Lecture 10

Introduction to dimensional analysis

Any investigation in physics ultimately comes down to determining certain quantity which may depend on various other quantities that characterizes the phenomena under consideration. The problem therefore reduces to establishing relationship which can always be represented in the form —

$$f = f(q_1, q_2, ..., q_n)$$

The quantities a_1, a_2, \ldots, a_n are called governing parameters which combines into a single term that bears the same dimension as f.

Dimensional and dimensionless quantities

If the numerical value of any quantity changes upon change of one unit of measurement to another unit of measurement, then the quantity is called a dimensional quantity.

Example: The angular frequency of a simple pendulum is given by, $\omega = \sqrt{8/L}$.

The numerical values of ω depends on a and L. If we change the unit of measurement from MKS to CGS, the numerical values of g, L and hence ω changes. These are dimensional quantities.

However, we can construct a quantity,

$$\Pi = \frac{\omega}{\sqrt{8/L}}$$

Upon changing the quantity unit of measurement from MKS to CGS, the numerical value of TI remains the same. So, TI is a dimensionless quantity.

Power-law monomial nature of dimensions

A function $f = f(a_1, a_2, ..., a_n)$ is called power law monomial is f can be expressed in terms of,

$$f = q_1^{\times} q_2^{\times} \dots q_n^{\times}$$

In fact, all Physical laws follow power law mono-

mial. There is no such instances where dimension function can be expressed as. It age as on $a_1^2 + a_5^{-3}$ etc.

This power law monomial nature actually helps us to deduce the exact nature of some function on quantity even without detailed information. Let's see an example.

The speed of travelling wave through a string depends on the clastic property of the medium and intertia. For a string, the clastic property in the tension on the string, and the inertia term in the mass density u.

$$\cdot \cdot \cdot \forall = \forall (T, u)$$

Power law monomial says.

$$[V] = [T]^{\alpha} [u]^{\beta}$$

$$\Rightarrow LT^{-1} = (MLT^{-2})^{\alpha} (ML^{-1})^{\beta} = M^{\alpha+\beta} \alpha - \beta - 2\alpha$$

Equating both sides we get.

$$\alpha + \beta = 0$$
 and $-2\alpha = -1$
 $\alpha - \beta = 1$ $\therefore \alpha = \frac{1}{2}$

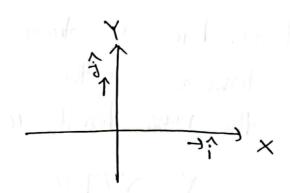
$$\beta = -\frac{1}{2}$$

$$[v] = [\tau]^{1/2} [u]^{-1/2}$$

Upto some dimensionless constant, the expression of the velocity is exact.

Reviews of coordinate systems and vectors Corresian coordinates and polar coordinates

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In 2D cartesian coordinates, we define two unit vectors— i and i along the x and Y axis. The unit vectors are typically in the directions of increasing x and Y.

If we move to polar coordinates, then the unit vectors are is and it. is typically away

From the origin in outward direction and changes of its perpendicular to it and its in the direction of increasing 0. The major difference between the Cortesian and polar coordinate in that and of its fixed in terms of directions. However, or and of continuously can change direction At any time, or and of can form the expressed as—

 $\hat{\gamma}(t) = \cos\theta \hat{i} + \sin\theta \hat{j}$

 $\hat{\theta}(t) = -\sin\theta \hat{\tau} + \cos\theta \hat{\theta}$

if the particle is moving in a circle. The position of any particle at some time is exactly given by,

Vector dot and cross product $\vec{A} = Ax\hat{i} + Ay\hat{J} + A_z\hat{k}$ $\vec{B} = B_x\hat{i} + B_y\hat{J} + B_z\hat{k}$

A.B = AzBz+AyBy +AzBz

$$\overrightarrow{A} \times \overrightarrow{B} = \begin{vmatrix} \widehat{1} & \widehat{2} & \widehat{k} \\ A_{x} & A_{y} & A_{z} \\ B_{x} & B_{y} & B_{z} \end{vmatrix}$$

$$\hat{t} = \hat{S}_{ij} = Kronecker delta$$

$$\hat{t} \times \hat{t} = Kronecker delta$$

$$S_{ij} = \begin{cases} 1 & i = j \\ 0 & i \neq j \end{cases}$$

In general, we can express dot product as,

$$\overrightarrow{A}.\overrightarrow{B}' = \sum A; B; = A; B; \sim |Einstein sum | Convention \rightarrow nepeated indices are summed over$$

$$(\overrightarrow{A} \times \overrightarrow{B}) = \mathcal{E}_{ijk} A_{j} \beta_{k}$$

$$= \sum_{k=1}^{3} \mathcal{E}_{ilk} (A_{j} \beta_{k}) + \sum_{k=1}^{3} \mathcal{E}_{ilk} (A_{j} \beta_{k}) + \sum_{k=1}^{3} \mathcal{E}_{ilk} (A_{j} \beta_{k})$$

$$= \left[0 + \left[0 + 0 + A_{j} \beta_{j} \right] + \left[0 + (-1) A_{j} \beta_{j} + 0 \right] \right]$$

$$= A_{j} \beta_{j} - A_{3} \beta_{j} = A_{j} \beta_{z} - A_{z} \beta_{y} \text{ and so on } .$$