

A cascade of two travelling waves for the two-layer model of “gravitational fingering”



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Based on joint work with Sergey Tikhomirov, Yalchin Efendiev :

“Propagating terrace in a two-tubes model of gravitational fingering”, 2024, ArXiv: 2401.05981

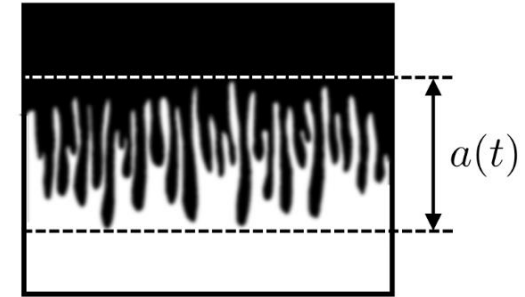


Outline

1. Motivation

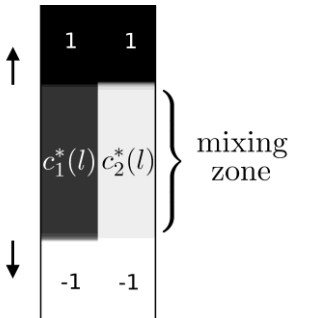
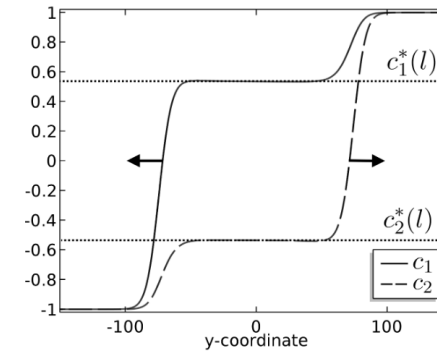
Miscible displacement in porous media

- viscous fingering
- gravitational fingering



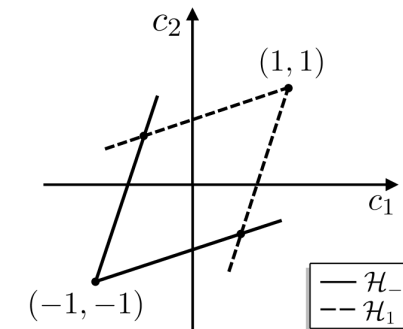
2. Problem statement

- Two-tubes model
- Main theorem



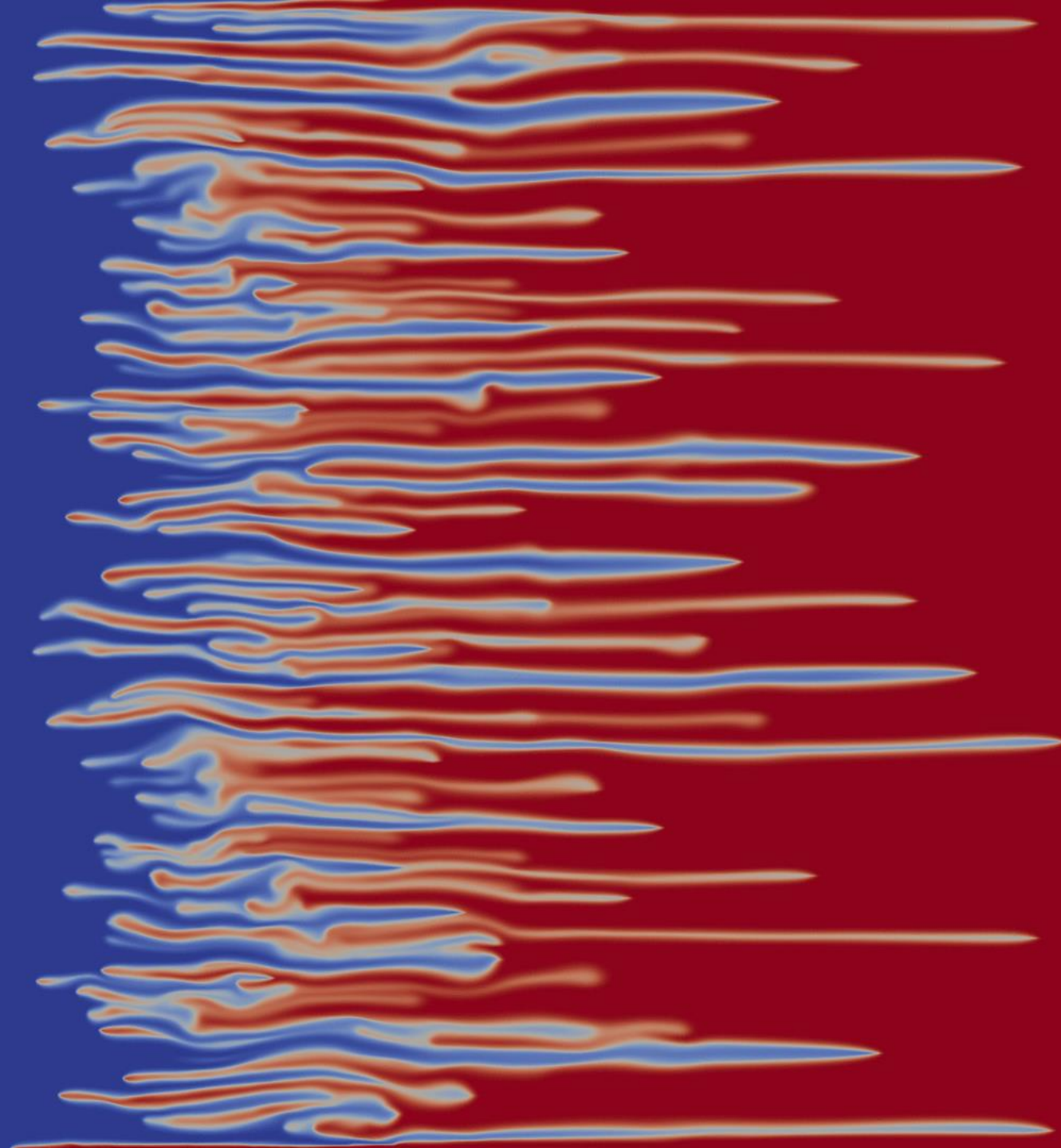
3. Sketch of proof:

- traveling waves
- slow-fast systems



"Miscible displacement in porous media"
Credit: Pavlov Dmitrii, St. Petersburg State University

Homsy , 1987 "Viscous Fingering in Porous Media"



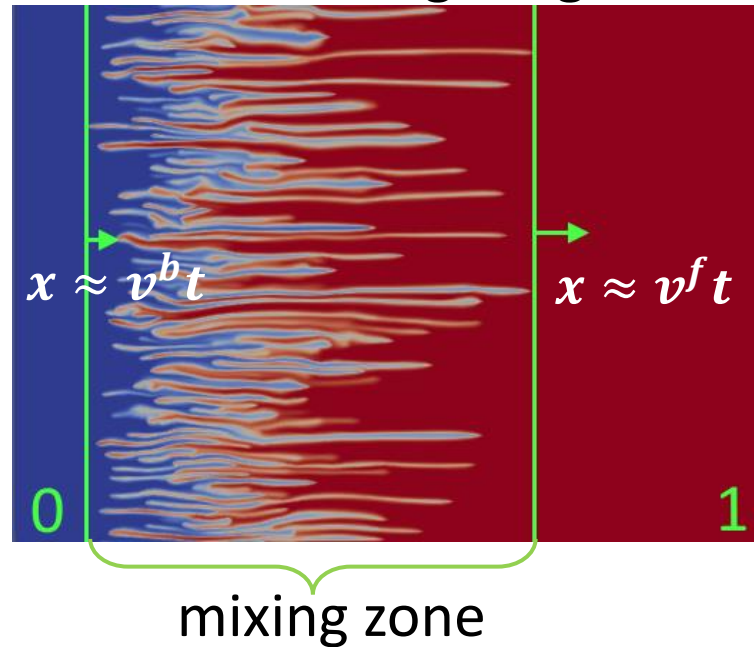
Viscous fingering phenomenon

water (blue color)

polymerized water (red color) 1

Incompressible Porous Medium eq – IPM, 2D (Two formulations)

Viscous fingering



$$c_t + \operatorname{div}(uc) = \varepsilon \cdot \Delta c$$

$$\operatorname{div}(u) = 0$$

(viscosity)

$$u = -m(c) K \nabla p$$

(gravity)

$$u = -\nabla p - (0, c)$$

$c = c(t, x, y)$ – concentration

$u = u(t, x, y)$ – velocity

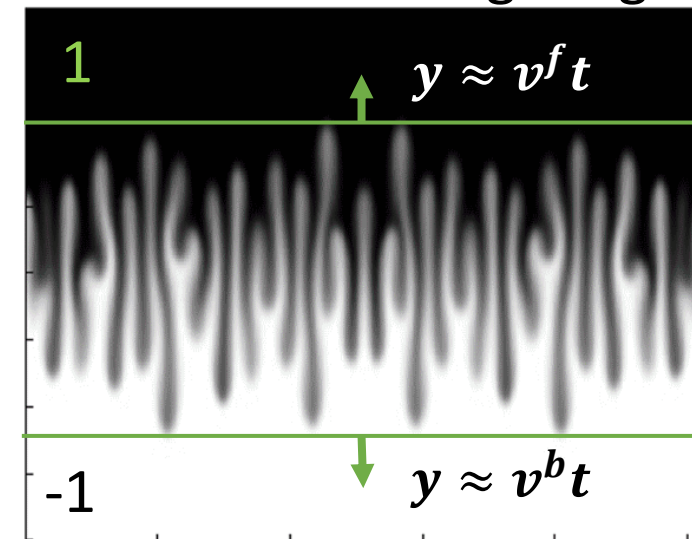
$p = p(t, x, y)$ – pressure

$\varepsilon \geq 0$ – diffusion

$m(c)$ – mobility

K – permeability

Gravitational fingering

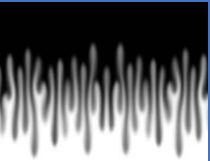


- many laboratory and numerical experiments show *linear growth of the mixing zone* ^{[1], [2]}

Question: how to find speeds v^b and v^f of propagation?

[1] Nijjer J., Hewitt D., and Neufeld J. The dynamics of miscible viscous fingering from onset to shutdown. *Journal of Fluid Mechanics*, 2018.

[2] Bakharev, F., Enin, A., Groman, A., Kalyuzhnyuk, A., Matveenko, S., **Petrova, Y.**, Starkov, I. and Tikhomirov, S., Velocity of viscous fingers in miscible displacement: Comparison with analytical models. *Journal of Computational and Applied Mathematics*, 2022.



IPM: $\varepsilon = 0$ (without diffusion)

Active scalar:

$$\begin{aligned} c_t + u \cdot \nabla c &= 0 \\ u &= A(c) \end{aligned}$$

$$u = \nabla^\perp (-\Delta)^{-1} \partial_1 c \quad (\text{Biot-Savart law})$$

Discontinuous initial data: free boundary problem (Muskat problem) – ill-posed for unstable stratification

2011 - A. Córdoba, D. Córdoba, F. Gancedo (Annals of Mathematics)

“Interface evolution: the Hele-Shaw and Muskat problems”

Existence: smooth initial data

2007 – D. Cordoba, F. Gancedo, R. Orive (JMP): local well-posedness for initial data H^s

global solution vs finite-time blow-up? open

2017 – T. Elgindi (ARMA): global solution for small perturbations of $c = -y$

2023 – S. Kiselev, Y. Yao (ARMA): if solutions stay “smooth” for all times, then there is blow-up at $t = +\infty$

Uniqueness: non-uniqueness of weak solutions – by convex integration

2011 – D. Córdoba, D. Faraco, F. Gancedo (ARMA)

2012 – L. Szekelyhidi Jr.

...and many others...

