1.

(a)

$$\langle \alpha | \beta \rangle = \exp\left(-\frac{|\alpha|^2 + |\beta|^2}{2}\right) \sum_{n} \frac{(\alpha^*)^n \beta^n}{n!}$$
$$= \exp\left(-\frac{|\alpha|^2 + |\beta|^2}{2} + \alpha^* \beta\right)$$

(b)

$$\langle n \int |\alpha\rangle \langle \alpha| \frac{\mathrm{d}^{2}\alpha}{\pi} |m\rangle = \mathrm{e}^{-|\alpha|^{2}} \int \frac{|\alpha|^{2n} \alpha^{m-n}}{\sqrt{n!m!}} |\alpha| \frac{\mathrm{d}|\alpha| \mathrm{d}\theta}{\pi}$$
$$= \delta_{mn} \mathrm{e}^{-|\alpha|^{2}} \int \frac{|\alpha|^{2n}}{n!} \mathrm{d}|\alpha|^{2}$$
$$= \delta_{mn}$$

(c)

$$D(\alpha)|0\rangle = \exp\left(\alpha a^{\dagger} - \alpha^* a\right)|0\rangle$$

$$= e^{\alpha a^{\dagger}} e^{-\alpha^* a} e^{\alpha \alpha^*} [a^{\dagger}, a]/2|0\rangle$$

$$= \exp\left(-\frac{|\alpha|^2}{2}\right) e^{\alpha a^{\dagger}}|0\rangle$$

$$= |\alpha\rangle$$

(d)

$$\langle E_x \rangle = E_0 \sin(kz) \langle a + a^{\dagger} \rangle$$

$$= E_0 \sin(kz) (\alpha + \alpha^*)$$

$$\langle E_x^2 \rangle = E_0^2 \sin^2(kz) \langle (a + a^{\dagger})^2 \rangle$$

$$= E_0^2 \sin^2(kz) \langle a^2 + 1 + 2a^{\dagger}a + a^{\dagger^2} \rangle$$

$$= E_0^2 \sin^2(kz) ((\alpha + \alpha^*)^2 + 1)$$

$$\sqrt{\langle \Delta E_x^2 \rangle} = E_0 \sin(kz)$$

The deviation of the field does not change (even for vacuum) because the coherent state is just a displacement of the vacuum state.

(e)

For number state

$$P(\phi) = \frac{1}{2\pi}$$

so the number state has equal probability of having any phase.

$$P(\phi, \theta) = \frac{1}{2\pi} \left| \sum_{n} \langle n|e^{-in\phi}| (|\alpha|e^{i\theta}) \rangle \right|^{2}$$
$$= \frac{1}{2\pi} \left| \sum_{n} e^{-in(\phi-\theta)} \langle n|(|\alpha|) \rangle \right|^{2}$$
$$= P(\phi - \theta, 0)$$

Therefore P rotates in the same way as the phase of α changes

$$P(\phi, 0) = \frac{1}{2\pi} \left| \sum_{n} \langle n | e^{-in\phi} | (|\alpha|) \rangle \right|^{2}$$
$$= \frac{e^{-|\alpha|^{2}}}{2\pi} \left| \sum_{n} e^{-in\phi} \frac{|\alpha|^{n}}{\sqrt{n!}} \right|^{2}$$

So $P(\phi,0)$ is symmetrically distributed and maximized at $\phi=0$ (therefore $P(\phi,\theta)$ is maximized when $\phi=\theta$) if $|\alpha|\neq 0$. Therefore the expected value of ϕ (when $\alpha\neq 0$) is θ .

2.

(a)

$$\begin{split} \frac{V_{\psi}}{V_{0}} &= \frac{1}{4} \langle \psi, 0 | \left(a^{\dagger} + b^{\dagger} \right) (a+b) \left(a^{\dagger} - b^{\dagger} \right) (a-b) | \psi, 0 \rangle \\ &= \frac{1}{4} \langle \psi, 0 | a^{\dagger} \left(a a^{\dagger} - a b^{\dagger} + b a^{\dagger} - b b^{\dagger} \right) a | \psi, 0 \rangle \\ &= \frac{1}{4} \langle a^{\dagger} a^{\dagger} a a \rangle \end{split}$$

(b)

$$\begin{split} g_{cl}^{(2)}(0) = & \frac{\left\langle I^2 \right\rangle}{\left\langle I \right\rangle^2} \\ = & 1 + \frac{\left\langle I^2 \right\rangle - \left\langle I \right\rangle^2}{\left\langle I \right\rangle^2} \\ = & 1 + \frac{\left\langle \left(I - \left\langle I \right\rangle\right)^2 \right\rangle}{\left\langle I \right\rangle^2} \\ \geqslant & 1 \end{split}$$

