**1.** 

(a)

In x, y, z basis

$$Q_{ij} = \begin{pmatrix} -2d^2q & & \\ & -2d^2q & \\ & 4d^2q \end{pmatrix}$$
$$= 2d^2q \begin{pmatrix} -1 & & \\ & -1 & \\ & 2 \end{pmatrix}$$

Potential

$$\Phi = \frac{1}{8\pi\varepsilon_0} \frac{Q_{ij}x_ix_j}{r^5}$$

$$= \frac{d^2q}{4\pi\varepsilon_0} \frac{-x^2 - y^2 + 2z^2}{r^5}$$

$$= \frac{d^2q}{4\pi\varepsilon_0} \frac{3\cos^2\theta - 1}{r^3}$$

Electric field

$$E_r = \frac{3d^2q}{4\pi\varepsilon_0} \frac{3\cos^2\theta - 1}{r^4}$$
 
$$E_\theta = \frac{d^2q}{4\pi\varepsilon_0} \frac{3\sin 2\theta}{r^4}$$

Field line

$$\frac{\mathrm{d}r}{\mathrm{d}\theta} = r \frac{3\cos^2\theta - 1}{\sin 2\theta}$$
$$\frac{\mathrm{d}r}{r} = \frac{3\cos^2\theta - 1}{\sin 2\theta} \mathrm{d}\theta$$
$$r = R_0 \sin \theta \sqrt{\cos \theta}$$

The field line is the furthest when  $\frac{\mathrm{d}r}{\mathrm{d}\theta}=0$ 

$$\theta = \arccos \frac{1}{\sqrt{3}}$$

(b)

Magnetic field

$$\begin{split} \vec{B} &= -\frac{\mu_0}{24\pi c^2 r} \hat{r} \times \left( \hat{r} \cdot \overset{..}{Q} \right) \\ &= -\frac{\mu_0}{24\pi c^2 r} \hat{r} \times \left( \hat{r} \cdot \frac{\mathrm{d}^3 2 d^2 q (3 \hat{z} \hat{z} - 1)}{\mathrm{d} t^3} \right) \\ &= -\frac{q \mu_0 d_0^2}{12\pi c^2 r} \hat{r} \times (3 (\hat{r} \cdot \hat{z}) \hat{z} - \hat{r}) \frac{\mathrm{d}^3 \cos^2 \omega t}{\mathrm{d} t^3} \\ &= \frac{q \mu_0 d_0^2 \omega^3}{\pi c^2 r} (\hat{r} \cdot \hat{z}) (\hat{r} \times \hat{z}) \sin 2\omega t \\ &= -\frac{q \mu_0 d_0^2 \omega^3}{2\pi c^2 r} \sin 2\theta \sin 2\omega t \hat{e}_{\phi} \end{split}$$

Electric field

$$\begin{split} \vec{E} = & c\vec{B} \times \vec{r} \\ = & -\frac{q\mu_0 d_0^2 \omega^3}{2\pi cr} \sin 2\theta \sin 2\omega t \hat{e}_{\theta} \end{split}$$

Power

$$\frac{\mathrm{d}P}{\mathrm{d}\Omega} = \frac{q^2 d_0^4 \omega^6}{8\pi^2 \varepsilon_0 c^5} \sin^2 2\theta$$

The radiation is the strongest for  $\theta = \frac{\pi}{4}$  and  $\theta = \frac{3\pi}{4}$ 

$$\begin{split} P &= \int_0^\pi \mathrm{d}\theta \sin\theta \int_0^{2\pi} \mathrm{d}\phi \frac{q^2 d_0^4 \omega^6}{8\pi^2 \varepsilon_0 c^5} \sin^2 2\theta \\ &= \frac{4q^2 d_0^4 \omega^6}{15\pi \varepsilon_0 c^5} \end{split}$$

The radiation appears at  $2\omega$ 

(c)

The radiation will be a dipole radiation appears at  $\omega$  with total power

$$P = \frac{q^2 d_0^2 \omega^4}{12\pi\varepsilon_0 c^3}$$

Ratio of power

$$\frac{P_{quad}}{P} = \frac{16d_0^2\omega^2}{5c^2}$$

$$\ll 1$$

2.

(a)

$$\begin{split} \vec{B} = & \nabla \times A \\ = & \frac{\mu_0}{4\pi} \int \mathrm{d}^3 \vec{r}' \nabla \times \left( \frac{1}{|\vec{r} - \vec{r}'|} \vec{j} \left( \vec{r}', t - \frac{|\vec{r} - \vec{r}'|}{c} \right) \right) \\ = & \frac{\mu_0}{4\pi} \int \mathrm{d}^3 \vec{r}' \frac{1}{|\vec{r} - \vec{r}'|} \nabla \times \vec{j} \left( \vec{r}', t - \frac{|\vec{r} - \vec{r}'|}{c} \right) + \left( \nabla \frac{1}{|\vec{r} - \vec{r}'|} \right) \times \vec{j} \left( \vec{r}', t - \frac{|\vec{r} - \vec{r}'|}{c} \right) \\ = & \frac{\mu_0}{4\pi} \int \mathrm{d}^3 \vec{r}' \vec{j} \left( \vec{r}', t - \frac{|\vec{r} - \vec{r}'|}{c} \right) \times \frac{\vec{r} - \vec{r}'}{|\vec{r} - \vec{r}'|^3} + \frac{\partial}{\partial t} \vec{j} \left( \vec{r}', t - \frac{|\vec{r} - \vec{r}'|}{c} \right) \times \frac{\vec{r} - \vec{r}'}{c|\vec{r} - \vec{r}'|^2} \end{split}$$