IPM: $\varepsilon > 0$ (with diffusion)

Question: Are those estimates sharp?

Estimates on the growth:

2005 – F. Otto, G. Menon. Proved estimates

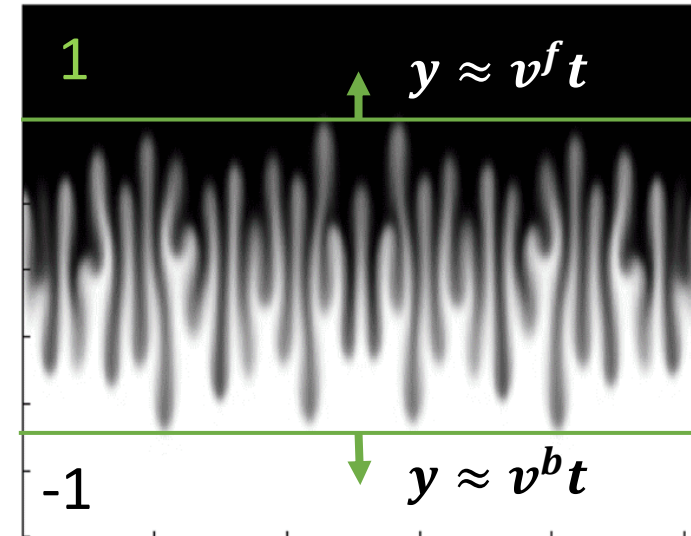
- Full model (IPM) $v^f \leq 2$
- Simplified model (TFE) $v^f \leq 1$

(gravity) Bofetta: $v^f \approx 0.67$
(viscosity) empirical models
numerics: NOT sharp

Transverse Flow Equilibrium = TFE

$$p(t, x, y) \approx p(t, y)$$

$$\begin{aligned} c_t + u \cdot \nabla c &= \varepsilon \Delta c \\ \operatorname{div}(u) &= 0 \\ u &= (u^1, u^2), \quad u^2 = \bar{c} - c \end{aligned}$$

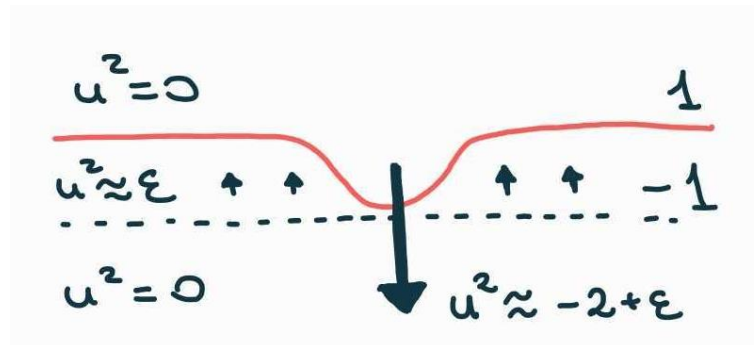


Why fingers appear?

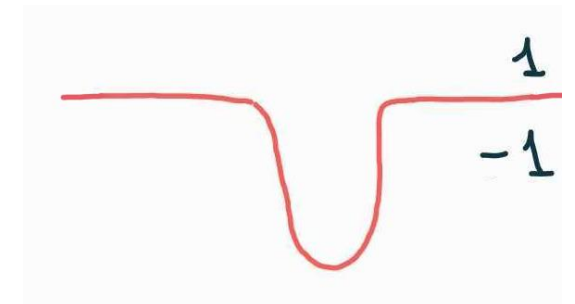
It is a hair-trigger effect!

$$\begin{array}{c} u^2 = 0 \\ \hline u^2 = 0 \end{array} \quad \begin{array}{c} 1 \\ -1 \end{array}$$

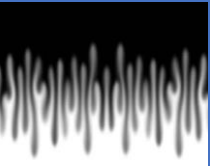
No flow



Velocity u changes
due to concentration c

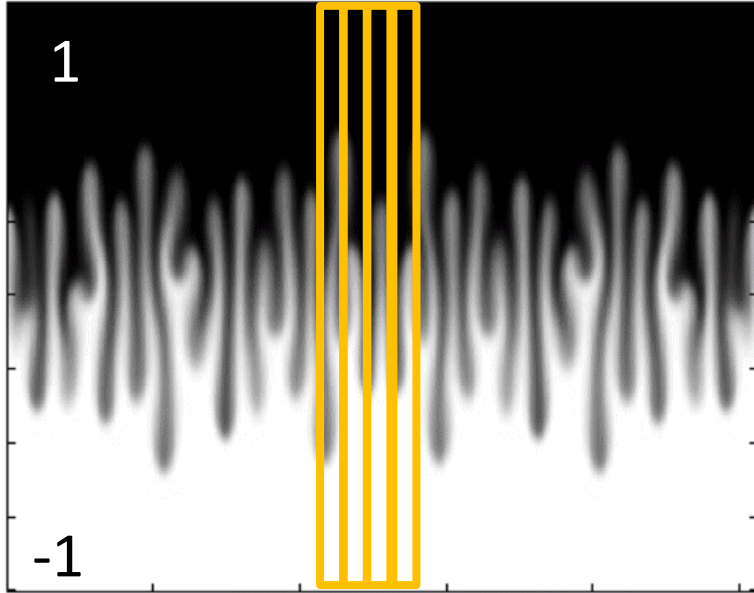


Concentration c changes
due to velocity u



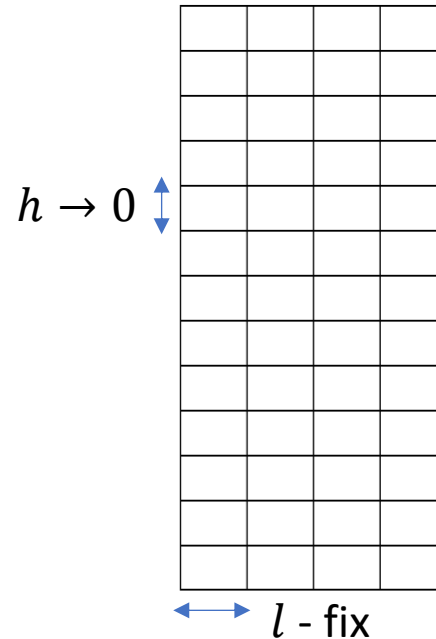
IDEA: semi-discrete model of gravitational fingering

