## 1. The Dressed Atom

- (a)
- (b)
- (c)
- (d)
- (e)
- (f)

## 2. Sideband Cooling

(a)

Interaction term

$$H_{I} = \hbar\Omega(\sigma_{+} + \sigma_{-})\cos(k\hat{x} - \omega t)$$

$$= \hbar\Omega(\sigma_{+} + \sigma_{-})\cos(\eta(a + a^{\dagger}) - \omega t)$$

$$= \frac{\hbar\Omega}{2}(\sigma_{+} + \sigma_{-})\left(e^{i(\eta(a + a^{\dagger}) - \omega t)} + e^{-i(\eta(a + a^{\dagger}) - \omega t)}\right)$$

In the Lamb-Dicke limit

$$\begin{split} H_{I} &= \frac{\hbar\Omega}{2} (\sigma_{+} + \sigma_{-}) \Big( \mathrm{e}^{\mathrm{i}\eta \left( a + a^{\dagger} \right)} \mathrm{e}^{-\mathrm{i}\omega t} + \mathrm{e}^{-\mathrm{i}\eta \left( a + a^{\dagger} \right)} \mathrm{e}^{\mathrm{i}\omega t} \Big) \\ &= \frac{\hbar\Omega}{2} (\sigma_{+} + \sigma_{-}) \Big( \Big( 1 + \mathrm{i}\eta \left( a + a^{\dagger} \right) \Big) \mathrm{e}^{-\mathrm{i}\omega t} + \Big( 1 - \mathrm{i}\eta \left( a + a^{\dagger} \right) \Big) \mathrm{e}^{\mathrm{i}\omega t} \Big) \\ &= \frac{\hbar\Omega}{2} (\sigma_{+} + \sigma_{-}) \Big( \mathrm{e}^{-\mathrm{i}\omega t} + \mathrm{e}^{\mathrm{i}\omega t} \Big) + \mathrm{i}\eta \frac{\hbar\Omega}{2} (\sigma_{+} + \sigma_{-}) \Big( a + a^{\dagger} \Big) \Big( \mathrm{e}^{-\mathrm{i}\omega t} - \mathrm{e}^{\mathrm{i}\omega t} \Big) \end{split}$$

In the interaction picture

$$\begin{split} \sigma'_{\pm} &= \mathrm{e}^{\mathrm{i} H_0 t/\hbar} \sigma_{\pm} \mathrm{e}^{-\mathrm{i} H_0 t/\hbar} \\ &= \sigma_{\pm} \mathrm{e}^{\mathrm{i} \omega_0 t} \\ a' &= \mathrm{e}^{\mathrm{i} H_0 t/\hbar} a \mathrm{e}^{-\mathrm{i} H_0 t/\hbar} \\ &= a \mathrm{e}^{-\mathrm{i} \nu t} \\ a'^{\dagger} &= \mathrm{e}^{\mathrm{i} H_0 t/\hbar} a^{\dagger} \mathrm{e}^{-\mathrm{i} H_0 t/\hbar} \\ &= a^{\dagger} \mathrm{e}^{\mathrm{i} \nu t} \end{split}$$

In the rotating wave approximation

$$\begin{split} H_I' = & \mathrm{e}^{\mathrm{i} H_0 t/\hbar} H_I \mathrm{e}^{-\mathrm{i} H_0 t/\hbar} \\ = & \frac{\hbar \Omega}{2} \left( \sigma_+ \mathrm{e}^{-\mathrm{i} \delta t} + \sigma_- \mathrm{e}^{\mathrm{i} \delta t} \right) + \mathrm{i} \eta \frac{\hbar \Omega}{2} \left( \sigma_+ a \mathrm{e}^{-\mathrm{i} (\delta + \nu) t} - \sigma_- a^\dagger \mathrm{e}^{\mathrm{i} (\delta + \nu) t} \right) \\ & + \mathrm{i} \eta \frac{\hbar \Omega}{2} \left( \sigma_+ a^\dagger \mathrm{e}^{-\mathrm{i} (\delta - \nu) t} - \sigma_- a \mathrm{e}^{\mathrm{i} (\delta - \nu) t} \right) \end{split}$$

When the light is in resonance with the sideband, the Rabi frequencies

$$\begin{split} \Omega_{n,n-1} &= \left| \langle n-1, e | \mathrm{i} \eta \frac{\hbar \Omega}{2} \left( \sigma_{+} a - \sigma_{-} a^{\dagger} \right) | n, g \rangle \right| \\ &= \frac{\eta \hbar \Omega}{2} \langle n-1 | a | n \rangle \\ &= \frac{\eta \hbar \Omega}{2} \sqrt{n} \\ \Omega_{n,n+1} &= \left| \langle n+1, e | \mathrm{i} \eta \frac{\hbar \Omega}{2} \left( \sigma_{+} a^{\dagger} - \sigma_{-} a \right) | n, g \rangle \right| \\ &= \frac{\eta \hbar \Omega}{2} \langle n+1 | a^{\dagger} | n \rangle \\ &= \frac{\eta \hbar \Omega}{2} \sqrt{n+1} \end{split}$$

