

Vanishing adsorption admissibility criterion for contact discontinuities in the polymer model

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H + M = Lopes²



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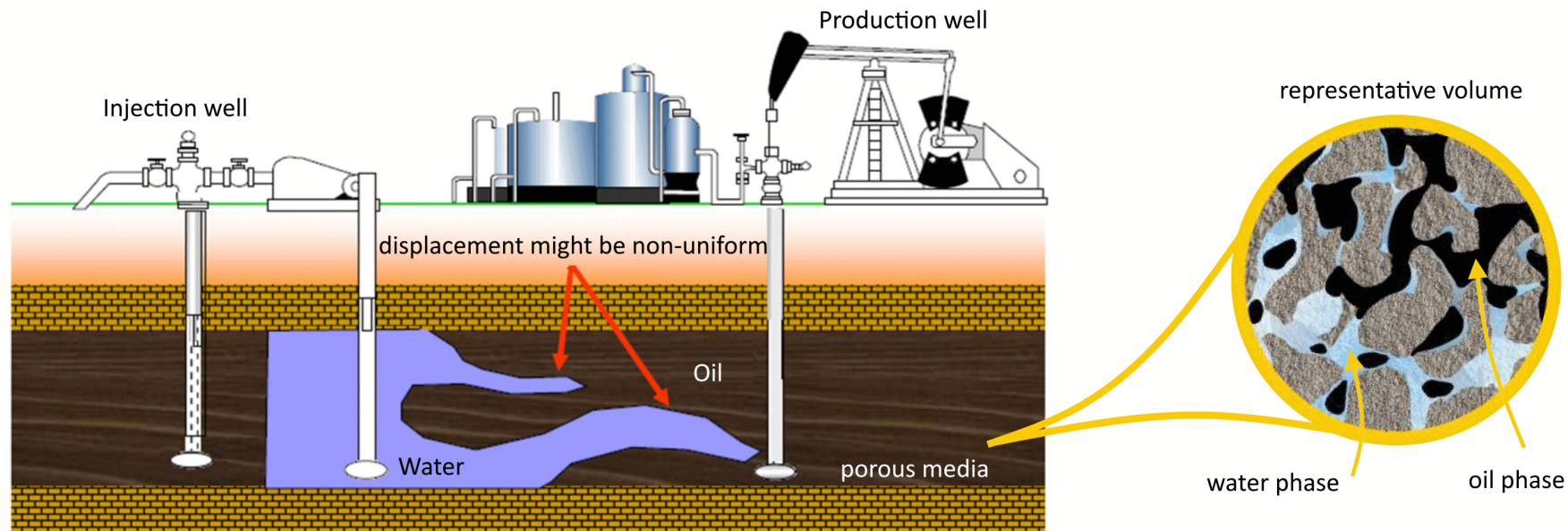
The talk is based on:

- **Y. Petrova**, B. Plohr, and Marchesin D. Vanishing adsorption admissibility criterion for contact discontinuities in the polymer model. arXiv:2211.10326.
- F. Bakharev, A. Enin, **Y. Petrova**, and N. Rastegaev. Impact of dissipation ratio on vanishing viscosity solutions of the Riemann problem for chemical flooding model. Journal of Hyperbolic Differential Equations, 20:1–26, 2023.

Motivation: enhanced oil recovery (EOR)

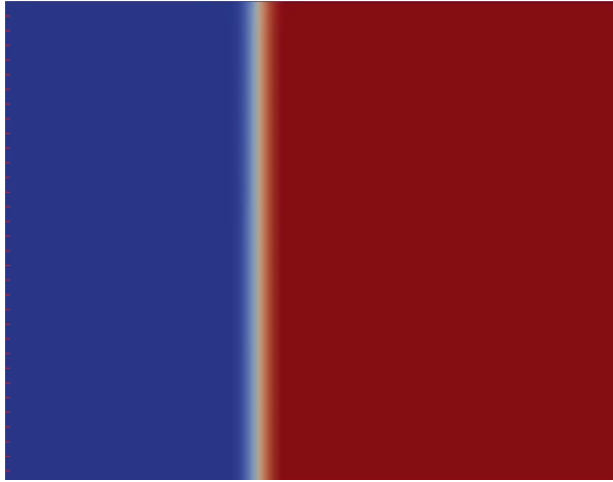
We are interested in the mathematical model of oil recovery.

- porous media (averaged models of flow)
- unknown variables: $s \in [0,1]$ - water saturation, $1 - s$ - oil concentration
- relatively small speeds (≈ 1 meter per day): Navier-Stokes \rightarrow Darcy's law
- multiphase flow: oil, water, gas
- applications to EOR methods: thermal, gas, chemical flooding



Two main directions of investigation

Stable displacement (1-dim)

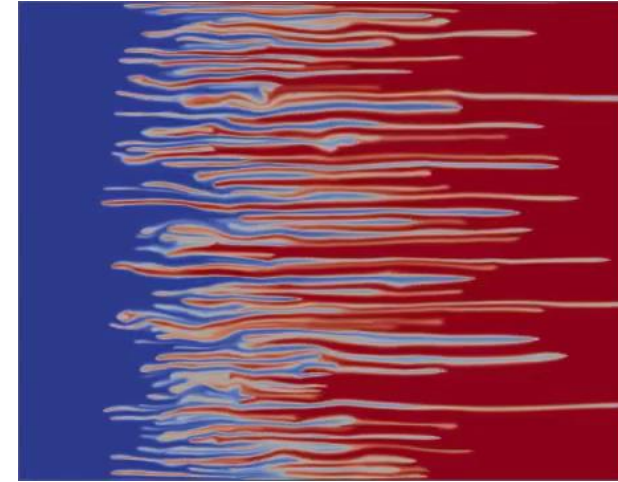


- hyperbolic conservation laws
- main question: find an exact solution for a Riemann problem

$$\begin{aligned} s_t + f(s, c)_x &= 0 \\ (sc)_t + (cf(s, c))_x &= 0 \end{aligned}$$

Example: polymer model

Unstable displacement (2-dim)



- viscous fingering phenomenon
- source of instability: water and oil/polymer have different viscosities

$$\begin{aligned} c_t + \operatorname{div}(uc) &= \Delta c \\ \operatorname{div}(u) &= 0, \quad u = -m(c) \nabla p \end{aligned}$$

Example: incompressible porous media equation (IPM)
[remember yesterday's talk of Sergey Tikhomirov] 2

