

Quantum Mechanics II

Lecture Notes 3

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The Dirac Notation:

In the often used Dirac notation vectors in a Hilbert space are denoted by a 'ket': $|\psi\rangle$. The addition of two vectors is written as:

$$|\psi\rangle = c_1|\psi_1\rangle + c_2|\psi_2\rangle, \text{ where } c_i \in \mathbb{C}.$$

Note that for a given vector $|\psi\rangle \in \mathcal{H}$, the inner product defines the map:

$$\langle\psi|\cdot\rangle \equiv (\psi, \cdot): \mathcal{H} \rightarrow \mathbb{C}.$$

This map is antilinear

$$((c_1\psi_1 + c_2\psi_2), \psi) = c_1^* (\psi_1, \psi) + c_2^* (\psi_2, \psi)$$

$$= c_1^* \langle\psi_1|\psi\rangle + c_2^* \langle\psi_2|\psi\rangle$$

$$\langle c_1\psi_1 + c_2\psi_2|\psi\rangle = (c_1^* \langle\psi_1| + c_2^* \langle\psi_2|)\psi\rangle$$

The number of linearly independent maps is equal to the dimension of the Hilbert space. Thus the space of antilinear maps from $H \rightarrow \mathbb{C}$ itself forms a vector space. This space is called a dual vector space H^* .

The inner product gives an isomorphism between H and H^* . We can express this by choosing a set of orthonormal basis $\{|\alpha\rangle\}$:

$$\langle \alpha | \alpha' \rangle = \delta_{\alpha \alpha'}$$

Operators in the bracket notation:

We write $(\psi, \hat{A} \phi) = \langle \psi | \hat{A} | \phi \rangle$ where \hat{A} is implicitly thought to act on $|\phi\rangle$.

$\langle \psi | \hat{A}$, on the other hand, is dual to $\hat{A}^\dagger | \psi \rangle$ since

$$\begin{aligned} (\psi, \hat{A} \phi) &= (\hat{A}^\dagger \psi, \phi) \\ &= (\phi, \hat{A}^\dagger \psi)^* \\ &= \langle \phi | \hat{A}^\dagger | \psi \rangle^* \end{aligned}$$

When \hat{A} is Hermitian:

$$(\psi, \hat{A}^\dagger \varphi) = (\psi, \hat{A} \varphi) \quad \forall \psi \in \varphi.$$

$$\Rightarrow \langle \psi | \hat{A}^\dagger | \varphi \rangle = \langle \psi | \hat{A} | \varphi \rangle$$

Projection Operators:

Suppose \hat{A} is a Hermitian operator. Let a_i be its eigenvalues and $|a_i\rangle$ the corresponding eigenvector. Then we can express any vector $|\psi\rangle$ as a linear combination of the eigenvectors of \hat{A} :

$$|\psi\rangle = \sum_i c_i |a_i\rangle$$

Now suppose we want to know what is the j -th component of $|\psi\rangle$ in the basis $\{|a_i\rangle\}$.

We then define an operator \hat{P}_j st.

$$\hat{P}_j |\alpha_j\rangle = |\alpha_j\rangle$$

$$\hat{P}_j |\alpha_i\rangle = 0 \quad \text{if } i \neq j.$$

These two equations can be expressed as

$$\hat{P}_i |\alpha_j\rangle = \delta_{ij} |\alpha_i\rangle \quad (\text{no sum})$$

\hat{P}_j is called a projection operator. Note that

$$\begin{aligned} \hat{P}_j |\psi\rangle &= \hat{P}_j \sum_i c_i |\alpha_i\rangle \\ &= \sum_i c_i \hat{P}_j |\alpha_i\rangle \\ &= \sum_i c_i \delta_{ij} |\alpha_j\rangle \\ &= c_j |\alpha_j\rangle \end{aligned}$$

Comments:

1. \hat{P}_j peels off the j -th component of $|\psi\rangle$

when expressed in the $\{|\alpha_i\rangle\}$ basis.

