1.

2.

(a)

Since  $J_z = L_z + S_z$  and  $[L_z, S_z] = 0$ , the  $J_z$  eigenstates are sum of  $L_z$  and  $S_z$  eigenstates that has the same sum of  $m_l$  and  $m_s$ . (The only variable on the RHS are  $m_l$  and  $m_s$  among which  $m_s$  can only be  $\pm \frac{1}{2}$  so the constraint above limit the decomposition to the given form.)

(b)

Hamiltonian

$$\begin{split} H = & \frac{L^2}{2ma^2} + V_0 + \frac{e^2 \vec{L} \cdot \vec{S}}{2mc^2 a^3} \\ = & \frac{L^2}{2ma^2} + V_0 + \frac{e^2}{4mc^2 a^3} \big( J^2 - L^2 - S^2 \big) \\ = & \frac{l(l+1)}{2ma^2} + V_0 + \frac{e^2}{4mc^2 a^3} \bigg( j(j+1) - l(l+1) - \frac{3}{4} \bigg) \end{split}$$

When l = 0 there's only one manifold instead of two.

**3.** 

$$e^{\lambda A}Be^{-\lambda A} = \left(\sum_{m=0}^{\infty} \frac{\lambda^m A^m}{m!}\right) B\left(\sum_{n=0}^{\infty} \frac{\lambda^n (-A)^n}{n!}\right)$$

$$= \sum_{m,n=0}^{\infty} \frac{\lambda^m A^m}{m!} B\frac{\lambda^n (-A)^n}{n!}$$

$$= \sum_{m=0}^{\infty} \sum_{n=0}^{m} \frac{\lambda^m A^{m-n} B(-A)^n}{(m-n)! n!}$$

$$= \sum_{m=0}^{\infty} \frac{\lambda^m}{m!} \sum_{n=0}^{m} \frac{m! A^{m-n} B(-A)^n}{(m-n)! n!}$$

So now we just need to show that

$$[A^{(m)}, B] = \sum_{n=0}^{m} \frac{m! A^{m-n} B(-A)^n}{(m-n)! n!}$$

where  $[A^{(0)}, B] \equiv B$ ,  $[A^{(n)}, B] \equiv [A, [A^{(n-1)}, B]]$ . For m = 0, this is trivially true. If it is true for m - 1,

$$\begin{split} \left[A^{(m)}, B\right] &= \left[A, \sum_{n=0}^{m-1} \frac{(m-1)! A^{m-n-1} B(-A)^n}{(m-n-1)! n!}\right] \\ &= \sum_{n=0}^{m-1} \frac{(m-1)! A^{m-n} B(-A)^n}{(m-n-1)! n!} + \sum_{n=1}^{m} \frac{(m-1)! A^{m-n} B(-A)^n}{(m-n)! (n-1)!} \end{split}$$

Extend the summation taking advantage of  $(-1)! = \infty$ 

$$\begin{split} \left[A^{(m)},B\right] &= \sum_{n=0}^{m} \frac{(m-1)!A^{m-n}B(-A)^n}{(m-n-1)!n!} + \frac{(m-1)!A^{m-n}B(-A)^n}{(m-n)!(n-1)!} \\ &= \sum_{n=0}^{m} \frac{(m)!A^{m-n}B(-A)^n}{(m-n)!n!} \left(\frac{m-n}{m} + \frac{n}{m}\right) \\ &= \sum_{n=0}^{m} \frac{(m)!A^{m-n}B(-A)^n}{(m-n)!n!} \end{split}$$

Therefore, the equation is true for all  $m \ge 0$ . If  $[A, B] = \gamma B$ ,  $[A^{(n)}, B] = \gamma^n B$ 

$$\begin{aligned} \mathrm{e}^{\lambda A} B \mathrm{e}^{-\lambda A} &= \sum_{m=0}^{\infty} \frac{\lambda^m}{m!} \Big[ A^{(m)}, B \Big] \\ &= \sum_{m=0}^{\infty} \frac{\lambda^m \gamma^m}{m!} B \\ &= \mathrm{e}^{\lambda \gamma} B \end{aligned}$$

4.

**5**.

(a)

Eigenvalue  $\lambda$ 

$$0 = (h - \lambda)^{2} - |g|^{2}$$
$$h - \lambda = \pm |g|$$
$$\lambda = h \pm |g|$$

Corresponding eigen vectors are  $\frac{1}{\sqrt{2}} \left( 1, \pm \frac{g}{|g|} \right)$ 

(b)

Initial state

$$|\psi_0\rangle = \frac{1}{\sqrt{2}}(|+\rangle + |-\rangle)$$

At time t,

$$\begin{split} |\psi_t\rangle &= \frac{1}{\sqrt{2}} \bigg( \exp\bigg(-\mathrm{i}\frac{h + |g|}{\hbar}t\bigg) |+\rangle + \exp\bigg(-\mathrm{i}\frac{h - |g|}{\hbar}t\bigg) |-\rangle\bigg) \\ &= \frac{\mathrm{e}^{-\mathrm{i}\hbar t/\hbar}}{\sqrt{2}} \bigg( \mathrm{e}^{-\mathrm{i}|g|t/\hbar}|+\rangle + \mathrm{e}^{\mathrm{i}|g|t/\hbar}|-\rangle\bigg) \\ &= \frac{\mathrm{e}^{-\mathrm{i}\hbar t/\hbar}}{2} \bigg( \mathrm{e}^{-\mathrm{i}|g|t/\hbar} \bigg( |1\rangle + \frac{g}{|g|} |2\rangle\bigg) + \mathrm{e}^{\mathrm{i}|g|t/\hbar} \bigg( |1\rangle - \frac{g}{|g|} |2\rangle\bigg) \bigg) \\ &= \mathrm{e}^{-\mathrm{i}\hbar t/\hbar} \bigg( \cos\bigg(\frac{|g|t}{\hbar}\bigg) |1\rangle - \mathrm{i}\frac{g}{|g|} \sin\bigg(\frac{|g|t}{\hbar}\bigg) |2\rangle\bigg) \end{split}$$

- 6.
- (a)
- (b)