

1.1 应力与应变

本文以 CFRC 悬臂板为研究对象，将 MFC 粘贴于 CFRC 板的根部作为致动器和传感器，CFRC 悬臂板的结构如图 1 所示。图一 (a) 中的蓝色区域为 CFRC-MFC 悬臂层合板的固定约束区域，图一 (b) 为 CFRC 板的碳纤维材料铺层示意图。

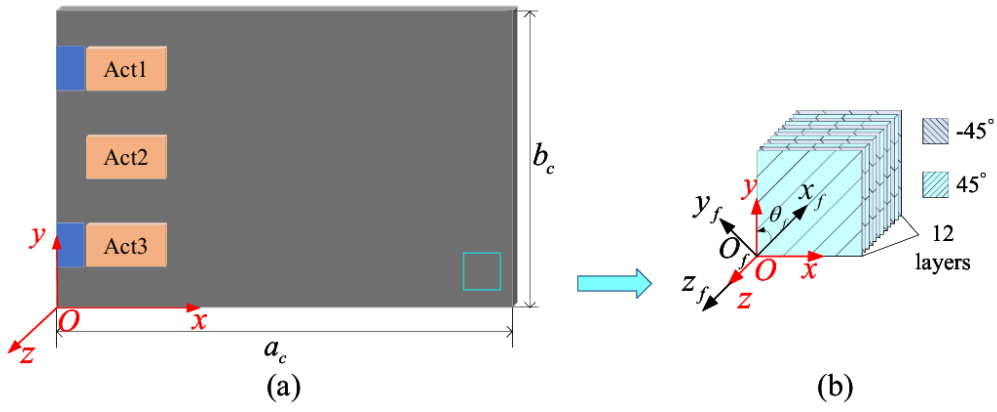


Fig1 CFRC-MFC 悬臂板示意图:(a) 层合板结构; (b) 铺层示意图

CFRC 板的单层碳纤维材料在纤维坐标系 $O_f - x_f y_f z_f$ 中描述的线性本构关系如下：

$$\boldsymbol{\sigma}_c^f = \begin{bmatrix} \sigma_{x_f x_f}^c \\ \sigma_{y_f y_f}^c \\ \tau_{x_f y_f}^c \end{bmatrix} = \begin{bmatrix} C_{x_f x_f} & C_{x_f y_f} & 0 \\ C_{y_f x_f} & C_{y_f y_f} & 0 \\ 0 & 0 & G_{x_f y_f}^c \end{bmatrix} \begin{bmatrix} \varepsilon_{x_f x_f}^c \\ \varepsilon_{y_f y_f}^c \\ \gamma_{x_f y_f}^c \end{bmatrix} = \mathbf{D}_c^f \boldsymbol{\varepsilon}_c^f \quad (1)$$

其中， $\boldsymbol{\sigma}_c^f$ ， $\boldsymbol{\varepsilon}_c^f$ 和 \mathbf{D}_c^f 分别为单层碳纤维材料的应力矢量，应变矢量和弹性系数矩阵。弹性系数表达式如下：

$$C_{x_f x_f}^c = \frac{E_1^c}{1 - \nu_{12}^c \nu_{21}^c}, C_{y_f y_f}^c = \frac{E_2^c}{1 - \nu_{12}^c \nu_{21}^c}, C_{y_f x_f}^c = \frac{\nu_{21}^c E_1^c}{1 - \nu_{12}^c \nu_{21}^c}, C_{x_f y_f}^c = \frac{\nu_{12}^c E_2^c}{1 - \nu_{12}^c \nu_{21}^c}, G_{x_f y_f}^c = G_{12}^c \quad (2)$$

其中， E_1^c 、 E_2^c 分别为碳纤维材料在纤维轴向 (x_f) 和纤维横向 (y_f) 的杨氏模量，

G_{12}^c 为剪切模量， ν_{12}^c 、 ν_{21}^c 为泊松比，1 代指 x_f ，2 代指 y_f ， $\nu_{21}^c = \nu_{12}^c E_2^c / E_1^c$ 。

而 MFC 的粘贴位置和结构如图 2 所示。图 2 (a) 和(b)分别为 CFRC-MFC 悬臂层合板的侧视图和俯视图，展示了 MFC 传感器和致动器的粘贴位置；图 2 (c) 和(d)分别为 d33 type MFC 和 d31 type MFC 的结构示意图。

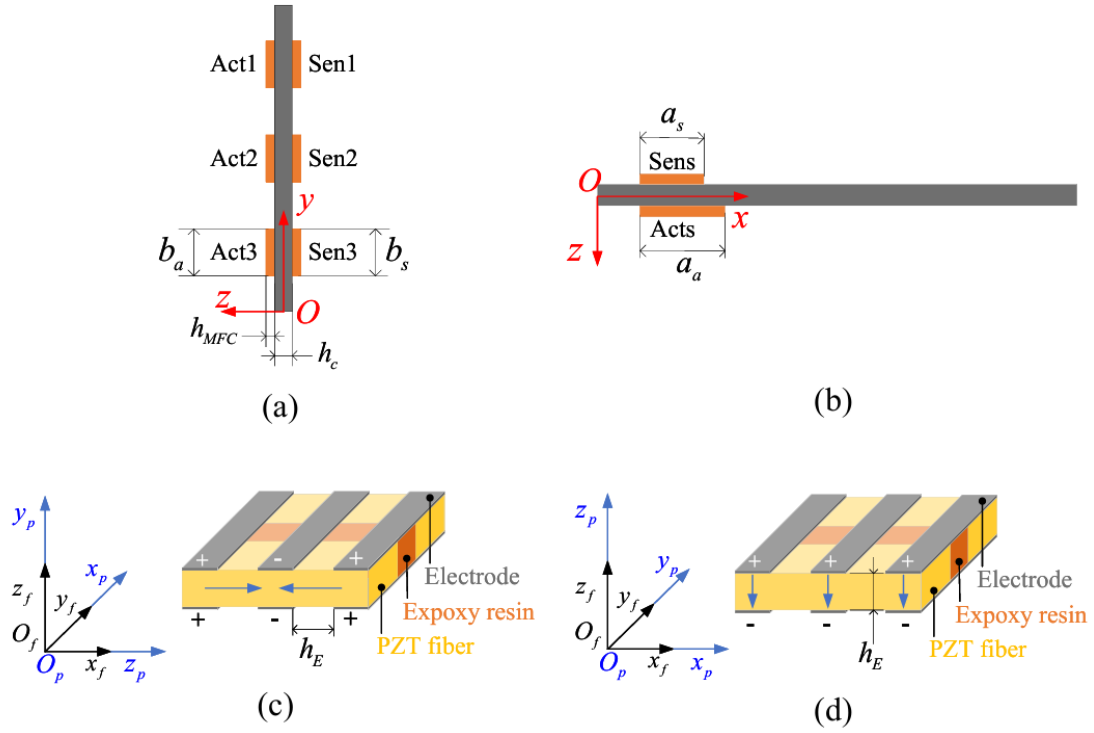


Fig2 MFC 的粘贴位置和结构示意图:(a) 侧视图; (b) 俯视图; (c) d33 type MFC;(d)d31 type MFC.

MFC 在纤维坐标系 $O_f - x_f y_f z_f$ 中描述的压电本构关系如下：

$$\begin{aligned} \boldsymbol{\sigma}_p^f &= \begin{bmatrix} \sigma_{x_f x_f}^p \\ \sigma_{y_f y_f}^p \\ \tau_{x_f y_f}^p \end{bmatrix} = \begin{bmatrix} P_{x_f x_f} & P_{x_f y_f} & 0 \\ P_{y_f x_f} & P_{y_f y_f} & 0 \\ 0 & 0 & G_{x_f y_f}^p \end{bmatrix} \begin{bmatrix} \varepsilon_{x_f x_f}^p \\ \varepsilon_{y_f y_f}^p \\ \gamma_{x_f y_f}^p \end{bmatrix} - \begin{bmatrix} e_{z_p x_f} \\ e_{z_p y_f} \\ 0 \end{bmatrix} E_{z_p} = \mathbf{D}_p^f \boldsymbol{\varepsilon}_p^f - \mathbf{e}^f E_{z_p} \\ D_{z_p} &= \begin{bmatrix} e_{z_p x_f} & e_{z_p y_f} & 0 \end{bmatrix} \begin{bmatrix} \varepsilon_{x_f x_f}^p \\ \varepsilon_{y_f y_f}^p \\ \gamma_{x_f y_f}^p \end{bmatrix} + \varsigma_{z_p z_p} E_{z_p} = (\mathbf{e}^f)^T \boldsymbol{\varepsilon}_p^f + \varsigma_{z_p z_p} E_{z_p} \end{aligned} \quad (3)$$

其中， $\boldsymbol{\sigma}_p^f$ 、 $\boldsymbol{\varepsilon}_p^f$ 、 \mathbf{D}_p^f 为 MFC 的应力矢量、应变矢量和弹性系数矩阵，下标

$p = a, s$ ，分别代表致动器和传感器。弹性系数表达式如下：

$$P_{x_f x_f}^p = \frac{E_1^p}{1 - \nu_{12}^p \nu_{21}^p}, P_{y_f y_f}^p = \frac{E_2^p}{1 - \nu_{12}^p \nu_{21}^p}, P_{y_f x_f}^p = \frac{\nu_{21}^p E_1^p}{1 - \nu_{12}^p \nu_{21}^p}, P_{x_f y_f}^p = \frac{\nu_{12}^p E_2^p}{1 - \nu_{12}^p \nu_{21}^p}, G_{x_f y_f}^p = G_{12}^p \quad (4)$$

其中, E_1^p 和 E_2^p 分别是 MFC 纤维主轴方向和纤维横向上的杨氏模量, G_{12}^p 是 MFC 的剪切模量, ν_{12}^p 、 ν_{21}^p 是 MFC 泊松比。而 \mathbf{e}^f 、 $\zeta_{z_p z_p}$ 、 E_{z_p} 、 D_{z_p} 分别为 MFC 的压电应力常数、介电常数、极化方向的电场强度和电位移矢量。值得注意的是, 由图 2 (c), (d) 可知, d33 型 MFC 和 d31 型 MFC 的极化坐标系并不一致, 因此本文采用极化坐标系与纤维坐标系相结合的表达方法, 获得了 MFC 压电本构方程的统一形式。此外, 压电应力常数 \mathbf{e}^f 各分量和电场强度 E_{z_p} 的表达式如下:

$$\begin{aligned} e_{z_p x_f} &= d_{z_p x_f} P_{x_f x_f} + d_{z_p y_f} P_{y_f x_f}, \\ e_{z_p y_f} &= d_{z_p x_f} P_{x_f y_f} + d_{z_p y_f} P_{y_f y_f}, \quad (5) \\ E_{z_p} &= \frac{\phi_{z_p}}{h_E} \end{aligned}$$

其中, $d_{z_p x_f}$ 和 $d_{z_p y_f}$ 为压电应变常数 \mathbf{d}^f 的分量, ϕ_{z_p} 为极化方向上的外加电压, h_E 为叉状电极正负极间的距离。smartstructure 公司发布的 mfc 的压电参数 d33 等的描述方式是基于极化坐标系 $O_p - x_p y_p z_p$ 的, 在本文中 d33 型 (P1 型) MFC, $d_{z_p x_f} = d_{33}$, 而对于 d31 型 (P2 型) MFC, $d_{z_p x_f} = d_{31}$ 。

