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Hopf fibration

In the mathematical field of <u>differential topology</u>, the **Hopf fibration** (also known as the **Hopf bundle** or **Hopf map**) describes a <u>3-sphere</u> (a <u>hypersphere</u> in <u>four-dimensional space</u>) in terms of <u>circles</u> and an ordinary <u>sphere</u>. Discovered by <u>Heinz Hopf</u> in 1931, it is an influential early example of a <u>fiber bundle</u>. Technically, Hopf found a many-to-one <u>continuous function</u> (or "map") from the 3-sphere onto the 2-sphere such that each distinct *point* of the 2-sphere comes from a distinct *circle* of the 3-sphere (<u>Hopf 1931</u>). Thus the 3-sphere is composed of fibers, where each fiber is a circle—one for each point of the 2-sphere.

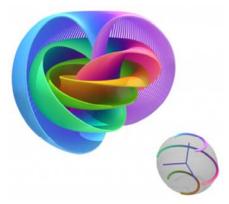
This fiber bundle structure is denoted

$$S^1 \hookrightarrow S^3 \overset{p}{\longrightarrow} S^2,$$

meaning that the fiber space S^1 (a circle) is <u>embedded</u> in the total space S^3 (the 3-sphere), and $p:S^3\to S^2$ (Hopf's map) projects S^3 onto the base space S^2 (the ordinary 2-sphere). The Hopf fibration, like any fiber bundle, has the important property that it is <u>locally</u> a <u>product space</u>. However it is not a *trivial* fiber bundle, i.e., S^3 is not *globally* a product of S^2 and S^1 although locally it is indistinguishable from it.

This has many implications: for example the existence of this bundle shows that the higher <u>homotopy</u> groups of spheres are not trivial in general. It also provides a basic example of a <u>principal bundle</u>, by identifying the fiber with the <u>circle</u> group.

Stereographic projection of the Hopf fibration induces a remarkable structure on ${\bf R}^3$, in which space is filled with nested <u>tori</u> made of linking <u>Villarceau circles</u>. Here each fiber projects to a <u>circle</u> in space (one of which is a line, thought of as a "circle through infinity"). Each torus is the stereographic projection of the <u>inverse image</u> of a circle of latitude of the 2-sphere. (Topologically, a torus is the product of two circles.) These tori are illustrated in the images at right. When ${\bf R}^3$ is compressed to a ball, some geometric structure is lost although the topological structure is retained (see <u>Topology and geometry</u>). The loops are <u>homeomorphic</u> to circles, although they are not geometric <u>circles</u>.



The Hopf fibration can be visualized using a stereographic projection of S^3 to ${\bf R}^3$ and then compressing R^3 to a ball. This image shows points on S^2 and their corresponding fibers with the same color.



Pairwise linked keyrings mimic part of the Hopf fibration.

There are numerous generalizations of the Hopf fibration. The unit sphere in <u>complex coordinate space</u> \mathbb{C}^{n+1} fibers naturally over the <u>complex projective space</u> $\mathbb{C}\mathbf{P}^n$ with circles as fibers, and there are also <u>real</u>, <u>quaternionic</u>, [1] and <u>octonionic</u> versions of these fibrations. In particular, the Hopf fibration belongs to a family of four fiber bundles in which the total space, base space, and fiber space are all spheres:

$$S^0 \hookrightarrow S^1 \rightarrow S^1, \ S^1 \hookrightarrow S^3 \rightarrow S^2, \ S^3 \hookrightarrow S^7 \rightarrow S^4, \ S^7 \hookrightarrow S^{15} \rightarrow S^8.$$

By Adams' theorem such fibrations can occur only in these dimensions.

The Hopf fibration is important in twistor theory.

