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### Exercise 12.1 (Schmidt decomposition and entanglement entropy)

As in tutorial 12, let  $|\psi\rangle$  be a pure state of a composite system, AB. The Schmidt decomposition of this state is denoted by  $|\psi\rangle = \sum_{i=1}^k \sigma_i |i_A\rangle |i_B\rangle$ .

(a) Verify that

$$\langle\psi|\psi\rangle = \sum_{i=1}^k \sigma_i^2.$$

In general, the *von Neumann entropy* of a density matrix  $\rho$  is defined as

$$\mathcal{S}(\rho) = -\text{tr}[\rho \log(\rho)],$$

with the logarithm interpreted as matrix function, and the convention  $0 \log(0) = \lim_{x \rightarrow 0} x \log(x) = 0$ .

In tutorial 12 we found the reduced density matrices of the subsystems, defined as  $\rho_1 = \text{tr}_2[|\psi\rangle\langle\psi|]$  and  $\rho_2 = \text{tr}_1[|\psi\rangle\langle\psi|]$ . We observed that  $\rho_1$  and  $\rho_2$  have the same eigenvalues  $(\sigma_i^2)_{i=1,\dots,k}$ . The *entanglement entropy* between the two subsystems is then given by

$$\mathcal{S}_{\text{ent}} = \mathcal{S}(\rho_1) = \mathcal{S}(\rho_2) = -\sum_{i=1}^k \sigma_i^2 \log(\sigma_i^2).$$

(You should convince yourself that  $\mathcal{S}(\rho_1)$  and  $\mathcal{S}(\rho_2)$  are indeed equal to the sum on the right.) Intuitively, the entanglement entropy measures how strongly the subsystems are intertwined.

(b) Which sets of singular values  $(\sigma_i)_{i=1,\dots,k}$  minimize and maximize the entanglement entropy, respectively, under the normalization condition  $\sum_{i=1}^k \sigma_i^2 = 1$ ? ( $k$  should be regarded as fixed.)

Hints: The smallest possible entanglement entropy is zero. Regarding maximization, you can take the normalization condition via a Lagrange multiplier into account.

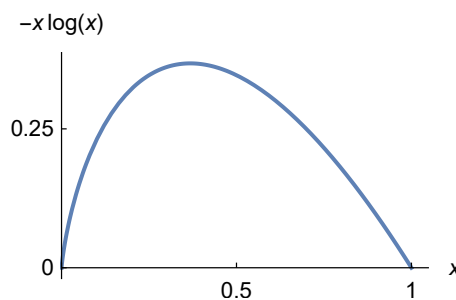
(c) Show that  $\mathcal{S}_{\text{ent}} = 0$  (completely unentangled case) implies that  $|\psi\rangle$  can be written as tensor product of a state from subsystem A and one from subsystem B.

### Solution

(a) Inserting the Schmidt decomposition of  $|\psi\rangle$  directly leads to:

$$\langle\psi|\psi\rangle = \sum_{i,j=1}^k \sigma_i \sigma_j \underbrace{\langle i_A | j_A \rangle}_{\delta_{ij}} \underbrace{\langle i_B | j_B \rangle}_{\delta_{ij}} = \sum_i \sigma_i^2.$$

(b) We know that singular values are (in general) real and non-negative. Moreover, due to the normalization condition,  $\sigma_i^2 \in [0, 1]$  for all  $i$ . The following figure visualizes  $-x \log(x)$ , which is non-negative for any  $x \in [0, 1]$ , and equal to 0 precisely if  $x = 0$  or  $x = 1$ .



By identifying  $x$  with  $\sigma_i^2$ , one concludes that the entanglement entropy is non-negative.  $\mathcal{S}_{\text{ent}} = 0$  is reached by setting the first singular values to 1 and the others to 0 (which satisfies the normalization condition).

Regarding maximization of the entanglement entropy, we take the normalization constraint by a Lagrange multiplier  $\lambda \in \mathbb{R}$  into account, and abbreviate  $\sigma_i^2 = x_i$  for convenience:

$$\mathcal{L}(x_1, \dots, x_k, \lambda) = - \sum_{i=1}^k x_i \log(x_i) - \lambda \left( \sum_{i=1}^k x_i - 1 \right).$$

Finding an extremum of  $\mathcal{L}$  by differentiation w.r.t.  $x_i$ , and using that  $\log'(x) = \frac{1}{x}$  for  $x > 0$ , gives

$$0 \stackrel{!}{=} \frac{\partial \mathcal{L}}{\partial x_i} = -\log(x_i) - 1 - \lambda \quad \rightsquigarrow \quad x_i = e^{-1-\lambda}.$$

In particular, all  $x_i$  take the same value; combined with the normalization condition, one arrives at  $x_i = \frac{1}{k}$  for all  $i = 1, \dots, k$ . This assignment indeed maximizes  $\mathcal{L}$  since  $-x \log(x)$  is concave. The corresponding singular values are  $\sigma_i = \frac{1}{\sqrt{k}}$  for  $i = 1, \dots, k$ , and

$$\max_{\sigma_1, \dots, \sigma_k} \mathcal{S}_{\text{ent}} = -\log(1/k) = \log(k).$$

- (c) As already mentioned,  $\mathcal{S}_{\text{ent}} = 0$  is reached by setting the first singular values to 1 and the others to 0, and this is actually the only case in which  $\mathcal{S}_{\text{ent}} = 0$  since  $-x \log(x) = 0$  implies  $x = 0$  or  $x = 1$ . In terms of the Schmidt decomposition  $|\psi\rangle = \sum_{i=1}^k \sigma_i |i_A\rangle |i_B\rangle$ , only the first term remains, i.e.,

$$|\psi\rangle = |1_A\rangle |1_B\rangle$$

is a tensor product of two basis states.