

1. A Zeeman Slower

(a)

The maximum deceleration is achieved when the scattering rate is maximized. Since the maximum population when pumping with a laser in such a two level system is 50% the maximum scattering rate is $\frac{\gamma}{2}$. Maximum deceleration

$$\begin{aligned} a_{max} &= \frac{\gamma}{2} \frac{p_{rec}}{m} \\ &= \frac{h\gamma}{2m\lambda} \\ &= 9.3 \cdot 10^5 \text{ m} \cdot \text{s}^{-2} \end{aligned}$$

(b)

Assuming the atom flux (and therefore the optical depth of the slower) is small so that the intensity of light is almost constant in the slower.

Length of the slower,

$$\begin{aligned} L &= \frac{v_{max}^2}{2a} \\ &= \frac{v_{max}^2}{2fa_{max}} \\ &= \frac{1}{f} \frac{k_B T \lambda}{h\gamma} \\ &= \frac{0.12}{f} \text{ m} \end{aligned}$$

Maximum velocity in the slower

$$v = \sqrt{2fa_{max}(L - x)}$$

Doppler shift

$$\begin{aligned} \delta_{doppler} &= \frac{v}{\lambda} \\ &= \frac{\sqrt{2fa_{max}(L - x)}}{\lambda} \end{aligned}$$

This should be canceled by the Zeeman shift

$$B = \frac{\sqrt{2fa_{max}(L - x)}}{g\mu_B\lambda}$$

(c)

Variance of Δx for each emission for three situations

$$\begin{aligned}\langle \Delta p^2 \rangle_0 &= p_{rec}^2 \begin{cases} 1 & \text{i} \\ \int_{-1}^1 x^2 dx & \text{ii} \\ \int_{-1}^1 x^4 dx & \text{iii} \end{cases} \\ &= p_{rec}^2 \begin{cases} 1 & \text{i} \\ \frac{1}{3} & \text{ii} \\ \frac{1}{5} & \text{iii} \end{cases}\end{aligned}$$

$$\begin{aligned}\mathcal{D} &= \frac{d\langle \Delta p^2 \rangle}{dt} \\ &= \langle \Delta p^2 \rangle_0 \Gamma_s \\ &= \Gamma_s p_{rec}^2 \begin{cases} 1 & \text{i} \\ \frac{1}{3} & \text{ii} \\ \frac{1}{5} & \text{iii} \end{cases}\end{aligned}$$

2. Slowing an atom with off-resonant light

(a)

Scattering rate

$$\begin{aligned}\Gamma_s &= \frac{\gamma}{2} \frac{s}{1 + s + \left(\frac{2\delta}{\gamma}\right)^2} \\ &= \frac{\gamma}{2} \frac{s}{1 + s + \left(\frac{2v}{\lambda\gamma}\right)^2} \\ \frac{dv}{dt} &= -\frac{p_{rec}}{m} \Gamma_s \\ &= -a_{max} \frac{s}{1 + s + \left(\frac{2v}{\lambda\gamma}\right)^2}\end{aligned}$$

Total time

$$\begin{aligned}T &= \int_0^{v_{max}} \frac{1}{sa_{max}} \left(1 + s + \left(\frac{2v}{\lambda\gamma}\right)^2\right) dv \\ &= \frac{v_{max}}{sa_{max}} \left(1 + s + \frac{4v_{max}^2}{3\lambda^2\gamma^2}\right) \\ &= 43.5 \frac{v_{max}}{a_{max}} \\ &= 21.9\text{ms}\end{aligned}$$

(b)

Let t be the time until the atom stops

$$\begin{aligned}
 t &= \frac{v}{sa_{max}} \left(1 + s + \frac{4v^2}{3\lambda^2\gamma^2} \right) \\
 &\approx \\
 v &\approx \sqrt[3]{\frac{3sa_{max}\lambda^2\gamma^2}{4}} t^{1/3} \\
 l &\approx \sqrt[3]{\frac{3sa_{max}\lambda^2\gamma^2}{4}} \int_0^T t^{1/3} dt \\
 &= \sqrt[3]{\frac{3sa_{max}\lambda^2\gamma^2}{4}} \frac{3}{4} T^{4/3} \\
 &\approx 7.8\text{m}
 \end{aligned}$$

(c)

This is much longer and less efficient than a Zeeman slower.

3. Density Limit in a MOT

(a)

The radiation from a single atom is

$$S_{single} = \frac{6\sigma_L I}{4\pi r^2} \hat{r}$$

Force

$$F_{single} = \frac{6\sigma_L \sigma_R I}{4\pi r^2 c} \hat{r}$$

From Gauss' law

$$\nabla \cdot F_R = \frac{6\sigma_L \sigma_R I n}{c}$$

(b)

For an isotropic light field, the light force is proportional to the gradient of light intensity. For the case of scattering force, each atom is acting as a source for the gradient while for attenuation force, each atom is acting as a drain for the gradient (thus the minus sign). Furthermore, the frequency of the light is the trap light so the absorption cross section should be replaced by σ_L .

(c)

$$\begin{aligned}
 \nabla \cdot \vec{F} &= \frac{6\sigma_L(\sigma_R - \sigma_L)In}{c} - 3\kappa \\
 n_{max} &= \frac{3\kappa c}{6\sigma_L(\sigma_R - \sigma_L)I}
 \end{aligned}$$

(d)

On average $\sigma_R > \sigma_L$ due to the resonance component of the scattered light (Mollow triplet).

(e)

This effectively replaces n by $fn < n$ in F_R and F_A so the both F_R and F_A decreases (with same n) while F_T remains the same (and might increase a little due to the lack of attenuation from the atoms). n_{max} will increase by f^{-1} .