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Definition and construction

For any <u>natural number</u> n, an n-dimensional sphere, or <u>n-sphere</u>, can be defined as the set of points in an (n+1)-dimensional <u>space</u> which are a fixed distance from a central <u>point</u>. For concreteness, the central point can be taken to be the <u>origin</u>, and the distance of the points on the sphere from this origin can be assumed to be a unit length. With this convention, the n-sphere, S^n , consists of the points $(x_1, x_2, ..., x_{n+1})$ in \mathbb{R}^{n+1} with $x_1^2 + x_2^2 + \cdots + x_{n+1}^2 = 1$. For example, the 3-sphere consists of the points (x_1, x_2, x_3, x_4) in \mathbb{R}^4 with $x_1^2 + x_2^2 + x_3^2 + x_4^2 = 1$.

The Hopf fibration $p: S^3 \to S^2$ of the 3-sphere over the 2-sphere can be defined in several ways.

Direct construction

Identify \mathbf{R}^4 with \mathbf{C}^2 and \mathbf{R}^3 with $\mathbf{C} \times \mathbf{R}$ (where \mathbf{C} denotes the complex numbers) by writing:

$$(x_1,x_2,x_3,x_4) \leftrightarrow (z_0,z_1) = (x_1+ix_2,x_3+ix_4)$$

and

$$(x_1,x_2,x_3) \leftrightarrow (z,x) = (x_1+ix_2,x_3).$$

Thus S^3 is identified with the <u>subset</u> of all (z_0, z_1) in \mathbb{C}^2 such that $|z_0|^2 + |z_1|^2 = 1$, and S^2 is identified with the subset of all (z, x) in $\mathbb{C} \times \mathbb{R}$ such that $|z|^2 + x^2 = 1$. (Here, for a complex number z = x + iy, $|z|^2 = z z^* = x^2 + y^2$, where the star denotes the <u>complex</u> conjugate.) Then the Hopf fibration p is defined by

$$p(z_0,z_1) = (2z_0z_1^*, |z_0|^2 - |z_1|^2).$$

The first component is a complex number, whereas the second component is real. Any point on the 3-sphere must have the property that $|z_0|^2 + |z_1|^2 = 1$. If that is so, then $p(z_0, z_1)$ lies on the unit 2-sphere in $\mathbb{C} \times \mathbb{R}$, as may be shown by squaring the complex and real components of p

$$2z_0z_1^*\cdot 2z_0^*z_1 + \left(|z_0|^2 - |z_1|^2\right)^2 = 4|z_0|^2|z_1|^2 + |z_0|^4 - 2|z_0|^2|z_1|^2 + |z_1|^4 = \left(|z_0|^2 + |z_1|^2\right)^2 = 1$$

Furthermore, if two points on the 3-sphere map to the same point on the 2-sphere, i.e., if $p(z_0, z_1) = p(w_0, w_1)$, then (w_0, w_1) must equal $(\lambda z_0, \lambda z_1)$ for some complex number λ with $|\lambda|^2 = 1$. The converse is also true; any two points on the 3-sphere that differ by a common complex factor λ map to the same point on the 2-sphere. These conclusions follow, because the complex factor λ cancels with its complex conjugate λ^* in both parts of p: in the complex $2z_0z_1^*$ component and in the real component $|z_0|^2 - |z_1|^2$.

Since the set of complex numbers λ with $|\lambda|^2 = 1$ form the unit circle in the complex plane, it follows that for each point m in S^2 , the inverse image $p^{-1}(m)$ is a circle, i.e., $p^{-1}m \cong S^1$. Thus the 3-sphere is realized as a disjoint union of these circular fibers.

A direct parametrization of the 3-sphere employing the Hopf map is as follows.^[2]

$$egin{aligned} z_0 &= e^{irac{\xi_1+\xi_2}{2}} \sin\eta \ z_1 &= e^{irac{\xi_2-\xi_1}{2}} \cos\eta. \end{aligned}$$

or in Euclidean R⁴

$$egin{aligned} x_1 &= \cos\left(rac{\xi_1+\xi_2}{2}
ight) \sin\eta \ x_2 &= \sin\left(rac{\xi_1+\xi_2}{2}
ight) \sin\eta \ x_3 &= \cos\left(rac{\xi_2-\xi_1}{2}
ight) \cos\eta \ x_4 &= \sin\left(rac{\xi_2-\xi_1}{2}
ight) \cos\eta \end{aligned}$$

Where η runs over the range 0 to $\pi/2$, and ζ_1 and ζ_2 can take any values between 0 and 2π . Every value of η , except 0 and $\pi/2$ which specify circles, specifies a separate <u>flat torus</u> in the 3-sphere, and one round trip (0 to 2π) of either ζ_1 or ζ_2 causes you to make one full circle of both limbs of the torus.