根据复合材料力学, 在纤维坐标系 $O_f - x_f y_f z_f$ 内描述的应力矢量和应变矢量与在整体坐标系 $O - xyz$ 内描述的应力矢量和应变矢量间的转换关系如下:

$$\boldsymbol{\sigma} = \begin{bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \tau_{xy} \end{bmatrix} = \begin{bmatrix} c^2 & s^2 & -2sc \\ s^2 & c^2 & 2sc \\ sc & -sc & c^2 - s^2 \end{bmatrix} \bigg|_{\theta=\theta_f} \begin{bmatrix} \sigma_{x_f x_f} \\ \sigma_{y_f y_f} \\ \tau_{y_f y_f} \end{bmatrix} = \mathbf{T}_\sigma \boldsymbol{\sigma}^f \quad (6)$$

$$\boldsymbol{\varepsilon}^f = \begin{bmatrix} \varepsilon_{x_f x_f} \\ \varepsilon_{y_f y_f} \\ \gamma_{x_f y_f} \end{bmatrix} = \begin{bmatrix} c^2 & s^2 & sc \\ s^2 & c^2 & -sc \\ -2sc & 2sc & c^2 - s^2 \end{bmatrix} \bigg|_{\theta=\theta_f} \begin{bmatrix} \varepsilon_{xx} \\ \varepsilon_{yy} \\ \gamma_{xy} \end{bmatrix} = \mathbf{T}_\varepsilon \boldsymbol{\varepsilon} \quad (7)$$

其中, $c = \cos(\theta)$, $s = \sin(\theta)$, $\theta = \theta_f$ 为纤维坐标系与整体坐标系间的偏转角度。结合式(6),(7)和式(1)可得在整体坐标系 $O - xyz$ 中描述的 CFRC 板单层

碳纤维材料的线性本构关系和 MFC 压电贴片的压电本构关系如下：

$$\boldsymbol{\sigma}_c^k = \mathbf{T}_\sigma^k \mathbf{D}_c^f \mathbf{T}_\varepsilon^k \boldsymbol{\varepsilon}_c^k = \bar{\mathbf{D}}_c^k \boldsymbol{\varepsilon}_c^k \quad (8)$$

$$\begin{aligned} \boldsymbol{\sigma}_p &= \mathbf{T}_\sigma^p \mathbf{D}_p^f \mathbf{T}_\varepsilon^p \boldsymbol{\varepsilon}_p - \mathbf{T}_\sigma^p \mathbf{e}^f E_{z_p} = \bar{\mathbf{D}}_p \boldsymbol{\varepsilon}_p - \bar{\mathbf{e}} E_{z_p} \\ D_{z_p} &= \left(\left(\mathbf{T}_\varepsilon^p \right)^T \mathbf{e}^f \right)^T \boldsymbol{\varepsilon}_p + \zeta_{z_p z_p} E_{z_p} = (\bar{\mathbf{e}})^T \boldsymbol{\varepsilon}_p + \zeta_{z_p z_p} E_{z_p} \end{aligned} \quad (9)$$

其中， $\boldsymbol{\sigma}_c^k$ 和 $\boldsymbol{\varepsilon}_c^k$ 分别为第 k 层碳纤维材料在整体坐标系 $O-xyz$ 内的应力矢量和应变矢量， $\boldsymbol{\sigma}_p$ 和 $\boldsymbol{\varepsilon}_p$ 分别为 MFC 压电片在整体坐标系 $O-xyz$ 内的应力矢量和应变矢量。 $\mathbf{T}_\sigma^k = \mathbf{T}_\sigma|_{\theta=\theta_f^k}$ ， $\mathbf{T}_\varepsilon^k = \mathbf{T}_\varepsilon|_{\theta=\theta_f^k}$ ， $\mathbf{T}_\sigma^p = \mathbf{T}_\sigma|_{\theta=\theta_f^p}$ ， $\mathbf{T}_\varepsilon^p = \mathbf{T}_\varepsilon|_{\theta=\theta_f^p}$ ， θ_f^k 为第 k 层碳纤维材料层的铺层角度， θ_f^p 为 MFC 传感器 ($p=s$) 或致动器 ($p=a$) 的粘贴角度。

接下来，本文基于经典板理论，结合 ANCF 理论建立 C F R C-MFC 层合板在整体坐标系 $O-xyz$ 的应变—位移关系：

$$\boldsymbol{\varepsilon} = \begin{bmatrix} \frac{1}{2} \left((\mathbf{r}_{,x})^T \mathbf{r}_{,x} - 1 \right) \\ \frac{1}{2} \left((\mathbf{r}_{,y})^T \mathbf{r}_{,y} - 1 \right) \\ (\mathbf{r}_{,x})^T \mathbf{r}_{,y} \end{bmatrix} \quad (10)$$

\mathbf{r} 为板单元内任一点的位置矢量

$$\mathbf{r} = \mathbf{r} + z \mathbf{n}_0 \quad (11)$$

\mathbf{r} 和 \mathbf{n}_0 分别为中面上任意点的位置矢量和单位法向量。

$$\mathbf{n}_0 = \frac{\mathbf{n}}{\bar{n}}, \mathbf{n} = \mathbf{r}_{,x} \times \mathbf{r}_{,y}, \bar{n} = \sqrt{\mathbf{n}^T \mathbf{n}} \quad (12)$$

\mathbf{n} 为板中面上任意点的法向量，其具体表达式详见于附录 A， \bar{n} 为 \mathbf{n} 的模长。

将式(12)代入式(11)，忽略 z^2 的高阶项，可推得：

$$\boldsymbol{\varepsilon} = \begin{bmatrix} \frac{1}{2} \left((\mathbf{r}_{,x})^T \mathbf{r}_{,x} - 1 \right) \\ \frac{1}{2} \left((\mathbf{r}_{,y})^T \mathbf{r}_{,y} - 1 \right) \\ (\mathbf{r}_{,x})^T \mathbf{r}_{,y} \end{bmatrix} + z \begin{bmatrix} (\mathbf{r}_{,x})^T \frac{\partial \mathbf{n}_0}{\partial x} \\ (\mathbf{r}_{,y})^T \left(\frac{\partial \mathbf{n}_0}{\partial y} \right) \\ (\mathbf{r}_{,x})^T \frac{\partial \mathbf{n}_0}{\partial y} + \left(\frac{\partial \mathbf{n}_0}{\partial x} \right)^T \mathbf{r}_{,y} \end{bmatrix} \quad (13)$$

由于 $(\mathbf{r}_{,x})^T \mathbf{n}_0 = 0, (\mathbf{r}_{,y})^T \mathbf{n}_0 = 0$, 对等式两边分别对 x, y 求导可得:

$$\begin{aligned} \frac{\partial \left((\mathbf{r}_{,x})^T \mathbf{n}_0 \right)}{\partial x} &= (\mathbf{n}_0)^T \frac{\partial (\mathbf{r}_{,x})}{\partial x} + (\mathbf{r}_{,x})^T \frac{\partial (\mathbf{n}_0)}{\partial x}, \quad \frac{\partial \left((\mathbf{r}_{,x})^T \mathbf{n}_0 \right)}{\partial y} = (\mathbf{n}_0)^T \frac{\partial (\mathbf{r}_{,x})}{\partial y} + (\mathbf{r}_{,x})^T \frac{\partial (\mathbf{n}_0)}{\partial y} \\ \frac{\partial \left((\mathbf{r}_{,y})^T \mathbf{n}_0 \right)}{\partial x} &= (\mathbf{n}_0)^T \frac{\partial (\mathbf{r}_{,y})}{\partial x} + (\mathbf{r}_{,y})^T \frac{\partial (\mathbf{n}_0)}{\partial x}, \quad \frac{\partial \left((\mathbf{r}_{,y})^T \mathbf{n}_0 \right)}{\partial y} = (\mathbf{n}_0)^T \frac{\partial (\mathbf{r}_{,y})}{\partial y} + (\mathbf{r}_{,y})^T \frac{\partial (\mathbf{n}_0)}{\partial y} \end{aligned} \quad (14)$$

式中的等式左端都是 0, 于是推得:

$$\begin{aligned} (\mathbf{r}_{,x})^T \frac{\partial (\mathbf{n}_0)}{\partial x} &= -(\mathbf{n}_0)^T \mathbf{r}_{,xx}, \quad (\mathbf{r}_{,x})^T \frac{\partial (\mathbf{n}_0)}{\partial y} = -(\mathbf{n}_0)^T \mathbf{r}_{,xy} \\ (\mathbf{r}_{,y})^T \frac{\partial (\mathbf{n}_0)}{\partial x} &= -(\mathbf{n}_0)^T \mathbf{r}_{,yx}, \quad (\mathbf{r}_{,y})^T \frac{\partial (\mathbf{n}_0)}{\partial y} = -(\mathbf{n}_0)^T \mathbf{r}_{,yy} \end{aligned} \quad (15)$$

将式(15)代入式(13), 于是推得:

$$\boldsymbol{\varepsilon} = \begin{bmatrix} \frac{1}{2} \left((\mathbf{r}_{,x})^T \mathbf{r}_{,x} - 1 \right) \\ \frac{1}{2} \left((\mathbf{r}_{,y})^T \mathbf{r}_{,y} - 1 \right) \\ (\mathbf{r}_{,x})^T \mathbf{r}_{,y} \end{bmatrix} - z \begin{bmatrix} (\mathbf{r}_{,xx})^T \mathbf{n}_0 \\ (\mathbf{r}_{,yy})^T \mathbf{n}_0 \\ 2(\mathbf{r}_{,xy})^T \mathbf{n}_0 \end{bmatrix} = \boldsymbol{\varepsilon}_m - z \boldsymbol{\kappa} \quad (16)$$

其中, $\boldsymbol{\varepsilon}_m$ 和 $\boldsymbol{\kappa}$ 分别为板中面的膜应变和弯曲曲率。

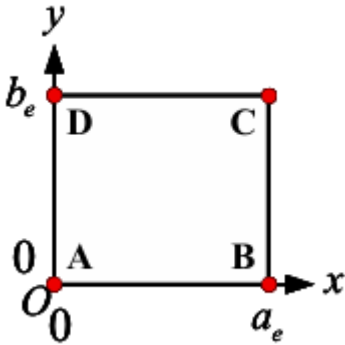


Fig3 单元节点示意图

采用有限元方法划分网格，单元内节点和局部坐标系如图 3 所示，其中 a_e 和 b_e 分别为单元长度和宽度。采用形函数对单元内中面上任意点的位置矢量 \mathbf{r} 进行离散可得：

$$\mathbf{r} = \mathbf{S}\mathbf{q}_e \quad (17)$$

其中， \mathbf{q}_e 为四节点单元位移矢量，表达式如下：

$$\mathbf{q}_e = \left[(\mathbf{q}_e^1)^T, (\mathbf{q}_e^2)^T, (\mathbf{q}_e^3)^T, (\mathbf{q}_e^4)^T \right]^T \quad (18)$$

$$\mathbf{q}_e^i = \left[\mathbf{r}^T, (\mathbf{r}_{,x})^T, (\mathbf{r}_{,y})^T \right]^T, i = 1, 2, 3, 4 \quad (19)$$

