

Dynamical Tides in Eccentric Massive Stellar Binaries

Group Meeting

Yubo Su

Jan 23, 2020

- Massive star with eccentric binary companion inducing dynamical tides.
- Primary difficulty: dynamical tides is typically messy, sum over many modes, hard to gain analytical intuition.
- Question: can we obtain a *simple closed form for dynamical tides* in this system?
- Dynamical tide in massive stars due to companion on *circular orbit* (Kushnir et. al. 2017).

$$\tau(\omega; r_c) = \beta_2 \frac{GM_2^2 r_c^5}{a^6} \frac{\rho_c}{\bar{\rho}_c} \left(1 - \frac{\rho_c}{\bar{\rho}_c}\right)^2 \times \left(\frac{\omega}{\sqrt{GM_c/r_c^3}}\right)^{8/3}.$$

- Eccentric forcing is just sum of many circular forcings (Fourier transform, e.g. Vick et. al. 2017)

$$\tau_{\text{tot}} = T_0 \sum_{N=-\infty}^{\infty} F_{N2}^2 \text{sgn}(\sigma) \tau(\omega = |\sigma|),$$

where $\sigma \equiv N\Omega - 2\Omega_s$ and F_{Nm} are the *Hansen coefficients*

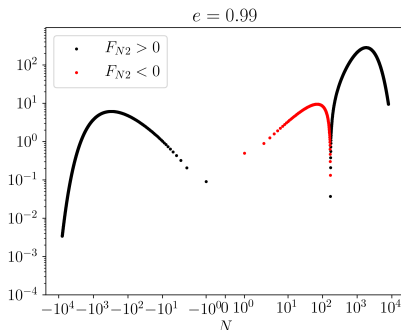
$$F_{Nm} = \frac{1}{\pi} \int_0^\pi \frac{\cos[N\mathcal{M}(E) - mf(E)]}{(1 - e \cos E)^2} dE,$$

where f , \mathcal{M} , and E are the true, mean, and eccentric anomalies.

- Thus, want to evaluate something of form

$$\tau_{\text{tot}} = \hat{T}(r_c, \Omega) \sum_{N=-\infty}^{\infty} F_{N2}^2(e) \operatorname{sgn}\left(N - 2\frac{\Omega_s}{\Omega}\right) \left|N - 2\frac{\Omega_s}{\Omega}\right|^{8/3}.$$

- The F_{N2} look like (note: FT is fastest to compute coefficients, below took ~ 2 s)



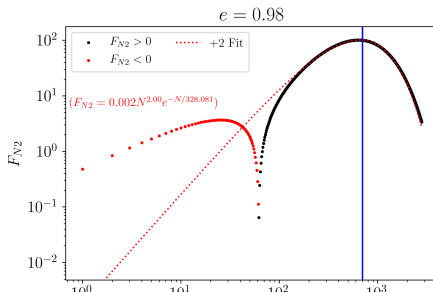
- Key insight: only one important hump ($\sim N_{\text{peri}}$), seek inspired fit.

$$\tau_{\text{tot}} = \hat{T}(r_c, \Omega) \sum_{N=-\infty}^{\infty} F_{N2}^2(e) \operatorname{sgn}\left(N - 2\frac{\Omega_s}{\Omega}\right) \left|N - 2\frac{\Omega_s}{\Omega}\right|^{8/3}.$$

- Criteria for approximate $F_{N2}(e)$:
 - Should only have one scale, N_{peri}
 - Should exponentially fall off for large N (smoothness)
 - $F_{02}(e) \approx 0$.
- Guess: maybe

$$F_{N2}(e) \simeq C(e) N^p(e) e^{-N/\eta(e)}. \quad (1)$$

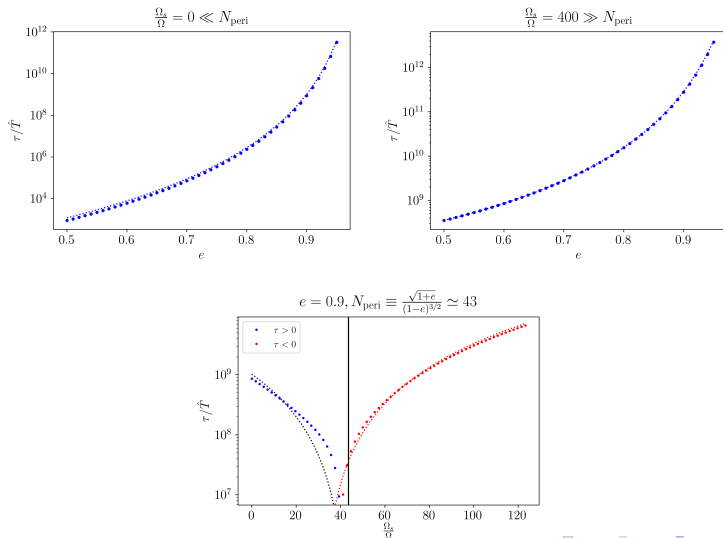
Turns out $p \approx 2$. Furthermore, $\operatorname{argmax}_N F_{N2}(e) = p\eta(e)$, so $\eta \simeq N_{\text{peri}}/2$. C is fixed by normalization (Parseval's).



Solution

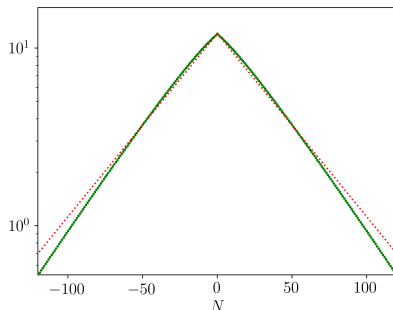
Tidal Torque

- Use resulting τ_{tot} in closed form (piecewise for $\Omega_s \gg N_{\text{peri}}\Omega/2$ or $\Omega_s \ll N_{\text{peri}}\Omega/2$), some small fudge factors, compare with explicit sum:



- For heating, same treatment for F_{N2} , need $m = 0$ Hansen coefficients F_{N0} too:

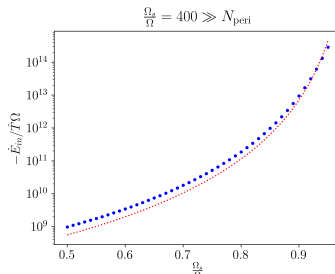
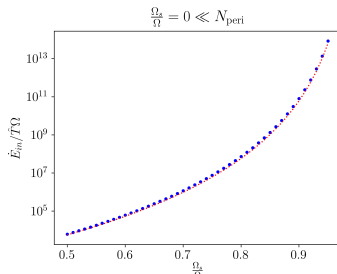
$$\dot{E}_{\text{in}} = \frac{1}{2} \hat{T}(r_c, \Omega) \sum_{N=-\infty}^{\infty} \left[N \Omega F_{N2}^2 \text{sgn}(\sigma) |\sigma|^{8/3} + \left(\frac{W_{20}}{W_{22}} \right)^2 \Omega F_{N0}^2 |N|^{11/3} \right]. \quad (2)$$



- Only characteristic scale is still N_{peri} .

- Guess $F_{N0} \simeq Ae^{-\frac{|N|}{N_{\text{peri}}}}$. Empirically, find

$$F_{N0} \approx Ae^{-\frac{|N|}{N_{\text{peri}}/\sqrt{2}}}, \quad (3)$$



- Application to J0045+7319?
 - Since e is known, τ_{tot} is very easy to evaluate without approximation.
- Tidal synchronization timescale as a function of e ?