Glimm-Isaacson model (KKIT)*

Two-phase oil-water flow with *polymer* in the water (1980)

$$\begin{aligned} s_t + f(s, c)_x &= 0 \\ (cs)_t + (cf(s, c))_x &= 0 \end{aligned}$$

- $s \in [0,1]$ – water saturation
- $c \in [0,1]$ – polymer concentration in water
- $f(s, c)$ – fractional flow function: affected by polymer
 - S-shaped in s (for fixed c)

Case 1:

$f'_c < 0$ (monotone in c)

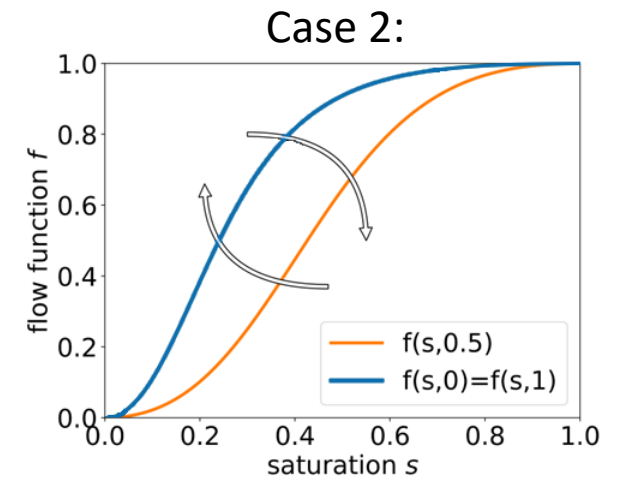
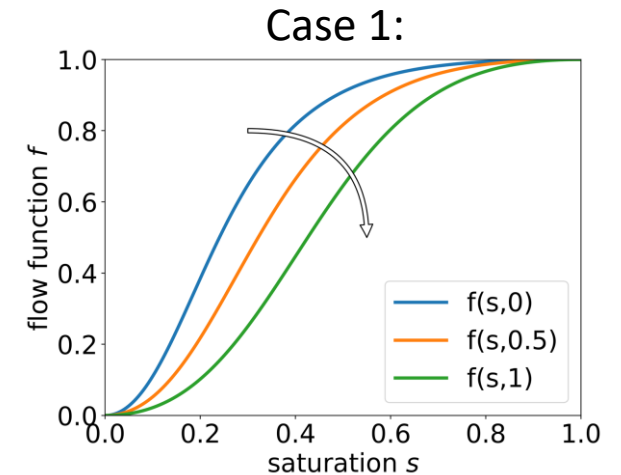
Case 2:

f changes monotonicity once

Initial data: $(s, c)(x, 0) = \begin{cases} (s_L, c_L), & x \leq 0 \\ (s_R, c_R), & x \geq 0 \end{cases}$

Question: find an exact solution $s(x, t)$ and $c(x, t)$ to any Riemann problem

NB: $s(x, t) = s(x/t)$ – self-similar



* KKIT = Keyfitz, Kranzer, Isaacson, Temple

Main idea

Polymer model (1980' E. Isaacson)

$$\begin{aligned} s_t + f(s, c)_x &= 0 \\ (sc)_t + (cf(s, c))_x &= 0 \end{aligned} \quad (M_0)$$

- Contact discontinuities (linearly degenerate) \Rightarrow non-uniqueness of solutions to a Riemann problem
- Vanishing viscosity criteria doesn't help
- Existing admissibility criteria needs to be justified from "physical" perspective

Aim: select unique physically admissible weak solution

Main idea: add small physical effect – adsorption of polymer on the rock (1987 ' T. Johansen, R. Winther)

$$\begin{aligned} s_t + f(s, c)_x &= 0 \\ (sc + \alpha a(c))_t + (cf(s, c))_x &= 0 \end{aligned} \quad (M_\alpha)$$

Here $\alpha > 0$ - small, $a(c)$ – strictly convex

- Contact discontinuities \rightsquigarrow rarefaction and shock waves \rightsquigarrow vanishing viscosity criteria is applicable

Vanishing adsorption criterion: the admissible contacts for M_0 are the L^1_{loc} limits of a family of admissible solutions for a Riemann problem for M_α as $\alpha \rightarrow 0$.

Hyperbolic systems of conservation laws*

$$G(U)_t + F(U)_x = 0$$

- $G(U) \in \mathbb{R}^n$ - accumulation function (conserved quantities)
- $F(U) \in \mathbb{R}^n$ - flux function (flux of conserved quantities)

$$U_t + F(U)_x = 0$$

$$U_t + DF(U) \cdot U_x = 0$$

Eigenvalues of $DF(U)$ play crucial role

Simplest example: wave equation

$$y_{tt} - c^2 y_{xx} = 0 \quad (\text{J. d'Alembert, 1750})$$

can be rewritten as a system of two first-order equations on the state-vector $u = \begin{pmatrix} y_x \\ y_t \end{pmatrix}$

$$u_t + Du_x = 0 \quad \text{with} \quad D = \begin{pmatrix} 0 & -1 \\ -c^2 & 0 \end{pmatrix}$$

- eigenvalues $\lambda_1 = c$ and $\lambda_2 = -c$ are **real**, the system is **hyperbolic (strictly)**.
Solution contains two **wave modes** that propagate at the **velocities** λ_1 and λ_2 .

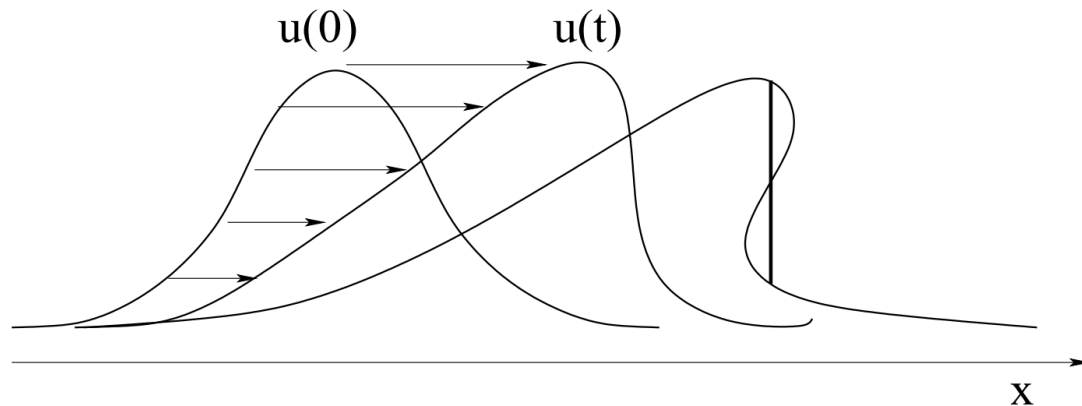
*A. Bressan "Hyperbolic conservation laws: an illustrated tutorial" , lecture notes from Cetraro, Italy 2009