(b)

The ratio of the two term is on the order of (replacing time derivative with T and replace  $|\vec{r} - \vec{r}'|$  with d)

$$\frac{d}{cT} \ll 1$$

3.

(a)

$$\begin{split} P = & \frac{e^4 E_0^2}{12\pi\varepsilon_0 m_e^2 c^3} \\ = & \langle S \rangle \frac{e^4}{6\pi\varepsilon_0^2 m_e^2 c^4} \\ = & \langle S \rangle \frac{8\pi}{3} r_e^2 \end{split}$$

(b)

The radiation (and therefore the stress tensor) has mirror symmetry along each axis.

(c)

$$\begin{split} F = & \frac{P}{c} \\ = & \langle S \rangle \frac{8\pi}{3c} r_e^2 \end{split}$$

4.

$$\begin{split} a_y &= -\frac{eE_0}{m_e}\cos\omega t - \frac{e^2}{6\pi\varepsilon_0 m_e c^3}\dot{a}_y\\ v_y &= -\frac{eE_0}{m_e\omega}\sin\omega t - \frac{e^2}{6\pi\varepsilon_0 m_e c^3}\dot{v}_y\\ \tilde{v}_y &= \mathrm{i}\frac{eE_0}{m_e\omega}\mathrm{e}^{\mathrm{i}\omega t} - \mathrm{i}\omega\frac{e^2}{6\pi\varepsilon_0 m_e c^3}\tilde{v}_y\\ \tilde{v}_y &= \mathrm{i}\frac{eE_0}{m_e\omega}\mathrm{e}^{\mathrm{i}\omega t} - \mathrm{i}\frac{2}{3}\omega\tau_e\tilde{v}_y\\ \tilde{v}_y &= \mathrm{i}\frac{eE_0}{m_e\omega}\frac{1}{1+\mathrm{i}\frac{2}{3}\omega\tau_e}\mathrm{e}^{\mathrm{i}\omega t} \end{split}$$

Force

$$\begin{split} \langle F \rangle = & \frac{1}{2} \Re \left( i e \frac{eE_0}{m_e c \omega} \frac{1}{1 + i \frac{2}{3} \omega \tau_e} e^{i \omega t} E_0 e^{-i \omega t} \right) \\ = & - \frac{1}{2} \Im \left( \frac{e^2 E_0^2}{m_e c \omega} \frac{1}{1 + i \frac{2 \omega r_e}{3c}} \right) \\ = & \frac{1}{2} \frac{e^2 E_0^2}{m_e c \omega} \frac{\frac{2 \omega r_e}{3c}}{1 + \frac{4 \omega^2 r_e^2}{9c^2}} \\ \approx & \frac{e^4 E_0^2}{12 \pi \varepsilon_0 m_e^2 c^4} \end{split}$$

**5**.

$$\frac{\sigma_T Lc}{4\pi} \leqslant GMm_p$$

$$L \leqslant \frac{4\pi GMm_p c}{\sigma_T}$$

$$= 1.2 \cdot 10^{31} J \cdot s^{-1} \qquad \text{(for the sun)}$$

6.

(a)

First term

$$\frac{1}{6\pi\varepsilon_0 c^2} \vec{a} \int d^3 x' \int d^3 x \frac{\rho \rho'}{R}$$

Second term

$$-\frac{1}{6\pi\varepsilon_0 c^3}\dot{\vec{a}}\int\mathrm{d}^3x'\int\mathrm{d}^3x\rho\rho'$$

The problem is that the first term (EM momentum) diverge.

(b)

$$E = \frac{Q}{\varepsilon_0 l^2}$$
 
$$U_E = \frac{Q^2 d}{2\varepsilon_0 l^2}$$

(c)

$$\vec{B} = \frac{\mu_0 QV}{l^2} \hat{z}$$

(d)

$$\begin{split} \vec{P} = & \varepsilon_0 \frac{Q}{\varepsilon_0 l^2} \frac{\mu_0 QV}{l^2} \hat{x} l^2 d \\ = & \frac{\mu_0 V Q^2 d}{l^2} \hat{x} \\ = & 2V \frac{U_e}{c^2} \hat{x} \end{split}$$

(e)

Field is pointing to -x in the plate on -y and +x in the plate on +y side.

$$\begin{split} E = & \frac{\mu_0 Qad}{2l^2} \\ F = & -\frac{\mu_0 Q^2 ad}{l^2} \hat{x} \\ = & -2a \frac{U_e}{c^2} \hat{x} \end{split}$$