A mapping of the above parametrization to the 2-sphere is as follows, with points on the circles parametrized by ζ_2 .

$$z = \cos(2\eta)$$

$$x = \sin(2\eta)\cos\xi_1$$

$$y = \sin(2\eta)\sin\xi_1$$

Geometric interpretation using the complex projective line

A geometric interpretation of the fibration may be obtained using the <u>complex projective line</u>, \mathbb{CP}^1 , which is defined to be the set of all complex one-dimensional <u>subspaces</u> of \mathbb{C}^2 . Equivalently, \mathbb{CP}^1 is the <u>quotient</u> of $\mathbb{C}^2 \setminus \{0\}$ by the <u>equivalence relation</u> which identifies (z_0, z_1) with $(\lambda z_0, \lambda z_1)$ for any nonzero complex number λ . On any complex line in \mathbb{C}^2 there is a circle of unit norm, and so the restriction of the <u>quotient map</u> to the points of unit norm is a fibration of S^3 over \mathbb{CP}^1 .

 ${\bf CP}^1$ is diffeomorphic to a 2-sphere: indeed it can be identified with the <u>Riemann sphere</u> ${\bf C}_{\infty} = {\bf C} \cup \{\infty\}$, which is the <u>one point compactification</u> of ${\bf C}$ (obtained by adding a <u>point at infinity</u>). The formula given for p above defines an explicit diffeomorphism between the complex projective line and the ordinary 2-sphere in 3-dimensional space. Alternatively, the point (z_0, z_1) can be mapped to the ratio z_1/z_0 in the Riemann sphere ${\bf C}_{\infty}$.

Fiber bundle structure

The Hopf fibration defines a <u>fiber bundle</u>, with bundle projection p. This means that it has a "local product structure", in the sense that every point of the 2-sphere has some <u>neighborhood</u> U whose inverse image in the 3-sphere can be <u>identified</u> with the product of U and a circle: $p^{-1}(U) \cong U \times S^1$. Such a fibration is said to be locally trivial.

For the Hopf fibration, it is enough to remove a single point m from S^2 and the corresponding circle $p^{-1}(m)$ from S^3 ; thus one can take $U = S^2 \setminus \{m\}$, and any point in S^2 has a neighborhood of this form.

Geometric interpretation using rotations

Another geometric interpretation of the Hopf fibration can be obtained by considering rotations of the 2-sphere in ordinary 3-dimensional space. The <u>rotation group SO(3)</u> has a <u>double cover</u>, the <u>spin group Spin(3)</u>, <u>diffeomorphic</u> to the 3-sphere. The spin group acts <u>transitively</u> on S^2 by rotations. The <u>stabilizer</u> of a point is isomorphic to the <u>circle group</u>. It follows easily that the 3-sphere is a principal circle bundle over the 2-sphere, and this is the Hopf fibration.

To make this more explicit, there are two approaches: the group Spin(3) can either be identified with the group Sp(1) of unit quaternions, or with the special unitary group SU(2).

In the first approach, a vector (x_1, x_2, x_3, x_4) in \mathbb{R}^4 is interpreted as a quaternion $q \in \mathbb{H}$ by writing

$$q = x_1 + \mathbf{i}x_2 + \mathbf{j}x_3 + \mathbf{k}x_4.$$

The 3-sphere is then identified with the <u>versors</u>, the quaternions of unit norm, those $q \in \mathbf{H}$ for which $|q|^2 = 1$, where $|q|^2 = q q^*$, which is equal to $x_1^2 + x_2^2 + x_3^2 + x_4^2$ for q as above.

