1.

(a)

$$\begin{split} H = &\hbar\omega b^{\dagger}b + \mathrm{i}\hbar\Lambda \left(b^{\dagger^2} + b^2 \right) \\ \frac{\mathrm{d}b}{\mathrm{d}t} = &\frac{\mathrm{i}}{\hbar}[H,b] + \frac{\partial b}{\partial t} \\ = &\left[\mathrm{i}\omega b^{\dagger}b - \Lambda b^{\dagger^2}, b \right] + \mathrm{i}\omega b \\ = &2\Lambda b^{\dagger} \\ \frac{\mathrm{d}b^{\dagger}}{\mathrm{d}t} = &2\Lambda b \end{split}$$

(b)

$$\begin{split} \vec{E} &= \mathrm{i} \mathcal{E}_{\omega} \vec{\varepsilon} \Big(b \mathrm{e}^{\mathrm{i} (\vec{k} \cdot \vec{r} - \omega t)} - b^{\dagger} \mathrm{e}^{-\mathrm{i} (\vec{k} \cdot \vec{r} - \omega t)} \Big) \\ &= \mathrm{i} \mathcal{E}_{\omega} \vec{\varepsilon} \Big(b \cos \Big(\vec{k} \cdot \vec{r} - \omega t \Big) + \mathrm{i} b \sin \Big(\vec{k} \cdot \vec{r} - \omega t \Big) - b^{\dagger} \cos \Big(\vec{k} \cdot \vec{r} - \omega t \Big) + \mathrm{i} b^{\dagger} \sin \Big(\vec{k} \cdot \vec{r} - \omega t \Big) \Big) \\ &= -2 \mathcal{E}_{\omega} \vec{\varepsilon} \Big(\frac{b - b^{\dagger}}{2 \mathrm{i}} \cos \Big(\vec{k} \cdot \vec{r} - \omega t \Big) + \frac{b + b^{\dagger}}{2} \sin \Big(\vec{k} \cdot \vec{r} - \omega t \Big) \Big) \\ &= -2 \mathcal{E}_{\omega} \vec{\varepsilon} \Big(b_Q \cos \Big(\vec{k} \cdot \vec{r} - \omega t \Big) + b_P \sin \Big(\vec{k} \cdot \vec{r} - \omega t \Big) \Big) \\ &= -2 \mathcal{E}_{\omega} \vec{\varepsilon} \Big(b_Q \cos \Big(\vec{k} \cdot \vec{r} - \omega t \Big) + b_P \sin \Big(\vec{k} \cdot \vec{r} - \omega t \Big) \Big) \\ &= -2 \mathcal{E}_{\omega} \vec{\varepsilon} \Big(b_Q \cos \Big(\vec{k} \cdot \vec{r} - \omega t \Big) + b_P \sin \Big(\vec{k} \cdot \vec{r} - \omega t \Big) \Big) \\ &= -2 \mathcal{E}_{\omega} \vec{\varepsilon} \Big(b_Q \cos \Big(\vec{k} \cdot \vec{r} - \omega t \Big) + b_P \sin \Big(\vec{k} \cdot \vec{r} - \omega t \Big) \Big) \\ &= -2 \mathcal{E}_{\omega} \vec{\varepsilon} \Big(b_Q \cos \Big(\vec{k} \cdot \vec{r} - \omega t \Big) + b_P \sin \Big(\vec{k} \cdot \vec{r} - \omega t \Big) \Big) \\ &= -2 \mathcal{E}_{\omega} \vec{\varepsilon} \Big(b_Q \cos \Big(\vec{k} \cdot \vec{r} - \omega t \Big) + b_P \sin \Big(\vec{k} \cdot \vec{r} - \omega t \Big) \Big) \\ &= -2 \mathcal{E}_{\omega} \vec{\varepsilon} \Big(b_Q \cos \Big(\vec{k} \cdot \vec{r} - \omega t \Big) + b_P \sin \Big(\vec{k} \cdot \vec{r} - \omega t \Big) \Big) \\ &= -2 \mathcal{E}_{\omega} \vec{\varepsilon} \Big(b_Q \cos \Big(\vec{k} \cdot \vec{r} - \omega t \Big) + b_P \sin \Big(\vec{k} \cdot \vec{r} - \omega t \Big) \Big) \end{aligned}$$