- Discretize in horizontal direction
- Take n tubes, $n = 2, 3, 4, \dots$



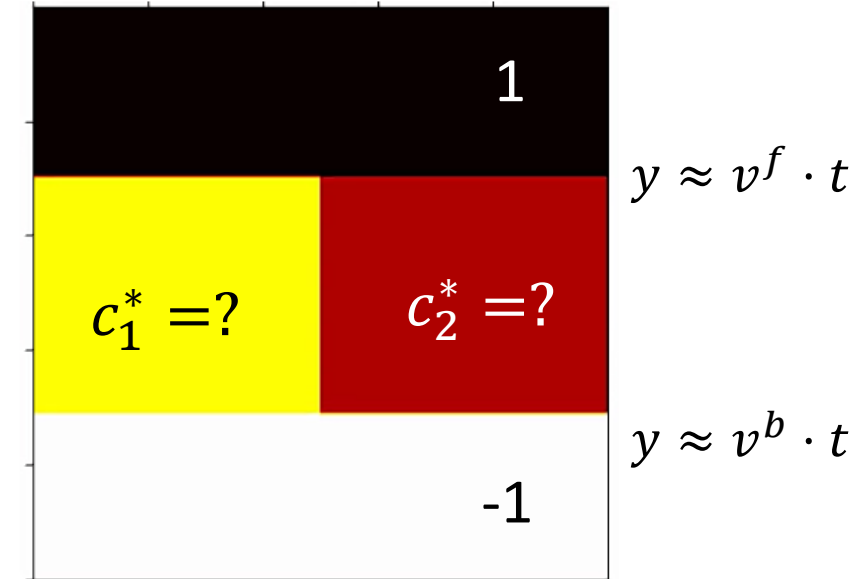
Tubes (layer, lane,...) models:

Limit of
numerical scheme



- Finite volume
- Upwind

- For simplicity, $n = 2$

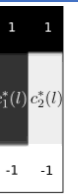


We observe two traveling waves:

$$c(y, t) = c(y - vt)$$

- 1995 — Y. Yortsos “A theoretical analysis of vertical flow equilibrium”
2019 — A. Armiti-Juber, C. Rohde “On Darcy- and Brinkman-type models for two-phase flow in asympt. flat domains”
2006 — J.C. Da Mota, S. Schechter “Combustion fronts in a porous medium with two layers”
2019 — H. Holden, N. Risebro “Models for dense multilane vehicular traffic”

Two-tubes model



1. Original equation on c :
Two-tubes equations on c :

$$c_t + \operatorname{div}(uc) - \Delta c = 0$$

$$\begin{aligned} \partial_t c_1 + \partial_y(u_1 c_1) - \partial_{yy} c_1 &= -B \\ \partial_t c_2 + \partial_y(u_2 c_2) - \partial_{yy} c_2 &= +B \end{aligned}$$

2. Original equation on p :
Two-tubes equations on p :

$$u = -\nabla p - (0, c)$$

$$u_1 = -\partial_y p_1 - c_1$$

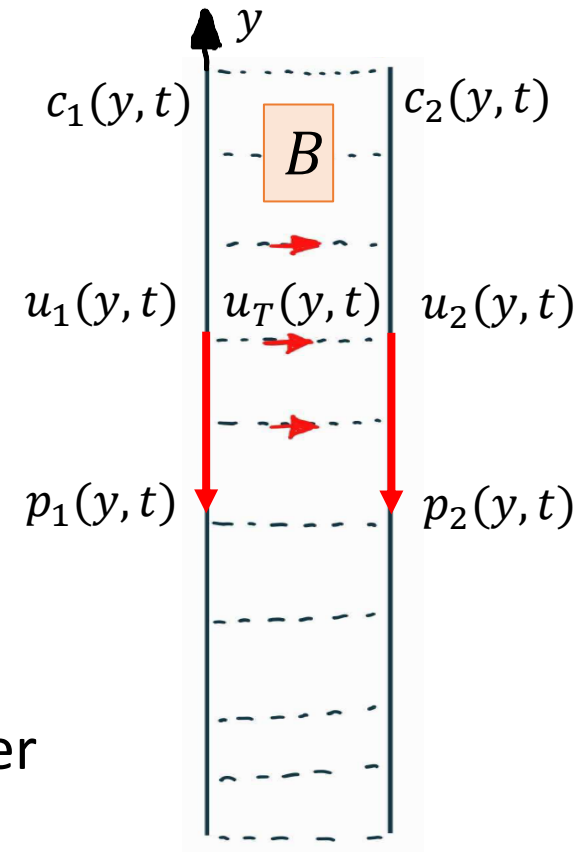
$$u_2 = -\partial_y p_2 - c_2$$

$$u_T = -\frac{p_2 - p_1}{l}$$

3. Original equation on u :
Two-tubes equations on u :

$$\operatorname{div}(u) = 0$$

$$\partial_y u_1 + \frac{u_T}{l} = 0$$



l - parameter

$$B = \begin{cases} \frac{u_T}{l} \cdot c_1, & u_T > 0, \\ \frac{u_T}{l} \cdot c_2, & u_T < 0 \end{cases}$$

Two-tubes model

$$\begin{matrix} 1 & 1 \\ c_1^*(l) & c_2^*(l) \\ -1 & -1 \end{matrix}$$

1. Original equation on c :
Two-tubes equations on c :

$$c_t + \operatorname{div}(uc) - \Delta c = 0$$

$$\begin{aligned} \partial_t c_1 + \partial_y(u_1 c_1) - \partial_{yy} c_1 &= -B \\ \partial_t c_2 + \partial_y(u_2 c_2) - \partial_{yy} c_2 &= +B \end{aligned}$$