On the other hand, a vector (y_1, y_2, y_3) in \mathbb{R}^3 can be interpreted as an imaginary quaternion

$$p = \mathbf{i}y_1 + \mathbf{j}y_2 + \mathbf{k}y_3.$$

Then, as is well-known since Cayley (1845), the mapping

$$p\mapsto qpq^*$$

is a rotation in \mathbf{R}^3 : indeed it is clearly an <u>isometry</u>, since $|q\ p\ q^*|^2 = q\ p\ q^*\ q\ p^*\ q^* = q\ p\ p^*\ q^* = |p|^2$, and it is not hard to check that it preserves orientation.

In fact, this identifies the group of <u>versors</u> with the group of rotations of \mathbb{R}^3 , modulo the fact that the versors q and -q determine the same rotation. As noted above, the rotations act transitively on S^2 , and the set of versors q which fix a given right versor p have the form q = u + v p, where u and v are real numbers with $u^2 + v^2 = 1$. This is a circle subgroup. For concreteness, one can take $p = \mathbf{k}$, and then the Hopf fibration can be defined as the map sending a versor ω to ω \mathbf{k} ω^* . All the quaternions ωq , where q is one of the circle of versors that fix k, get mapped to the same thing (which happens to be one of the two 180° rotations rotating k to the same place as ω does).

Another way to look at this fibration is that every versor ω moves the plane spanned by $\{1, k\}$ to a new plane spanned by $\{\omega, \omega k\}$. Any quaternion ωq , where q is one of the circle of versors that fix k, will have the same effect. We put all these into one fibre, and the fibres can be mapped one-to-one to the 2-sphere of 180° rotations which is the range of $\omega k \omega^*$.

This approach is related to the direct construction by identifying a quaternion $q = x_1 + \mathbf{i} x_2 + \mathbf{j} x_3 + \mathbf{k} x_4$ with the 2×2 matrix:

$$\begin{bmatrix} x_1 + \mathbf{i}x_2 & x_3 + \mathbf{i}x_4 \\ -x_3 + \mathbf{i}x_4 & x_1 - \mathbf{i}x_2 \end{bmatrix}.$$

This identifies the group of versors with SU(2), and the imaginary quaternions with the skew-hermitian 2×2 matrices (isomorphic to $C\times R$).

Explicit formulae

The rotation induced by a unit quaternion $q = w + \mathbf{i} x + \mathbf{j} y + \mathbf{k} z$ is given explicitly by the orthogonal matrix

$$egin{bmatrix} 1-2(y^2+z^2) & 2(xy-wz) & 2(xz+wy) \ 2(xy+wz) & 1-2(x^2+z^2) & 2(yz-wx) \ 2(xz-wy) & 2(yz+wx) & 1-2(x^2+y^2) \end{bmatrix}.$$

Here we find an explicit real formula for the bundle projection. For, the fixed unit vector along the z axis, (0,0,1), rotates to another unit vector,

$$\Big(2(xz+wy),2(yz-wx),1-2(x^2+y^2)\Big),$$

which is a continuous function of (w, x, y, z). That is, the image of q is where it aims the z axis. The fiber for a given point on S^2 consists of all those unit quaternions that aim there.

To write an explicit formula for the fiber over a point (a, b, c) in S^2 , we may proceed as follows. Multiplication of unit quaternions produces composition of rotations, and

$$q_{\theta} = \cos \theta + \mathbf{k} \sin \theta$$

is a rotation by 2θ around the z axis. As θ varies, this sweeps out a great circle of S^3 , our prototypical fiber. So long as the base point, (a, b, c), is not the antipode, (0, 0, -1), the quaternion

$$q_{(a,b,c)}=rac{1}{\sqrt{2(1+c)}}(1+c-\mathbf{i}b+\mathbf{j}a)$$

will aim there. Thus the fiber of (a, b, c) is given by quaternions of the form $q_{(a, b, c)}q_{\theta}$, which are the S^3 points

$$rac{1}{\sqrt{2(1+c)}} \Big((1+c)\cos(heta), a\sin(heta) - b\cos(heta), a\cos(heta) + b\sin(heta), (1+c)\sin(heta) \Big).$$

