are central, conservative forces.

While the force is only radial, the motion of a particle may be in any direction. In particular, it is interesting to consider the component of motion \perp to the radial direction.

The angular momentum of a particle located a distance \vec{r} from a given origin & moving w/ momentum $\vec{p} = m\vec{v}$ is defined as

$$\vec{J} = \vec{r} \times \vec{p}$$

The time derivative is: $\dot{\vec{J}} = \dot{\vec{r}} \times \vec{p} + \vec{r} \times \dot{\vec{p}}$

But $\dot{\vec{r}} \times \vec{p} = \dot{\vec{r}} \times (m\dot{\vec{r}}) = 0$, $\dot{\vec{J}} = \vec{r} \times \dot{\vec{p}} = \vec{r} \times \vec{F}$ from Newton's 2nd Law

The cross product $\vec{G} = \dot{\vec{J}}$ is the moment of the force, or torque. For central forces, $\vec{r} \parallel \vec{F}$, so $\dot{\vec{J}} = 0$ & \vec{J} is constant w/ time.