Hyperbolic systems of conservation laws

- Inviscid Burgers equation (1948)

$$u_t + \left(\frac{u^2}{2}\right)_x = 0$$

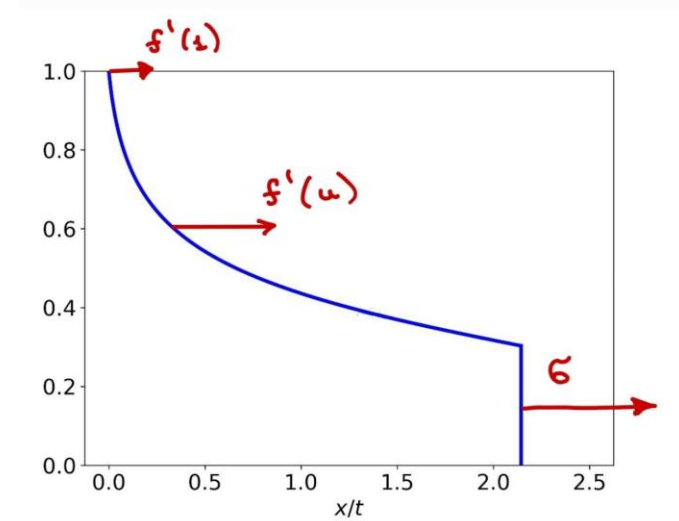
- non-linearity implies **wave speed** $\lambda(u) = u$ **depends on state** u
- so the wave can spread (**rarefaction wave**) or concentrate (**shock wave**)



- General scalar conservation law

$$u_t + f(u)_x = 0$$

- existence, uniqueness of solution of Cauchy problem was established by Olga Oleinik (1957)



Shock speed σ between the states u_- and u_+ is defined by the Rankine-Hugoniot condition $\sigma = \frac{f(u_-) - f(u_+)}{u_- - u_+}$

Riemann problem (1858)

Riemann solved the initial-value problem with data having a **single jump**

$$U(x, 0) = \begin{cases} U_L, & x \leq 0 \\ U_R, & x \geq 0 \end{cases}$$

Solution to a Riemann problem is important because:

- it appears in a long-term behaviour of Cauchy problem
- helps to prove the existence of solutions to Cauchy problem (Glimm's method)
- helps to construct numerical solution (Godunov method)

$$U_t + F(U)_x = 0$$

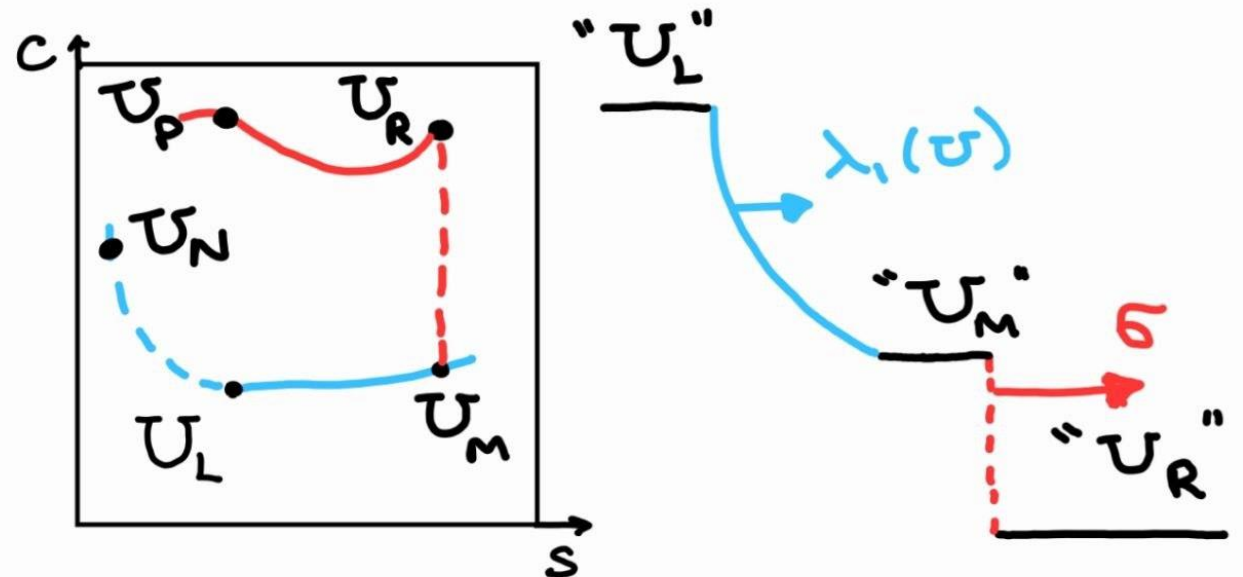
Representation of solution $U = (s, c)$ for 2×2 systems:

Eigenvalues of $DF(U)$: $\lambda_1(U) < \lambda_2(U)$

Take U_L and $\lambda_1(U)$

Take U_R and $\lambda_2(U)$

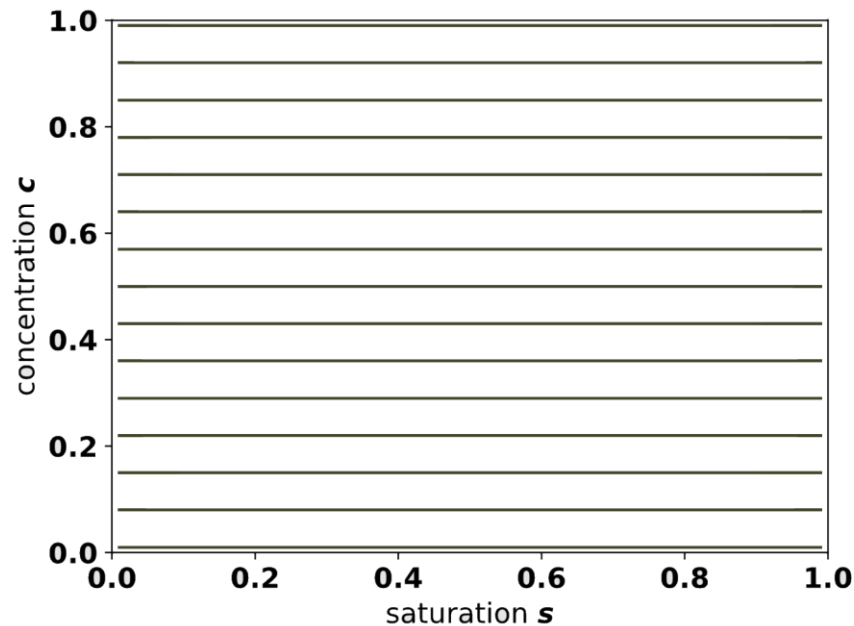
We have constructed a solution!