Since multiplication by $q_{(a,b,c)}$ acts as a rotation of quaternion space, the fiber is not merely a topological circle, it is a geometric circle. The final fiber, for (0,0,-1), can be given by using $q_{(0,0,-1)} = \mathbf{i}$, producing

$$(0,\cos(\theta),-\sin(\theta),0),$$

which completes the bundle.

Thus, a simple way of visualizing the Hopf fibration is as follows. Any point on the 3-sphere is equivalent to a <u>quaternion</u>, which in turn is equivalent to a <u>particular rotation</u> of a <u>Cartesian coordinate frame</u> in three dimensions. The set of all possible quaternions produces the set of all possible rotations, which moves the tip of one unit vector of such a coordinate frame (say, the **z** vector) to all possible points on a unit 2-sphere. However, fixing the tip of the **z** vector does not specify the rotation fully; a further rotation is possible about the **z**-axis. Thus, the 3-sphere is mapped onto the 2-sphere, plus a single rotation.

Fluid mechanics

If the Hopf fibration is treated as a vector field in 3 dimensional space then there is a solution to the (compressible, non-viscous) Navier-Stokes equations of fluid dynamics in which the fluid flows along the circles of the projection of the Hopf fibration in 3 dimensional space. The size of the velocities, the density and the pressure can be chosen at each point to satisfy the equations. All these quantities fall to zero going away from the centre. If a is the distance to the inner ring, the velocities, pressure and density fields are given by:

$$egin{aligned} \mathbf{v}(x,y,z) &= Aig(a^2+x^2+y^2+z^2ig)^{-2}ig(2(-ay+xz),2(ax+yz),a^2-x^2-y^2+z^2ig)\ p(x,y,z) &= -A^2Big(a^2+x^2+y^2+z^2ig)^{-3},\
ho(x,y,z) &= 3Big(a^2+x^2+y^2+z^2ig)^{-1} \end{aligned}$$

for arbitrary constants A and B. Similar patterns of fields are found as soliton solutions of magnetohydrodynamics:^[3]

Generalizations

The Hopf construction, viewed as a fiber bundle $p: S^3 \to \mathbb{CP}^l$, admits several generalizations, which are also often known as Hopf fibrations. First, one can replace the projective line by an n-dimensional <u>projective space</u>. Second, one can replace the complex numbers by any (real) division algebra, including (for n = 1) the octonions.

Real Hopf fibrations

A real version of the Hopf fibration is obtained by regarding the circle S^1 as a subset of \mathbb{R}^2 in the usual way and by identifying antipodal points. This gives a fiber bundle $S^1 \to \mathbb{RP}^1$ over the <u>real projective line</u> with fiber $S^0 = \{1, -1\}$. Just as \mathbb{CP}^1 is diffeomorphic to a sphere, \mathbb{RP}^1 is diffeomorphic to a circle.

More generally, the *n*-sphere S^n fibers over real projective space \mathbb{RP}^n with fiber S^o .

Complex Hopf fibrations

The Hopf construction gives circle bundles $p: S^{2n+1} \to \mathbb{CP}^n$ over <u>complex projective space</u>. This is actually the restriction of the <u>tautological line bundle</u> over \mathbb{CP}^n to the unit sphere in \mathbb{C}^{n+1} .

Quaternionic Hopf fibrations

Similarly, one can regard S^{4n+3} as lying in \mathbf{H}^{n+1} (quaternionic n-space) and factor out by unit quaternion (= S^3) multiplication to get the quaternionic projective space \mathbf{HP}^n . In particular, since $S^4 = \mathbf{HP}^1$, there is a bundle $S^7 \to S^4$ with fiber S^3 .