2. Original equation on p :
Two-tubes equations on p :

$$u = -\nabla p - (0, c)$$

$$u_1 = -\partial_y p_1 - c_1$$

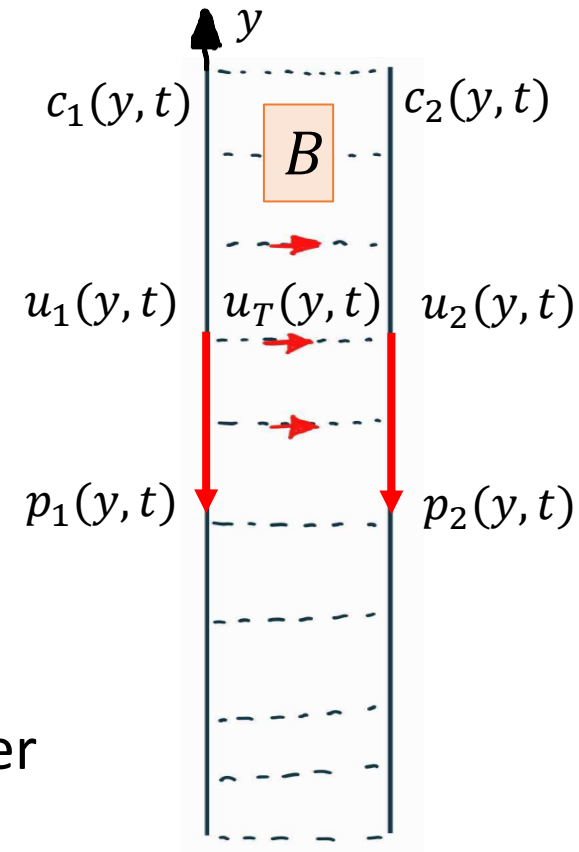
$$u_2 = -\partial_y p_2 - c_2$$

$$\boxed{\frac{u_T}{l}} = -\frac{p_2 - p_1}{l^2}$$

3. Original equation on u :
Two-tubes equations on u :

$$\operatorname{div}(u) = 0$$

$$\partial_y u_1 + \boxed{\frac{u_T}{l}} = 0$$



l - parameter

$$\boxed{B} = \begin{cases} \boxed{\frac{u_T}{l}} \cdot c_1, & u_T > 0, \\ \boxed{\frac{u_T}{l}} \cdot c_2, & u_T < 0 \end{cases}$$



Main result

Questions?

(*)
$$\begin{cases} \partial_t c_1 + \partial_y(u_1 c_1) - \partial_{yy} c_1 = -B \\ \partial_t c_2 + \partial_y(u_2 c_2) - \partial_{yy} c_2 = B \\ u_1 = -\partial_y p_1 - c_1 \\ u_2 = -\partial_y p_2 - c_2 \\ \partial_y u_1 = -\partial_y u_2 = \frac{p_2 - p_1}{l^2} \end{cases}$$

$$B = \begin{cases} -\partial_y u_1 \cdot c_1, & \partial_y u_1 < 0, \\ +\partial_y u_2 \cdot c_2, & \partial_y u_1 > 0 \end{cases}$$

Remark: $\lim_{l \rightarrow 0} c_1^*(l) = -0.5$ $\lim_{l \rightarrow 0} v^b(l) = -0.25$
 $\lim_{l \rightarrow 0} c_2^*(l) = +0.5$ $\lim_{l \rightarrow 0} v^f(l) = +0.25$

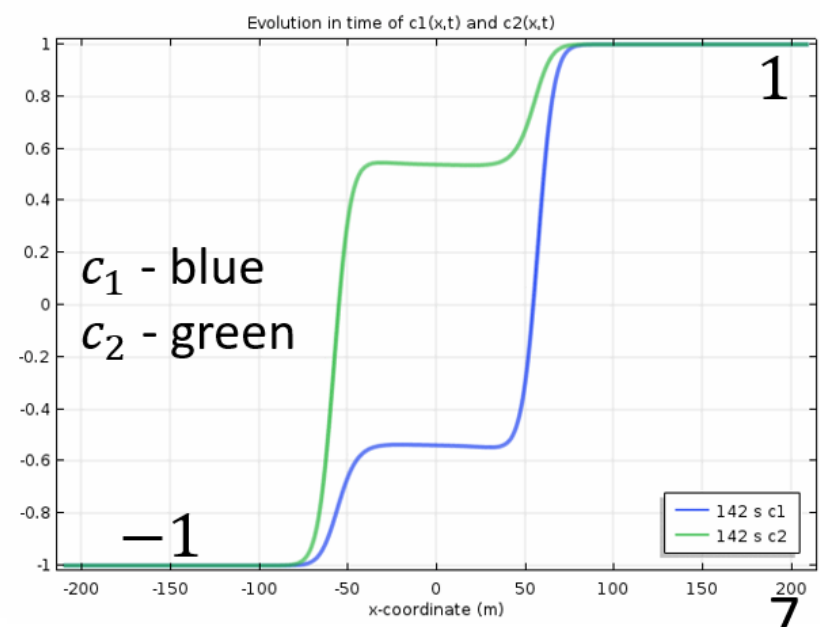
Theorem (Efendiev, P., Tikhomirov, 2024, arXiv: 2401.05981)

Consider a two-tube model with gravity (*).

Then for all $l > 0$ *sufficiently small* there exists $c_1^*(l), c_2^*(l)$ such that there exist two traveling waves (TW):

TW1 with speed $v^b(l)$: $(-1, -1) \rightarrow (c_1^*(l), c_2^*(l))$
TW2 with speed $v^f(l)$: $(c_1^*(l), c_2^*(l)) \rightarrow (1, 1)$.