其中， \mathbf{q}_e^i 为单元内第 i 个节点的位移矢量，单元位移矢量与整体位移矢量间的关系为：

$$\mathbf{q}_e = \mathbf{B}_e \mathbf{q} \quad (20)$$

其中， \mathbf{B}_e 为单元 e 的布尔矩阵。

\mathbf{S} 为单元的形函数，表达式如下：

$$\mathbf{S} = \left[S_1 I_{3 \times 3} \quad S_2 I_{3 \times 3} \quad \mathbf{L} \quad S_{12} I_{3 \times 3} \right] \quad (21)$$

$$\begin{aligned}
S_1 &= -(\xi - 1)(\eta - 1)(2\eta^2 - \eta + 2\xi^2 - \xi - 1) \\
S_2 &= -a_e \xi (\xi - 1)^2 (\eta - 1) \\
S_3 &= -b_e \eta (\xi - 1)(\eta - 1)^2 \\
S_4 &= \xi (\eta - 1)(2\eta^2 - \eta + 2\xi^2 - 3\xi) \\
S_5 &= -a_e \xi^2 (\xi - 1)(\eta - 1) \\
S_6 &= b_e \xi \eta (\eta - 1)^2 \\
S_7 &= -\xi \eta (2\eta^2 - 3\eta + 2\xi^2 - 3\xi + 1) \\
S_8 &= a_e \xi^2 \eta (\xi - 1) \\
S_9 &= b_e \xi \eta^2 (\eta - 1) \\
S_{10} &= \eta (\xi - 1)(2\eta^2 - 3\eta + 2\xi^2 - \xi) \\
S_{11} &= a_e \xi \eta (\xi - 1)^2 \\
S_{12} &= -b_e \eta^2 (\xi - 1)(\eta - 1)
\end{aligned} \tag{22}$$

其中, a_e 和 b_e 分别代表未变形状态下板单元的长度和宽度; $I_{3 \times 3}$ 是一个 3×3

的单位矩阵。 $\xi = \frac{x}{a_e}, \eta = \frac{y}{b_e}$ 。 (x, y) 为未变形时中面上任意点在板单元坐标

系内的坐标。形函数顺序对应的单元节点顺序为 $A(0, 0), B(a_e, 0), C(a_e, b_e), D(0, b_e)$, 这决定了单元节点坐标在整体坐标中的对应位置。

将式(17)代入式(16)中, 获得面内膜应变 $\boldsymbol{\varepsilon}_m$ 和弯曲曲率 $\boldsymbol{\kappa}$ 的表达式如下:

$$\boldsymbol{\varepsilon}_m = \frac{1}{2} \begin{bmatrix} (\mathbf{r}_{,x})^T \mathbf{r}_{,x} - 1 \\ (\mathbf{r}_{,y})^T \mathbf{r}_{,y} - 1 \\ 2(\mathbf{r}_{,x})^T \mathbf{r}_{,y} \end{bmatrix} = \frac{1}{2} \begin{bmatrix} (\mathbf{S}_{,x} \mathbf{q}_e)^T \mathbf{S}_{,x} \mathbf{q}_e \\ (\mathbf{S}_{,y} \mathbf{q}_e)^T \mathbf{S}_{,y} \mathbf{q}_e \\ 2(\mathbf{S}_{,x} \mathbf{q}_e)^T \mathbf{S}_{,y} \mathbf{q}_e \end{bmatrix} - \frac{1}{2} \begin{bmatrix} 1 \\ 1 \\ 0 \end{bmatrix} \tag{23}$$

$$\mathbf{\kappa} = \begin{bmatrix} \frac{(\mathbf{r}_{,xx})^T \mathbf{n}}{\bar{n}} \\ \frac{(\mathbf{r}_{,yy})^T \mathbf{n}}{\bar{n}} \\ \frac{(\mathbf{r}_{,xy})^T \mathbf{n}}{\bar{n}} \end{bmatrix} = \begin{bmatrix} (\mathbf{S}_{,xx} \mathbf{q}_e)^T \\ (\mathbf{S}_{,yy} \mathbf{q}_e)^T \\ (2\mathbf{S}_{,xy} \mathbf{q}_e)^T \end{bmatrix} \frac{\mathbf{n}}{\bar{n}} \quad (24)$$

综上所述, 本小节建立了 CFRC-MFC 层合板在整体坐标系 $O-xyz$ 内的应力-应变关系和应变-位移关系, 接下来在下一小节将基于 Hamilton 原理建立 CFRC-MFC 层合板的动力学方程。

1.2 ANCF 动力学方程推导

首先, Hamilton 原理是常用的动力学建模方法, 其表达式如下:

$$\delta \int_{t_1}^{t_2} L dt = 0 \quad (25)$$

其中, L 为拉格朗日量, 其表达式如下:

$$L = (T_c + T_p) - (U_c + U_p) + (H_F - H_p) \quad (26)$$

其中, T_c 、 T_p 、 U_c 、 U_p 、 H_F 和 H_p 分别为 CFRC 层合板的动能、MFC 压电片的动能、CFRC 层合板的应变能、MFC 压电片的应变能、集中力所做功和 MFC 压电片外加电压所做的功, $p = a, s$ 。

对式 (25) 进行变分推导可得拉格朗日方程如下:

$$-\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\mathbf{q}}} \right) + \frac{\partial L}{\partial \mathbf{q}} = 0 \quad (27)$$

上式中的 \mathbf{q} 代表独立变量, 实际上包含 \mathbf{q} (整体位移矢量)、 ϕ_a (MFC 致动器电压) 和 ϕ_s (MFC 传感器电压), 分别获得如下表达式:

$$\begin{cases} -\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{\mathbf{q}}}\right) + \frac{\partial L}{\partial \mathbf{q}} = 0 \\ -\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{\phi}_a}\right) + \frac{\partial L}{\partial \phi_a} = 0 \\ -\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{\phi}_s}\right) + \frac{\partial L}{\partial \phi_s} = 0 \end{cases} \quad (28)$$

将式(26)代入式(27)可得

$$\begin{cases} -\frac{d}{dt}\left(\frac{\partial T_c}{\partial \dot{\mathbf{q}}} + \frac{\partial T_p}{\partial \dot{\mathbf{q}}}\right) + \left(-\frac{\partial U_c}{\partial \mathbf{q}} - \frac{\partial U_p}{\partial \mathbf{q}} + \frac{\partial H_F}{\partial \mathbf{q}} - \frac{\partial H_p}{\partial \mathbf{q}}\right) = 0 \\ -\frac{\partial U_a}{\partial \phi_a} - \frac{\partial H_a}{\partial \phi_a} = 0 \\ \frac{\partial U_s}{\partial \phi_s} = 0 \end{cases} \quad (29)$$

由式 (29) 可知推导动力学方程需要能量方程及其对独立变量的导数, CFRC 层合板的动能和应变能表达式如下:

$$T_c = \frac{1}{2} \iiint_{V_c} \rho_c \dot{\mathbf{u}}^T \dot{\mathbf{u}} dV = \frac{1}{2} (\dot{\mathbf{q}})^T \sum_e^{N_e} (\mathbf{B}_e)^T \iint_{S_e} I_{c1} (\mathbf{S})^T \mathbf{S} dx dy \mathbf{B}_e \dot{\mathbf{q}} = \frac{1}{2} (\dot{\mathbf{q}})^T \mathbf{M}_c \dot{\mathbf{q}} \quad (30)$$

$$\begin{aligned} U_c &= \frac{1}{2} \iiint_{V_c} (\boldsymbol{\varepsilon}_c)^T \boldsymbol{\sigma}_c dV = \frac{1}{2} \sum_e^{N_e} \iint_{S_e} \sum_{k=1}^{N_f} \int_{h_b^k}^{h_u^k} (\boldsymbol{\varepsilon}_c^k)^T \bar{\mathbf{D}}_c^k \boldsymbol{\varepsilon}_c^k dz dx dy \\ &= \frac{1}{2} \sum_e^{N_e} \iint_{S_e} ((\boldsymbol{\varepsilon}_m)^T \mathbf{D}_{c1} \boldsymbol{\varepsilon}_m + 2(\boldsymbol{\kappa})^T \mathbf{D}_{c2} \boldsymbol{\varepsilon}_m + (\boldsymbol{\kappa})^T \mathbf{D}_{c3} \boldsymbol{\kappa}) dx dy \end{aligned} \quad (31)$$

其中, N_e^c 是 CFRC 板的单元数量, N_f 为碳纤维材料层数, ρ_c 为复合材料密度, \mathbf{D}_{c1} 、 \mathbf{D}_{c2} 、 \mathbf{D}_{c3} 和 I_{c1} 等刚度和惯性系数矩阵的具体内容见于附录 A。k 表示碳纤维材料层的编号, h_u^k 、 h_b^k 分别为第 k 层碳纤维材料的上界和下界。将式(30),(31)代入式(29)的求导过程可以获得:

$$\frac{d}{dt}\left(\frac{\partial T_c}{\partial \dot{\mathbf{q}}}\right) = (\dot{\mathbf{q}})^T \mathbf{M}_c \quad (32)$$

$$\begin{aligned}
\frac{\partial U_c}{\partial \mathbf{q}} &= \frac{1}{2} \sum_e^{N_e^c} \iint_{S_e} \left(\frac{\partial (\boldsymbol{\varepsilon}_m)^T \mathbf{D}_{c1} \boldsymbol{\varepsilon}_m}{\partial \mathbf{q}_e} + 2 \frac{\partial (\boldsymbol{\kappa})^T \mathbf{D}_{c2} \boldsymbol{\varepsilon}_m}{\partial \mathbf{q}_e} + \frac{\partial (\boldsymbol{\kappa})^T \mathbf{D}_{c3} \boldsymbol{\kappa}}{\partial \mathbf{q}_e} \right) \frac{\partial \mathbf{q}_e}{\partial \mathbf{q}} dx dy \\
&= \sum_e^{N_e^c} \iint_{S_e} \left((\boldsymbol{\varepsilon}_m)^T \mathbf{D}_{c1} \frac{\partial \boldsymbol{\varepsilon}_m}{\partial \mathbf{q}_e} + (\boldsymbol{\kappa})^T \mathbf{D}_{c2} \frac{\partial \boldsymbol{\varepsilon}_m}{\partial \mathbf{q}_e} + (\boldsymbol{\varepsilon}_m)^T \mathbf{D}_{c2} \frac{\partial \boldsymbol{\kappa}}{\partial \mathbf{q}_e} + (\boldsymbol{\kappa})^T \mathbf{D}_{c3} \frac{\partial \boldsymbol{\kappa}}{\partial \mathbf{q}_e} \right) \mathbf{B}_e dx dy \\
&= \sum_e^{N_e^c} \iint_{S_e} \left((\boldsymbol{\varepsilon}_m)^T \mathbf{D}_{c1} \boldsymbol{\varepsilon}_{m,q_e} + (\boldsymbol{\kappa})^T \mathbf{D}_{c2} \boldsymbol{\varepsilon}_{m,q_e} + (\boldsymbol{\varepsilon}_m)^T \mathbf{D}_{c2} \boldsymbol{\kappa}_{,q_e} + (\boldsymbol{\kappa})^T \mathbf{D}_{c3} \boldsymbol{\kappa}_{,q_e} \right) \mathbf{B}_e dx dy \\
&= (\mathbf{F}_c^{ela})^T
\end{aligned} \tag{33}$$

其中, \mathbf{M}_c 为 CFRC 层合板的质量阵, \mathbf{F}_c^{ela} 为 CFRC 层合板的非线性弹性力矢量, $\boldsymbol{\varepsilon}_{m,q_e}$ 和 $\boldsymbol{\kappa}_{,q_e}$ 的表达式见于附录 A。