Therefore

$$\begin{split} b_P = & \mathrm{e}^{2\Lambda t} b_{P0} \\ b_Q = & \mathrm{e}^{-2\Lambda t} b_{Q0} \\ b = & b_P + \mathrm{i} b_Q \\ = & \mathrm{e}^{2\Lambda t} b_{P0} + \mathrm{i} \mathrm{e}^{-2\Lambda t} b_{Q0} \\ = & b_0 \cosh 2\Lambda t + b_0^\dagger \sinh 2\Lambda t \\ b^\dagger = & b_0^\dagger \cosh 2\Lambda t + b_0 \sinh 2\Lambda t \end{split}$$

(c)

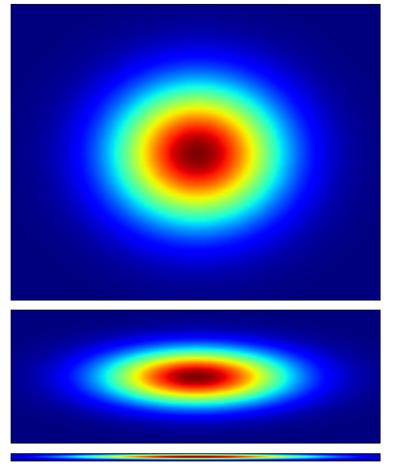
$$\begin{split} \langle N \rangle = & \langle 0|b^{\dagger}b|0 \rangle \\ = & \langle 0| \left(b_0^{\dagger} \cosh 2\Lambda t + b_0 \sinh 2\Lambda t \right) \left(b_0 \cosh 2\Lambda t + b_0^{\dagger} \sinh 2\Lambda t \right) |0 \rangle \\ = & \sinh^2 2\Lambda t \\ \Delta b_P = & \mathrm{e}^{2\Lambda t} \Delta b_{P0} \\ = & \frac{1}{2} \mathrm{e}^{2\Lambda t} \\ \Delta b_Q = & \mathrm{e}^{-2\Lambda t} \Delta b_{Q0} \\ = & \frac{1}{2} \mathrm{e}^{-2\Lambda t} \end{split}$$

The state is squeezed in Q direction while the product of the uncertainty in P and Q remains the same.

(d)

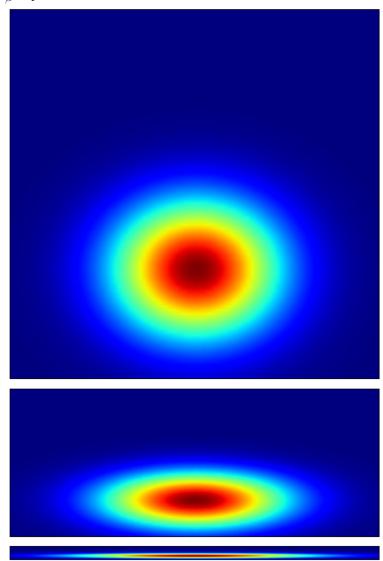
(e)

i.



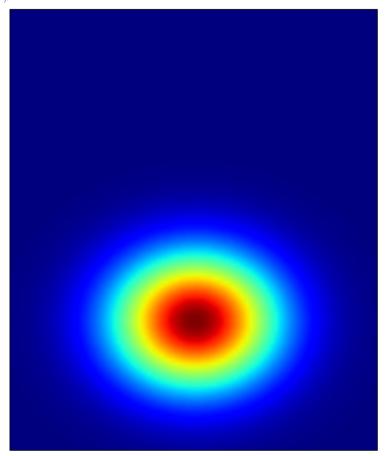
ii.

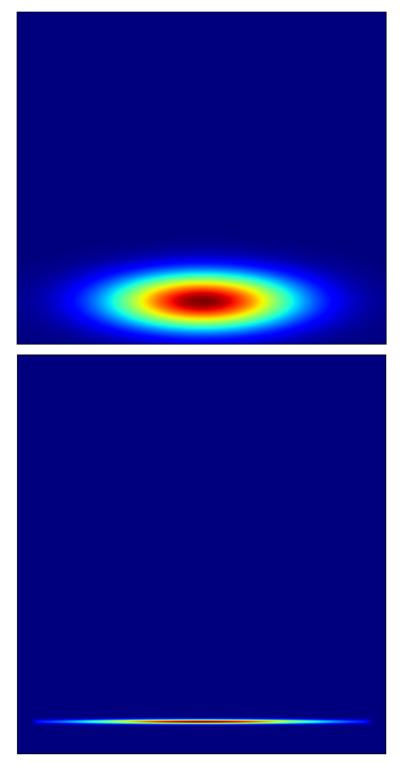
 $\beta = i$



iii.

 $\beta = i$





Weither a phase or number squeezed state is created depends on the initial state.

- 2.
- (a)
- (b)
- (c)
- i.
- ii.
- iii.
- iv.
- 3.
- (a)
- (b)
- (c)
- (d)