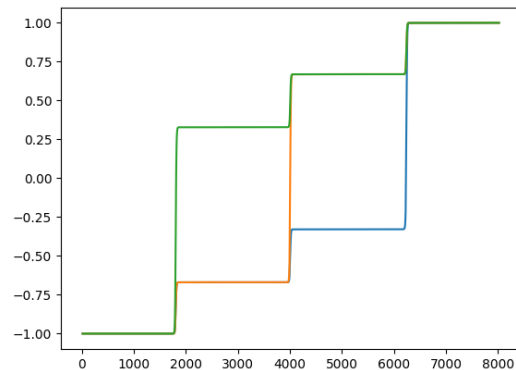
As $t \rightarrow \infty$ we observe:



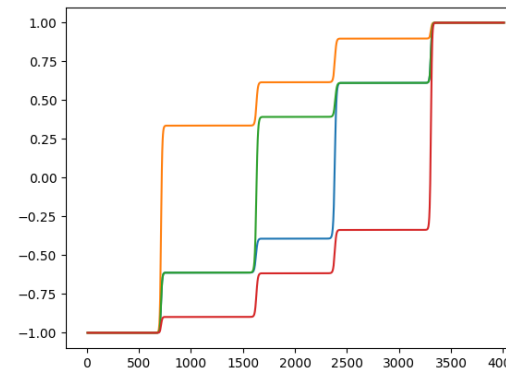
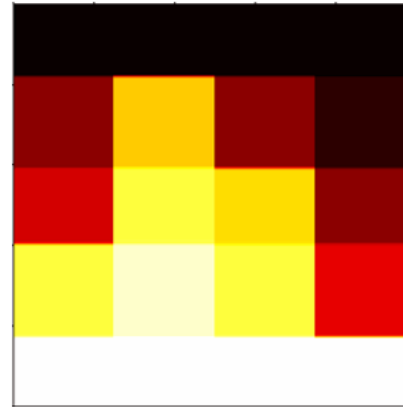
Many tubes: numerics



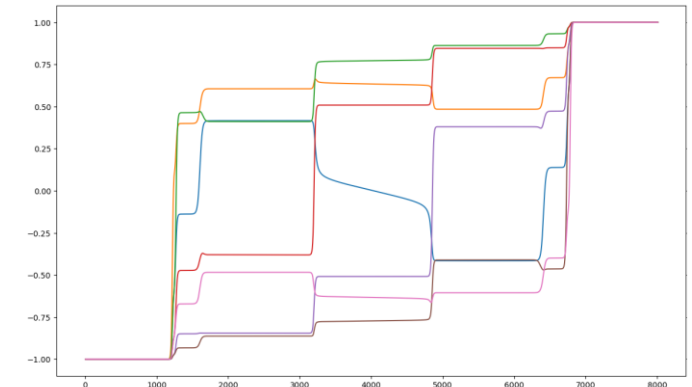
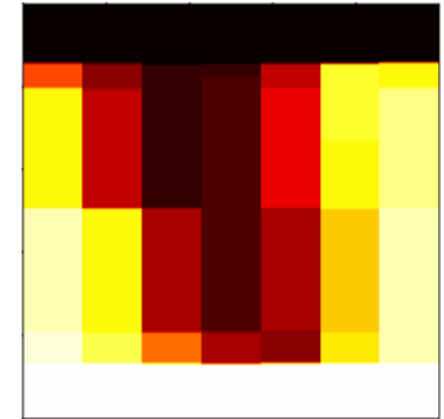
3 tubes



4 tubes



7 tubes



Questions:
(open)

- (1) explain the structure of “asymptotic solutions” for n tubes and study their stability
- (2) find speed of growth of the mixing zone
- (3) understand the behaviour as $n \rightarrow \infty$. Do we approximate 2-dim IPM?
- (4) can we repeat this story for the many tubes viscous fingering model?

Scheme of proof: step 1 - TW

Travelling wave (TW) ansatz with fixed v :

$$c_1(t, y) = c_1(y - vt)$$

$$c_2(t, y) = c_2(y - vt)$$

$$u_1(t, y) = u_1(y - vt)$$

$$u_2(t, y) = u_2(y - vt)$$

$$p_1(t, y) = p_1(y - vt)$$

$$p_2(t, y) = p_2(y - ct)$$



System of ODEs in \mathbb{R}^6 :

$$\begin{cases} \dot{X} = F_v(X, Y) \\ l \cdot \dot{Y} = AY - BX \end{cases}$$

Here:

$$\bullet X = \begin{pmatrix} c_1 \\ c_2 \\ \partial_\xi c_1 \\ \partial_\xi c_2 \end{pmatrix} \in \mathbb{R}^4, \quad Y = \begin{pmatrix} u_1 \\ p_1 - p_2 \end{pmatrix} \in \mathbb{R}^2$$

$$\bullet A = \begin{pmatrix} 0 & 1 \\ 2 & 0 \end{pmatrix}, \quad B \in M^{2 \times 4}, \quad l \ll 1$$

With condition at $+\infty$:

$$c_1(+\infty) = 1$$

$$c_2(+\infty) = 1$$

$$u_1(+\infty) = 0$$

$$u_2(+\infty) = 0$$

$$(p_1 - p_2)(+\infty) = 0$$

Observation:

for $l \rightarrow 0$ this system has a special ``slow-fast'' structure.

Key tool: **geometric singular perturbation theory (GSPT)**

by Fenichel (JDE, 1979)

