

Lesson 16:

- I. Lagrange's Equations (cont'd) (MLS 4.2)
 - II. Inertial Properties of Rigid Bodies (MLS 4.2.2)
 - III. Lagrangian of Open-Chain Robot (MLS 4.3.1)
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I. Lagrange's Equations (cont'd)

• Recall: $L(q, \dot{q}) := K(q, \dot{q}) - P(q)$

\uparrow \uparrow \uparrow
Lagrangian kinetic energy potential energy

The EOM of a mechanical system is given by:

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} = \gamma$$

• Generalized momentum: $p_i = \frac{\partial L}{\partial \dot{q}_i}$

$$\Rightarrow \frac{d}{dt} p_i = \frac{\partial L}{\partial q_i} + \gamma_i$$

\uparrow
conservative forces and kinetic terms

, if RHS is zero, then
conservation of momentum
($p_i = \text{const.}$)

- Preview of HW 4: Hamilton's Method uses

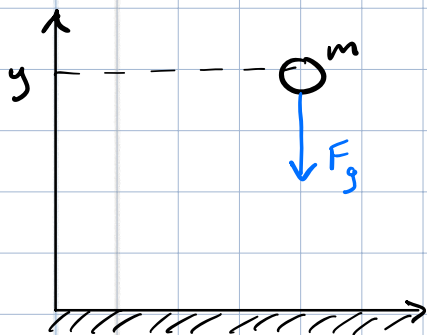
$$x = \begin{bmatrix} q \\ p \end{bmatrix} \in \mathbb{R}^{2n} \quad \text{with set of } 2n \text{ 1st order ODEs}$$

Def: The Hamiltonian function $H(q, p) \in \mathbb{R}$ given by

$$H(q, p) = \underbrace{K(q, p)}_{\text{kinetic energy}} + \underbrace{P(q)}_{\text{potential energy}}$$

→ See HW 4 for Hamilton's equations.

Ex: Falling mass



Find equation of motion (Eom):

• Recall Newton's law

$$p = m \cdot v \xrightarrow{\frac{d}{dt}} \dot{p} = ma = F$$

$$m\ddot{y} = F_g = -mg$$

$$\ddot{y} = -g \quad \checkmark$$

• Now let's try with E-L:

$$K = \frac{1}{2} m \dot{y}^2 \quad P = mgy$$

$$\Rightarrow L = K - P = \frac{1}{2} m \dot{y}^2 - mgy$$

$$\text{E-L: } \frac{d}{dt} \frac{\partial L}{\partial \dot{y}} - \frac{\partial L}{\partial y} = 0$$

$$\frac{d}{dt} (m \dot{y}) - (-mg) = 0$$

$$m\ddot{y} + mg = 0$$

$$\ddot{y} = -g \quad \checkmark$$

* Same process for multi-body kinematic chains,
but need to derive K and P first.

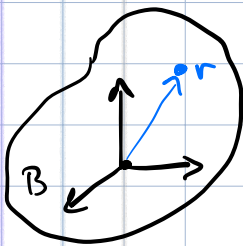
II. Inertial Properties of Rigid Bodies

$V \subset \mathbb{R}^3$ is volume occupied by a body.

For position $r = \begin{bmatrix} x \\ y \\ z \end{bmatrix} \in V$, density $\rho(r)$ (mass distribution)

$$\text{mass } m = \int_V \rho(r) dV = \int_{x_1}^{x_2} \int_{y_1}^{y_2} \int_{z_1}^{z_2} \rho(x, y, z) dz dy dx$$

$$\text{center of mass (com)} \quad \bar{r} = \frac{1}{m} \int_V \rho(r) r dV \in \mathbb{R}^3$$



Let $r \in \mathbb{R}^3$ be coords. of a
body point relative to the
body frame.

Inertia Tensor
(constant 3×3)

$$\tilde{I} = \begin{bmatrix} I_{xx} & I_{xy} & I_{xz} \\ & I_{yy} & I_{yz} \\ & & I_{zz} \end{bmatrix} = \tilde{I}^T$$

$$= - \int_V \rho(r) \hat{r}^2 dV$$

where $\hat{r}^2 = \begin{bmatrix} -(z^2+y^2) & xy & xz \\ & -(x^2+z^2) & yz \\ & & -(x^2+y^2) \end{bmatrix} = (\hat{r}^2)^T$

Then,

$$I_{xx} = - \int_V \rho(r) (-z^2 - y^2) dx dy dz \quad (\text{Moment of Inertia})$$

$$= \int_V \rho(r) (z^2 + y^2) dx dy dz$$

$$\vdots \quad \left(\begin{array}{l} \tau_x = I_{xx} \dot{\omega}_x \text{ for ang. vel.} \\ \omega_x \text{ about } x\text{-axis, if } x\text{-axis} \\ \text{is principal axis.} \end{array} \right)$$

$$I_{xy} = - \int_V \rho(r) xy dx dy dz \quad (\text{Product of Inertia})$$

\vdots

Note: Inertia tensor is real, symmetric

\Rightarrow diagonalizable \Rightarrow can find principal axes.

