

Homework 10

①

Problem A It's a classic problem. Let's use the superposition principle

no symmetries almost!

drill

superposition

$|\vec{J}| = I_0 / \pi R^2$

$|\vec{J}| = I_0 / (\pi R^2 - \pi a^2)$

less material to distribute I

It's ok if you kept $\vec{J} = \frac{I}{\pi R^2}$

$\vec{J}_1 = \vec{J}$

$\vec{J}_2 = -\vec{J}$

Ampère

$\vec{B}_1 = B_1(s_1) \hat{\phi}_1$

$\oint \vec{B}_1 \cdot d\vec{\ell} = \mu_0 I_{\text{through}}$

$$B_1(s_1) \cdot 2\pi s_1 = \mu_0 J \cdot \pi s_1^2$$

$$\rightarrow \boxed{B_1(s_1) = \frac{\mu_0 J s_1}{2} \hat{\phi}_1}$$



$$\vec{B}_2 = B_2(s_2) \hat{\phi}_2$$

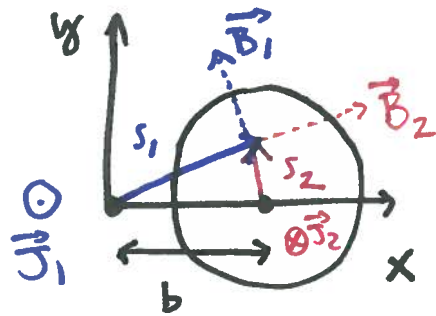
$$\oint \vec{B}_2 \cdot d\vec{\ell} = \mu_0 I_{\text{through}}$$

different coordinate systems!

$$\rightarrow \boxed{B_2(s_2) = -\frac{\mu_0 J s_2}{2} \hat{\phi}_2}$$

In the hole see that \vec{B}_1 and \vec{B}_2 have different directions ??

(2)



Look at that $\vec{B}_1 = \frac{\mu_0 J s_1}{2} \hat{\phi}_1 \equiv \frac{\mu_0}{2} \vec{J} \times \vec{s}_1 = \frac{\mu_0}{2} \vec{J} \times \vec{s}_1$

can you see it?

$$\vec{B}_2 = -\frac{\mu_0 J s_2}{2} \hat{\phi}_2 \equiv \frac{\mu_0}{2} \vec{J} \times \vec{s}_2 = -\frac{\mu_0}{2} \vec{J} \times \vec{s}_2$$

$$\text{So } \vec{B}_1 + \vec{B}_2 = \frac{\mu_0}{2} (\vec{J} \times \vec{s}_1 + \vec{J} \times \vec{s}_2) = \frac{\mu_0}{2} (\vec{J} \times (\vec{s}_1 - \vec{s}_2))$$

But from drawing: $\vec{b} = \vec{s}_1 - \vec{s}_2$ and therefore

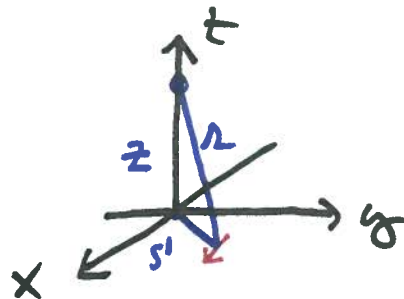
$$\vec{B} = \vec{B}_1 + \vec{B}_2 = \frac{\mu_0}{2} \vec{J} \times \vec{b}$$

The field is uniform in the hole!!!