(b)

For spontanious decay, when the emitted photon is  $\theta$  from the axis of the trap, the probability to go from n to n+1 is  $\eta^2 \cos^2 \theta(n+1)$  Assuming isotropic emission patter (i.e. non-polarized) the averate probability is

$$p_{n,n+1} = \frac{1}{4\pi} \int_0^{\pi} d\theta \int_0^{2\pi} \sin\theta d\phi \eta^2 \cos^2\theta (n+1)$$
$$= \frac{\eta^2 (n+1)}{2} \int_{-1}^1 dz z^2$$
$$= \frac{\eta^2 (n+1)}{3}$$

Similarly, the probability of going to n-1 is

$$p_{n,n-1} = \frac{\eta^2 n}{3}$$

To the lowest order in  $\eta$ , the rate at which  $|n,g\rangle \to |n,e\rangle \to |n-1,g\rangle$  happens is

$$\frac{n\Omega^2\eta^2}{3}\frac{\Gamma}{\Gamma^2+4\delta^2}$$

For  $|n,g\rangle \to |n-1,e\rangle \to |n-1,g\rangle$ 

$$n\Omega^2 \eta^2 \frac{\Gamma}{\Gamma^2 + 4(\delta + \nu)^2}$$
$$A_- = \Omega^2 \Gamma \eta^2 \left( \frac{1}{\Gamma^2 + 4(\delta + \nu)^2} + \frac{1}{3} \frac{1}{\Gamma^2 + 4\delta^2} \right)$$

Similarly

$$A_{+} = \Omega^{2} \Gamma \eta^{2} \left( \frac{1}{\Gamma^{2} + 4(\delta - \nu)^{2}} + \frac{1}{3} \frac{1}{\Gamma^{2} + 4\delta^{2}} \right)$$

(c)

$$\frac{\mathrm{d}\langle n \rangle}{\mathrm{d}t} = \sum_{n=0}^{\infty} n \frac{\mathrm{d}p_n}{\mathrm{d}t}$$

$$= \sum_{n=0}^{\infty} n(nA_+p_{n-1} + (n+1)A_-p_{n+1} - (n+1)A_+p_n - nA_-p_n)$$

$$= \sum_{n=0}^{\infty} \left( n^2A_+p_{n-1} + n(n+1)A_-p_{n+1} - n(n+1)A_+p_n - n^2A_-p_n \right)$$

$$= \sum_{n=0}^{\infty} \left( (n+1)^2A_+p_n + (n-1)nA_-p_n - n(n+1)A_+p_n - n^2A_-p_n \right)$$

$$= \sum_{n=0}^{\infty} \left( (n+1)A_+p_n - nA_-p_n \right)$$

$$= A_+ - (A_- - A_+)\langle n \rangle$$

The solution is a exponential decay to  $\langle n \rangle = \frac{A_+}{A_- - A_+}$  with decay rate  $A_- - A_+$ . For driving cooling sideband in the resolved limit

$$A_{-} \approx \frac{\Omega^{2} \eta^{2}}{\Gamma}$$

$$A_{+} \approx \frac{7\Omega^{2} \Gamma \eta^{2}}{48\nu^{2}}$$

Final temperature

$$\begin{split} \langle n \rangle_{\infty} \approx & \frac{A_{+}}{A_{-}} \\ &= \frac{7\Gamma^{2}}{48\nu^{2}} \\ T_{\infty} = & \langle n \rangle_{\infty} \frac{\hbar\nu}{k_{B}} \\ \approx & \frac{7\hbar\Gamma^{2}}{48k_{B}\nu} \end{split}$$

Decay time

$$\tau \approx \frac{1}{A_{-}}$$
 
$$\approx \frac{\Gamma}{\Omega^{2} \eta^{2}}$$

(d)

For the narrow line

$$\begin{split} \eta = & k \sqrt{\frac{\hbar}{2m\nu}} \\ = & 0.022 \\ T_{\infty} \approx & \frac{7\hbar\Gamma^2}{48k_B\nu} \\ = & 0.11\text{aK} \\ \tau \approx & \frac{\Gamma}{\Omega^2\eta^2} \\ = & 22\text{hr} \end{split}$$

For broadened line

$$T_{\infty} \approx \frac{7\hbar\Gamma^2}{48k_B\nu}$$

$$= 7.0 \text{nK}$$

$$\tau \approx \frac{\Gamma}{\Omega^2 \eta^2}$$

$$= 0.32 \text{s}$$

## 3. Optical dipole trap

- (a)
- (b)
- (c)