Optimizing Liquid Scintillators for Low Energy Events via Wavelength-Shifter and Quantum-Dot Doping

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ABSTRACT: The composition of liquid scintillators are extremely versatile. This adaptablility allows liquid scintillators to be tuned for low energy events, in partilcular, optimization of the light yield and transparency of the scintillator to Cherenkov radiation. In this study, the light yield of seven potential solvents are presented: Toluene, Pseudocumene (PC), Linear Alkyl Benzene (LAB), Phenylcyclohexane (PCH), Phenyl-o-Xylyl Ethane (PXE) and two grades of Di-Isopropyl Napthalene (DIN). The wavelength-shifter 2,5-Diphenyloxazole (PPO) is utilized in various concentrations to study the light yield of binary scintillator mixtures. The experimental data is then fit to a model for the light yield dependence on the PPO concentration. The fit results provide insight about the energy transfer processes in the liquid scintillator mixtures and determine the optimal PPO concentration for each solvent. In addition, the optical properties of 2 g/L PPO in toluene doped with quantum-dots (QDs) are explored. QDs ability to re-sculpt the absorption and emission spectrum of liquid scintillators demonstrate that QDs are viable canidates for conserving the Cherenkov photons produced by interactions in the liquid scintillator.

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1. Introduction

The aim of this study is to theorize an ideal liquid scintillator for observing low energy events. This is achieved by maximizing the light yield by choosing the scintillator liquid and the 2,5-Diphenyloxazole (PPO) concentration while maintaining control over the emission and absorption spectrum of a scintillator with quantum-dots (QDs) to preserve the directional information inherent in Cherenkov radiation produced in scintillator interactions.

Table 1. Solvents and solute names or abbreviations used in this paper, the corresponding IUPAC names, CAS numbers and information on the source of the chemicals.

Name	IUPAC Name	CAS Number	Source Information
Toluene	Methylbenzene	108-88-03	Sigma Aldrich [20] (Chromasolv Plus)
Pseudocumene	1,2,4-Trimethylbenzene	95-63-6	Aldrich (98 %) [20]
(PC)			
Linear Alkyl Benzene	various chain lengths	67774-74-7	Cepsa/Petresa [21]
(LAB)			PETRELAB® 550-Q
Phenyl-o-Xylylethane	1,2-Dimethyl-4-	6196-95-8	Dixie Chemical
(PXE)	(1-Phenyl-Ethyl)-Benzene		Company [22]
Di-Isopropylnapthalene	Isomer mixture	38640-62-9	courtesy of
(DIN)			PerkinElmer [23]
Di-Isopropylnapthalene,	Isomer mixture	38640-62-9	courtesy of
high purity (DIN HP)			PerkinElmer [23]
Phenylcyclohexane	Cyclohexylbenzene	827-52-1	Aldrich (≥ 97 %) [20]
(PCH)			
Diphenyloxazole	2,5-Diphenyloxazole	92-71-7	Aldrich (99 %) [20]
(PPO)			suitable for scint.

1.1 Liquid Scintillators

As charged particles propagate through a scintillator, electrons in the π -bonds of the aromatic rings are readily excited [17]. Since solvent molecules are present in high abundance, interaction with solvent molecules is likely. The absorption and emission spectra of a single-molecule based scintillator overlaps. In order to prevent degradation of efficiency due to the re-absorption of scintillation light by the scintillator, the scintillator is mixed with a wavelength-shifter, solute, that absorbs the energy from the excited electron and subsequently releases light of a longer wavelength to which the scintillator is more transparent. Optimizing the energy transfer between solvent and solute molecules maximizes the light yield. The light yield is the number of photons produced per deposited energy. Light yield is a key property of the scintillator because it is directly related to the energy resolution of the detector.

In addition to testing the light yield of the various liquid scintillator candidates listed in Table 1, the concentration of PPO, solute, in each sample is also varied; the relative light yields are compared. There are two main effects important in the description of the PPO concentration dependency. First, the efficiency of energy transfer from solvent to solute increases with increasing PPO concentration. At the studied PPO concentrations the energy transfers are predominantly non-radiative Förster energy transfers [29, 30], but the emission and re-absorption of real photons or collision energy transfer and others [31] can also contribute to the overall effective energy transfer rate. Second, there is the competing effect of self-quenching [17] which becomes more important

at higher PPO concentrations. In this process two PPO molecules interact with each other and no light is emitted. The light yield $I(c_{PPO})$ as a function of PPO concentration c_{PPO} can be described by

$$I(c_{PPO}) = p_1 \cdot \frac{1}{1 + p_2/c_{PPO}} \cdot \frac{1}{1 + p_3 \cdot c_{PPO}}$$
(1.1)

with three solvent-specific parameters p_1 , p_2 and p_3 . For details on the derivation of these types of equations see [32, 33] and references therein. The maximum light yield without self-quenching is characterized by p_1 , the effectiveness of energy transfer from solvent molecules to solute molecules corresponds to p_2 and the effect of self-quenching is described by p_3 .

