1. Classical Model of the Light Force

(a) Time averaged force

$$\begin{split} \left\langle \vec{F} \right\rangle = & \langle (\hat{p} \cdot \hat{\epsilon})(u\cos(\omega t + \theta) - v\sin(\omega t + \theta)) \nabla (E_0\cos(\omega t + \theta)) \rangle \\ = & \langle (\hat{p} \cdot \hat{\epsilon})(u\cos(\omega t + \theta) - v\sin(\omega t + \theta))(\cos(\omega t + \theta) \nabla E_0 - E_0\sin(\omega t + \theta) \nabla \theta) \rangle \\ = & \frac{1}{2}(\hat{p} \cdot \hat{\epsilon})(u\nabla E_0 + vE_0\nabla \theta) \end{split}$$

(b) The potential picture

$$\langle U \rangle = -\left\langle \vec{p} \cdot \vec{E} \right\rangle$$

$$= -\left\langle (\hat{p} \cdot \hat{\epsilon})(u\cos(\omega t + \theta) - v\sin(\omega t + \theta))E_0\cos(\omega t + \theta) \right\rangle$$

$$= -\frac{1}{2}(\hat{p} \cdot \hat{\epsilon})uE_0$$

$$\left\langle \vec{F} \right\rangle = \frac{1}{2}(\hat{p} \cdot \hat{\epsilon})u\nabla E_0$$

(c) Dipole moment of electron

Let $\vec{r} = \hat{\epsilon} \tilde{r}_0 e^{i\omega t}$

$$-e\hat{\epsilon}E_0e^{\mathrm{i}\omega t + \mathrm{i}\theta}$$

$$= (-m\omega^2 + \mathrm{i}\omega\gamma + m\omega_0^2)\hat{\epsilon}\tilde{r}_0e^{\mathrm{i}\omega t}$$

$$\tilde{r}_0 = \frac{eE_0e^{\mathrm{i}\theta}}{m\omega^2 - \mathrm{i}\omega\gamma - m\omega_0^2}$$

$$\tilde{\vec{p}} = -\hat{\epsilon}e^{\mathrm{i}\omega t}\frac{e^2E_0e^{\mathrm{i}\theta}}{m\omega^2 - \mathrm{i}\omega\gamma - m\omega_0^2}$$

Real part

$$\begin{split} \vec{p} &= -\hat{\epsilon} \Re \left(\mathrm{e}^{\mathrm{i}\omega t + \mathrm{i}\theta} \frac{e^2 E_0}{m\omega^2 - \mathrm{i}\omega\gamma - m\omega_0^2} \right) \\ &= -\hat{\epsilon} \left(\cos\left(\omega t + \theta\right) \Re \left(\frac{e^2 E_0}{m\omega^2 - \mathrm{i}\omega\gamma - m\omega_0^2} \right) - \sin\left(\omega t + \theta\right) \Im \left(\frac{e^2 E_0}{m\omega^2 - \mathrm{i}\omega\gamma - m\omega_0^2} \right) \right) \end{split}$$

Where

$$\begin{split} &\frac{e^2 E_0}{m\omega^2 - \mathrm{i}\omega\gamma - m\omega_0^2} \\ = &e^2 E_0 \frac{1}{m(\omega^2 - \omega_0^2) - \mathrm{i}\omega\gamma} \\ \approx &\frac{e^2 E_0}{m\omega_0} \frac{1}{2\delta - \mathrm{i}\Gamma} \\ = &\frac{e^2 E_0}{m\omega_0} \frac{2\delta + \mathrm{i}\Gamma}{4\delta^2 + \Gamma^2} \end{split}$$

Compare to the definition of u and v

$$\begin{split} u &= -\Re\left(\frac{e^2 E_0}{m\omega^2 - \mathrm{i}\omega\gamma - m\omega_0^2}\right) \\ &= -\frac{e^2 E_0}{m\omega_0} \frac{2\delta}{4\delta^2 + \Gamma^2} \\ v &= -\Im\left(\frac{e^2 E_0}{m\omega^2 - \mathrm{i}\omega\gamma - m\omega_0^2}\right) \\ &= -\frac{e^2 E_0}{m\omega_0} \frac{\Gamma}{4\delta^2 + \Gamma^2} \end{split}$$

