## 1.

(a)

Since  $i(\lambda a^{\dagger} - \lambda^* a)$  is Hermitian,  $S_{\lambda} \equiv \exp(\lambda a^{\dagger} - \lambda^* a)$  is unitary and  $|\lambda\rangle \equiv S_{\lambda}|0\rangle$  is normalized. Since  $[a, a^{\dagger}] = 1$  commutes with both a and  $a^{\dagger}$ 

$$\begin{split} |\lambda\rangle &= \exp\left(\lambda a^{\dagger} - \lambda^* a\right) |0\rangle \\ &= \exp\left(\lambda a^{\dagger}\right) \exp\left(-\lambda^* a\right) \exp\left(-\frac{1}{2} \left[\lambda a^{\dagger}, -\lambda^* a\right]\right) |0\rangle \\ &= \exp\left(\lambda a^{\dagger}\right) \exp\left(-\lambda^* a\right) \exp\left(-\frac{|\lambda|^2}{2}\right) |0\rangle \\ &= \exp\left(-\frac{|\lambda|^2}{2}\right) \exp\left(\lambda a^{\dagger}\right) |0\rangle \end{split}$$

$$\begin{aligned} a|\lambda\rangle &= \exp\left(-\frac{|\lambda|^2}{2}\right) a \exp\left(\lambda a^{\dagger}\right) |0\rangle \\ &= \exp\left(-\frac{|\lambda|^2}{2}\right) \left(\exp\left(\lambda a^{\dagger}\right) a + \left[a, \exp\left(\lambda a^{\dagger}\right)\right]\right) |0\rangle \\ &= \exp\left(-\frac{|\lambda|^2}{2}\right) \left[a, a^{\dagger}\right] \lambda \exp\left(\lambda a^{\dagger}\right) |0\rangle \\ &= \lambda |\lambda\rangle \end{aligned}$$

(b)

$$x = z_0(a + a^{\dagger}), p = i\frac{\hbar}{2z_0}(a^{\dagger} - a)$$

$$\langle x \rangle = z_0\langle a + a^{\dagger} \rangle$$

$$= z_0(\lambda + \lambda^*)$$

$$\langle x^2 \rangle = z_0^2 \langle (a + a^{\dagger})^2 \rangle$$

$$= z_0^2 \langle a^2 + a^{\dagger} + aa^{\dagger} + a^{\dagger} a \rangle$$

$$= z_0^2 \langle a^2 + a^{\dagger} + 2a^{\dagger} a + 1 \rangle$$

$$= z_0^2 \langle (\lambda + \lambda^*)^2 + 1 \rangle$$

$$\langle \Delta x^2 \rangle = \langle x^2 \rangle - \langle x \rangle^2$$

$$= z_0^2 \langle a^2 + a^2 +$$

$$\langle p \rangle = i \frac{\hbar}{2z_0} \langle a^{\dagger} - a \rangle$$

$$= i \frac{\hbar}{2z_0} (\lambda^* - \lambda)$$

$$\langle p^2 \rangle = -\frac{\hbar^2}{4z_0^2} \langle (a^{\dagger} - a)^2 \rangle$$

$$= -\frac{\hbar^2}{4z_0^2} \langle a^2 + a^{\dagger^2} - aa^{\dagger} - a^{\dagger}a \rangle$$

$$= -\frac{\hbar^2}{4z_0^2} \langle a^2 + a^{\dagger^2} - 2a^{\dagger}a - 1 \rangle$$

$$= -\frac{\hbar^2}{4z_0^2} \langle (\lambda^* - \lambda)^2 - 1 \rangle$$

$$\langle \Delta p^2 \rangle = \langle p^2 \rangle - \langle p \rangle^2$$

$$= \frac{\hbar^2}{4z_0^2}$$

$$\langle \Delta p^2 \rangle \langle \Delta x^2 \rangle = \frac{\hbar^2}{4}$$

$$= \frac{1}{4} |\langle [x, p] \rangle|^2$$

(c)

$$|\lambda\rangle = \exp\left(-\frac{|\lambda|^2}{2}\right) \exp\left(\lambda a^{\dagger}\right)|0\rangle$$
$$= \exp\left(-\frac{|\lambda|^2}{2}\right) \sum_{n=0}^{\infty} \frac{\left(\lambda a^{\dagger}\right)^n}{n!}|0\rangle$$
$$= \exp\left(-\frac{|\lambda|^2}{2}\right) \sum_{n=0}^{\infty} \frac{\lambda^n}{\sqrt{n!}}|n\rangle$$

Therefore,

$$P(n) = \exp^{-|\lambda|^2} \frac{|\lambda|^{2n}}{n!}$$

$$[n]_{av} = \sum_{n=0}^{\infty} nP(n)$$

$$= \exp^{-|\lambda|^2} \sum_{n=0}^{\infty} \frac{|\lambda|^{2n}}{(n-1)!}$$

$$= |\lambda|^2 \exp^{-|\lambda|^2} \sum_{n=0}^{\infty} \frac{|\lambda|^{2n}}{n!}$$

$$= |\lambda|^2$$

$$[E_n]_{av} = \hbar\omega \left(|\lambda|^2 + \frac{1}{2}\right)$$

(d)

$$\langle n^2 \rangle = \sum_{n=0}^{\infty} n^2 P(n)$$

$$= \exp^{-|\lambda|^2} \sum_{n=0}^{\infty} n \frac{|\lambda|^{2n}}{(n-1)!}$$

$$= \exp^{-|\lambda|^2} \left( \sum_{n=0}^{\infty} \frac{|\lambda|^{2n}}{(n-2)!} + \sum_{n=0}^{\infty} \frac{|\lambda|^{2n}}{(n-1)!} \right)$$

$$= \exp^{-|\lambda|^2} \left( |\lambda|^4 \sum_{n=0}^{\infty} \frac{|\lambda|^{2n}}{n!} + |\lambda|^2 \sum_{n=0}^{\infty} \frac{|\lambda|^{2n}}{n!} \right)$$

$$= |\lambda|^4 + |\lambda|^2$$

$$\Delta n = |\lambda|$$

$$\Delta E = \hbar \omega |\lambda|$$

$$\frac{\Delta E}{(E_n)_{av}} = \frac{1}{|\lambda|}$$

so the relative uncertainty goes to 0 at large n limit.

- 2.
- (a)
- (b)
- (c)
- (d)
- 3.

The probability is

$$\begin{aligned} \left| \langle u | \chi \rangle \right|^2 &= \left| c_1 \right|^2 \\ \left| \langle u | \chi \rangle \right|^2 &= \langle u | \chi \rangle \langle \chi | u \rangle \\ &= \langle u | \rho | u \rangle \\ &= \operatorname{Tr}(\langle u | \rho | u \rangle) \\ &= \operatorname{Tr}(\rho | u \rangle \langle u |) \end{aligned}$$

- **4.**
- (a)
- (b)
- (c)
- (d)
- (e)
- **(f)**
- (g)
- **5.**
- (a)
- (b)