1.2 Quantum-Dots

Quantum-Dots are semiconducting nanocrystals. The size of a QD is inversely proportional to its band gap and therefore proportional to the wavelength of light it absorbs and emits. Doping the scintillator with QDs enables the absorbed and emitted wavelength to be tuned. These features of the QDs can be exploited improve liquid scintillators as follows.

Cherenkov radiation is emitted whenever a charged particle travels through a medium at a speed greater than the phase velocity of light in that medium. Unlike scintillation light that is emitted isotropically, Cherenkov light is radiated in a cone which is directionally relevant to the incident particle. The Cherenkov photons are not monochromatic and have an energy distribution. Doping the scintillator with QDs that absorb and emit only high energy photons enables the longer wavelengths of Cherenkov light to travel toward the detector unimpeded by scintillation process. Because long wavelengths travel faster than short wavelengths of light in a medium due to chromatic dispersion, and the scintillation absorption and emission process takes time, the long wavelengths of Cherenkov light will reach the detector before the scintillation light. The difference in arrival times allows the two types of light to be distinguished. Furthermore, adding QDs to the scintillator solvent generally narrows the absorption spectrum of the liquid scintillator. This is vital for low energy events. For example, a 1 MeV electron emits only 83 Cherenkov photons between 370-550 nm[18]; therefore, it is imperative that the scintillator absorption spectrum is minimized as much as possible. The measured Cherenkov light can be used for reconstructing trajectory of particles cite[18], thus eliminating background signals.

Unfortunately, doping liquid scintillators with QDs lowers the light yield as shown in Section 3.1.3 and discussed Section 3.2.2. Fortunately, this can be at least partially countered by choosing the emission peak to be at an optimal wavelength for the quantum efficiency of the photomultiplier tube (PMT)[34] The QDs of various emission and absorption spectra that undergo study for suitability are listed in Table 2.

2. Experimental

The light yield of candidate liquid scintillators with varying amounts of PPO and different QDs are measured. This is achieved by collecting the charge produced by a gamma ray irradiated scintillator with a PMT. In addition, the emission and absorption spectrum of the QDs are quantified with two different spectrophotometers. The details of the experimental setup are explained below followed by a description of the sample handling.

Table 2. Quantum-Dot composition and source information.

Composition	Source Information			
(core/shell)				
CdS	NN-Labs-CS360[24]			
CdS	NN-Labs-CS400[24]			
CdS	MKN-CdS-T360[25]			
CdS	MKN-CdS-T380[25]			
CdS	MKN-CdS-T400[25]			
CdS/ZnS	Ocean NanoTech QZR-400-0010[26]			
CdS/ZnS	Ocean NanoTech QZR-425-0010 [26]			
CdSeS/ZnS	Crystalplex NC-450-A [27]			

2.1 Light Yield Setup

2.1.1 Setup

The setup is illustrated in Figure 1. A $(1 \text{ cm} \times 1 \text{ cm} \times 3.5 \text{ cm})$ UV-transparent quartz cell (Starna 21-Q-10 [28]) holds the scintillator mixtures. The cell is coupled to the PMT with transparent silicone optical grease (Saint Gobain BC-630) with a similar index of refraction to the cell and PMT glass so that light losses due to reflection are minimized. The quartz cell is then secured in a reflective white Teflon block to further increase the light collection efficiency. The scintillator is excited by a 137 Cs source with an activity of $\approx 1~\mu\text{Ci}$ (=37 kBq). The source is attached to the outside of the Teflon block. Isotope 137 Cs undergoes beta decay with the subsequent emission of a single 662 keV gamma. The emitted gamma rays enter the quartz cell, then typically Compton scatter with electrons, which excite the scintillator. Cherenkov light is also produced for electrons with energies above a given threshold as discussed in Section 1.2 . A fraction of the emitted optical photons hit the photocathode of a Hamamatsu R1828-01 PMT [34]. A Hamamatsu E2979-500 base [35] is used, and the high voltage of -1675 V is provided by a LeCroy 1454 high voltage system. The charge incident on the photocathode produce pulses which are recorded by an AlazarTech ATS9870 PCI waveform digitizer [36].

