

# Analytical Predictions of Explanet Obliquities Generated by Planet-Disk Interactions

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## ABSTRACT

Abstract here

**Key words:** planet–star interactions

## 1 INTRODUCTION

Separatrix crossing was studied by (Henrard 1982).

## 2 THEORY AND EQUATIONS

We study an oblate planet orbiting around a star hosting a further out planet. The equations of motion in the absence of dissipation are:

We define  $s_c$  to be the critical spin where the precession frequencies  $\omega_{s1}$  and  $\omega_{21}$  are equal, or:

Throughout this paper, we often set the orbital frequency of the inner planet  $\Omega_1 = 1$ .

### 2.1 Cassini States

Equilibria of spin dynamics are *Cassini States* (CSs). Refer to Su & Lai 2020 for more detailed discussion. In the notation of the previous paper,

$$\eta = s_c/s. \tag{1}$$

Plots of CS locations as always.

Note that  $\phi$  is conjugate to  $\cos \theta$ , and thus the spin dynamics exhibit Hamiltonian:

Plot of level curves of Hamiltonian, including  $C_{\pm}$  notation.

## 2.2 Weak Tidal Friction

We use the weak tidal friction model (Lai 2012). The EOM for  $\theta, s$  become:

The phase portrait for these EOM in  $(s, \theta)$  space is shown in Fig. flow-diagram.

TODO overplot CS1, CS2 for two values of  $s_c$ . TODO show on plot the location where tides are so strong that CS2 becomes unstable (is the only  $\epsilon$ -dependent result of the paper).

## 2.3 Stable Equilibria of Tidal Friction

Generally, these EOM have equilibria at the points that are both a CS and satisfy  $\dot{s} = 0$  under weak tidal friction. These are stable as  $\epsilon \rightarrow 0$ , see Appendix A. On Fig. flow-diagram, these are the intersection of the locations of CS1 and CS2 with the line  $\dot{s} = 0$ . It is clear that the locations of these equilibria depend on the value of  $s_c$ . Call these generalized equilibria *tidal Cassini States* (tCS), and number them tCS1 and tCS2 depending on whether they are CS1 or CS2 states.

Furthermore, if tides are too strong (large  $\epsilon$ ), CS2 can become unstable (cite Fabrycky). This can be quantified as:

# 3 PROBABILITY DISTRIBUTION OF OUTCOMES

The general question of the dynamics is then as follows: given problem parameters (including  $s_c$ ) and an initial planet spin and obliquity  $(s, \theta)$ , what are the possible outcomes and their associated probabilities? We study this first at fixed  $s_c$  as a function of  $\theta$  in Section 3.1, then as a function of  $s_c$  for an isotropic initial spin  $\hat{s}$  in Section 3.2.

## 3.1 Distribution As a Function of $\theta$

As a result of subsection 2.3, the tCS are the only possible final outcomes. We generate initial spin vectors  $\hat{s}$  for isotropic initial distribution  $(\cos \theta, \phi)$ , and marginalize over  $\phi$  (which generally cannot be fixed by any formation mechanism) in histograms:

The governing mechanism for these probabilities is probabilistic separatrix crossing, which is covered in Appendix B3. The calculation is difficult analytically, but can be performed semi-analytically (describe procedure here, calculate  $\eta_\star$  numerically etc.). The agreement of this curve with the histograms is examined for the same three  $s_c$  values below:

### 3.2 Distribution As a Function of $s_c$

Assuming that  $\hat{s}$  is drawn from an isotropic distribution, we may then calculate the probabilities of going to either tCS as a function of  $s_c$ . This is shown below:

Note that while only the results for an isotropic distribution of initial  $\hat{s}$  are shown, in principle arbitrary distributions  $P(\theta_i)$  can be convolved against the  $P_{tCS2}(\theta_i)$  distributions shown in [subsection 3.1](#).

## 4 SUMMARY AND DISCUSSION

### REFERENCES

Henrard J., 1982, *Celestial Mechanics and Dynamical Astronomy*, 27, 3  
 Lai D., 2012, *Monthly Notices of the Royal Astronomical Society*, 423, 486

### APPENDIX A: PROOF OF STABILITY OF CASSINI STATES 1 & 2

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### APPENDIX B: SEPARATRIX CROSSING DYNAMICS

#### B1 Theory

##### *B1.1 Application of Adiabatic Invariance: Henrard Theory*

Review Henrard result, is already very successful e.g. MMR capture.

##### *B1.2 Melnikov Integral*

The first substantial new result.

#### B2 Example: Constant $s$

“Toy problem 1”, the nice  $P_c \propto \eta^{3/2}$  result. Application of Section [B1.2](#).

#### B3 Separatrix Crossing Probability: Tidal Friction

Application of the full formula presented in Section [B1](#). The key result is that one integrates

The capture probabilities are then

An analytical form that holds when  $s \gg s_c$  is:

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