MFC 压电片的动能和应变能的表达式如下:

$$\begin{aligned}
T_p &= \frac{1}{2} \iiint_{V_p} \rho_p \mathbf{\dot{\boldsymbol{\Phi}}}^T \mathbf{\dot{\boldsymbol{\Phi}}} dV = \frac{1}{2} (\boldsymbol{\Phi})^T \sum_e^{N_e^p} (\mathbf{B}_e)^T \iint_{S_e} I_{p1} (\mathbf{S})^T \mathbf{S} dx dy \mathbf{B}_e \boldsymbol{\Phi} = \frac{1}{2} (\boldsymbol{\Phi})^T \mathbf{M}_p \boldsymbol{\Phi} \tag{34} \\
U_p &= \frac{1}{2} \iiint_{V_p} (\boldsymbol{\varepsilon}_p)^T \boldsymbol{\sigma}_p dV - \frac{1}{2} \iiint_{V_p} D_{z_p} E_{z_p} dV \\
&= \frac{1}{2} \sum_e^{N_e^p} \iint_{S_e} (\boldsymbol{\varepsilon}_m)^T \mathbf{D}_{p1} \boldsymbol{\varepsilon}_m + 2 (\boldsymbol{\kappa})^T \mathbf{D}_{p2} \boldsymbol{\varepsilon}_m + (\boldsymbol{\kappa})^T \mathbf{D}_{p3} \boldsymbol{\kappa} dx dy \\
&\quad - \sum_e^{N_e^p} \iint_{S_e} \phi_p (\mathbf{e}_{p1})^T \frac{1}{h_E} \boldsymbol{\varepsilon}_m + \phi_p (\mathbf{e}_{p2})^T \frac{1}{h_E} \boldsymbol{\kappa} dx dy \\
&\quad - \frac{1}{2} \sum_e^{N_e^p} \iint_{S_e} \zeta_{z_p z_p}^1 \left(\frac{\phi_p}{h_E} \right)^2 dx dy
\end{aligned} \tag{35}$$

其中, N_e^p 是单个 MFC 传感器 ($p = s$) 或致动器 ($p = a$) 的单元数量, ρ_p 为 MFC 的密度。 I_{p1} 、 \mathbf{D}_{p1} 、 \mathbf{D}_{p2} 、 \mathbf{D}_{p3} 、 \mathbf{e}_{p1} 、 \mathbf{e}_{p2} 、 $\zeta_{z_p z_p}^1$ 等参数矩阵可以通过在厚度方向上进行积分运算获得, 具体内容见于附录 A。将式(34),(35)代入式(29)的求导过程可以获得:

$$\frac{d}{dt} \left(\frac{\partial T_p}{\partial \boldsymbol{\Phi}} \right) = (\boldsymbol{\Phi})^T \mathbf{M}_p \tag{36}$$

$$\begin{aligned} \frac{\partial U_p}{\partial \mathbf{q}} &= \frac{1}{2} \sum_e^{N_e^p} \iint_{S_e} \left(\frac{\partial (\boldsymbol{\varepsilon}_m)^T \mathbf{D}_{p1} \boldsymbol{\varepsilon}_m}{\partial \mathbf{q}_e} + 2 \frac{\partial (\boldsymbol{\kappa})^T \mathbf{D}_{p2} \boldsymbol{\varepsilon}_m}{\partial \mathbf{q}_e} + \frac{\partial (\boldsymbol{\kappa})^T \mathbf{D}_{p3} \boldsymbol{\kappa}}{\partial \mathbf{q}_e} \right) \frac{\partial \mathbf{q}_e}{\partial \mathbf{q}} dx dy \\ &- \sum_e^{N_e^p} \iint_{S_e} \left(\phi_p (\mathbf{e}_{p1})^T \frac{1}{h_E} \frac{\partial \boldsymbol{\varepsilon}_m}{\partial \mathbf{q}_e} + \phi_p (\mathbf{e}_{p2})^T \frac{1}{h_E} \frac{\partial \boldsymbol{\kappa}}{\partial \mathbf{q}_e} \right) \frac{\partial \mathbf{q}_e}{\partial \mathbf{q}} dx dy \end{aligned} \quad (37)$$

$$\begin{aligned} &= \sum_e^{N_e^p} \iint_{S_e} \left((\boldsymbol{\varepsilon}_m)^T \mathbf{D}_{p1} \boldsymbol{\varepsilon}_{m,q_e} + (\boldsymbol{\kappa})^T \mathbf{D}_{p2} \boldsymbol{\varepsilon}_{m,q_e} + (\boldsymbol{\varepsilon}_m)^T \mathbf{D}_{p2} \boldsymbol{\kappa}_{,q_e} + (\boldsymbol{\kappa})^T \mathbf{D}_{p3} \boldsymbol{\kappa}_{,q_e} \right) \mathbf{B}_e dx dy \\ &- \sum_e^{N_e^p} \iint_{S_e} \left(\phi_p (\mathbf{e}_{p1})^T \frac{1}{h_E} \boldsymbol{\varepsilon}_{m,q_e} + \phi_p (\mathbf{e}_{p2})^T \frac{1}{h_E} \boldsymbol{\kappa}_{,q_e} \right) \mathbf{B}_e dx dy \end{aligned}$$

$$\begin{aligned} &= (\mathbf{F}_p^{ela})^T - \mathbf{K}_{cp} \phi_p \\ \frac{\partial U_p}{\partial \phi_p} &= - \sum_e^{N_e^p} \iint_{S_e} (\mathbf{e}_{p1})^T \frac{1}{h_E} \boldsymbol{\varepsilon}_m + (\mathbf{e}_{p2})^T \frac{1}{h_E} \boldsymbol{\kappa} dx dy \\ &- \sum_e^{N_e^p} \iint_{S_e} \zeta_{z_p z_p}^1 \left(\frac{1}{h_E} \right)^2 \phi_p dx dy \end{aligned} \quad (38)$$

$$= -K_{pc} - K_{pp} \phi_p$$

其中, \mathbf{M}_p 为 MFC 压电片的质量阵, \mathbf{F}_p^{ela} 为 MFC 压电片的非线性弹性力矢量, 而 \mathbf{K}_{cp} 和 K_{pc} 分别是 MFC 压电片的机电耦合系数矩阵和等效机电耦合系数, K_{pp} 是 MFC 压电片的等效介电系数。

集中力所做的功表达式如下:

$$H_F = \mathbf{F} \mathbf{r}(x_F, y_F, t) = \mathbf{F} \mathbf{S}_F \mathbf{q} \quad (39)$$

其中, $\mathbf{F} = [0 \quad 0 \quad F_{ext}]$ 为外部集中力, F_{ext} 为集中力幅值。

$\mathbf{S}_F = \mathbf{S}(x_F, y_F)$, (x_F, y_F) 为集中力施加位置所在单元内局部坐标系下的坐标。

等式两边都对 \mathbf{q} 求导获得:

$$\frac{\partial H_F}{\partial \mathbf{q}} = \mathbf{F} \mathbf{S}_F \quad (40)$$

在 MFC 致动器外部施加的电压所做的功的表达式如下:

$$H_p = \frac{1}{2} \left(\frac{h_E}{h_p} \right)^2 K_{aa} (\phi_a)^2 \quad (41)$$

等式两边都对 ϕ_a 求导获得:

$$\frac{\partial H_p}{\partial \phi_a} = \left(\frac{h_E}{h_p} \right)^2 K_{aa} \phi_a^c \quad (42)$$

在实际控制中，式 (42) 右边的 ϕ_a 是指前一时间步内保持的 $\phi_a(t - dt)$ ，即在电压 ϕ_a 变化的微时间段 dt 内，认为 $\frac{\partial H_p}{\partial \phi_a}$ 不变。或可以将式 (42) 右边的 ϕ_a 理解为是施加的控制电压 ϕ_a^c ，施加了之后就变成了节点上的独立变量 ϕ_a ，并在控制电压再次施加前随应变变化而产生变化 $d\phi_a$ 。

将式 (32)、(33)、(36)、(37)、(38)、(40)、(42) 代入式 (29) 推得如下非线性动力学方程组：

$$(\mathbf{M}_c + \mathbf{M}_a + \mathbf{M}_s) \ddot{\mathbf{q}} + \mathbf{F}_c^{ela} + \mathbf{F}_a^{ela} + \mathbf{F}_s^{ela} - (\mathbf{K}_{ca})^T \phi_a - (\mathbf{K}_{cs})^T \phi_s - (\mathbf{F} \mathbf{S}_F)^T = 0 \quad (43)$$

$$K_{ac} + K_{aa} \phi_a - \left(\frac{h_E}{h_p} \right)^2 K_{aa} \phi_a^c = 0 \quad (44)$$

$$K_{sc} + K_{ss} \phi_s = 0 \quad (45)$$

将式 (44)、(45) 代入式(43),并引入阻尼矩阵可得 MFC-CFRC 层合板的整体动力学方程：

$$\mathbf{M} \ddot{\mathbf{q}} + \mathbf{C} \dot{\mathbf{q}} + \mathbf{F}^{ela} = (\mathbf{F} \mathbf{S}_F)^T + \mathbf{K}_{con} \phi_a^c \quad (46)$$

式中， \mathbf{F}^{ela} 为 CFRC-MFC 层合板的非线性弹性力列矢量， \mathbf{K}_{con} 为 MFC 致动器控制力输入矩阵， \mathbf{C} 是采用瑞丽阻尼模型建立的阻尼矩阵，其表达式如下：

$$\mathbf{C} = c_0 \mathbf{M} + c_1 \frac{\partial \mathbf{F}^{ela}}{\partial \mathbf{q}} \bigg|_{q=q_0} \quad (47)$$

其中， c_0 和 c_1 是瑞丽阻尼参数，表达式如下：

$$c_0 = 2 \left(\frac{\xi_2}{\omega_2} - \frac{\xi_1}{\omega_1} \right) / \left(\frac{1}{(\omega_2)^2} - \frac{1}{(\omega_1)^2} \right), c_1 = 2(\xi_2 \omega_2 - \xi_1 \omega_1) / ((\omega_2)^2 - (\omega_1)^2) \quad (48)$$

其中, ω_1 、 ω_2 、 ξ_1 、 ξ_2 分别为 CFRC-MFC 层合板的前两阶自然频率及对应的模态阻尼比。将式 (45) 整理可得 MFC 传感器输出方程:

$$\phi_s = \mathbf{K}_{out} \mathbf{q} \quad (49)$$

其中, \mathbf{K}_{out} 为 MFC 传感器电压输出矩阵。此外, 式 (46) 和 (49) 中的各个矩阵详见附录 B。

CFRC 悬臂板的约束区域如图 1 所示, 边界约束方程如下:

$$\Phi(\mathbf{q}, t) = 0 \quad (50)$$

采用拉格朗日乘子 λ 将边界约束方程引入动力学方程可得:

$$\mathbf{M}\ddot{\mathbf{q}} + \mathbf{C}\dot{\mathbf{q}} + \mathbf{K}\mathbf{q} + (\Phi_{,q})^T \lambda = (\mathbf{F}\mathbf{S}_F)^T + \mathbf{K}_{con}\phi_a^c \quad (51)$$

$$\Phi(\mathbf{q}, t) = 0$$

其中, $\Phi_{,q} = \mathbf{G}$ 。

1.3 数值求解算法-基于加速度

$$\mathbf{M}\ddot{\mathbf{q}} + \mathbf{C}\dot{\mathbf{q}} + \mathbf{F}^{ela} + (\Phi_{,\ddot{\mathbf{q}}})^T \lambda = (\mathbf{F}\mathbf{S}_F)^T + \mathbf{K}_{con}\phi_a^c \quad (52)$$

$$\Phi(\mathbf{q}, t) = 0$$

对于基于加速度的求解, 令 $\Phi(\mathbf{q}, t) = \mathbf{G}\mathbf{q}$, $\Phi_{,\ddot{\mathbf{q}}} = \mathbf{G}$, 矩阵 \mathbf{G} 用于标记受约束的自由度。