2.1.2 Photomultiplier Tube Calibration

2.1.3 Sample Handling and Procedure

In order to ensure reproducible and accurate measurements, a sample handling procedure is defined. In particular, contamination of the samples with dust, residual chemicals and oxygen is avoided.

The bottles which store the samples are rinsed with isopropanol, cyclohexane, dried and finally rinsed with the solvent. In addition, the bottles are sealed under N_2 to reduce the amount of oxygen in contact with the sample. Oxygen interferes with the light production processes (oxygen

quenching [17]) and leads to a loss of light output. The lids of the bottles are sealed with Teflon tape and electric tape to further limit the amount of oxygen entering the bottle.

As for the measurement, the quartz cell is cleaned thoroughly when samples are changed. The optical grease on the outside is removed with ethanol and on the inside it is washed with isopropanol then dried with N_2 twice, washed with cyclohexane, and again dried with N_2 . The amount of liquid scintillator sample in the cell is kept constant to avoid effects due to differences in light collection. Before each measurement, a pipet is inserted into the liquid to slowly purge the samples with nitrogen for about 10 minutes to actively remove remaining oxygen.

The coupling of the cell to the PMT is done consistently, and the PMT is aligned in the same way in order to avoid effects due to the earth's magnetic field. Section 3.1 presents the studies that determine how reproducible the measurements are when these procedures are applied. The quartz cell containing the liquid scintillator sample is then placed in the dark box setup as described in the previous section, and a measurement is performed for a duration of 30 minutes.

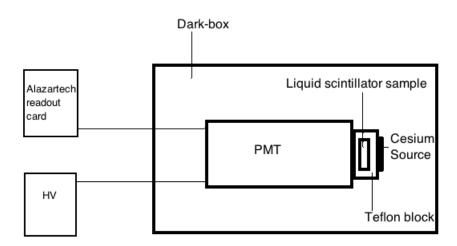


Figure 1. Schematic view of the light yield measurement setup.

2.2 Absorption and Emission Setup

2.2.1 Setup

Two spectrophotometers are utilized in this experiment: a Perkin Elmer Lambda 25 UV/Vis for measuring the emission spectrum and a Shimadzu UV3101 for measuring absorption spectrum of

the QD doped liquid scintillators. The instruments work via a double-beam setup; the beams are comprised of white light produced by a deuterium UV lamp and a halogen lamp. [?]

2.2.2 Sample Handling and Procedure

To measure the absorption spectrum of each quantum-dot sample, the following procedure is executed. A (1 cm×1 cm×3.5 cm) UV-transparent quartz cell (Starna 21-Q-10 [28]) is rinsed with toluene to dissolve any residue on the wall. Next, the cell is soaked in hydrochloric acid for 5 minutes to clean the cell and destroy any remaining QD. To wash away the hydrochloric acid and the residue QD constituents, the quartz cell is rinsed with distilled water and is then subjected to three iterations of liquid isopropyl alcohol evaporated by nitrogen gas. Since all of the QD samples are dissolved in toluene, a background absorption spectrum analysis is required to isolate the quantum-dot spectrum. Thus, the quartz vile is filled with toluene. The cell is then placed into the Shimadzu UV3101, and a background scan is performed. Next, the toluene is removed from the cell and replaced by the QD sample of choice. Finally, The sample is inserted back into the spectrophotometer, and a scan is initiated.

As for the emission spectrum measurements of the QD samples, the cleaning procedure is exactly the same as for the absorption trials; the only difference is that the emission spectrum is measured by a Perkin Elmer Lambda 25 UV/Vis which does not require a background toluene scan.

3. Results

This section presents the analysis methodology as well as the main results; namely the light yield measurements with varying PPO concentrations for each of the solvents listed in Table 1 and for the different QDs listed in Table 2 dissolved in toluene with 2 g/L of PPO added. The emission and absorption spectrum of the QD scintillators are also explicitly examined.

3.1 Light Yield

3.1.1 Signal Analysis

- Baseline: At the beginning of each waveform, 100 samples are averaged to get the baseline value.
- Pulse finder: Pulses are defined as a consecutive series of samples below the rounded baseline value.
- Pulse times: For each pulse the start time is defined as the time where the waveform falls below the rounded baseline value. The stop time is set once the pulse reaches this value again.
- Charge: The integral of the baseline-subtracted waveform between the start and stop times.

