

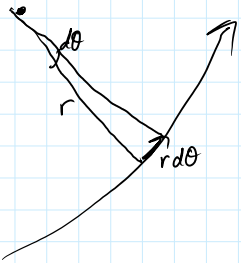
An important consequences of this result:

The vectors \vec{r} & \vec{v} define a plane in space. \vec{J} is always \perp to this plane (e.g. let \vec{r} & \vec{v} lie in the xy plane, then \vec{J} is in the \hat{k} direction). So, if \vec{J} is constant, it must be constant in both mag. and direction. \therefore Orientation of the plane of motion is fixed in space. So the motion of any body whose net force is central is simply motion in a plane.

The meaning of the constancy of $|\vec{J}|$ is easiest to see in polar coordinates

$$\vec{v} = \dot{r}\hat{r} + r\dot{\theta}\hat{\theta}$$

$$\begin{aligned}\text{So, } |\vec{J}| &= |\vec{r} \times m\vec{v}| = |r\hat{r} \times m(\dot{r}\hat{r} + r\dot{\theta}\hat{\theta})| \\ &= mr^2\dot{\theta}|\hat{r} \times \hat{\theta}| = mr^2\dot{\theta}\end{aligned}$$



So, for central forces, $|\vec{J}| = mr^2\dot{\theta} = \text{constant}$

Consider the area swept out by the radius vector when the particle moved by $d\theta$

$$dA = \frac{1}{2}r^2d\theta$$

So, the rate of area being swept is

$$\frac{dA}{dt} = \frac{1}{2}r^2\dot{\theta} = \frac{J}{2m} = \text{constant}$$

Thus, the rate of area being swept out is constant \Rightarrow Kepler's 2nd Law in the context of bodies moving in a grav. field. Although it applies to any motion subject to a central force (even non-cen. central force)

Orbit of a Particle in Central Force Field

Work in polar coordinates, as interested in the trace

of (r, θ) of the particle. $\therefore m\ddot{\vec{r}} = f(r)\hat{r}$

where $f(r)$ is the central force acting on the particle of mass m

From the 1st week: $\ddot{\vec{r}} = (\ddot{r} - r\dot{\theta}^2, 2\dot{r}\dot{\theta} + r\ddot{\theta})$

Thus, the 2 e.o.m are

$$m(\ddot{r} - r\dot{\theta}^2) = f(r)$$

$$m(2\dot{r}\dot{\theta} + r\ddot{\theta}) = 0$$

x by r $\rightarrow m \frac{d}{dt}(r^2\dot{\theta}) = 0$

$$\text{so, } r^2\dot{\theta} = \text{constant} = \frac{J}{m} \text{ where } J = |\vec{J}|$$

specific ang momentum

To find the path in (r, θ) plane (the orbit), re-write these ode's w/o time. Easiest to define a new variable $u = \frac{1}{r} \rightarrow r = \frac{1}{u}$

$$\text{then } \dot{r} = -\frac{1}{u^2} \dot{u} = -\frac{1}{u^2} \frac{du}{dt} = -\frac{1}{u^2} \frac{du}{d\theta} \frac{d\theta}{dt} = -\frac{1}{u^2} \frac{J}{mr^2} \frac{du}{d\theta} = -\frac{J}{m} \frac{du}{d\theta}$$

$$\rightarrow \ddot{r} = -\frac{J}{m} \frac{d}{dt} \left(\frac{du}{d\theta} \right) = -\frac{J}{m} \frac{d\theta}{dt} \left(\frac{d^2 u}{d\theta^2} \right) = -\frac{J^2}{m^2 r^2} \left(\frac{d^2 u}{d\theta^2} \right) = -\frac{J^2 u^2}{m^2} \frac{d^2 u}{d\theta^2}$$

Sub. this into the r-component of our e.o.m.

$$m \left(-\frac{J^2 u^2}{m^2} \frac{d^2 u}{d\theta^2} - \frac{1}{u} \frac{J^2 u^4}{m^2} \right) = f(u^{-1})$$

$$-\frac{J^2 u^2}{m} \frac{d^2 u}{d\theta^2} - \frac{J^2 u^3}{m} = f(u^{-1})$$

x by $\left(\frac{-m}{J^2 u^2} \right) \rightarrow \boxed{\frac{d^2 u}{d\theta^2} + u = -\frac{mf(u^{-1})}{J^2 u^2}}$

Diff. eqn of the orbit of a particle moving under a central force.

Ex: A particle in a central field moves in a spiral orbit

" $r = c\theta^4$. Determine the force function.

$$u = \frac{1}{r} = \frac{1}{c\theta^2}, \text{ so } \frac{du}{d\theta} = -\frac{2}{c}\theta^{-3}; \quad \frac{d^2u}{d\theta^2} = \frac{6}{c}\theta^{-4} = \frac{6}{c\theta^4} = 6cu^2$$

Sub. into o.d.e:

$$6cu^2 + u = -\frac{mf(u^{-1})}{J^2u^2} \rightarrow f(u^{-1}) = -\frac{J^2u^2}{m}6cu^2 - \frac{J^2u^3}{m} \\ = -\frac{J^2}{m}(6cu^4 + u^3)$$

$$\text{or } f(r) = -\frac{J^2}{m}\left(\frac{6c}{r^4} + \frac{1}{r^3}\right)$$