采用 newmark-beta 法基于加速度求解

其中, \mathbf{q}_{i+1} 和 \mathbf{q}_{i+1} 为第 $i+1$ 个时间步的预测值, 他们的表达式如下:

$$\mathbf{q}_{i+1} = \mathbf{q}_i + \Delta t \dot{\mathbf{q}}_i + \frac{\Delta t^2}{2} ((1-2\beta)\ddot{\mathbf{q}}_i + 2\beta\ddot{\mathbf{q}}_{i+1}) \quad (53)$$

$$\dot{\mathbf{q}}_{i+1} = \dot{\mathbf{q}}_i + \Delta t ((1-\gamma)\ddot{\mathbf{q}}_i + \gamma\ddot{\mathbf{q}}_{i+1})$$

其中,当 $\gamma \geq \frac{1}{2}, \beta \geq \frac{\gamma}{2}$ 时, 求解过程是无条件稳定的。第 $i+1$ 个时间步的动力

学方程残差 \mathbf{R}_d 和约束方程残差 \mathbf{R}_c 如下:

$$\begin{aligned}\mathbf{R}_d(\mathbf{q}_{i+1}^k, \lambda_{i+1}^k) &= \mathbf{M}\mathbf{q}_{i+1}^k + \mathbf{C}\mathbf{q}_{i+1}^k + \mathbf{F}^{ela}(\mathbf{q}_{i+1}^k) + \mathbf{G}^T \lambda_{i+1}^k - (\mathbf{F}(t_{i+1}) \mathbf{S}_F)^T \\ \mathbf{R}_c(\mathbf{q}_{i+1}^k, \lambda_{i+1}^k) &= \mathbf{G}\mathbf{q}_{i+1}^k\end{aligned}\quad (54)$$

根据式 (53), 式中只有 $\mathbf{q}_{i+1}^k, \lambda_{i+1}^k$ 是未知独立变量。

接下来采用牛顿-拉斐逊迭代法迭代求解 \mathbf{q}_{i+1}^k 和 λ_{i+1}^k , 我们定义 $(q_{i+1}^k, \mathbf{q}_{i+1}^k, \lambda_{i+1}^k)$ 是在第 k 次迭代中获得的 $(q_{i+1}, \mathbf{q}_{i+1}, \lambda_{i+1})$ 的近似解, 那么第 $k+1$ 次迭代的残差方程为:

$$\begin{aligned}\mathbf{R}_d(\mathbf{q}_{i+1}^{k+1}, \lambda_{i+1}^{k+1}) &\approx \mathbf{R}_d(\mathbf{q}_{i+1}^k, \lambda_{i+1}^k) + \left. \frac{\partial \mathbf{R}_d}{\partial \mathbf{q}} \right|_{\mathbf{q}=\mathbf{q}_{i+1}^k} \Delta \mathbf{q}_{i+1}^k + \left. \frac{\partial \mathbf{R}_d}{\partial \lambda} \right|_{\lambda=\lambda_{i+1}^k} \Delta \lambda_{i+1}^k \\ \mathbf{R}_c(\mathbf{q}_{i+1}^{k+1}, \lambda_{i+1}^{k+1}) &\approx \mathbf{R}_c(\mathbf{q}_{i+1}^k, \lambda_{i+1}^k) + \left. \frac{\partial \mathbf{R}_c}{\partial \mathbf{q}} \right|_{\mathbf{q}=\mathbf{q}_{i+1}^k} \Delta \mathbf{q}_{i+1}^k + \left. \frac{\partial \mathbf{R}_c}{\partial \lambda} \right|_{\lambda=\lambda_{i+1}^k} \Delta \lambda_{i+1}^k\end{aligned}\quad (55)$$

迭代目标是使得第 $k+1$ 次迭代的残差趋近于零, 因此假设第 $k+1$ 次迭代的残差为零, 式 (55) 可写成下式:

$$\mathbf{S}^k \begin{bmatrix} \Delta \mathbf{q}_{i+1}^k \\ \Delta \lambda_{i+1}^k \end{bmatrix} = -\mathbf{R}^k \quad (56)$$

\mathbf{S}^k 和 \mathbf{R}^k 具体表达式见于附录 C。

于是

$$\begin{bmatrix} \Delta \mathbf{q}_{i+1}^k \\ \Delta \lambda_{i+1}^k \end{bmatrix} = -(\mathbf{S}^k)^{-1} \mathbf{R}^k \quad (57)$$

$$\begin{bmatrix} \mathbf{q}_{i+1}^{k+1} \\ \lambda_{i+1}^{k+1} \end{bmatrix} = \begin{bmatrix} \mathbf{q}_{i+1}^k \\ \lambda_{i+1}^k \end{bmatrix} + \begin{bmatrix} \Delta \mathbf{q}_{i+1}^k \\ \Delta \lambda_{i+1}^k \end{bmatrix} \quad (58)$$

将近似解 \mathbf{q}_{i+1}^{k+1} 代入式 (53) 获得 \mathbf{q}_{i+1}^{k+1} 、 \mathbf{f}_{i+1}^{k+1} ，并重新计算残差 \mathbf{R}^{k+1} 。若是残差满足 $\|\mathbf{R}^{k+1}\| \leq \Xi$ ，则认为此次迭代的近似解满足精度要求，进入下一时间步，直到求解结束。采用上述流程，即可获得一段时间内的非线性振动响应。

值得注意的是，每个时间步的迭代求解过程都要更新 \mathbf{q}_{i+1}^k 对应的 $\mathbf{F}^{ela}(\mathbf{q}_{i+1})$ 和 $\frac{\partial \mathbf{F}^{ela}}{\partial \mathbf{q}}$ ，为了进一步提高计算效率，本文采用了高斯-勒让德数值积分算法，以实现上述迭代过程中对非线性弹性力及其雅可比矩阵的快速更新。以非线性弹性

力的高斯数值积分公式为例，其表达式如下：

$$\begin{aligned} \mathbf{F}^{ela}(\mathbf{q}_{i+1}^k) &= \int_0^{b_e} \int_0^{a_e} \mathbf{f}^{ela}(x, y, \mathbf{q}_{i+1}^k) dx dy \\ &= \frac{a_e b_e}{2} \int_{-1}^1 \int_{-1}^1 \mathbf{f}^{ela} \left(\frac{a_e}{2} \xi + \frac{a_e}{2}, \frac{b_e}{2} \eta + \frac{b_e}{2}, \mathbf{q}_{i+1}^k \right) d\xi d\eta \quad (59) \\ &= \frac{a_e b_e}{2} \sum_{i=1}^m \sum_{j=1}^n w_{\xi_i} w_{\eta_j} \mathbf{f}^{ela} \left(\frac{a_e}{2} \xi_i + \frac{a_e}{2}, \frac{b_e}{2} \eta_j + \frac{b_e}{2}, \mathbf{q}_{i+1}^k \right) \end{aligned}$$

式中， $\mathbf{f}^{ela}(x, y, \mathbf{q}_{i+1}^k) = (\mathbf{B}_e)^T \left((\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{c1} \boldsymbol{\varepsilon}_m + (\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{c2} \boldsymbol{\kappa} + (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{c2} \boldsymbol{\varepsilon}_m + (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{c3} \boldsymbol{\kappa} \right)$ ， ξ_i 、 η_j 、 w_{ξ_i} 、 w_{η_j} ($i=1, L, m$; $j=1, L, n$) 分别为 x 、 y 方向的高斯点及其对应的权重因子， m 、 n 分别 x 、 y 方向的高斯点数量。 ξ 、 η 为积分区间转换后的积分坐标。

$\mathbf{f}^{ela}(x, y, \mathbf{q}_{i+1}^k) = (\mathbf{B}_e)^T \left((\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{c1} \boldsymbol{\varepsilon}_m + (\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{c2} \boldsymbol{\kappa} + (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{c2} \boldsymbol{\varepsilon}_m + (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{c3} \boldsymbol{\kappa} \right)$ 为非线性弹性力积分前的矩阵函数。高斯点的选取是高斯数值积分的重点，构造高斯点的方法有很多，其中最常用的是使用 Legendre 多项式函数构造高斯点。

首先通过迭代方法构造 Legendre 多项式函数：

$$\begin{aligned}
L_0(x) &= 1 \\
L_1(x) &= x \\
L_n(x) &= \frac{2n-1}{n} x L_{n-1}(x) - \frac{n-1}{n} L_{n-2}(x) \quad (n \geq 2)
\end{aligned} \tag{60}$$

使用 matlab 求解上述 Legendre 多项式 $L_n(x)$ 的零点 x_k 作为高斯点。 n 为 Legendre 多项式的阶数，同时也是零点个数。根据求解的精度需求可以选择合适高斯点数，从而确定 Legendre 多项式的阶数。

高斯点对应的权重因子的计算公式如下：

$$w_k = \frac{2}{(1 - (x_k)^2) \left(\frac{\partial L_n(x)}{\partial x} \Big|_{x=x_k} \right)^2} \tag{61}$$

附录 A

法向量表达式如下

$$\begin{aligned}
\mathbf{n} = \mathbf{r}_{,x} \times \mathbf{r}_{,y} &= \begin{bmatrix} 0 & -\mathbf{r}_{3,x} & \mathbf{r}_{2,x} \\ \mathbf{r}_{3,x} & 0 & -\mathbf{r}_{1,x} \\ -\mathbf{r}_{2,x} & \mathbf{r}_{1,x} & 0 \end{bmatrix} \begin{bmatrix} \mathbf{r}_{1,y} \\ \mathbf{r}_{2,y} \\ \mathbf{r}_{3,y} \end{bmatrix} \\
&= \begin{bmatrix} \mathbf{r}_{2,x}\mathbf{r}_{3,y} - \mathbf{r}_{3,x}\mathbf{r}_{2,y} \\ \mathbf{r}_{3,x}\mathbf{r}_{1,y} - \mathbf{r}_{1,x}\mathbf{r}_{3,y} \\ \mathbf{r}_{1,x}\mathbf{r}_{2,y} - \mathbf{r}_{2,x}\mathbf{r}_{1,y} \end{bmatrix} = \begin{bmatrix} (\mathbf{S}_{2,x}\mathbf{q}_e)^T \mathbf{S}_{3,y}\mathbf{q}_e - (\mathbf{S}_{3,x}\mathbf{q}_e)^T \mathbf{S}_{2,y}\mathbf{q}_e \\ (\mathbf{S}_{3,x}\mathbf{q}_e)^T \mathbf{S}_{1,y}\mathbf{q}_e - (\mathbf{S}_{1,x}\mathbf{q}_e)^T \mathbf{S}_{3,y}\mathbf{q}_e \\ (\mathbf{S}_{1,x}\mathbf{q}_e)^T \mathbf{S}_{2,y}\mathbf{q}_e - (\mathbf{S}_{2,x}\mathbf{q}_e)^T \mathbf{S}_{1,y}\mathbf{q}_e \end{bmatrix}
\end{aligned} \tag{A.1}$$