The duration of each run is around thirty minutes and during this time on the order of $3 \cdot 10^5$ waveforms are recorded. Each pulse corresponds to one gamma ray that interacted within the liquid scintillator sample. The amount of light collected by the PMT is dependent on the scattering angle of the electron that is produced by Compton scattering of the 0.662 MeV gamma ray. The

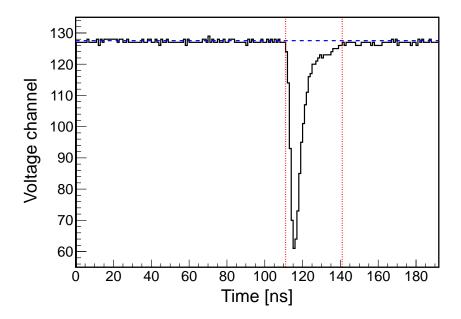


Figure 2. Example waveform recorded by the AlazarTech digitizer card. For each waveform, 192 samples are taken with a rate of 1Gs/s. The blue, dashed line is the baseline which is calculated for each waveform. The red, dotted lines show the start and stop time for the main pulse in this waveform. Integration of the baseline subtracted waveform between start and stop yields 520. The trigger threshold is set at voltage channel 116.

result of the charge calculation is placed in a histogram, Figure 3. Since the Compton effect is the dominant mechanism, this histogram is a Compton spectrum. The events where the Compton scattered electron obtained close to the maximal energy, 478 keV for the 662 keV gammas from ¹³⁷Cs, is associated with events at the right end of this spectrum. Multiple pulses are found for each waveform, including baseline fluctuations. The small fake pulses from baseline fluctuations did not affect analysis. In order to compare different scintillator mixtures, the charge where the number of events dropped to half of the Compton edge maximum is compared, Figure 3. This quantity is proportional to the scintillator light yield. For practical reasons and because this study only looks at relative light yield, arbitrary charge units are used instead of real charge units.

3.1.2 Investigation of Uncertainties

This section focuses on the reproducibility of a measurement utilizing the sample handle procedure defined in Section 2.1 in order to estimate the errors of a single measurement.

First, the error due to differences in the procedure of coupling the scintillator cell to the PMT glass is checked. For this test, the same sample, LAB with 5 g/L PPO, is sealed in the cell and only the coupling is renewed. The RMS divided by the sample mean for several trials is 0.30 %. This number is used as a rough estimate for the relative error of a single measurement due to differences in the coupling. The small uncertainty in the determination of the charge value from the histograms, Figure 3, is already included here. The relative spread of 0.30 % is negligible compared to other uncertainties discussed below.

The effect of oxygen is also studied quantitatively. A comparatively old sample, 24 days after mixing, of DIN HP with 5 g/L PPO is purged with nitrogen bubbles for 10 minutes as in the standard measurement procedure and is measured. Then the sample is opened and left in contact with oxygen for 15 minutes. It is then re-measured without nitrogen purging; the charge value dropped by 2.04 %. After an additional 3 hours of contact with oxygen the same is measured. The relative difference between the first and the last measurement is about 6.92 %. This shows that the nitrogen purging and reproducible procedures in preventing oxygen from entering the scintillator are important.

In order to estimate the total error of a single light yield measurement, the same sample with the full sample change protocol, as described in Section 2.1, is measured repeatedly. For these measurements an older sample, 7 weeks after mixing, of DIN HP is used as well as a standard DIN sample, 3-4 weeks old. For the older samples, 10 minutes of nitrogen purging is not sufficient to remove oxygen. The effect of differences in the oxygen removal is included in these reproducibility measurements. However, it is difficult to assess exactly how much oxygen affects the measurements when the full measurement protocol is applied. The RMS for these data sets is regarded as a rough estimate of the relative uncertainty of a single light yield measurement. It is about 3.70 %.

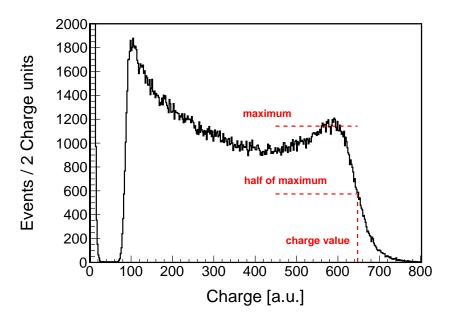


Figure 3. Example charge spectrum. To compare different scintillators, the charge at a characteristic point close to the true Compton edge is determined.

