

Tutorial 2 (Dirac notation and inner products)

The Dirac notation (also called bra-ket notation), which you have seen being used in the lecture, uses “kets”, such as $|\psi\rangle$, to represent a quantum state. For our purposes, a ket is always a complex (column) vector.¹ ψ is usually the actual vector itself, or can be an identifier or index for the quantum state, as for $|0\rangle$ and $|1\rangle$.

The corresponding “bra” $\langle\psi|$ is then the conjugate-transposed $|\psi\rangle$, i.e., a row vector with complex-conjugated entries of $|\psi\rangle$. A motivation for this notation is that “bras” are linear maps from quantum states to complex numbers via the inner product. Namely, given $\phi \in \mathbb{C}^n$:

$$\langle\phi| : \mathbb{C}^n \rightarrow \mathbb{C}, \quad |\psi\rangle \mapsto \langle\phi|\psi\rangle = \sum_{j=1}^n \phi_j^* \psi_j.$$

(a) Write down the matrix representation of the following expressions:

- $|0\rangle\langle 1|$
 - $|0\rangle\langle 0| + |1\rangle\langle 1|$
 - $|+\rangle\langle 0|$, with $|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$
- Handwritten notes:* $\langle 0|0\rangle = \langle 1|1\rangle = 1$, $\langle 0|1\rangle = \langle 1|0\rangle = 0$

(b) Express the Hadamard gate H using Dirac notation in the computational basis (i.e. $\{|0\rangle, |1\rangle\}$).

(c) Given the qubit state $|\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$, compute $H|\psi\rangle$ using only the bra-ket notation.

(d) For any $\psi, \phi \in \mathbb{C}^n$ and $A \in \mathbb{C}^{n \times n}$, verify that

$$\sum_{j=1}^n \phi_j^* (A\psi)_j = \sum_{j=1}^n (A^\dagger \phi)_j^* \psi_j \Leftrightarrow (\cdot) = \sum_{j=1}^n (A^\dagger \phi)_j^* \psi_j$$

$$\langle\phi|A\psi\rangle = \langle A^\dagger \phi|\psi\rangle, \quad \text{transpose on scalar} \quad (\cdot) = \sum_{j=1}^n \psi_j^* A_{ji}^* \phi_i$$

with $A^\dagger = (A^*)^T$ denoting the conjugate transpose (adjoint) of A .

$$\langle AB|AB\rangle = \langle A^\dagger A B|B\rangle$$

(e) Prove that unitary matrices are norm-preserving, i.e., $\|U\psi\| = \|\psi\|$ for all unitary $U \in \mathbb{C}^{n \times n}$ and $\psi \in \mathbb{C}^n$.

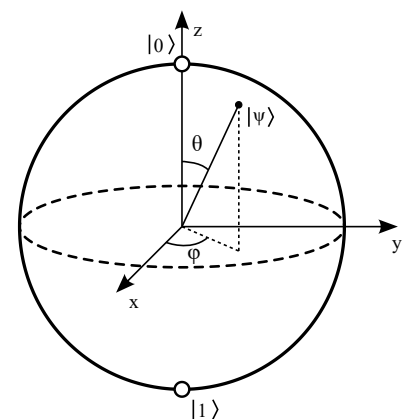
Hint: Use that $\|\psi\|^2 = \langle\psi|\psi\rangle$ and part (d). *Handwritten notes:* $\|U\psi\| = \sqrt{\langle U\psi|U\psi\rangle} = \sqrt{\sum (U\psi)_j^* (U\psi)_j} = \sqrt{\sum \psi_j^* \psi_j} = \sqrt{\langle\psi|\psi\rangle} = \|\psi\|$ // also by $\langle 0\psi|0\psi\rangle = \langle \psi^* 0\psi\rangle$

Exercise 2.1 (Bloch sphere and single qubit rotation gates)

Recall from the lecture that an arbitrary single qubit quantum state can be parametrized as

$$|\psi\rangle = e^{i\gamma} \left(\cos \frac{\theta}{2} |0\rangle + e^{i\varphi} \sin \frac{\theta}{2} |1\rangle \right)$$

where θ , φ and γ are real numbers, which can be chosen such that $0 \leq \theta \leq \pi$ and $0 \leq \varphi \leq 2\pi$. The angles θ and φ define the Bloch sphere representation of $|\psi\rangle$, as shown on the right.



https://commons.wikimedia.org/wiki/File:Bloch_sphere.svg

(a) Find the Bloch angles θ and φ of $|\psi\rangle = \frac{i}{2}|0\rangle - \frac{\sqrt{3}}{2}|1\rangle$, and the corresponding Bloch vector

$$\vec{r} = (\cos(\varphi) \sin(\theta), \sin(\varphi) \sin(\theta), \cos(\theta)).$$

For a real unit vector $\vec{v} \in \mathbb{R}^3$, the rotation by an angle ω about the \vec{v} axis is defined as

$$R_{\vec{v}}(\omega) = \exp(-i\omega \vec{v} \cdot \vec{\sigma}/2) = \cos(\omega/2)I - i \sin(\omega/2)(\vec{v} \cdot \vec{\sigma}),$$

where $\vec{\sigma} = (\sigma_1, \sigma_2, \sigma_3)$ is the Pauli vector. The rotations R_x , R_y , R_z about the standard axes correspond to the special cases $\vec{v} = (1, 0, 0)$, $\vec{v} = (0, 1, 0)$ and $\vec{v} = (0, 0, 1)$, respectively.

(b) Compute $R_x(\frac{2\pi}{3})|\psi\rangle$ for the state $|\psi\rangle$ defined in (a), and visualize this operation on the Bloch sphere.

$$\text{Hint: } \cos(\frac{\pi}{3}) = \frac{1}{2} \text{ and } \sin(\frac{\pi}{3}) = \frac{\sqrt{3}}{2}.$$

¹In general, quantum states can also be complex-valued functions (e.g., electronic orbitals of atoms), but these will not play a role in this course.

- (c) The *Z-Y decomposition* theorem states the following: given any unitary 2×2 matrix U , there exist real numbers $\alpha, \beta, \gamma, \delta \in \mathbb{R}$ such that

$$U = e^{i\alpha} R_z(\beta) R_y(\gamma) R_z(\delta).$$

Find the Z-Y decomposition of the Hadamard gate $H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$.

Hint: There exists a solution with $\beta = 0$.

Exercise 2.2 (Basic single qubit gates)

Imagine you are playing a game against an adversary. The game consists of multiple trials through which the adversary performs one of the following with equal probability:

1. They flip a coin and send you $|0\rangle$ or $|1\rangle$ depending on the outcome.

OR

2. They send you the state $\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$.

Your goal is to decide which of the two they performed, and you win if you can decide correctly for $\frac{3}{4}$ of the trials on average.

- (a) Before you make your guess (based on a quantum measurement on the qubit), you are allowed to perform **one** of the gates X , Y , Z or H . Compute the outputs you would obtain in each situation with each of these gates.
- (b) Which of the gates would allow you to win the game? Explain your strategy.