Characteristic families: s and c-waves

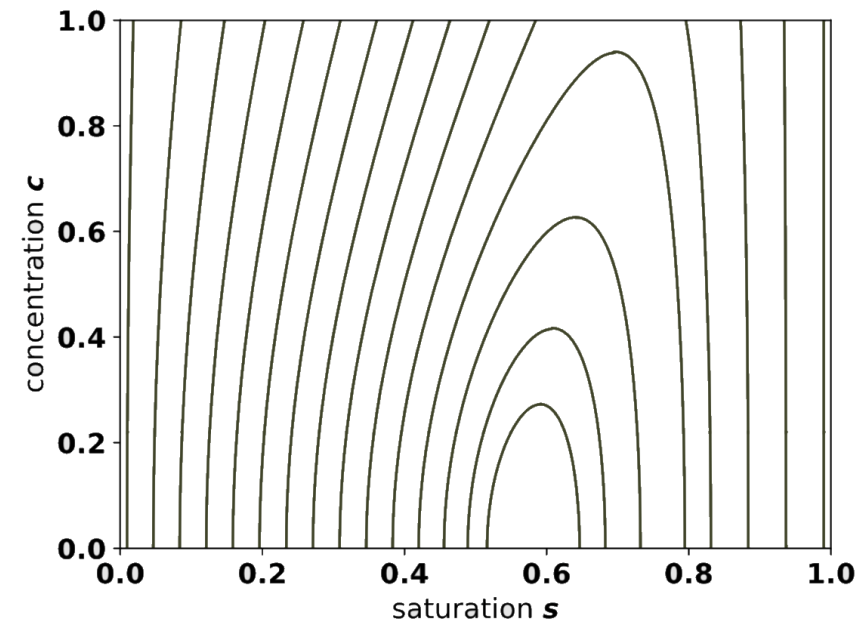
s-waves

- $\lambda^s = f'_s$
- Solve the Buckley-Leverett equation $c = \text{const}$
- Riemann invariant $c = \text{const}$
- “line” family



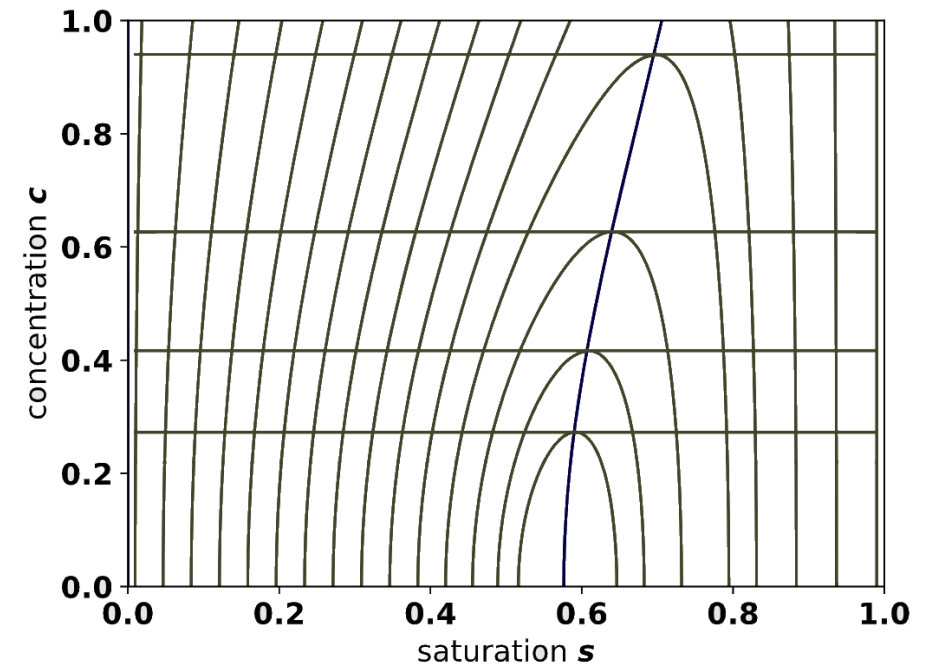
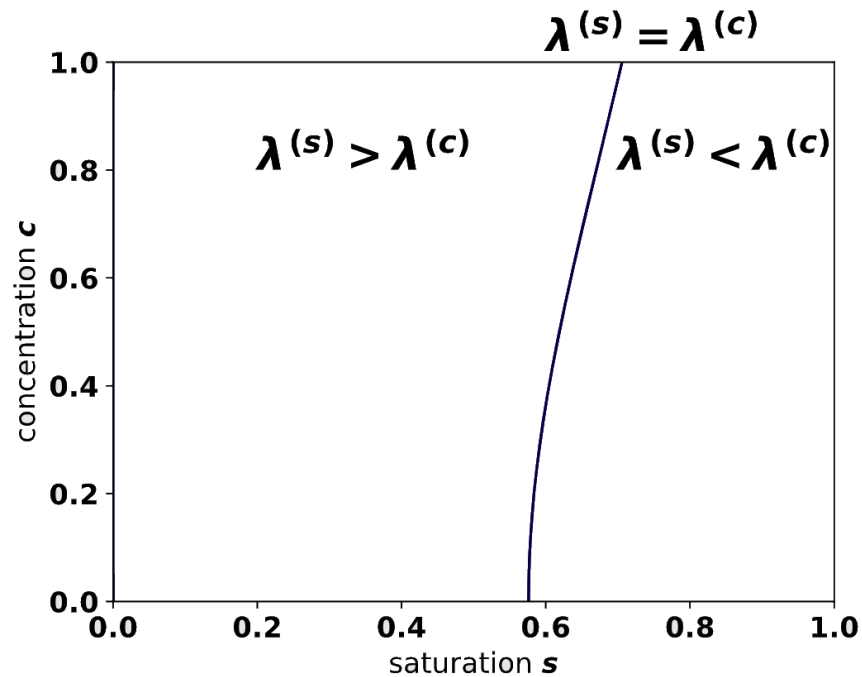
c-waves