Problem B

① The object has finite extent so $\vec{A}(\vec{r}) = \frac{\mu_0}{4\pi} \iint \frac{\vec{K}(\vec{r}') da'}{r}$ seems appropriate

here $\vec{K} = K_0 \hat{x}$



$da' = s' d\phi' ds'$
cylindrical

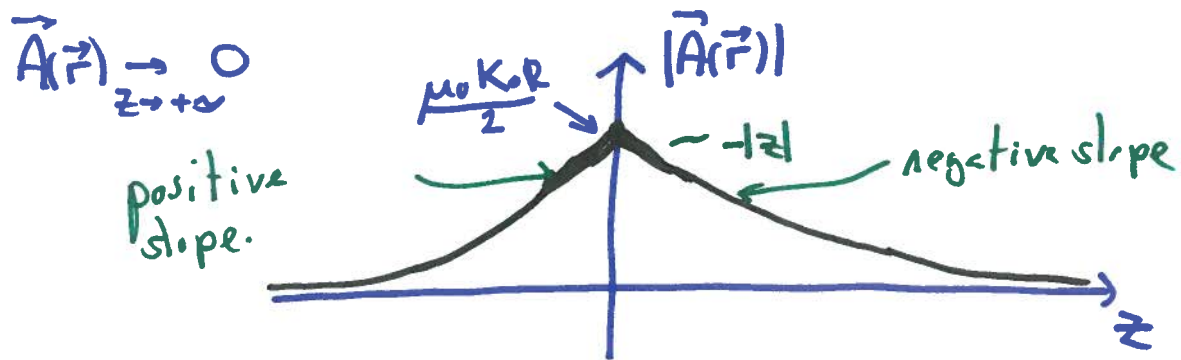
$r = \sqrt{s'^2 + z^2}$
see drawing

$$\text{then } \vec{A}(\vec{r}) = \frac{\mu_0}{4\pi} \int_0^R \int_0^{2\pi} \frac{K_0 s' ds' d\phi'}{\sqrt{s'^2 + z^2}} \hat{x} = \frac{K_0 \mu_0}{4\pi} \cdot 2\pi \cdot \hat{x} \cdot \int_0^R \frac{s' ds'}{\sqrt{s'^2 + z^2}} = \frac{K_0 \mu_0}{2} \left[\sqrt{s'^2 + z^2} \right]_0^R \hat{x}$$

$$\text{therefore } \vec{A}(\vec{r}) = \frac{K_0 \mu_0}{2} \left(\left[\sqrt{R^2 + z^2} - \sqrt{z^2} \right] \right) \hat{x} = \frac{\mu_0 K_0}{2} \hat{x} \cdot \left[\sqrt{R^2 + z^2} - |z| \right] \quad (3)$$

$\uparrow \Delta \sqrt{z^2} = |z|$

$$\vec{A}(\vec{r}) \xrightarrow{z \rightarrow 0} \frac{\mu_0 K_0}{2} \hat{x} \cdot \left[R \sqrt{1 + \frac{z^2}{R^2}} - |z| \right] \approx \frac{\mu_0 K_0}{2} \hat{x} \left[R + \frac{z^2}{2R} - |z| \right] \approx \frac{\mu_0 K_0}{2} \hat{x} [R - |z|]$$



This tells us there is a break in the slope of $|A(\vec{r})|$ close to $z=0$
 $\sim R - |z| \ll R \quad \forall z!$

② Theorem: $\frac{\partial \vec{A}^{\text{above}}}{\partial m} - \frac{\partial \vec{A}^{\text{below}}}{\partial m} = -\mu_0 \vec{K}$ at a current sheet.

here $z \equiv m$

$$\left. \frac{\partial \vec{A}^{\text{above}}}{\partial z} \right|_{z \geq 0} = \frac{1}{2} \mu_0 K_0 \frac{\partial}{\partial z} \left(\sqrt{R^2 + z^2} - \sqrt{z^2} \right) \hat{x} = \frac{1}{2} \mu_0 K_0 \left(\frac{z}{\sqrt{R^2 + z^2}} - \frac{z}{\sqrt{z^2}} \right) \hat{x} = \frac{1}{2} \mu_0 K_0 \left(\frac{z}{\sqrt{R^2 + z^2}} - 1 \right) \hat{x}$$

here $\sqrt{z^2} = |z| = z$
 as $z \geq 0, \frac{z}{|z|} = 1$

(4)

$$\left. \frac{\partial \vec{A}^{\text{below}}}{\partial z} \right|_{z \leq 0} = \frac{1}{2} \mu_0 K_0 \hat{x} \left(\frac{z}{\sqrt{R^2 + z^2}} - \frac{z}{\sqrt{z^2}} \right) = \frac{1}{2} \mu_0 K_0 \left(\frac{z}{\sqrt{R^2 + z^2}} + 1 \right) \hat{x}$$