$\boldsymbol{\varepsilon}_{m,q_e}$ 和 $\boldsymbol{\kappa}_{q_e}$ 的表达式如下

$$\boldsymbol{\varepsilon}_{m,q_e} = \begin{bmatrix} (\mathbf{S}_{,x}\mathbf{q}_e)^T \mathbf{S}_{,x} \\ (\mathbf{S}_{,y}\mathbf{q}_e)^T \mathbf{S}_{,y} \\ (\mathbf{S}_{,x}\mathbf{q}_e)^T \mathbf{S}_{,y} + (\mathbf{S}_{,y}\mathbf{q}_e)^T \mathbf{S}_{,x} \end{bmatrix} \tag{A.2}$$

$$\begin{aligned}
\boldsymbol{\kappa}_{,q_e} = \frac{\partial \boldsymbol{\kappa}}{\partial \mathbf{q}_e} &= \begin{bmatrix} \frac{\partial \left((\mathbf{r}_{,xx})^T \mathbf{n} \right)}{\partial \mathbf{q}_e} \bar{n} - (\mathbf{r}_{,xx})^T \mathbf{n} \frac{\partial (\bar{n})}{\partial \mathbf{q}_e}}{\bar{n}^2} \\ \frac{\partial \left((\mathbf{r}_{,yy})^T \mathbf{n} \right)}{\partial \mathbf{q}_e} \bar{n} - (\mathbf{r}_{,yy})^T \mathbf{n} \frac{\partial (\bar{n})}{\partial \mathbf{q}_e}}{\bar{n}^2} \\ \frac{\partial \left((\mathbf{r}_{,xy})^T \mathbf{n} \right)}{\partial \mathbf{q}_e} \bar{n} - (\mathbf{r}_{,xy})^T \mathbf{n} \frac{\partial (\bar{n})}{\partial \mathbf{q}_e}}{\bar{n}^2} \end{bmatrix} = \begin{bmatrix} \left(\mathbf{n}^T \frac{\partial (\mathbf{r}_{,xx})}{\partial \mathbf{q}_e} + (\mathbf{r}_{,xx})^T \frac{\partial (\mathbf{n})}{\partial \mathbf{q}_e} \right) \bar{n} - (\mathbf{r}_{,xx})^T \mathbf{n} \frac{\partial (\bar{n})}{\partial \mathbf{q}_e}}{\bar{n}^2} \\ \left(\mathbf{n}^T \frac{\partial (\mathbf{r}_{,yy})}{\partial \mathbf{q}_e} + (\mathbf{r}_{,yy})^T \frac{\partial (\mathbf{n})}{\partial \mathbf{q}_e} \right) \bar{n} - (\mathbf{r}_{,yy})^T \mathbf{n} \frac{\partial (\bar{n})}{\partial \mathbf{q}_e}}{\bar{n}^2} \\ \left(\mathbf{n}^T \frac{\partial (\mathbf{r}_{,xy})}{\partial \mathbf{q}_e} + (\mathbf{r}_{,xy})^T \frac{\partial (\mathbf{n})}{\partial \mathbf{q}_e} \right) \bar{n} - (\mathbf{r}_{,xy})^T \mathbf{n} \frac{\partial (\bar{n})}{\partial \mathbf{q}_e}}{\bar{n}^2} \end{bmatrix} \\
&= \begin{bmatrix} \frac{(\mathbf{n}^T \mathbf{S}_{,xx} + (\mathbf{r}_{,xx})^T \mathbf{n}_{,q_e}) \bar{n} - (\mathbf{r}_{,xx})^T \mathbf{n} \bar{n}_{,q_e}}{\bar{n}^2} \\ \frac{(\mathbf{n}^T \mathbf{S}_{,yy} + (\mathbf{r}_{,yy})^T \mathbf{n}_{,q_e}) \bar{n} - (\mathbf{r}_{,yy})^T \mathbf{n} \bar{n}_{,q_e}}{\bar{n}^2} \\ \frac{(\mathbf{n}^T 2\mathbf{S}_{,xy} + (\mathbf{r}_{,xy})^T \mathbf{n}_{,q_e}) \bar{n} - (\mathbf{r}_{,xy})^T \mathbf{n} \bar{n}_{,q_e}}{\bar{n}^2} \end{bmatrix} = \begin{bmatrix} \frac{\mathbf{n}^T \mathbf{S}_{,xx}}{\bar{n}} \\ \frac{\mathbf{n}^T \mathbf{S}_{,yy}}{\bar{n}} \\ \frac{\mathbf{n}^T 2\mathbf{S}_{,xy}}{\bar{n}} \end{bmatrix} + \begin{bmatrix} \frac{(\mathbf{r}_{,xx})^T \mathbf{n}_{,q_e}}{\bar{n}} \\ \frac{(\mathbf{r}_{,yy})^T \mathbf{n}_{,q_e}}{\bar{n}} \\ \frac{(\mathbf{r}_{,xy})^T \mathbf{n}_{,q_e}}{\bar{n}} \end{bmatrix} - \boldsymbol{\kappa} \frac{\bar{n}_{,q_e}}{\bar{n}} \quad (\text{A.3})
\end{aligned}$$

其中

$$\begin{aligned}
\mathbf{r}_{,xx} &= \mathbf{S}_{,xx} \mathbf{q}_e \\
\mathbf{r}_{,yy} &= \mathbf{S}_{,yy} \mathbf{q}_e \\
\mathbf{r}_{,xy} &= 2\mathbf{S}_{,xy} \mathbf{q}_e
\end{aligned} \quad (\text{A.4})$$

$$\mathbf{n}_{,q_e} = \frac{\partial \mathbf{n}}{\partial \mathbf{q}_e} = \begin{bmatrix} (\mathbf{S}_{2,x} \mathbf{q}_e)^T \mathbf{S}_{3,y} + (\mathbf{S}_{3,y} \mathbf{q}_e)^T \mathbf{S}_{2,x} - (\mathbf{S}_{3,x} \mathbf{q}_e)^T \mathbf{S}_{2,y} - (\mathbf{S}_{2,y} \mathbf{q}_e)^T \mathbf{S}_{3,x} \\ (\mathbf{S}_{3,x} \mathbf{q}_e)^T \mathbf{S}_{1,y} + (\mathbf{S}_{1,y} \mathbf{q}_e)^T \mathbf{S}_{3,x} - (\mathbf{S}_{1,x} \mathbf{q}_e)^T \mathbf{S}_{3,y} - (\mathbf{S}_{3,y} \mathbf{q}_e)^T \mathbf{S}_{1,x} \\ (\mathbf{S}_{1,x} \mathbf{q}_e)^T \mathbf{S}_{2,y} + (\mathbf{S}_{2,y} \mathbf{q}_e)^T \mathbf{S}_{1,x} - (\mathbf{S}_{2,x} \mathbf{q}_e)^T \mathbf{S}_{1,y} - (\mathbf{S}_{1,y} \mathbf{q}_e)^T \mathbf{S}_{2,x} \end{bmatrix} \quad (\text{A.5})$$

$$\bar{n}_{,q_e} = \frac{\partial \bar{n}}{\partial \mathbf{q}_e} = \frac{\partial \bar{n}}{\partial \mathbf{n}} \frac{\partial \mathbf{n}}{\partial \mathbf{q}_e} = \frac{\mathbf{n}^T}{\bar{n}} \mathbf{n}_{,q_e} \quad (\text{A.6})$$

惯性、刚度和压电等参数矩阵如下：

$$I_{c1} = \rho_c (h_u - h_b) \quad (\text{A.7})$$

$$\mathbf{D}_{c1} = \sum_{k=1}^{N_f} (h_u^k - h_b^k) \bar{\mathbf{D}}_c^k$$

$$\mathbf{D}_{c2} = \sum_{k=1}^{N_f} \left(\frac{(h_u^k)^2}{2} - \frac{(h_b^k)^2}{2} \right) \bar{\mathbf{D}}_c^k$$

$$\mathbf{D}_{c3} = \sum_{k=1}^{N_f} \left(\frac{(h_u^k)^3}{3} - \frac{(h_b^k)^3}{3} \right) \bar{\mathbf{D}}_c^k$$

$$I_{p1} = \rho_p (h_p^u - h_p^b)$$

$$\mathbf{D}_{p1} = (h_p^u - h_p^b) \bar{\mathbf{D}}_p$$

$$\mathbf{D}_{p2} = \left(\frac{(h_p^u)^2}{2} - \frac{(h_p^b)^2}{2} \right) \bar{\mathbf{D}}_p$$

$$\mathbf{D}_{p3} = \left(\frac{(h_p^u)^3}{3} - \frac{(h_p^b)^3}{3} \right) \bar{\mathbf{D}}_p$$

$$\mathbf{e}_{p1} = (h_p^u - h_p^b) \bar{\mathbf{e}}$$

$$\mathbf{e}_{p2} = \left(\frac{(h_p^u)^2}{2} - \frac{(h_p^b)^2}{2} \right) \bar{\mathbf{e}}$$

$$\varsigma_{33}^1 = (h_p^u - h_p^b) \varsigma_{33}$$

附录 B

$$(\mathbf{M}_c + \mathbf{M}_a + \mathbf{M}_s) \mathbf{q} + \mathbf{F}_c^{ela} + \mathbf{F}_a^{ela} + \mathbf{F}_s^{ela} - (\mathbf{K}_{ca})^T \boldsymbol{\phi}_a - (\mathbf{K}_{cs})^T \boldsymbol{\phi}_s - (\mathbf{F} \mathbf{S}_F)^T = 0$$

(B.1)

$$K_{ac} + K_{aa} \phi_a - \left(\frac{h_E}{h_p} \right)^2 K_{aa} \phi_a^c = 0 \quad (\text{B.2})$$

$$K_{sc} + K_{ss} \phi_s = 0 \quad (\text{B.3})$$

由 (B.2) 和 (B.3) 可得

$$\phi_a = (K_{aa})^{-1} \left(\left(\frac{h_E}{h_p} \right)^2 K_{aa} \phi_a^c - K_{ac} \right) \quad (\text{B.4})$$

$$\phi_s = (K_{ss})^{-1} (-K_{sc}) \quad (B.5)$$

将 (B.4) 和 (B.5) 代入式 (B.1)

$$\begin{aligned} & (\mathbf{M}_c + \mathbf{M}_a + \mathbf{M}_s) \mathbf{F} \\ & + \mathbf{F}_c^{ela} + \mathbf{F}_a^{ela} + \mathbf{F}_s^{ela} \\ & - (\mathbf{K}_{ca})^T (K_{aa})^{-1} \left(\left(\frac{h_E}{h_p} \right)^2 K_{aa} \phi_a^c - K_{ac} \right) \\ & - (\mathbf{K}_{cs})^T (K_{ss})^{-1} (-K_{sc}) \\ & = (\mathbf{F}_F)^T \end{aligned} \quad (B.6)$$

进一步推得

$$\begin{aligned} & (\mathbf{M}_c + \mathbf{M}_a + \mathbf{M}_s) \mathbf{F} \\ & + \mathbf{F}_c^{ela} + \mathbf{F}_a^{ela} + \mathbf{F}_s^{ela} + (\mathbf{K}_{ca})^T (K_{aa})^{-1} K_{ac} + (\mathbf{K}_{cs})^T (K_{ss})^{-1} K_{sc} \\ & = (\mathbf{F}_F)^T + (\mathbf{K}_{ca})^T \left(\frac{h_E}{h_p} \right)^2 \phi_a^c \end{aligned} \quad (B.7)$$

因此

$$\mathbf{M} = \mathbf{M}_c + \mathbf{M}_a + \mathbf{M}_s \quad (B.8)$$

$$\mathbf{F}^{ela} = \mathbf{F}_c^{ela} + \mathbf{F}_a^{ela} + \mathbf{F}_s^{ela} + (\mathbf{K}_{ca})^T (K_{aa})^{-1} K_{ac} + (\mathbf{K}_{cs})^T (K_{ss})^{-1} K_{sc} \quad (B.9)$$

$$\mathbf{K}_{con} = (\mathbf{K}_{ca})^T \left(\frac{h_E}{h_p} \right)^2 \quad (B.10)$$