Scheme of proof: step 2 – cascade of 2 TW

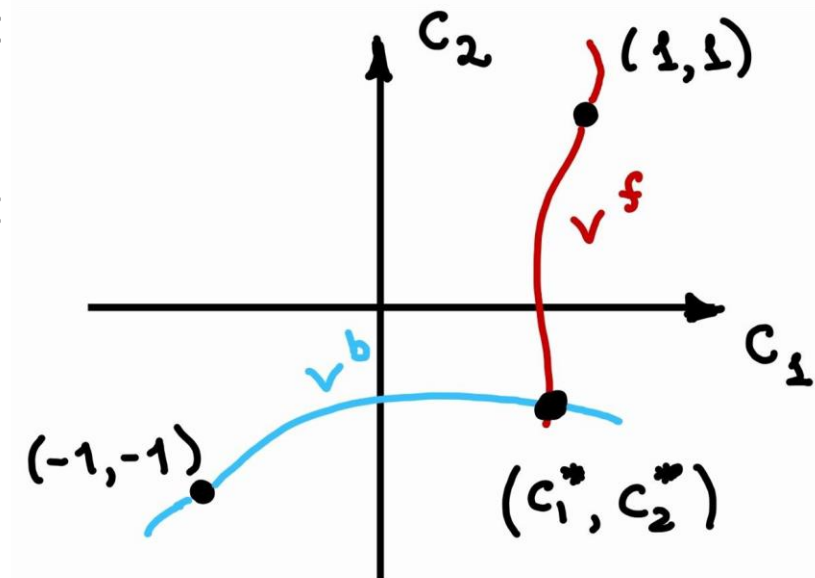
1) For each $v^f \in \mathbb{R}$ we find all points s.t. there exists a TW:

$$(c_1, c_2) \rightarrow (1, 1)$$

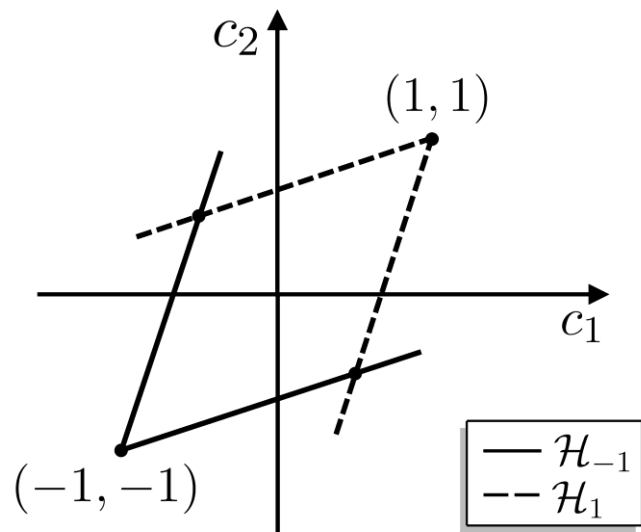
2) For each $v^b \in \mathbb{R}$ we find all points s.t. there exists a TW:

$$(-1, -1) \rightarrow (c_1, c_2)$$

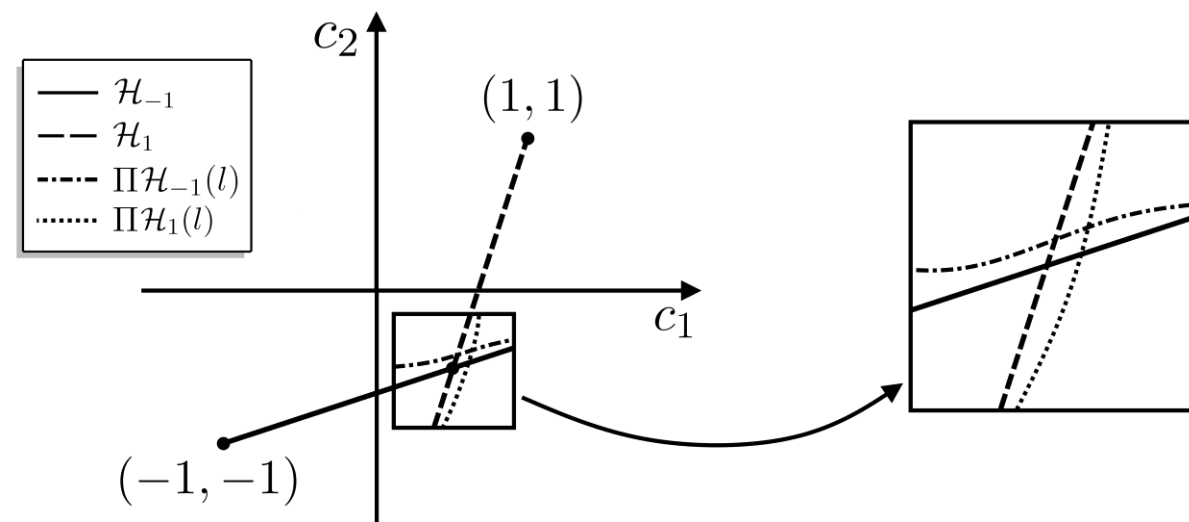
3) Find the intersection points of these two curves



$l = 0$ – these curves are just straight lines



$0 < l \ll 1$ – perturbation argument



But what is a singular limit $l = 0$?

$$\begin{aligned}
 & \partial_t c_1 + \partial_y(u_1 c_1) - \partial_{yy} c_1 = -B \\
 & \partial_t c_2 + \partial_y(u_2 c_2) - \partial_{yy} c_2 = B \\
 & u_1 = \frac{c_2 + c_1}{2} - c_1 = \bar{c} - c_1 \\
 & B = \begin{cases} -\partial_y u_1 \cdot c_1, & \partial_y u_1 < 0, \\ +\partial_y u_2 \cdot c_2, & \partial_y u_1 > 0 \end{cases}
 \end{aligned}$$

$l = 0$ corresponds to the two-tubes TFE equations (**) !!!

NB: (**) can be seen a hyperbolic system in non-conservative form (for fixed choice of B):

$$C_t + A(C)C_y = 0$$

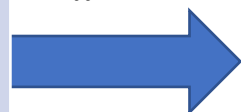
$$C(0, y) = \begin{cases} (+1, +1), & y > 0 \\ (-1, -1), & y \leq 0 \end{cases}$$

Question: TFE as a limit of IPM as $k_y/k_x \rightarrow \infty$?
(open) Can we use the connection to prove the linear growth in 2D IPM?