(c)

$$g_{\alpha}^{(2)} = \frac{|a|^4}{|a|^4}$$
=1
$$g_{n=2}^{(2)} = \frac{2 \cdot 3}{2^2}$$

$$= \frac{3}{2}$$

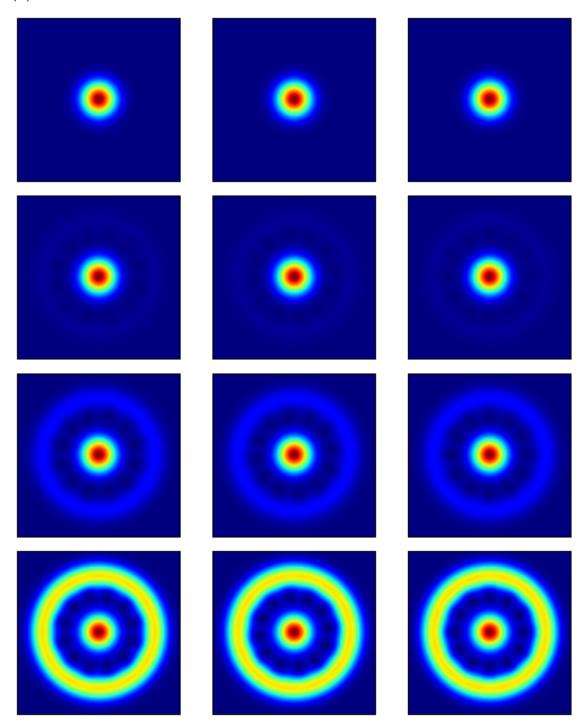
(d)

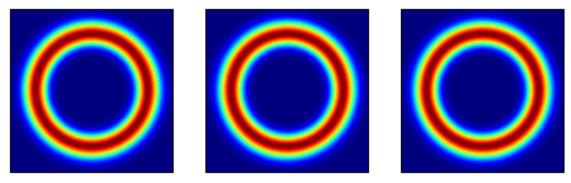
$$\begin{split} \langle \psi_3 | \psi_3 \rangle &= \left(1 + \mathrm{e}^{2|\alpha|^2} \right) \\ \langle \psi_4 | \psi_4 \rangle &= \left(1 - \mathrm{e}^{2|\alpha|^2} \right) \\ g_{3,4}^{(2)} &= \frac{\langle \psi_{3,4} | a^\dagger a^\dagger a a | \psi_{3,4} \rangle \langle \psi_{3,4} | \psi_{3,4} \rangle}{\langle \psi_{3,4} | a^\dagger a | \psi_{3,4} \rangle^2} \\ &= \frac{\left(\langle \alpha | \pm \langle -\alpha | \rangle a^\dagger a^\dagger a a (|\alpha \rangle \pm | -\alpha \rangle) \left(1 \pm \mathrm{e}^{2|\alpha|^2} \right)}{\left(\langle \alpha | \pm \langle -\alpha | \rangle a^\dagger a (|\alpha \rangle \pm | -\alpha \rangle) \left(1 \pm \mathrm{e}^{2|\alpha|^2} \right)} \\ &= \frac{|\alpha|^4 (\langle \alpha | \pm \langle -\alpha | \rangle (|\alpha \rangle \pm | -\alpha \rangle) \left(1 \pm \mathrm{e}^{2|\alpha|^2} \right)}{|\alpha|^4 (\langle \alpha | \mp \langle -\alpha | \rangle (|\alpha \rangle \mp | -\alpha \rangle)^2} \\ &= \left(\frac{1 \pm \mathrm{e}^{2|\alpha|^2}}{1 \mp \mathrm{e}^{2|\alpha|^2}} \right)^2 \end{split}$$

Therefore $g^{(2)}$ can be smaller than 1 for ψ_4 .

3.