Octonionic Hopf fibrations

A similar construction with the <u>octonions</u> yields a bundle $S^{15} \to S^8$ with fiber S^7 . But the sphere S^{31} does not fiber over S^{16} with fiber S^{15} . One can regard S^8 as the <u>octonionic projective lime</u> \mathbf{OP}^1 . Although one can also define an <u>octonionic projective plane</u> \mathbf{OP}^2 , the sphere S^{23} does not fiber over \mathbf{OP}^2 with fiber S^7 . [4][5]

Fibrations between spheres

Sometimes the term "Hopf fibration" is restricted to the fibrations between spheres obtained above, which are

- $S^1 \rightarrow S^1$ with fiber S^0
- $S^3 \rightarrow S^2$ with fiber S^1
- $S^7 \rightarrow S^4$ with fiber S^3
- $S^{15} \rightarrow S^8$ with fiber S^7

As a consequence of <u>Adams' theorem</u>, fiber bundles with <u>spheres</u> as total space, base space, and fiber can occur only in these dimensions. Fiber bundles with similar properties, but different from the Hopf fibrations, were used by <u>John Milnor</u> to construct exotic spheres.

Geometry and applications

The Hopf fibration has many implications, some purely attractive, others deeper. For example, stereographic projection $S^3 \to \mathbf{R}^3$ induces a remarkable structure in \mathbf{R}^3 , which in turn illuminates the topology of the bundle (Lyons 2003). Stereographic projection preserves circles and maps the Hopf fibers to geometrically perfect circles in \mathbf{R}^3 which fill space. Here there is one exception: the Hopf circle containing the projection point maps to a straight line in \mathbf{R}^3 — a "circle through infinity".

The fibers over a circle of latitude on S^2 form a <u>torus</u> in S^3 (topologically, a torus is the product of two circles) and these project to nested <u>toruses</u> in \mathbb{R}^3 which also fill space. The individual fibers map to linking <u>Villarceau circles</u> on these tori, with the exception of the circle through the projection point and the one through its <u>opposite point</u>: the former maps to a straight line, the latter to a unit circle perpendicular to, and centered on, this line, which may be viewed as a degenerate torus whose radius has shrunken to



The fibers of the Hopf fibration stereographically project to a family of Villarceau circles in **R**³.

zero. Every other fiber image encircles the line as well, and so, by symmetry, each circle is linked through *every* circle, both in \mathbb{R}^3 and in S^3 . Two such linking circles form a Hopf link in \mathbb{R}^3

Hopf proved that the Hopf map has <u>Hopf invariant</u> 1, and therefore is not <u>null-homotopic</u>. In fact it generates the <u>homotopy</u> group $\pi_3(S^2)$ and has infinite order.

In <u>quantum mechanics</u>, the Riemann sphere is known as the <u>Bloch sphere</u>, and the Hopf fibration describes the topological structure of a quantum mechanical <u>two-level system</u> or <u>qubit</u>. Similarly, the topology of a pair of entangled two-level systems is given by the Hopf fibration

$$S^3 \hookrightarrow S^7 \to S^4$$
.

(Mosseri & Dandoloff 2001).

The Hopf fibration is equivalent to the fiber bundle structure of the Dirac monopole. [6]

Notes

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External links

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- Dimensions Math (http://dimensions-math.org/Dim_reg_AM.htm) Chapters 7 and 8 illustrate the Hopf fibration with animated computer graphics.
- YouTube animation showing dynamic mapping of points on the 2-sphere to circles in the 3-sphere, by Professor Niles
 Johnson. (https://www.youtube.com/watch?v=AKotMPGFJYk/)
- YouTube animation of the construction of the 120-cell (https://www.youtube.com/watch?v=MFXRRW9goTs/)
 By Gian Marco Todesco shows the Hopf fibration of the 120-cell.
- Video of one 30-cell ring of the 600-cell (http://page.math.tu-berlin.de/~gunn/Movies/600cell.mp4) from http://page.math.tu-berlin.de/~gunn/.

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