• Let $v^b \in \mathbb{R}^3$ be linear component and $\omega^b \in \mathbb{R}^3$ be the angular component of the body velocity coords:

$$\underline{V}^b = \begin{bmatrix} v^b \\ \omega^b \end{bmatrix} = \begin{bmatrix} R^T \dot{p} \\ (R^T \dot{R})^v \end{bmatrix} \in \mathbb{R}^6 \iff \hat{\underline{V}}^b = g^{-1} \dot{g} \in \mathfrak{se}(3)$$

from path $g(t) = (R(t), p(t)) \in \text{SE}(3)$

• Then kinetic energy

$$K = \underbrace{\frac{1}{2} m \|v^b\|^2}_{\text{linear}} + \underbrace{\frac{1}{2} (\omega^b)^T \mathcal{I} \omega^b}_{\text{angular}}$$

$$= \frac{1}{2} (\mathbb{V}^b)^T \underbrace{\begin{bmatrix} m\mathbf{I}_{3 \times 3} & \mathbf{0}_{3 \times 3} \\ \mathbf{0}_{3 \times 3} & \mathcal{I} \end{bmatrix}}_{\mathcal{M}} \mathbb{V}^b$$

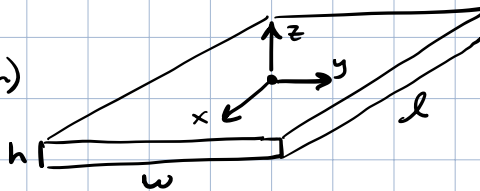
$$= \frac{1}{2} (\mathbb{V}^b)^T \mathcal{M} \mathbb{V}^b$$

where $\mathcal{M} \in \mathbb{R}^{6 \times 6}$ is the generalized inertia matrix of rigid body.

- real, symmetric, constant, and positive definite
 $(x^T \mathcal{M} x > 0, \forall x \neq 0)$

Ex: Homogeneous bar

$$\rho = \frac{m}{lwh} \text{ (uniform)}$$



(frame is at center of bar)

Moments of inertia:

$$\begin{aligned} I_{xx} &= \int_{\mathbb{V}} \rho (y^2 + z^2) d\mathbb{V} = \rho \int_{-\frac{h}{2}}^{\frac{h}{2}} \int_{-\frac{w}{2}}^{\frac{w}{2}} \int_{-\frac{l}{2}}^{\frac{l}{2}} (y^2 + z^2) dx dy dz \\ &= \rho l (w^3 h + h^3 w) / 12 \\ &= \frac{m}{12} (w^2 + h^2) \end{aligned}$$

I_{yy} and I_{zz} are similar...

Products of inertia:

$$\begin{aligned}
 I_{xy} &= -\rho \int_{-\frac{h}{2}}^{\frac{h}{2}} \int_{-\frac{w}{2}}^{\frac{w}{2}} \int_{-\frac{l}{2}}^{\frac{l}{2}} xy \, dx \, dy \, dz \\
 &= -\rho \int_{-\frac{h}{2}}^{\frac{h}{2}} \int_{-\frac{w}{2}}^{\frac{w}{2}} \underbrace{\left(\frac{1}{2} x^2 y \right) \Big|_{x=-\frac{l}{2}}^{x=\frac{l}{2}}}_{=0} dy \, dz = 0 = I_{xz} = I_{yz}
 \end{aligned}$$

* Products of inertia are always zero when axes are principal (here, located at COM and aligned with the dimensions of body)

$$\Rightarrow \mathcal{I} = \begin{bmatrix} \frac{m}{12}(w^2 + h^2) & 0 & 0 \\ 0 & \frac{m}{12}(l^2 + h^2) & 0 \\ 0 & 0 & \frac{m}{12}(l^2 + w^2) \end{bmatrix}$$

$$\Rightarrow \mathcal{M} = \begin{bmatrix} m I_{3 \times 3} & 0_{3 \times 3} \\ 0_{3 \times 3} & \mathcal{I} \end{bmatrix} = m \begin{bmatrix} I_{3 \times 3} & 0_{3 \times 3} \\ 0_{3 \times 3} & \frac{1}{12} \begin{bmatrix} w^2 + h^2 & 0 & 0 \\ 0 & l^2 + h^2 & 0 \\ 0 & 0 & l^2 + w^2 \end{bmatrix} \end{bmatrix}$$

III. Lagrangian of Open-Chain Robot (n -DOF/ n -links)

- Obtain KE by summing the KE of each link.