2D IPM model

$$\begin{aligned}
 c_t + u \cdot \nabla c &= \varepsilon \Delta c \\
 \operatorname{div} u &= 0 \\
 u &= - \begin{pmatrix} k_x & 0 \\ 0 & k_y \end{pmatrix} \nabla p - (0, c)
 \end{aligned}$$

$\frac{k_y}{k_x} \rightarrow \infty$



2D TFE model

$$\begin{aligned}
 c_t + u \cdot \nabla c &= \varepsilon \Delta c \\
 \operatorname{div} u &= 0 \\
 u &= (u^x, u^y) \\
 u^y &= \bar{c} - c
 \end{aligned}$$

- Use vanishing viscosity criteria to define admissible shocks
- For 2 and 3 tubes this is a “Temple-like” system (rarefaction and shock curves coincide and are linear)

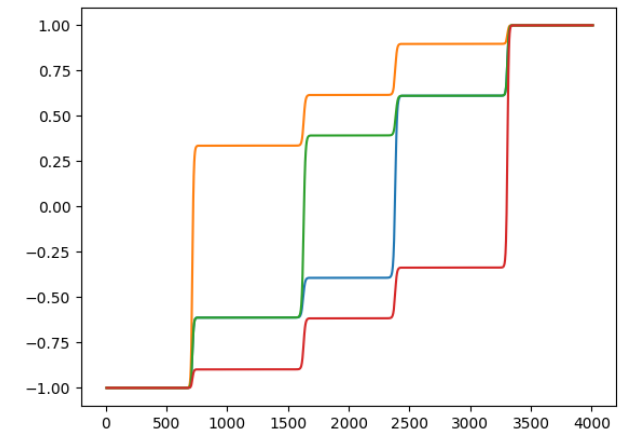
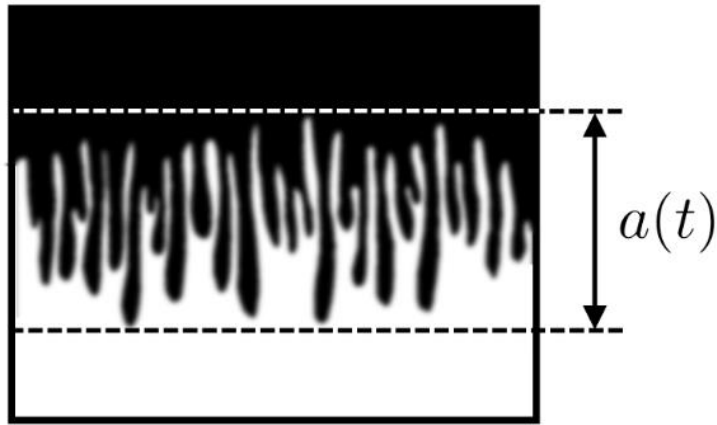
Thank you for your attention!

Muito obrigada!

¡Muchas gracias!

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<https://yulia-petrova.github.io/>



For more details see arXiv:2401.05981 (two-tubes model)

Any questions, comments and ideas are very welcome!

Own works on the topic of the talk:

1. **Yu. Petrova**, S. Tikhomirov, Ya. Efendiev, *“Propagating terrace in a two-tubes model of gravitational fingering”*
Submitted. ArXiv: 2401.05981; 2024.
2. Bakharev, F., Enin, A., Groman, A., Kalyuzhnyuk, A., Matveenko, S., **Petrova, Yu.**, Starkov, I. and Tikhomirov, S.,
“Velocity of viscous fingers in miscible displacement: Comparison with analytical models”.
Journal of Computational and Applied Mathematics, 402, p.113808; 2022.
3. Bakharev, F., Pavlov D., Enin, A., Matveenko, S., **Petrova, Yu.**, Rastegaev N., and Tikhomirov, S.,
“Velocity of viscous fingers in miscible displacement: Intermediate concentration”
Journal of Computational and Applied Mathematics, 451:116107, 2024.

Other references:

Dynamics of viscous fingering:

1. Nijjer J., Hewitt D., and Neufeld J. The dynamics of miscible viscous fingering from onset to shutdown. Journal of Fluid Mechanics 837 (2018): 520-545.
2. Menon, G. and Otto, F., 2006. Diffusive slowdown in miscible viscous fingering. Communications in Mathematical Sciences, 4(1), pp.267-273.
3. Menon, G. and Otto, F., 2005. Dynamic scaling in miscible viscous fingering. Communications in mathematical physics, 257, pp.303-317.
4. Homsy, G.M., 1987. Viscous fingering in porous media. Annual review of fluid mechanics, 19(1), pp.271-311.

Geometric singular perturbation theory (GSPT):

1. Fenichel, N., 1979. Geometric singular perturbation theory for ordinary differential equations. *Journal of differential equations*, 31(1), pp.53-98.
2. Wechselberger, M., 2020. *Geometric singular perturbation theory beyond the standard form* (Vol. 6). New York: Springer.
3. Kuehn, C., 2015. *Multiple time scale dynamics* (Vol. 191). Berlin: Springer.

Well-posedness for IPM:

1. Kiselev, A. and Yao, Y., 2023. Small scale formations in the incompressible porous media equation. *Archive for Rational Mechanics and Analysis (ARMA)*, 247(1), p.1.
2. A. Castro, D. Cordoba and D. Lear, Global existence of quasi-stratified solutions for the confined IPM equation, *Archive for Rational Mechanics and Analysis (ARMA)*, 232 (2019), no. 1, 437–471.
3. T. Elgindi, On the asymptotic stability of stationary solutions of the inviscid incompressible porous medium equation, *Archive for Rational Mechanics and Analysis (ARMA)*, 225 (2017), no. 2, 573–599.

Non-uniqueness for IPM:

1. D. Cordoba, D. Faraco and F. Gancedo, Lack of uniqueness for weak solutions of the incompressible porous media equation, *Archive for Rational Mechanics and Analysis (ARMA)* 200 (2011), no. 3, 725–746.
2. Shvydkoy, R.: Convex integration for a class of active scalar equations. *J. Am. Math. Soc.* 24(4), 1159–1174 (2011).
3. L. Szekelyhidi, Jr. Relaxation of the incompressible porous media equation, *Ann. Sci. de l'Ecole Norm. Superieure* (4) 45 (2012), no. 3, 491–509.