

Two photon coupling to atomic relative motional continuum

As the reverse process of the Raman transition that forms the molecule, the ground molecular state can be driven to the atomic state as well. If the energy for such atomic state is above the trap depth of the tweezer, this can couple the molecule to a continuum and therefore lead to loss. Physically, this loss corresponds to a photodissociation process where the two atoms flies apart with high kinetic energy.

The photodissociation rate is determined by the Fermi's golden rule,

$$\Gamma = 2\pi\Omega_{if}^2 \rho_f$$

where Ω_{if} is the matrix element between the initial and final state and ρ_f is the density of state near the final state. We calculate this by first discretizing the continuum and the rate becomes

$$\Gamma = 2\pi \frac{\Omega_{if}'^2}{\delta}$$

where Ω_{if}' is the Rabi frequency between the initial and final state and δ is the spacing between the discretized final states.

The Rabi frequency Ω_{if}' is proportional to the square root of FCF, which we calculate for the atomic ground motional states by exactly diagonalization of the atomic wavefunction. However, due to the high energy (≈ 1 GHz) of the final state, calculating FCF for the target state using this method requires an unrealistic number of states to be included. Instead, we approximate it using the results for the atomic ground state by assuming the FCF to be proportional to the probability density of non-interacting atomic wavefunction at zero relative distance so that,

$$\Omega_{if}' = \Omega_0 \frac{\psi_f(r_{rel} = 0)}{\psi_0(r_{rel} = 0)}$$

where Ω_0 is the Rabi frequency between the excited state and the atomic motional ground state and ψ_f and ψ_0 are the wavefunctions of the final and ground motional atomic state without the molecular potential. This approximation is justified because,

1. The molecular potential is short ranged and the excited molecular state is small.
2. Only the final atomic wavefunction within the molecular potential contributes to the FCF. For relatively low motional energy this wavefunction is proportional to the value of the wavefunction at the edge of the molecular potential which is well approximated by the atomic wavefunction without the molecular potential.

This approximation remains valid until the de Broglie wavelength for the atomic state is comparable to the size of the molecular potential which corresponds to a motional energy of ≈ 40 GHz. For atomic motional ground state, we have,

$$\begin{aligned} \psi_0(r_{rel} = 0) &= \prod_{i=1}^3 \left(\frac{\mu\omega_i}{\pi} \right)^{1/4} \\ &= \left(\frac{\mu^3\omega_1\omega_2\omega_3}{\pi^3} \right)^{1/4} \end{aligned}$$

where μ is the reduced mass $\mu \equiv m_{\text{Na}}m_{\text{Cs}}/(m_{\text{Na}} + m_{\text{Cs}})$, and ω_i 's are the relative trapping frequencies along the three trap axis. We discretize the continuum state by adding an infinitely deep spherical potential well of radius R around the center. The radial wavefunctions of the eigenstates within the well with quantum number $n = 1, 2, \dots$ is,

$$\psi_n = \frac{1}{\sqrt{4\pi r_{rel}}} \sqrt{\frac{2}{R}} \sin\left(\frac{\pi n r_{rel}}{R}\right)$$

and the corresponding energy,

$$E_n = \frac{\pi^2 n^2}{2\mu R^2}$$

The energy gap between neighboring states is

$$\delta_n \approx \frac{\pi^2 n}{\mu R^2}$$

and the wavefunction value at $r_{rel} = 0$ is,

$$\begin{aligned} \psi_n(r_{rel} = 0) &= \sqrt{\frac{1}{2\pi R}} \frac{d}{dr_{rel}} \sin\left(\frac{\pi n r_{rel}}{R}\right) \Big|_{r_{rel}=0} \\ &= \frac{\pi n}{\sqrt{2\pi R R}} \end{aligned}$$

Substituting into Γ and taking the limit of $R \rightarrow \infty$ we have

$$\begin{aligned} \Gamma &= 2\pi \frac{\Omega_0^2}{\delta} \frac{\psi_f^2(r_{rel} = 0)}{\psi_0^2(r_{rel} = 0)} \\ &= \Omega_0^2 \frac{n}{R} \sqrt{\frac{\pi^3}{\mu \omega_1 \omega_2 \omega_3}} = \Omega_0^2 \sqrt{\frac{2\pi E}{\omega_1 \omega_2 \omega_3}} \end{aligned}$$

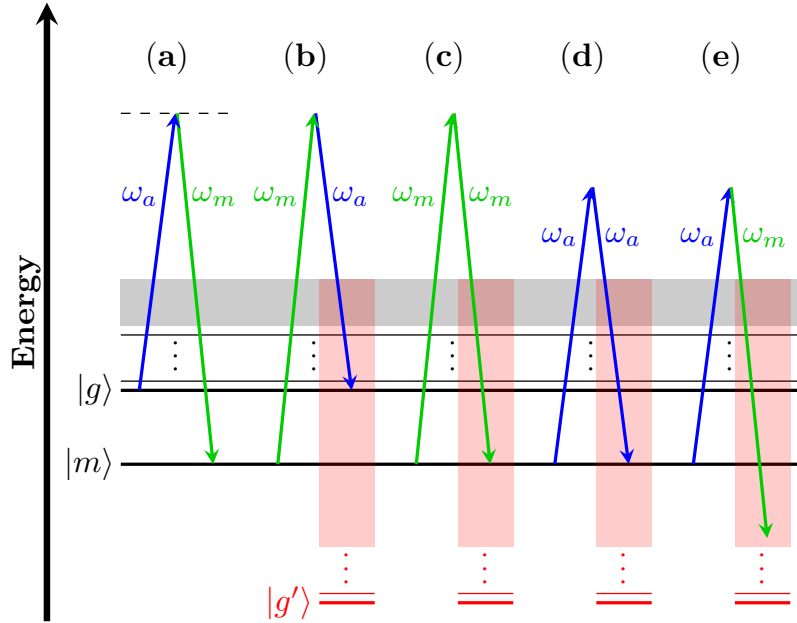


FIG. S1: Effect of two-photon coupling to atomic continuum on molecular state. (a) Raman transition from the atomic ground state $|g\rangle$ to the molecular bound state $|m\rangle$ using two Raman beams with frequencies ω_1 and ω_2 . (b-e) Two photon transitions starting from the molecular state $|m\rangle$ with different frequency combinations without additional atomic spin states. (b) This is the reverse Raman transition and reaches the initial atomic state. (c,d) The frequency pairs couples back to $|m\rangle$. (e) This couples to an energy below $|m\rangle$. Note that none of the two-photon transitions from $|m\rangle$ in (b-e) ends up in the atomic motional continuum (gray) for the spin state $|g\rangle$. However, it is coupled to the motional continuum for the spin state with lower energy $|g'\rangle$ and may be affected by the corresponding dissociative scattering.

As shown in Fig. S1, since we only have the two Raman frequencies in the experiment, the energy of the final state is never higher than that of the initial atomic state and well below the corresponding trap depth. Therefore, in order to couple to the continuum, the final state must belong to a different atomic spin state with lower energy. For our initial spin state of Na(2,2), Cs(3,3), the only candidate for such state is therefore Na(1,1), Cs(3,3), which is lower in energy by ≈ 1.77 GHz.