其中

$$\mathbf{M}_c = \sum_e^{N_e^c} \iint_{S_e} (\mathbf{B}_e)^T (I_{c1}(\mathbf{S})^T \mathbf{S}) \mathbf{B}_e dx dy \quad (B.11)$$

$$\mathbf{M}_p = \sum_e^{N_e^p} \iint_{S_e} (\mathbf{B}_e)^T (I_{p1}(\mathbf{S})^T \mathbf{S}) \mathbf{B}_e dx dy \quad (B.12)$$

$$\mathbf{F}_c^{ela} = \sum_e^{N_e^c} \iint_{S_e} (\mathbf{B}_e)^T \left((\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{c1} \boldsymbol{\varepsilon}_m + (\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{c2} \boldsymbol{\kappa} + (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{c2} \boldsymbol{\varepsilon}_m + (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{c3} \boldsymbol{\kappa} \right) dx dy \quad (B.13)$$

$$\mathbf{F}_p^{ela} = \sum_e^{N_e^p} \iint_{S_e} (\mathbf{B}_e)^T \left((\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{p1} \boldsymbol{\varepsilon}_m + (\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{p2} \boldsymbol{\kappa} + (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{p2} \boldsymbol{\varepsilon}_m + (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{p3} \boldsymbol{\kappa} \right) dx dy \quad (B.14)$$

$$\mathbf{K}_{cp} = \sum_e^{N_e^p} \iint_{S_e} \left((\mathbf{e}_{p1})^T \frac{1}{h_E} \boldsymbol{\varepsilon}_{m,q_e} + (\mathbf{e}_{p2})^T \frac{1}{h_E} \boldsymbol{\kappa}_{,q_e} \right) \mathbf{B}_e dx dy \quad (\text{B.15})$$

$$K_{pc} = \sum_e^{N_e^p} \iint_{S_e} (\mathbf{e}_{p1})^T \frac{1}{h_E} \boldsymbol{\varepsilon}_m + (\mathbf{e}_{p2})^T \frac{1}{h_E} \boldsymbol{\kappa} dx dy \quad (\text{B.16})$$

$$K_{pp} = \sum_e^{N_e^p} \iint_{S_e} \zeta_{z_p z_p}^{-1} \left(\frac{1}{h_E} \right)^2 \phi_p dx dy \quad (\text{B.17})$$

附录 C

$$\mathbf{S}^k = \begin{bmatrix} \left. \frac{\partial \mathbf{R}_d}{\partial \boldsymbol{\Phi}} \right|_{\boldsymbol{\Phi}=\boldsymbol{\Phi}_{i+1}^k} & \left. \frac{\partial \mathbf{R}_d}{\partial \boldsymbol{\lambda}} \right|_{\boldsymbol{\lambda}=\boldsymbol{\lambda}_{i+1}^k} \\ \left. \frac{\partial \mathbf{R}_c}{\partial \boldsymbol{\Phi}} \right|_{\boldsymbol{\Phi}=\boldsymbol{\Phi}_{i+1}^k} & \left. \frac{\partial \mathbf{R}_c}{\partial \boldsymbol{\lambda}} \right|_{\boldsymbol{\lambda}=\boldsymbol{\lambda}_{i+1}^k} \end{bmatrix} = \begin{bmatrix} \left. \frac{\partial \mathbf{R}_d}{\partial \boldsymbol{\Phi}} \right|_{\boldsymbol{\Phi}=\boldsymbol{\Phi}_{i+1}^k} & \mathbf{G}^T \\ \mathbf{G} & 0 \end{bmatrix} \quad (\text{C.1})$$

$$\mathbf{R}^k = \begin{bmatrix} \mathbf{R}_d(\boldsymbol{\Phi}_{i+1}^k, \boldsymbol{\lambda}_{i+1}^k) \\ \mathbf{R}_c(\boldsymbol{\Phi}_{i+1}^k, \boldsymbol{\lambda}_{i+1}^k) \end{bmatrix}$$

$$\frac{\partial \mathbf{R}_d}{\partial \boldsymbol{\Phi}} = \mathbf{M} + \mathbf{C} \frac{\partial \boldsymbol{\Phi}}{\partial \boldsymbol{\Phi}} + \frac{\partial \mathbf{F}^{ela}}{\partial \mathbf{q}} \frac{\partial \mathbf{q}}{\partial \boldsymbol{\Phi}} \quad (\text{C.2})$$

根据式 (53), 其中 \mathbf{q}_i , $\boldsymbol{\Phi}_i$ 和 $\boldsymbol{\lambda}_i$ 在 $i+1$ 时刻都是已知的, 由此可得

$$\frac{\partial \boldsymbol{\Phi}}{\partial \boldsymbol{\Phi}} = \Delta t \boldsymbol{\gamma}, \frac{\partial \mathbf{q}}{\partial \boldsymbol{\Phi}} = \frac{\Delta t^2}{2} 2\boldsymbol{\beta} \quad (\text{C.3})$$

$$\begin{aligned}
\frac{\partial \mathbf{F}^{ela}}{\partial \mathbf{q}} &= \frac{\partial \left(\mathbf{F}_c^{ela} + \mathbf{F}_a^{ela} + \mathbf{F}_s^{ela} + (\mathbf{K}_{ca})^T (K_{aa})^{-1} K_{ac} + (\mathbf{K}_{cs})^T (K_{ss})^{-1} K_{sc} \right)}{\partial \mathbf{q}} \\
&= \frac{\partial \mathbf{F}_c^{ela}}{\partial \mathbf{q}} + \frac{\partial \mathbf{F}_a^{ela}}{\partial \mathbf{q}} + \frac{\partial \mathbf{F}_s^{ela}}{\partial \mathbf{q}} + \frac{\partial \left((\mathbf{K}_{ca})^T (K_{aa})^{-1} K_{ac} \right)}{\partial \mathbf{q}} + \frac{\partial \left((\mathbf{K}_{cs})^T (K_{ss})^{-1} K_{sc} \right)}{\partial \mathbf{q}} \quad (C.4) \\
&= \frac{\partial \mathbf{F}_c^{ela}}{\partial \mathbf{q}} + \frac{\partial \mathbf{F}_a^{ela}}{\partial \mathbf{q}} + \frac{\partial \mathbf{F}_s^{ela}}{\partial \mathbf{q}} \\
&\quad + (\mathbf{K}_{ca})^T (K_{aa})^{-1} \frac{\partial (K_{ac})}{\partial \mathbf{q}} + (K_{aa})^{-1} K_{ac} \frac{\partial (\mathbf{K}_{ca})^T}{\partial \mathbf{q}} \\
&\quad + (\mathbf{K}_{cs})^T (K_{ss})^{-1} \frac{\partial (K_{sc})}{\partial \mathbf{q}} + (K_{ss})^{-1} K_{sc} \frac{\partial (\mathbf{K}_{cs})^T}{\partial \mathbf{q}}
\end{aligned}$$

其中

$$\begin{aligned}
\frac{\partial \mathbf{F}_c^{ela}}{\partial \mathbf{q}} &= \frac{\partial \left(\sum_e^{N_e} \iint_{S_e} (\mathbf{B}_e)^T \left((\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{c1} \boldsymbol{\varepsilon}_m + (\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{c2} \boldsymbol{\kappa} + (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{c2} \boldsymbol{\varepsilon}_m + (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{c3} \boldsymbol{\kappa} \right) dx dy \right)}{\partial \mathbf{q}} \\
&= \sum_e^{N_e} \iint_{S_e} (\mathbf{B}_e)^T \left(\frac{\partial (\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{c1} \boldsymbol{\varepsilon}_m}{\partial \mathbf{q}_e} + \frac{\partial (\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{c2} \boldsymbol{\kappa}}{\partial \mathbf{q}_e} + \frac{\partial (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{c2} \boldsymbol{\varepsilon}_m}{\partial \mathbf{q}_e} + \frac{\partial (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{c3} \boldsymbol{\kappa}}{\partial \mathbf{q}_e} \right) \frac{\partial \mathbf{q}_e}{\partial \mathbf{q}} dx dy \\
&= \sum_e^{N_e} \iint_{S_e} (\mathbf{B}_e)^T \left(\begin{aligned} &(\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{c1} \boldsymbol{\varepsilon}_{m,q_e} + \mathbf{D}_{c1} (1,:) \boldsymbol{\varepsilon}_m \boldsymbol{\varepsilon}_{m1,q,q_e} + \mathbf{D}_{c1} (2,:) \boldsymbol{\varepsilon}_m \boldsymbol{\varepsilon}_{m2,q,q_e} + \mathbf{D}_{c1} (3,:) \boldsymbol{\varepsilon}_m \boldsymbol{\varepsilon}_{m3,q,q_e} \\ &+ (\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{c2} \boldsymbol{\kappa}_{,q_e} + \mathbf{D}_{c2} (1,:) \boldsymbol{\kappa} \boldsymbol{\varepsilon}_{m1,q,q_e} + \mathbf{D}_{c2} (2,:) \boldsymbol{\kappa} \boldsymbol{\varepsilon}_{m2,q,q_e} + \mathbf{D}_{c2} (3,:) \boldsymbol{\kappa} \boldsymbol{\varepsilon}_{m3,q,q_e} \\ &+ (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{c2} \boldsymbol{\varepsilon}_{m,q_e} + \mathbf{D}_{c2} (1,:) \boldsymbol{\varepsilon}_m \boldsymbol{\kappa}_{1,q,q_e} + \mathbf{D}_{c2} (2,:) \boldsymbol{\varepsilon}_m \boldsymbol{\kappa}_{2,q,q_e} + \mathbf{D}_{c2} (3,:) \boldsymbol{\varepsilon}_m \boldsymbol{\kappa}_{3,q,q_e} \\ &+ (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{c3} \boldsymbol{\kappa}_{,q_e} + \mathbf{D}_{c3} (1,:) \boldsymbol{\kappa} \boldsymbol{\kappa}_{1,q,q_e} + \mathbf{D}_{c3} (2,:) \boldsymbol{\kappa} \boldsymbol{\kappa}_{2,q,q_e} + \mathbf{D}_{c3} (3,:) \boldsymbol{\kappa} \boldsymbol{\kappa}_{3,q,q_e} \end{aligned} \right) \mathbf{B}_e dx dy \quad (C.5)
\end{aligned}$$

