

# Probability

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## 1 Probability Spaces

### 1.1 Classical Probability Spaces

Textbook probability theory [1, 2, 3] is defined using the notions of a *sample space*  $\Omega$ , a space of *events*  $\mathcal{F}$ , and a *probability measure*  $\mu$ . In this paper, we will only consider *finite* sample spaces: we therefore define a sample space  $\Omega$  as an arbitrary non-empty finite set and the space of events  $\mathcal{F}$  as,  $2^\Omega$ , the powerset of  $\Omega$ . A *probability measure* is a function  $\mu : \mathcal{F} \rightarrow [0, 1]$  such that:

- $\mu(\Omega) = 1$ , and
- for a collection of pairwise disjoint events  $E_i$ , we have  $\mu(\bigcup E_i) = \sum \mu(E_i)$ .

*Example 1* (Two coin experiment). Consider an experiment that tosses two coins. We have four possible outcomes that constitute the sample space  $\Omega = \{HH, HT, TH, TT\}$ . The event that the first coin is “heads” is  $\{HH, HT\}$ ; the event that the two coins land on opposite sides is  $\{HT, TH\}$ ; the event that at least one coin is tails is  $\{HT, TH, TT\}$ . Depending on the assumptions regarding the coins, we can define several probability measures. Here is a possible one:

$\mu(\emptyset)$	$= 0$	$\mu(\{HT, TH\})$	$= 2/3$
$\mu(\{HH\})$	$= 1/3$	$\mu(\{HT, TT\})$	$= 0$
$\mu(\{HT\})$	$= 0$	$\mu(\{TH, TT\})$	$= 2/3$
$\mu(\{TH\})$	$= 2/3$	$\mu(\{HH, HT, TH\})$	$= 1$
$\mu(\{TT\})$	$= 0$	$\mu(\{HH, HT, TT\})$	$= 1/3$
$\mu(\{HH, HT\})$	$= 1/3$	$\mu(\{HH, TH, TT\})$	$= 1$
$\mu(\{HH, TH\})$	$= 1$	$\mu(\{HT, TH, TT\})$	$= 2/3$
$\mu(\{HH, TT\})$	$= 1/3$	$\mu(\{HH, HT, TH, TT\})$	$= 1$

### 1.2 Quantum Probability Spaces

A classical model decides the occurrence or non-occurrence of all events simultaneously which is inconsistent with quantum mechanics. Indeed, in the quantum world, there are (non-commuting) events which cannot happen simultaneously. To accommodate this situation, we completely abandon the sample space  $\Omega$  and define and reason directly about events. Thus a quantum probability space will consist of just two components: a set of events  $\mathcal{A}$  and a probability measure  $\phi : \mathcal{A} \rightarrow [0, 1]$ . These components are defined as follows.

We first assume an ambient Hilbert space  $\mathcal{H}$  and define the set of events  $\mathcal{A}$  as *projections* on  $\mathcal{H}$ . Similarly to the classical case, a probability measure is a function  $\phi : \mathcal{A} \rightarrow [0, 1]$  satisfying:

- $\phi(\mathbb{1}) = 1$ , and
- for all  $A \in \mathcal{A}$ , we have  $\phi(A^*A) \geq 0$ .

As an example, let  $P_1, P_2, \dots, P_k$  be mutually orthogonal projections on  $\mathcal{H}$  with sum  $\mathbb{1}$  and define the event space  $\mathcal{A}$  to be the linear span of these operators:

$$\mathcal{A} = \left\{ \sum_{j=1}^k \lambda_j P_j \mid \lambda_1, \dots, \lambda_k \in \mathbb{C} \right\}.$$

Yu-Tsung says: If we want to follow Swart [4]... Maybe we should define the set of projection and the

Each state  $\psi$  of the Hilbert space induces a probability measure  $\phi_\psi : \mathcal{A} \rightarrow [0, 1]$  defined as follows:

$$\phi_\psi(A) = \langle \psi \mid A\psi \rangle$$

### 1.3 Plan

We see that a quantum probability space embeds several infinite precision assumptions, both in the Hilbert space itself and in the range of the probability measure. We will now study the consequences of restricting each occurrence of  $\mathbb{R}$  or  $\mathbb{C}$  to some finite structure.

## References

- [1] William G. Faris. Appendix: Probability in quantum mechanics. In *The infamous boundary : seven decades of controversy in quantum physics*. Boston : Birkhauser, 1995.
- [2] R.L. Graham, D.E. Knuth, and O. Patashnik. *Concrete Mathematics: A Foundation for Computer Science*. A foundation for computer science. Addison-Wesley, 1994.
- [3] V.K. Rohatgi and A.K.M.E. Saleh. *An Introduction to Probability and Statistics*. Wiley Series in Probability and Statistics. Wiley, 2011.
- [4] JM Swart. Introduction to quantum probability. *Lecture Notes*, 2013.