Force

$$\begin{split} \left\langle \vec{F} \right\rangle &= -\frac{e^2}{2m\omega_0} \frac{2E_0 \nabla E_0 \delta + \Gamma E_0^2 \nabla \theta}{4\delta^2 + \Gamma^2} \\ &= -\frac{e^2}{2m\omega_0} \frac{\delta \nabla E_0^2 + \Gamma E_0^2 \nabla \theta}{4\delta^2 + \Gamma^2} \end{split}$$

(d) Force on a two-level atom

Since $\omega_R \propto E_0$

$$\left\langle \vec{F} \right\rangle \propto -\frac{\delta \nabla \omega_R^2 + \Gamma \omega_R^2 \nabla \theta}{4\delta^2 + \Gamma^2}$$
$$\propto F_{quantum}$$

2. Master equation for a damped optical cavity

(a)

$$\langle \tilde{\psi}_1 | \tilde{\psi}_1 \rangle = \Gamma dt \langle \psi | a^{\dagger} a | \psi \rangle$$

$$= dn$$

Since the imaginary and real part of H commute

$$\begin{split} \langle \tilde{\psi}_0 | \tilde{\psi}_0 \rangle &= \langle \psi | \mathrm{e}^{\mathrm{i} H^\dagger \mathrm{d} t / \hbar} \mathrm{e}^{-\mathrm{i} H \mathrm{d} t / \hbar} | \psi \rangle \\ &= \langle \psi | \mathrm{e}^{-\Gamma \mathrm{d} t a^\dagger a} | \psi \rangle \\ &= \exp \left(-\Gamma \mathrm{d} t \langle \psi | a^\dagger a | \psi \rangle \right) \\ &= \mathrm{e}^{-\mathrm{d} p} \\ &= 1 - \mathrm{d} p \end{split}$$

(b)

$$\begin{split} \rho + \mathrm{d}\rho = & |\tilde{\psi}_1\rangle\langle\tilde{\psi}_1| + |\tilde{\psi}_0\rangle\langle\tilde{\psi}_0| \\ = & \Gamma \mathrm{d}t a |\psi\rangle\langle\psi|a^{\dagger} + \left(1 - \frac{\mathrm{i}H^{\dagger}\mathrm{d}t}{\hbar}\right)|\psi\rangle\langle\psi|\left(1 + \frac{\mathrm{i}H\mathrm{d}t}{\hbar}\right) \\ = & \rho + \Gamma \mathrm{d}t a \rho a^{\dagger} + \rho \frac{\mathrm{i}H^{\dagger}\mathrm{d}t}{\hbar} - \frac{\mathrm{i}H\mathrm{d}t}{\hbar}\rho \\ = & \rho + \Gamma \mathrm{d}t a \rho a^{\dagger} - \frac{\mathrm{i}\mathrm{d}t}{\hbar}[H_0, \rho] - \mathrm{d}t \frac{\Gamma}{2}\rho a^{\dagger}a - \mathrm{d}t \frac{\Gamma}{2}a^{\dagger}a\rho \end{split}$$

(c)

$$\begin{split} \frac{\mathrm{d}\rho}{\mathrm{d}t} = & \Gamma a \rho a^{\dagger} - \frac{\mathrm{i}}{\hbar} [H_0, \rho] - \frac{\Gamma}{2} \rho a^{\dagger} a - \frac{\Gamma}{2} a^{\dagger} a \rho \\ = & -\frac{\mathrm{i}}{\hbar} [H_0, \rho] - \frac{1}{2} \left(\Gamma \rho a^{\dagger} a - 2 \Gamma a \rho a^{\dagger} + \Gamma a^{\dagger} a \rho \right) \\ = & -\frac{\mathrm{i}}{\hbar} [H_0, \rho] - \frac{1}{2} \left(\rho C^{\dagger} C - 2 C \rho C^{\dagger} + C^{\dagger} C \rho \right) \end{split}$$

where $C = \sqrt{\Gamma}a$