3.1.3 Light Yield Results

Table 3 contains the light yield results for the different liquid scintillator candidates described in Table 1.

In Figure 3, each plot shows the solvent's characteristic charge at the Compton edge as a function of PPO concentration in g/L. Concentrations of 0.5 g/L, 1 g/L, 2 g/L, 5 g/L, 10 g/L and 50

Table 3. Listed below are the fit results for each solvent's data set. The fit parameters and their errors are given and their respective χ^2 values, and the probabilities to get a worse χ^2 than the observed one. Since we have 6 data points and 3 parameters in the fit the number of degrees of freedom is 3.

Solvent	p ₁ [a.u.]	p ₂ [g/L]	p_3 [g/L]	χ^2	prob. %	max. LY
Toluene	681 ± 10	0.09 ± 0.014	0.019 ± 0.001	0.55	91	627 (at 2.2 g/L)
PC	578 ± 47	2.09 ± 0.37	0.014 ± 0.004	5.30	15	422 (at 12.2 g/L)
LAB	525 ± 24	0.36 ± 0.01	0.013 ± 0.003	4.33	23	462 (at 5.3 g/L)
PXE	694 ± 25	0.48 ± 0.96	0.013 ± 0.002	2.55	47	596 (at 6.1 g/L)
DIN	668 ± 19	0.27 ± 0.01	0.007 ± 0.001	2.34	50	613 (at 6.2 g/L)
DIN HP	673 ± 40	0.19 ± 0.08	0.005 ± 0.003	10.6	1	631 (at 6.1 g/L)
PCH	634 ± 29	0.34 ± 0.07	0.012 ± 0.003	6.11	11	562 (at 5.4 g/L)

g/L have been used. Tables 4 and 5 include the fit results; the parameters p1, p2 and p3 describe the light yield normalization, the energy transfer from solvent to solute and the self-quenching effect, respectively.

As clearly stated in Table 5, toluene produces the highest light yield around 2.2 g/L. Since all of the QDs are dissolved in a toluene solution, and maximizing light yield a main objective, the QD scintillators are doped with 2 g/L PPO for comparison with 2 g/L PPO in toluene. The light yield as a function of QD emission wavelength is plotted in Figure 6.

3.2 Emission and Aborption Spectrum of Quantum-Dot Doped Liquid Scintillators

3.2.1 Signal Analysis

Various trials are conducted to test the output signal of both the Perkin Elmer Lambda 25 UV/Vis and the Shimadzu UV3101 spectrophotometers[19]. The error is defined in terms of the relative intensities of the beam before transiting the quartz cell, I(0), and after it passes through the cell, I(x); the relationship is given by equation 3.1.

$$A(x) = \log_{10}\left(\frac{I(x)}{I(0)}\right) \tag{3.1}$$

The errors for the Perkin and Shimadzu spectrophotometers are measured to be $A = \pm 0.00075$ and $A = \pm 0.002$ respectfully. These errors are in close agreement with the manufactures specifications[?].

3.2.2 Absorption and Emission Results

The absorption and emission spectrum of the QDs were extensively probed, and the results are graphed in Figures 7 and 8 respectfully. As stated in Section 1.2 and proved in 3.1.3, the light yield of the QD doped scintillators is much lower than that of just toluene with PPO. This is due to the overlap in emission and absorption QD spectrum coupled with low quantum efficiency [18]. (I'm creating a figure that shows the overlap for emission and absorption for each quantum dot). The absorption spectrum of the QD doped liquid scintillators are clearly narrower than that of toluene

doped with 2 g/L PPO. (How can we back this? Need to take an absorption and emission scan of toluene + 2g/L PPO.)

4. Conclusion

Detection of low energy events in liquid scintillators require high light yield and preservation of Cherenkov radiation. This study successfully measures and models the light yield of seven scintillator liquids as a function of wavelength-shifter concentration in order to find the maximum light yield of each candidate. Furthermore, an extensive study on the effects of quantum-dot doping of toluene with 2 g/L of PPO in solution are analyzed. QDs are able to reduce the width and tune the location of the absorption spectrum peak. This allows the quantum efficiency of the photocathode to be matched and minimized the amount of Cherenkov radiation absorbed by the liquid scintillator.