- $\lambda^c = f/s$
- Are contact discontinuities
- Riemann invariant $f/s = \text{const}$
- “contact” family



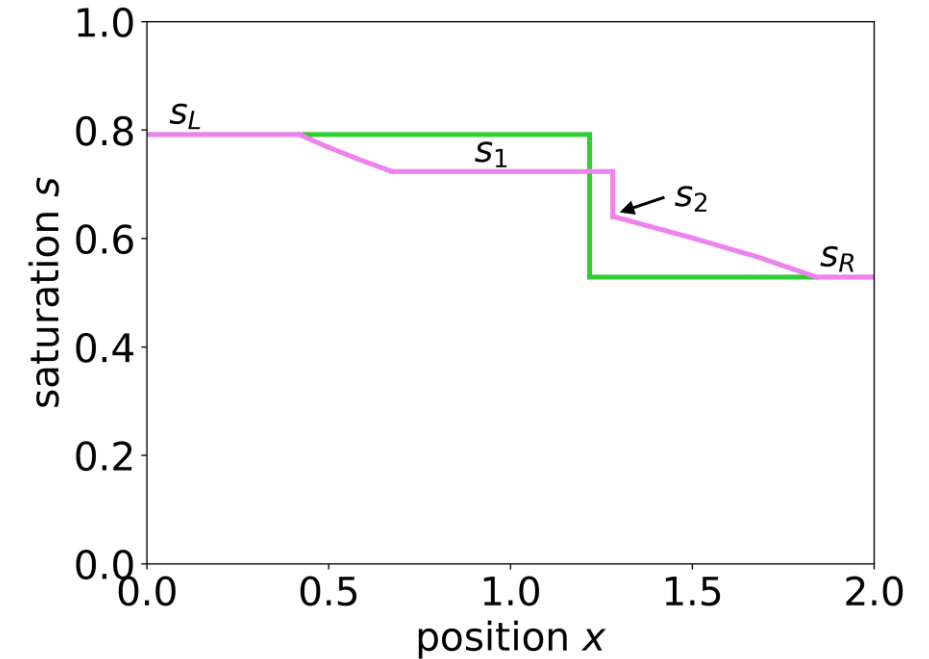
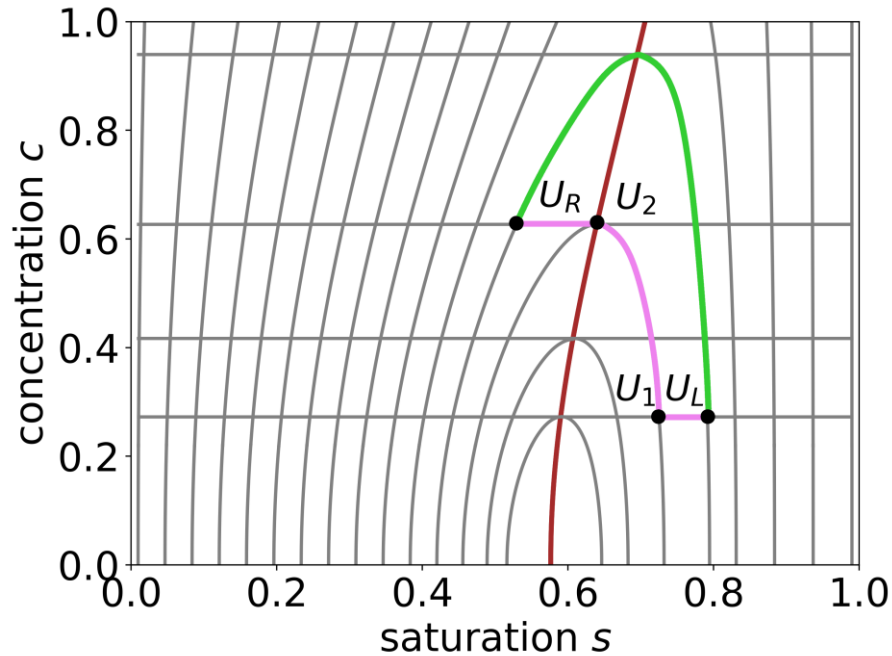
- For both families, the rarefaction and shock curves coincide! But in a different way (Temple'1983)
- Any solution for a Riemann problem is a combination of s and c waves

Non-strictly hyperbolic system



The coordinate system of wave curves is singular and the wave speeds coincide on a co-dimension one curve (coincidence locus): $\lambda^s = f'_s = f/s = \lambda^c$
 s and c waves are tangent on coincidence curve

Non-uniqueness of solutions



A contact discontinuity between U_- and U_+ is admissible if and only if:

Criterion 1 (Isaacson): either $U_-, U_+ \in \{\lambda^s \geq \lambda^c\}$ or $U_-, U_+ \in \{\lambda^s \leq \lambda^c\}$

Criterion 2 (de Souza-Marchesin): c is continuous and monotone along the sequence of contact curves, connecting U_- and U_+

What is the (physical) motivation of these criteria?

Main result

Polymer model

$$\begin{aligned} s_t + f(s, c)_x &= 0 \\ (sc + \alpha a(c))_t + (cf(s, c))_x &= 0 \end{aligned} \quad (M_\alpha)$$

Criterion 3: vanishing adsorption (Petrova-Marchesin-Plohr):

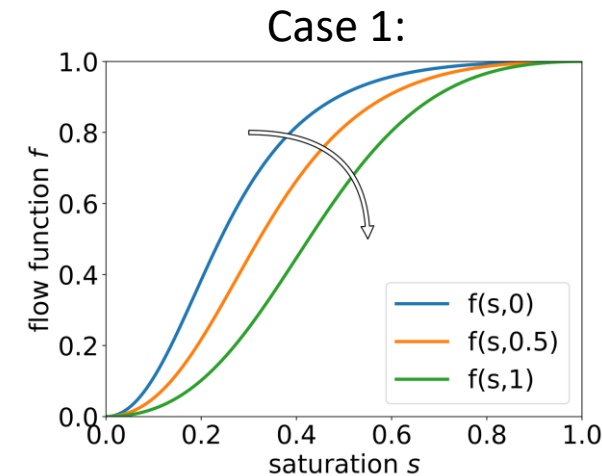
A contact discontinuity between U_- and U_+ for M_0 is admissible if and only if it is the L^1_{loc} limit of a family of admissible solutions for a Riemann problem for M_α as $\alpha \rightarrow 0$.