2. Note that the projection operators are basis dependent.

It is conventional to express \hat{P}_j in the Dirac notation as

$$\hat{P}_j = |\alpha_j\rangle\langle\alpha_j|$$

$$\text{Then } \langle\alpha_i|\alpha_j\rangle = \delta_{ij} \Rightarrow$$

$$\hat{P}_j |\psi\rangle = |\alpha_j\rangle\langle\alpha_j| \sum_i c_i |\alpha_i\rangle$$

$$= |\alpha_j\rangle \sum_i c_i \langle\alpha_j|\alpha_i\rangle$$

$$= |\alpha_j\rangle \sum_i c_i \delta_{ij}$$

$$= c_j |\alpha_j\rangle.$$

Note that

$$\hat{P}_i^2 = \hat{P}_i \quad (\text{Show})$$

$$\text{and} \quad \hat{P}_i \hat{P}_j = \delta_{ij} \hat{P}_j \quad (\text{no sum})$$

Finally, when we sum over all the projectors associated with a basis, we should get $\hat{1}$:

$$\sum_i \hat{P}_i = \sum_i |\alpha_i\rangle\langle\alpha_i| = \hat{1} \quad (\text{Show}).$$

When we express $|\psi\rangle$ in the basis $\{|\alpha_i\rangle\}$:

$$|\psi\rangle = \sum_i c_i |\alpha_i\rangle$$

$$\Rightarrow c_i = \langle\alpha_i|\psi\rangle$$

Previously we wrote this as

$$c_i = (\alpha_i, \psi)$$

One of the most useful formulas :

$$\sum_i P_i = \hat{1}$$

can expressed as:

$$\sum_i |\alpha_i\rangle\langle\alpha_i| = \hat{1}$$

The Hermitian matrix \hat{A} whose eigenvectors are $\{|\alpha_i\rangle\}$ can be expressed as

$$\hat{A} = \sum_i \alpha_i |\alpha_i\rangle\langle\alpha_i|$$

If \hat{B} is another operator such that it commutes with \hat{A} then \hat{B} can be expressed as:

$$\hat{B} = \sum_i \beta_i |\alpha_i\rangle\langle\alpha_i|.$$

Since $\hat{A} \nless \hat{B}$ commute, $\{|\alpha_i\rangle\}$ form simultaneous eigenstates.

$$[\hat{A}, \hat{B}] = \sum_{i,j} \alpha_i \beta_j |\alpha_i\rangle \langle \alpha_i| \alpha_j \langle \alpha_j|$$

$$- \sum_{i,j} \beta_i \alpha_j |\alpha_j\rangle \langle \alpha_j| \alpha_i \langle \alpha_i|$$

$$= \sum_{i,j} \alpha_i \beta_j |\alpha_i\rangle \delta_{ij} \langle \alpha_j|$$

$$- \sum_{i,j} \beta_i \alpha_j |\alpha_j\rangle \delta_{ji} \langle \alpha_i|$$

$$= \sum_i \alpha_i \beta_i |\alpha_i\rangle \langle \alpha_i|$$

$$- \sum_i \beta_i \alpha_i |\alpha_i\rangle \langle \alpha_i| = 0$$

The state $|\psi\rangle$ of a system in the position representation is defined to be the wave function:

$$\psi(x) = \langle x | \psi \rangle$$

where $\hat{X} |x\rangle = x |x\rangle$

and $\langle x | x' \rangle = \delta(x - x')$

Then a generic equation:

$$\hat{O} |\psi\rangle = |\varphi\rangle$$

becomes in the position representation:

$$\langle x | \hat{O} |\psi\rangle = \langle x | \varphi \rangle$$

$$\Rightarrow \int dx' \langle x | \hat{O} | x' \rangle \psi(x') = \varphi(x)$$

where we have used

$$\hat{1} = \int dx |x\rangle \langle x|$$

$$\Rightarrow \int dx \hat{O}_{xx'} \psi(x') = \psi(x)$$

$$\text{If } \hat{O} = \hat{p}$$

$$\hat{p}_{xx'} = i\hbar \frac{\partial}{\partial x'} \delta(x-x')$$

and we get :

$$-i\hbar \frac{\partial}{\partial x} \psi(x) = \psi(x).$$