

## Damped Oscillator

If there is energy loss during the oscillation, there will be a frictional term in the e.o.m. Again, consider small displacements from equilibrium, so  $x$  &  $\dot{x}$  will be small and therefore any non-linear terms ( $x^2$ ,  $x\dot{x}$  or  $\dot{x}^2$ ) can be neglected. Thus, a damped harmonic oscillator has

$$F = -Kx - \lambda \dot{x} \quad \text{where } \lambda \text{ is a } ^{+ve} \text{ constant}$$

$$\therefore \text{e.o.m. is } m\ddot{x} + \lambda\dot{x} + Kx = 0 \quad (\text{LRC circuits, atomic physics})$$

Consider solutions of the form  $x = e^{pt}$  & substitute into e.o.m.

$$\rightarrow mp^2 + \lambda p + K = 0$$

solving this quadratic equation

$$p = -\gamma \pm \sqrt{\gamma^2 - \omega_0^2}$$

$$\text{where } \gamma = \frac{\lambda}{2m} \quad ; \quad \omega_0 = \sqrt{\frac{K}{m}} \quad (\text{the freq. of the undamped oscillator})$$

### 3 Possible Cases

#### 1) Overdamping ( $\gamma > \omega_0$ )

If  $\lambda$  is large enough then both roots are real & negative

$$p = -\gamma_{\pm}, \quad \gamma_{\pm} = \gamma \pm \sqrt{\gamma^2 - \omega_0^2}$$

$$\text{and } x = \frac{1}{2} A e^{-\gamma_+ t} + \frac{1}{2} B e^{-\gamma_- t} \quad \text{where } A \text{ & } B \text{ are constants ;}$$

the  $\frac{1}{2}$ 's are going to be useful later

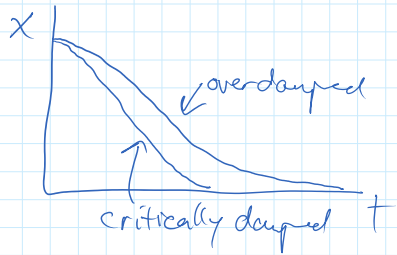
Displacement exponentially drops to zero, dominated by the  $K$  term

#### 2) Critical Damping ( $\gamma = \omega_0$ )

In this case  $p = -\gamma$  ; we only have  $x = Ae^{-\gamma t}$  so we need another solution in order to have 2 arbitrary constants.

Verify that  $x = te^{-\gamma t}$  is also a solution, so the general solution is  $x = (a + bt)e^{-\gamma t}$

Motion is again non-oscillatory ; displacement goes to zero faster than the overdamped case.



Critical damping is often ideal in measurement devices ; in suspension systems.

3) Underdamped. If  $\gamma$  is small, so that  $\gamma < \omega_0$ , the roots (for  $p$ ) are complex conjugates.

(See review of complex numbers)

$$\text{so, } p = -\gamma \pm \sqrt{\gamma^2 - \omega_0^2} \rightarrow p = -\gamma \pm i\omega \text{ where}$$

$$\omega = \sqrt{\omega_0^2 - \gamma^2}$$

so, the general solution is

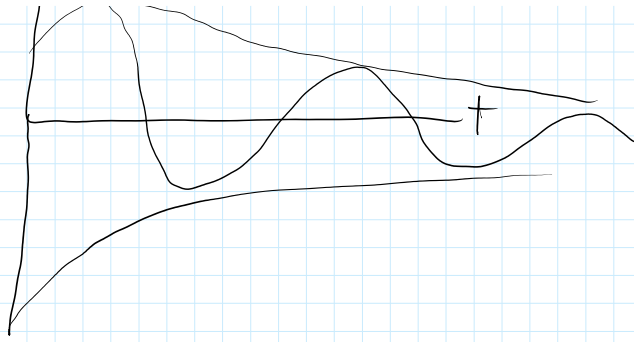
$$x = \frac{1}{2}Ae^{i\omega t - \gamma t} + \frac{1}{2}Be^{-i\omega t - \gamma t} \quad (A = ae^{-i\theta} ; B = ae^{i\theta})$$

Since these are complex conjugates

$$x = \text{Re}(Ae^{i\omega t - \gamma t}) = ae^{-\gamma t} \cos(\omega t - \theta)$$

These solutions are oscillatory w/ an exponentially decreasing amplitude ; ang. freq.  $\omega < \omega_0$





The time in which the amplitude is reduced by a factor of  $1/e$  is called the relaxation time.  $\gamma t_{rel} = 1 \rightarrow t_{rel} = \frac{1}{\gamma}$   
 or  $t_{rel} = \frac{2m}{\lambda}$

Another useful quantity is the quality factor  $Q$  of the oscillator.  $Q = \frac{m\omega_0}{\lambda} = \frac{\omega_0}{2\gamma}$ . If the damping is small then  $Q$  is large.

Why is  $Q$  useful?