Theorem 1 (P., Marchesin, Plohr, 2022, arxiv:2211.10326)

If f satisfies the monotonicity assumption $f'_c < 0$, then the set of admissible Riemann solutions for M_0 is the same for criteria 1, 2 and 3.

Corollary: any solution to a Riemann problem for M_0 exists and is unique.

Question: what happens when f is non-monotone in c ?



What happens when f is non-monotone in c ?

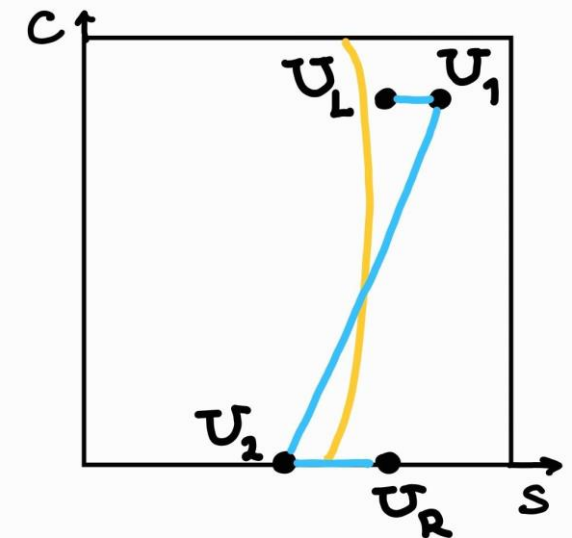
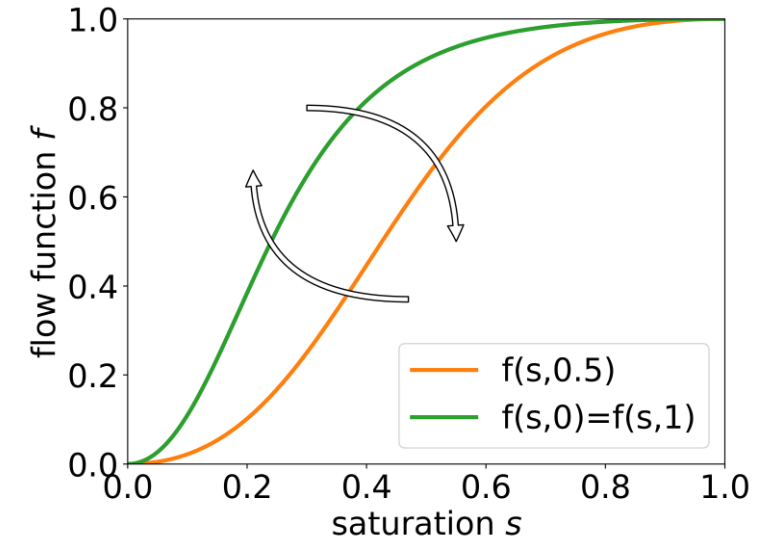
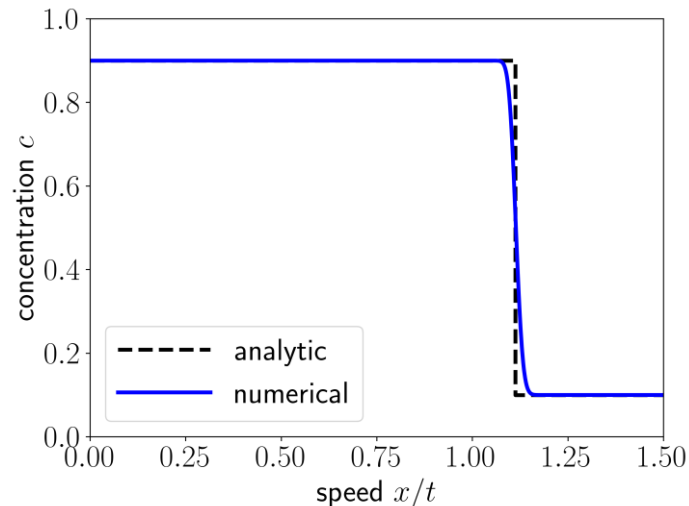
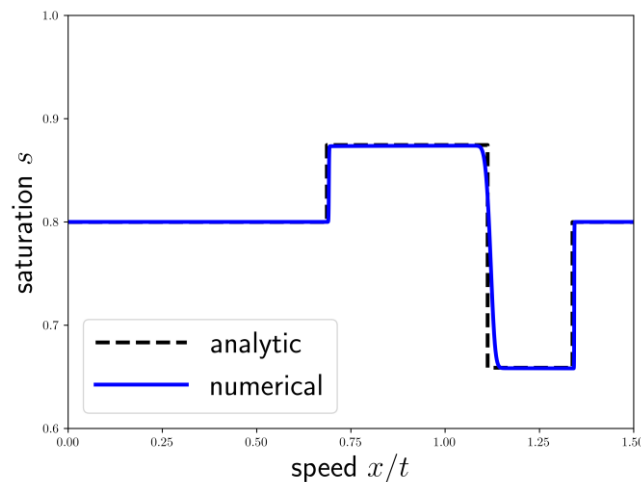
$$\begin{aligned} s_t + f(s, c)_x &= 0 \\ (sc)_t + (cf(s, c))_x &= 0 \end{aligned}$$

Example: “boomerang”

$$f(s, c) = \frac{s^2}{s^2 + \mu(c)(1-s)^2} \quad \text{with} \quad \mu(c) = 1 + 4c(1-c)$$