(a)

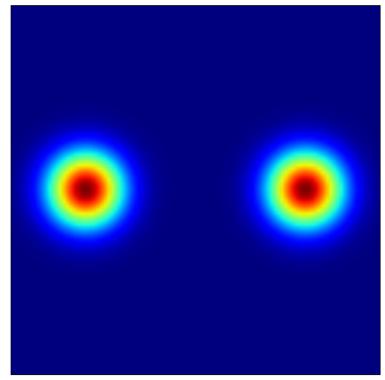




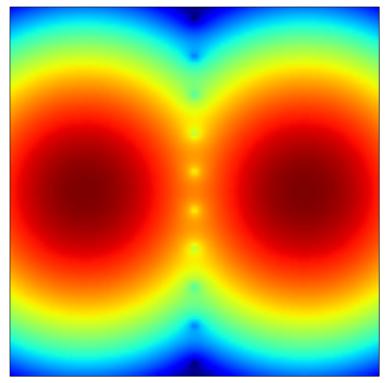
These are not minimum uncertainty states (except when it's $|0\rangle$) and not squeezed states.

(b)

 $\alpha=3\mathrm{i}$ Linear scale

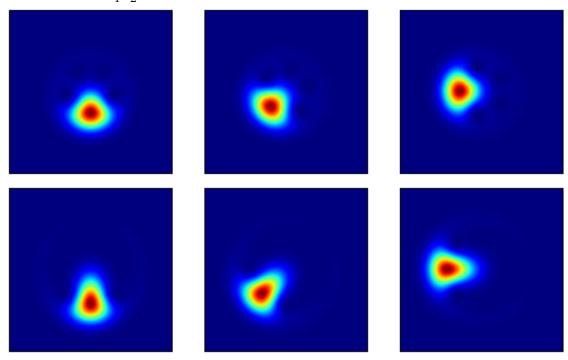


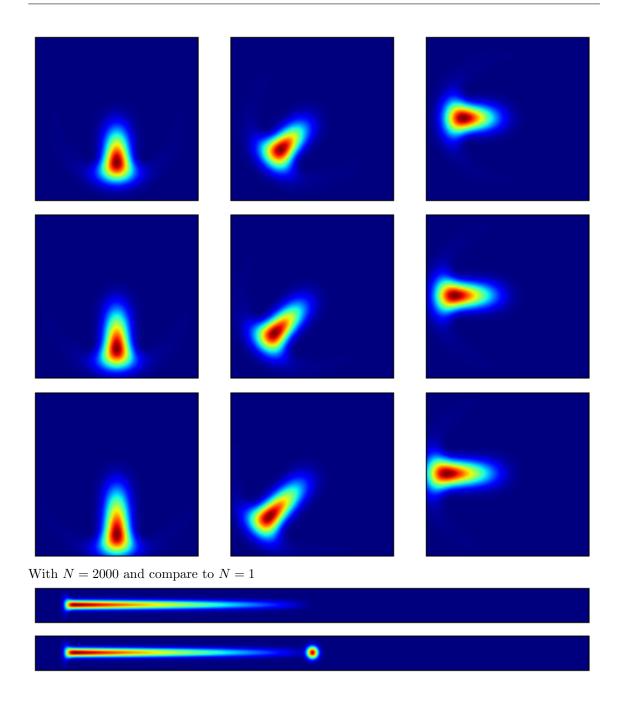
Log scale



Interference fringes appears because the two wavefunction overlaps.

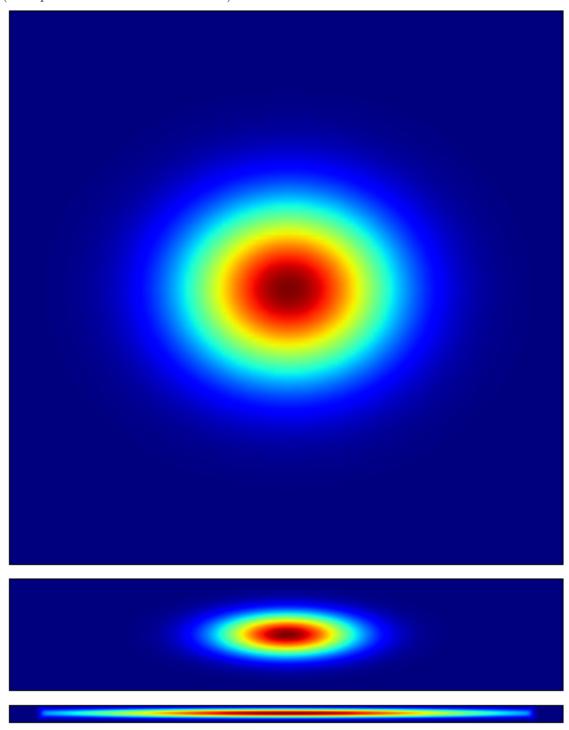
States with
$$\phi = 0, \frac{\pi}{4}, \frac{\pi}{2}$$
 and $N = 5, 10, 15, 20, 15$





(d)

(Transposed and not at the same scale)



(e)