here $\frac{z}{|z|} = -1$

$$\text{Therefore } \left. \frac{\partial \vec{A}}{\partial z} \right|_{z \neq 0} - \left. \frac{\partial \vec{A}}{\partial z} \right|_{z \neq 0} = -\frac{1}{2} \mu_0 K_0 \hat{x} - \frac{1}{2} \mu_0 K_0 \hat{x} = -\mu_0 K_0 \hat{x}!$$

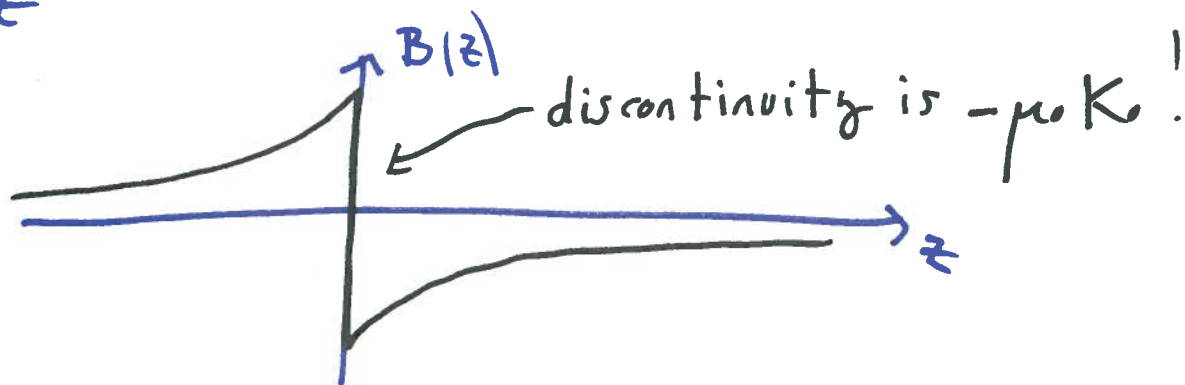
theorem is right

(3)

$$\text{By the way } \vec{B}(\vec{r}) = \vec{\nabla} \times \vec{A}(\vec{r}) = \vec{\nabla} \times (\vec{A}_x(z) \hat{x}) = \frac{\partial A_x(z)}{\partial z} \hat{y}$$

$$\text{So } \vec{B}^{\text{above}} = \frac{\partial A_x^{\text{above}}(z > 0)}{\partial z} \hat{y} = \frac{1}{2} \mu_0 K_0 \left(\frac{z}{\sqrt{R^2 + z^2}} - 1 \right) \hat{y}$$

$$\vec{B}^{\text{below}} = \frac{\partial A_x^{\text{below}}(z \leq 0)}{\partial z} \hat{y} = \frac{1}{2} \mu_0 K_0 \left(\frac{z}{\sqrt{R^2 + z^2}} + 1 \right) \hat{y}$$



④ $\lim_{R \rightarrow +\infty} \vec{A}(\vec{z}) = \frac{1}{z} \mu_0 K_0 (\sqrt{\infty^2 + z^2} - \sqrt{z^2}) \hat{x} \rightarrow +\infty$ ⑤
 i.e. \vec{A} diverges in the $R \rightarrow +\infty$ limit.

But $\left. \frac{\partial \vec{A}^{\text{above}}}{\partial z} \right|_{z=0} - \left. \frac{\partial \vec{A}^{\text{below}}}{\partial z} \right|_{z=0} = \left(-\frac{1}{z} \mu_0 K_0 - \frac{1}{z} \mu_0 K_0 \right) \hat{x} = -\mu_0 K_0 \hat{x}$

discontinuity remains finite and so does \vec{B} .

$\vec{B} \xrightarrow{R \rightarrow +\infty} + \frac{1}{z} \mu_0 K_0 \hat{y}$
 above
 below