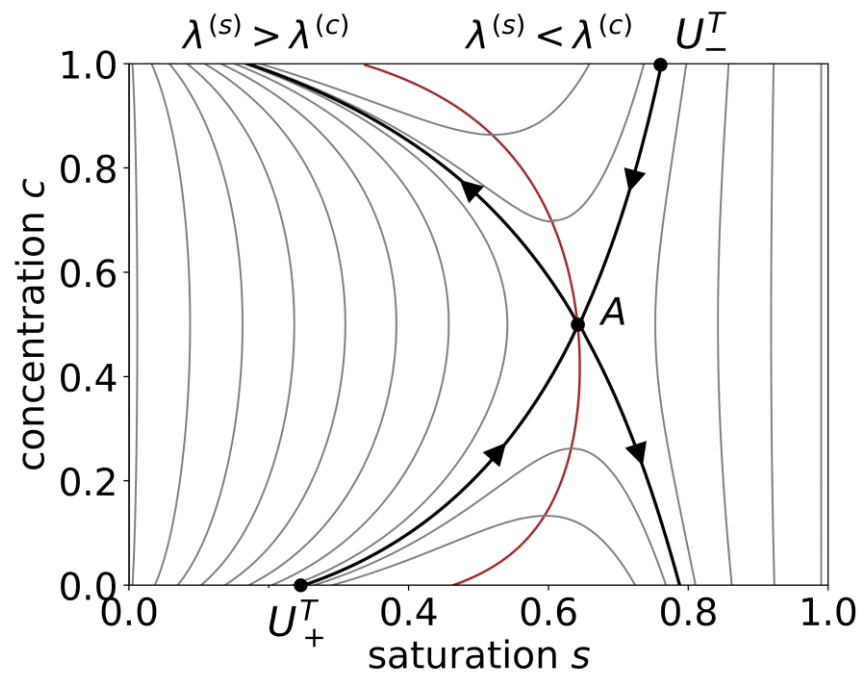
Riemann problem: $s_L = s_R = 0.82$, $c_L = 1$, $c_R = 0$

Results of numerical modelling:



...a more careful look...

$$\begin{aligned} s_t + f(s, c)_x &= 0 \\ (sc)_t + (cf(s, c))_x &= 0 \end{aligned}$$



Let's call the corresponding contact discontinuity undercompressive.

Is this contact discontinuity admissible by the existing criteria?

Criteria 1: NO

Criteria 2: YES

Criteria 3: YES

Theorem 2 (P., Marchesin, Plohr, 2022, arxiv:2211.10326)

The undercompressive contact discontinuities satisfy the vanishing adsorption admissibility criterion

Main step in proof of Thm 2

$$\begin{aligned} s_t + f(s, c)_x &= 0 \\ (sc + \alpha a(c))_t + (cf(s, c))_x &= 0 \end{aligned}$$

Add diffusion terms



$$\begin{aligned} s_t + f(s, c)_x &= \varepsilon_1 s_{xx} \\ (sc + \alpha a(c))_t + (cf(s, c))_x &= \varepsilon_1 (cs_x)_x + \varepsilon_2 c_{xx} \end{aligned}$$

$$\varepsilon_1, \varepsilon_2 \rightarrow 0, \quad k = \frac{\varepsilon_1}{\varepsilon_2}$$

Travelling wave ansatz



$$s = s(x - \sigma t) = s(\xi), \quad c = c(x - \sigma t) = c(\xi)$$

$$\begin{cases} \alpha \cdot s_\xi = f - \sigma(s + d_1) \\ c_\xi = \frac{\varepsilon_1}{\varepsilon_2} \cdot \sigma \cdot (d_1 c - d_2 - a(c)) \end{cases}$$

Theorem 3 (Bakharev, Enin, P., Rastegaev, 2023, JHDE)

For any $k = \frac{\varepsilon_1}{\varepsilon_2} > 0$ there exists $s_-(k)$ and $s_+(k)$ and velocity $\sigma(k)$ such that there exists a travelling wave, connecting two saddle points $(s_-(k), 1)$ and $(s_+(k), 0)$ with velocity $\sigma(k)$.

Own works:

1. F. Bakharev, A. Enin, Y. Petrova, and N. Rastegaev. Impact of dissipation ratio on vanishing viscosity solutions of the Riemann problem for chemical flooding model. *Journal of Hyperbolic Differential Equations*, 20:1–26, 2023.
2. Y. Petrova, B. Plohr, and Marchesin D. Vanishing adsorption admissibility criterion for contact discontinuities in the polymer model. [arXiv:2211.10326](https://arxiv.org/abs/2211.10326).

Other references:

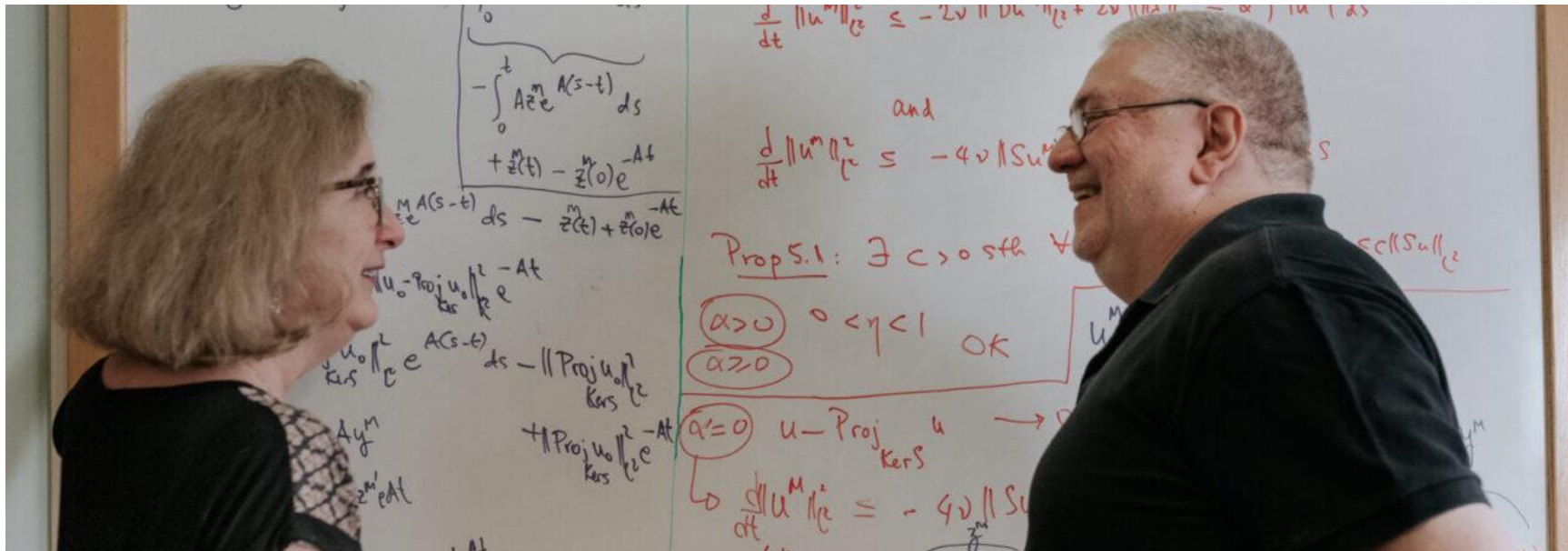
Polymer model:

1. W. Shen. On the Cauchy problems for polymer flooding with gravitation. *Journal of Differential Equations*, 261(1):627–653, 2016.
2. W. Shen. On the uniqueness of vanishing viscosity solutions for Riemann problems for polymer flooding. *Nonlinear Differential Equations and Applications NoDEA*, 24(4):37, 2017.
3. B. Temple. Global solution of the Cauchy problem for a class of 2×2 nonstrictly hyperbolic conservation laws. *Advances in Applied Mathematics*, 3(3):335–375, 1982.
4. T. Johansen and R. Winther. The solution of the Riemann problem for a hyperbolic system of conservation laws modeling polymer flooding. *SIAM Journal on Mathematical Analysis*, 19(3):541–566, 1988.
5. Keyfitz, B.L. and Kranzer, H.C., 1980. A system of non-strictly hyperbolic conservation laws arising in elasticity theory. *Archive for Rational Mechanics and Analysis*, 72(3), pp.219-241.
6. E. L. Isaacson, Global solution of a Riemann problem for a non-strictly hyperbolic system of conservation laws arising in enhanced oil recovery, Rockefeller University, New York, NY, preprint (1981).

Feliz aniversario, Helena e Milton!!!

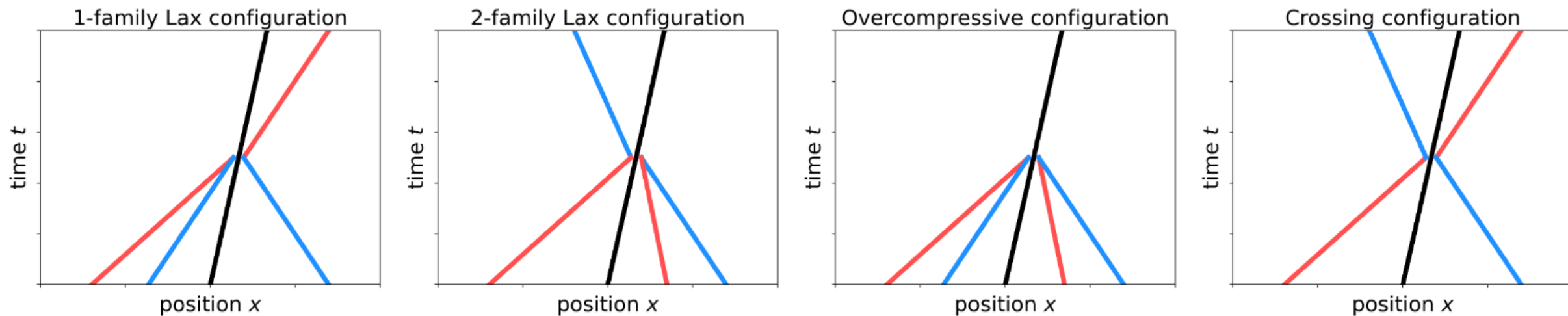
Parabéns pra vocês
Nesta data querida
Muitas felicidades
Muitos anos de vida

É pique, é pique
É pique, é pique, é pique
É hora, é hora
É hora, é hora, é hora
Ra-tim-bum!



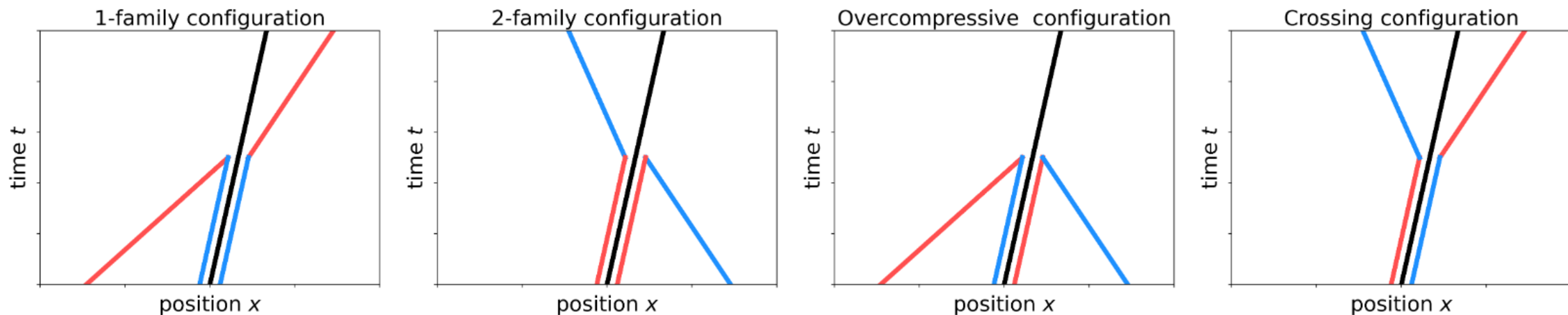
$$H + M = \text{Lopes}^2$$

Types of shocks



- 1-family Lax: $\lambda_1(U_-) > \sigma > \lambda_1(U_+)$ along with $\sigma < \lambda_2(U_-)$ and $\sigma < \lambda_2(U_+)$
- 2-family Lax: $\lambda_2(U_-) > \sigma > \lambda_2(U_+)$ along with $\sigma > \lambda_1(U_-)$ and $\sigma > \lambda_1(U_+)$
- overcompressive: $\lambda_1(U_-) > \sigma > \lambda_1(U_+)$ and $\lambda_2(U_-) > \sigma > \lambda_2(U_+)$
- crossing: $\lambda_2(U_-) > \sigma > \lambda_1(U_-)$ and $\lambda_1(U_+) < \sigma < \lambda_2(U_+)$.

Types of contact discontinuities



- 1-family: $\lambda_1(U_-) = \sigma = \lambda_1(U_+)$ along with $\sigma < \lambda_2(U_-)$ and $\sigma < \lambda_2(U_+)$;
- 2-family: $\lambda_2(U_-) = \sigma = \lambda_2(U_+)$ along with $\sigma > \lambda_1(U_-)$ and $\sigma > \lambda_1(U_+)$;
- overcompressive: $\lambda_1(U_-) = \sigma > \lambda_1(U_+)$ and $\lambda_2(U_-) = \sigma > \lambda_2(U_+)$;
- crossing: $\lambda_2(U_-) > \sigma = \lambda_1(U_-)$ and $\lambda_1(U_+) = \sigma < \lambda_2(U_+)$.