同理

$$\begin{aligned}
\frac{\partial \mathbf{F}_p^{ela}}{\partial \mathbf{q}} &= \sum_e^{N_p} \iint_{S_e} (\mathbf{B}_e)^T \left(\begin{aligned} &(\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{p1} \boldsymbol{\varepsilon}_{m,q_e} + \mathbf{D}_{p1} (1,:) \boldsymbol{\varepsilon}_m \boldsymbol{\varepsilon}_{m1,q,q_e} + \mathbf{D}_{p1} (2,:) \boldsymbol{\varepsilon}_m \boldsymbol{\varepsilon}_{m2,q,q_e} + \mathbf{D}_{p1} (3,:) \boldsymbol{\varepsilon}_m \boldsymbol{\varepsilon}_{m3,q,q_e} \\ &+ (\boldsymbol{\varepsilon}_{m,q_e})^T \mathbf{D}_{p2} \boldsymbol{\kappa}_{,q_e} + \mathbf{D}_{p2} (1,:) \boldsymbol{\kappa} \boldsymbol{\varepsilon}_{m1,q,q_e} + \mathbf{D}_{p2} (2,:) \boldsymbol{\kappa} \boldsymbol{\varepsilon}_{m2,q,q_e} + \mathbf{D}_{p2} (3,:) \boldsymbol{\kappa} \boldsymbol{\varepsilon}_{m3,q,q_e} \\ &+ (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{p2} \boldsymbol{\varepsilon}_{m,q_e} + \mathbf{D}_{p2} (1,:) \boldsymbol{\varepsilon}_m \boldsymbol{\kappa}_{1,q,q_e} + \mathbf{D}_{p2} (2,:) \boldsymbol{\varepsilon}_m \boldsymbol{\kappa}_{2,q,q_e} + \mathbf{D}_{p2} (3,:) \boldsymbol{\varepsilon}_m \boldsymbol{\kappa}_{3,q,q_e} \\ &+ (\boldsymbol{\kappa}_{,q_e})^T \mathbf{D}_{p3} \boldsymbol{\kappa}_{,q_e} + \mathbf{D}_{p3} (1,:) \boldsymbol{\kappa} \boldsymbol{\kappa}_{1,q,q_e} + \mathbf{D}_{p3} (2,:) \boldsymbol{\kappa} \boldsymbol{\kappa}_{2,q,q_e} + \mathbf{D}_{p3} (3,:) \boldsymbol{\kappa} \boldsymbol{\kappa}_{3,q,q_e} \end{aligned} \right) \mathbf{B}_e dx dy \quad (C.6) \\
\frac{\partial (K_{pc})}{\partial \mathbf{q}} &= \frac{\partial \left(\sum_e^{N_p} \iint_{S_e} (\mathbf{e}_{p1})^T \frac{1}{h_E} \boldsymbol{\varepsilon}_m + (\mathbf{e}_{p2})^T \frac{1}{h_E} \boldsymbol{\kappa} dx dy \right)}{\partial \mathbf{q}} \quad (C.7) \\
&= \sum_e^{N_p} \iint_{S_e} \left((\mathbf{e}_{p1})^T \frac{1}{h_E} \boldsymbol{\varepsilon}_{m,q_e} + (\mathbf{e}_{p2})^T \frac{1}{h_E} \boldsymbol{\kappa}_{,q_e} \right) \mathbf{B}_e dx dy
\end{aligned}$$

$$\begin{aligned}
\frac{\partial (\mathbf{K}_{cp})^T}{\partial \mathbf{q}} &= \frac{\partial \sum_e^{N_e^p} \iint_{S_e} (\mathbf{B}_e)^T \left((\boldsymbol{\varepsilon}_{m,q_e})^T \frac{1}{h_E} \mathbf{e}_{p1} + (\boldsymbol{\kappa}_{,q_e})^T \frac{1}{h_E} \mathbf{e}_{p2} \right) dx dy}{\partial \mathbf{q}} \\
&= \sum_e^{N_e^p} \iint_{S_e} (\mathbf{B}_e)^T \left(\frac{1}{h_E} \mathbf{e}_{p1} (1) \boldsymbol{\varepsilon}_{m1,q_e} + \frac{1}{h_E} \mathbf{e}_{p1} (2) \boldsymbol{\varepsilon}_{m2,q_e} + \frac{1}{h_E} \mathbf{e}_{p1} (3) \boldsymbol{\varepsilon}_{m3,q_e} \right. \\
&\quad \left. + \frac{1}{h_E} \mathbf{e}_{p2} (1) \boldsymbol{\kappa}_{1,q_e} + \frac{1}{h_E} \mathbf{e}_{p2} (2) \boldsymbol{\kappa}_{2,q_e} + \frac{1}{h_E} \mathbf{e}_{p2} (3) \boldsymbol{\kappa}_{3,q_e} \right) \mathbf{B}_e dx dy
\end{aligned} \tag{C.8}$$

其中

$$\begin{aligned}
\boldsymbol{\varepsilon}_{m1,q_e} &= (\mathbf{S}_{,x})^T \mathbf{S}_{,x} \\
\boldsymbol{\varepsilon}_{m2,q_e} &= (\mathbf{S}_{,y})^T \mathbf{S}_{,y} \\
\boldsymbol{\varepsilon}_{m3,q_e} &= (\mathbf{S}_{,x})^T \mathbf{S}_{,y} + (\mathbf{S}_{,y})^T \mathbf{S}_{,x}
\end{aligned} \tag{C.9}$$

$$\begin{aligned}
\boldsymbol{\kappa}_{1,q_e} &= \frac{\partial \boldsymbol{\kappa}_{1,q_e}}{\partial \mathbf{q}_e} = \frac{\partial}{\partial \mathbf{q}_e} \left(\frac{\mathbf{n}^T \mathbf{S}_{,xx}}{\bar{n}} + \frac{(\mathbf{r}_{,xx})^T \mathbf{n}_{,q_e}}{\bar{n}} - \boldsymbol{\kappa}_1 \frac{\bar{n}_{,q_e}}{\bar{n}} \right) \\
&= \frac{(\mathbf{S}_{,xx})^T \mathbf{n}_{,q_e} \bar{n} - (\mathbf{n}^T \mathbf{S}_{,xx})^T \bar{n}_{,q_e}}{\bar{n}^2} \\
&\quad + \frac{\left((\mathbf{n}_{,q_e})^T \mathbf{S}_{,xx} + \mathbf{r}_{1,xx} \mathbf{n}_{1,q_e} + \mathbf{r}_{2,xx} \mathbf{n}_{2,q_e} + \mathbf{r}_{3,xx} \mathbf{n}_{3,q_e} \right) \bar{n} - \left((\mathbf{r}_{,xx})^T \mathbf{n}_{,q_e} \right)^T \bar{n}_{,q_e}}{\bar{n}^2} \\
&\quad - \boldsymbol{\kappa}_1 \frac{\bar{n}_{,q_e} \bar{n} - (\bar{n}_{,q_e})^T \bar{n}_{,q_e}}{\bar{n}^2} - \left(\frac{\bar{n}_{,q_e}}{\bar{n}} \right)^T \boldsymbol{\kappa}_{1,q_e}
\end{aligned} \tag{C.10}$$

$$\begin{aligned}
\boldsymbol{\kappa}_{2,q_e} &= \frac{\partial \boldsymbol{\kappa}_{2,q_e}}{\partial \mathbf{q}_e} = \frac{\partial}{\partial \mathbf{q}_e} \left(\frac{\mathbf{n}^T \mathbf{S}_{,yy}}{\bar{n}} + \frac{(\mathbf{r}_{,yy})^T \mathbf{n}_{,q_e}}{\bar{n}} - \boldsymbol{\kappa}_1 \frac{\bar{n}_{,q_e}}{\bar{n}} \right) \\
&= \frac{(\mathbf{S}_{,yy})^T \mathbf{n}_{,q_e} \bar{n} - (\mathbf{n}^T \mathbf{S}_{,yy})^T \bar{n}_{,q_e}}{\bar{n}^2} \\
&\quad + \frac{\left((\mathbf{n}_{,q_e})^T \mathbf{S}_{,yy} + \mathbf{r}_{1,yy} \mathbf{n}_{1,q_e} + \mathbf{r}_{2,yy} \mathbf{n}_{2,q_e} + \mathbf{r}_{3,yy} \mathbf{n}_{3,q_e} \right) \bar{n} - \left((\mathbf{r}_{,yy})^T \mathbf{n}_{,q_e} \right)^T \bar{n}_{,q_e}}{\bar{n}^2} \\
&\quad - \boldsymbol{\kappa}_2 \frac{\bar{n}_{,q_e} \bar{n} - (\bar{n}_{,q_e})^T \bar{n}_{,q_e}}{\bar{n}^2} - \left(\frac{\bar{n}_{,q_e}}{\bar{n}} \right)^T \boldsymbol{\kappa}_{2,q_e}
\end{aligned} \tag{C.11}$$

$$\begin{aligned}
\kappa_{3,q_e,q_e} &= \frac{\partial \kappa_{3,q_e}}{\partial \mathbf{q}_e} = \frac{\partial}{\partial \mathbf{q}_e} \left(\frac{\mathbf{n}^T 2\mathbf{S}_{,xy}}{\bar{n}} + \frac{(\mathbf{r}_{,xy})^T \mathbf{n}_{,q_e}}{\bar{n}} - \kappa_1 \frac{\bar{n}_{,q_e}}{\bar{n}} \right) \\
&= \frac{(2\mathbf{S}_{,xy})^T \mathbf{n}_{,q_e} \bar{n} - (\mathbf{n}^T 2\mathbf{S}_{,xy})^T \bar{n}_{,q_e}}{\bar{n}^2} \\
&\quad + \frac{\left((\mathbf{n}_{,q_e})^T 2\mathbf{S}_{,xy} + \mathbf{r}_{1,xy} \mathbf{n}_{1,q_e,q_e} + \mathbf{r}_{2,xy} \mathbf{n}_{2,q_e,q_e} + \mathbf{r}_{3,xy} \mathbf{n}_{3,q_e,q_e} \right) \bar{n} - \left((\mathbf{r}_{,xy})^T \mathbf{n}_{,q_e} \right)^T \bar{n}_{,q_e}}{\bar{n}^2} \\
&\quad - \kappa_3 \frac{\bar{n}_{,q_e,q_e} \bar{n} - (\bar{n}_{,q_e})^T \bar{n}_{,q_e}}{\bar{n}^2} - \left(\frac{\bar{n}_{,q_e}}{\bar{n}} \right)^T \kappa_{3,q_e}
\end{aligned} \tag{C.12}$$

其中,

$$\bar{n}_{,q_e,q_e} = \frac{\partial \bar{n}_{,q_e}}{\partial \mathbf{q}_e} = \frac{\left((\mathbf{n}_{,q_e})^T \mathbf{n}_{,q_e} + \mathbf{n}_1 \mathbf{n}_{1,q_e,q_e} + \mathbf{n}_2 \mathbf{n}_{2,q_e,q_e} + \mathbf{n}_3 \mathbf{n}_{3,q_e,q_e} \right) \bar{n} - (\mathbf{n}^T \mathbf{n}_{,q_e})^T \bar{n}_{,q_e}}{\bar{n}^2} \tag{C.13}$$

$$\begin{aligned}
\mathbf{n}_{1,q_e,q_e} &= \frac{\partial \mathbf{n}_{1,q_e}}{\partial \mathbf{q}_e} = (\mathbf{S}_{2,x})^T \mathbf{S}_{3,y} + (\mathbf{S}_{3,y})^T \mathbf{S}_{2,x} - (\mathbf{S}_{3,x})^T \mathbf{S}_{2,y} - (\mathbf{S}_{2,y})^T \mathbf{S}_{3,x} \\
\mathbf{n}_{2,q_e,q_e} &= \frac{\partial \mathbf{n}_{2,q_e}}{\partial \mathbf{q}_e} = (\mathbf{S}_{3,x})^T \mathbf{S}_{1,y} + (\mathbf{S}_{1,y})^T \mathbf{S}_{3,x} - (\mathbf{S}_{1,x})^T \mathbf{S}_{3,y} - (\mathbf{S}_{3,y})^T \mathbf{S}_{1,x} \\
\mathbf{n}_{3,q_e,q_e} &= \frac{\partial \mathbf{n}_{3,q_e}}{\partial \mathbf{q}_e} = (\mathbf{S}_{1,x})^T \mathbf{S}_{2,y} + (\mathbf{S}_{2,y})^T \mathbf{S}_{1,x} - (\mathbf{S}_{2,x})^T \mathbf{S}_{1,y} - (\mathbf{S}_{1,y})^T \mathbf{S}_{2,x}
\end{aligned} \tag{C.14}$$