Overall, the results of the light yield trials are immediately useful. Regrettable, the quantum-dots reduce the light yield significantly which means that the Cherenkov photon recovered from utilization of the QD are not worth it; however, new quantum-dots with higher quantum efficiency and larger separation of absorption and emission peaks are developed constantly.

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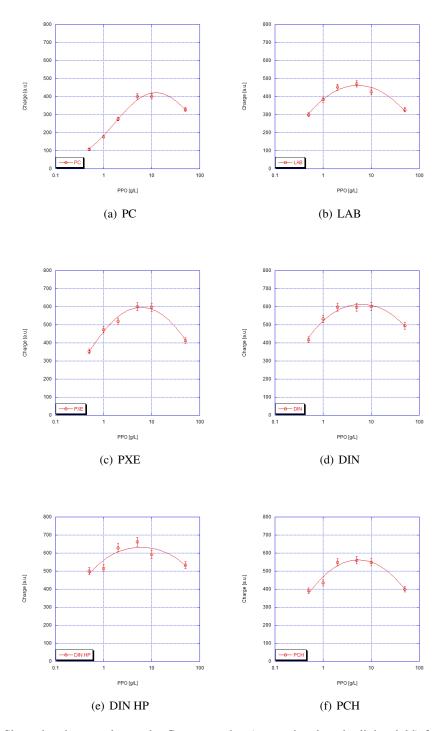


Figure 4. Show the charge value at the Compton edge (proportional to the light yield) for each of the 6 concentrations of PPO for six of the liquid scintillator candidates. A relative error of 3.7 % was used for each measurement. The red lines show the fit of equation 1.1. Note that the x-axis is logarithmic; the PPO concentrations span two orders of magnitude.

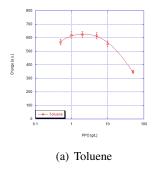


Figure 5. Show the charge value at the Compton edge (proportional to the light yield) for each of the 6 concentrations of PPO for toluene. A relative error of 3.7 % was used for each measurement. The red lines show the fit of equation 1.1. Note that the x-axis is logarithmic; the PPO concentrations span two orders of magnitude.

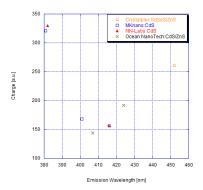


Figure 6. The light yield of quantum-dot doped toluene scintillators with 2 g/L of PPO versus emission wavelength of the QDs.

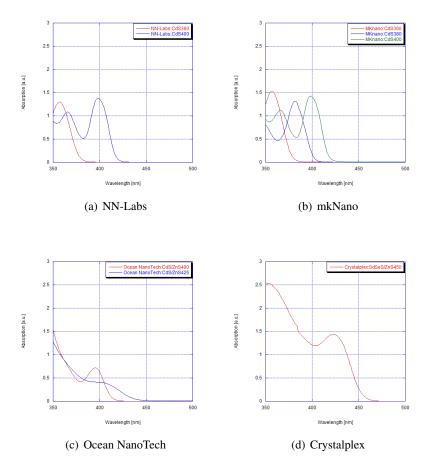


Figure 7. Absorption spectrum of the quantum-dots listed in Table 2 dissolved in toluene and grouped by manufacturer.

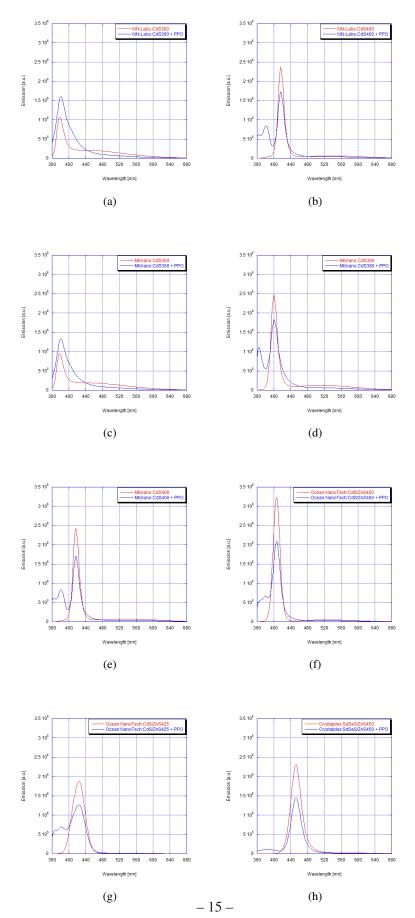


Figure 8. Emission spectrum of the quantum-dots listed in Table 2 dissolved in toluene with and without 2 g/L of PPO added.