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**Exercise 11.1** (von Neumann equation and time evolution with density operators)

- (a) Based on the Schrödinger equation (cf. tutorial 3), derive the following *von Neumann equation* for a density matrix  $\rho(t) = \sum_j p_j |\psi_j(t)\rangle \langle \psi_j(t)|$ :

$$i\hbar \frac{d}{dt} \rho(t) = [H, \rho(t)].$$

Here  $[\cdot, \cdot]$  is the matrix commutator.

Hint: Use the product rule for computing the time derivative of each term  $|\psi_j(t)\rangle \langle \psi_j(t)|$ .

- (b) What is the formal solution for  $\rho(t)$  expressed in terms of the time evolution operator  $U(t) = e^{-iHt/\hbar}$ ?
- (c) We consider the specific single-qubit Hamiltonian operator

$$H = JX$$

with parameter  $J \in \mathbb{R}$ . Compute the time-dependent density matrix  $\rho(t)$  starting from the initial state  $\rho_0 = \begin{pmatrix} 2/3 & 0 \\ 0 & 1/3 \end{pmatrix}$  at  $t = 0$ . For simplicity, you can set  $\hbar = 1$ .

- (d) Since the map  $\rho \mapsto [H, \rho]$  is linear, we can represent it as matrix-vector multiplication after “vectorizing”  $\rho$ , i.e., collecting its entries in a vector, denoted  $\vec{\rho}$  in the following. For the commutator, this leads to

$$\vec{\rho} \mapsto \text{vec}([H, \rho]) = (H \otimes I - I \otimes H^T) \vec{\rho},$$

where the identity matrix has the same dimension as  $H$ . Thus we can represent the von Neumann equation equivalently in the “superoperator” form

$$i\hbar \frac{d}{dt} \vec{\rho}(t) = \mathcal{H} \vec{\rho}(t), \quad \mathcal{H} = H \otimes I - I \otimes H^T.$$

Write down the formal solution of this differential equation, and determine  $\mathcal{H}$  for the Hamiltonian from (c).

**Solution**

- (a) By the product rule,

$$\begin{aligned} i\hbar \frac{d}{dt} \rho(t) &= i\hbar \frac{d}{dt} \sum_j p_j |\psi_j(t)\rangle \langle \psi_j(t)| = \sum_j p_j \left( \frac{i\hbar d}{dt} |\psi_j(t)\rangle \langle \psi_j(t)| + |\psi_j(t)\rangle \frac{i\hbar d}{dt} \langle \psi_j(t)| \right) \\ &= \sum_j p_j \left( (H |\psi_j(t)\rangle) \langle \psi_j(t)| - |\psi_j(t)\rangle (\langle \psi_j(t)| H^\dagger) \right) = H \rho(t) - \rho(t) H = [H, \rho(t)]. \end{aligned}$$

Here we have used that the Hamiltonian  $H$  is Hermitian, that is,  $H^\dagger = H$ . The minus sign stems from taking the conjugate-transpose of the Schrödinger equation.

- (b) The formal solution for a pure state is  $|\psi(t)\rangle = U(t) |\psi(0)\rangle$ , and thus

$$\rho(t) = U(t) \rho_0 U(t)^\dagger = U(t) \rho_0 U(-t).$$

This  $\rho(t)$  indeed solves the von Neumann equation, since  $\frac{d}{dt} U(t) = -\frac{i}{\hbar} H U(t)$ .

- (c) We first determine the unitary time evolution operator  $U(t)$ , using the formula for the  $R_x$  rotation operator:

$$U(t) = e^{-iHt} = e^{-iJXt} = \cos(Jt)I - i \sin(Jt)X = \begin{pmatrix} \cos(Jt) & -i \sin(Jt) \\ -i \sin(Jt) & \cos(Jt) \end{pmatrix}.$$

Inserted into the equation from (b) leads to

$$\begin{aligned} \rho(t) &= U(t) \rho_0 U(-t) = \begin{pmatrix} \cos(Jt) & -i \sin(Jt) \\ -i \sin(Jt) & \cos(Jt) \end{pmatrix} \begin{pmatrix} 2/3 & 0 \\ 0 & 1/3 \end{pmatrix} \begin{pmatrix} \cos(Jt) & i \sin(Jt) \\ i \sin(Jt) & \cos(Jt) \end{pmatrix} \\ &= \begin{pmatrix} \frac{2}{3} \cos^2(Jt) + \frac{1}{3} \sin^2(Jt) & \frac{i}{3} \cos(Jt) \sin(Jt) \\ -\frac{i}{3} \cos(Jt) \sin(Jt) & \frac{1}{3} \cos^2(Jt) + \frac{2}{3} \sin^2(Jt) \end{pmatrix} \\ &= \begin{pmatrix} \frac{1}{3} + \frac{1}{3} \cos^2(Jt) & \frac{i}{6} \sin(2Jt) \\ -\frac{i}{6} \sin(2Jt) & \frac{1}{3} + \frac{1}{3} \sin^2(Jt) \end{pmatrix} \\ &= \begin{pmatrix} \frac{1}{2} + \frac{1}{6} \cos(2Jt) & \frac{i}{6} \sin(2Jt) \\ -\frac{i}{6} \sin(2Jt) & \frac{1}{2} - \frac{1}{6} \cos(2Jt) \end{pmatrix} = \frac{1}{2} \left( I - \frac{1}{3} \sin(2Jt)Y + \frac{1}{3} \cos(2Jt)Z \right). \end{aligned}$$

- (d) The superoperator differential equation can be identified with a Schrödinger equation, with analogous formal solution

$$\vec{\rho}(t) = e^{-i\mathcal{H}t/\hbar} \vec{\rho}_0.$$

For the Hamiltonian from (c), we note that  $X^T = X$ , and obtain

$$\mathcal{H} = J(X \otimes I - I \otimes X) = J \begin{pmatrix} 0 & -1 & 1 & 0 \\ -1 & 0 & 0 & 1 \\ 1 & 0 & 0 & -1 \\ 0 & 1 & -1 & 0 \end{pmatrix}.$$