

Diffuse Relativistic Black Hole in the Cornell Box

Andrew Aikawa

Yuan Zhou

Weylan Wang

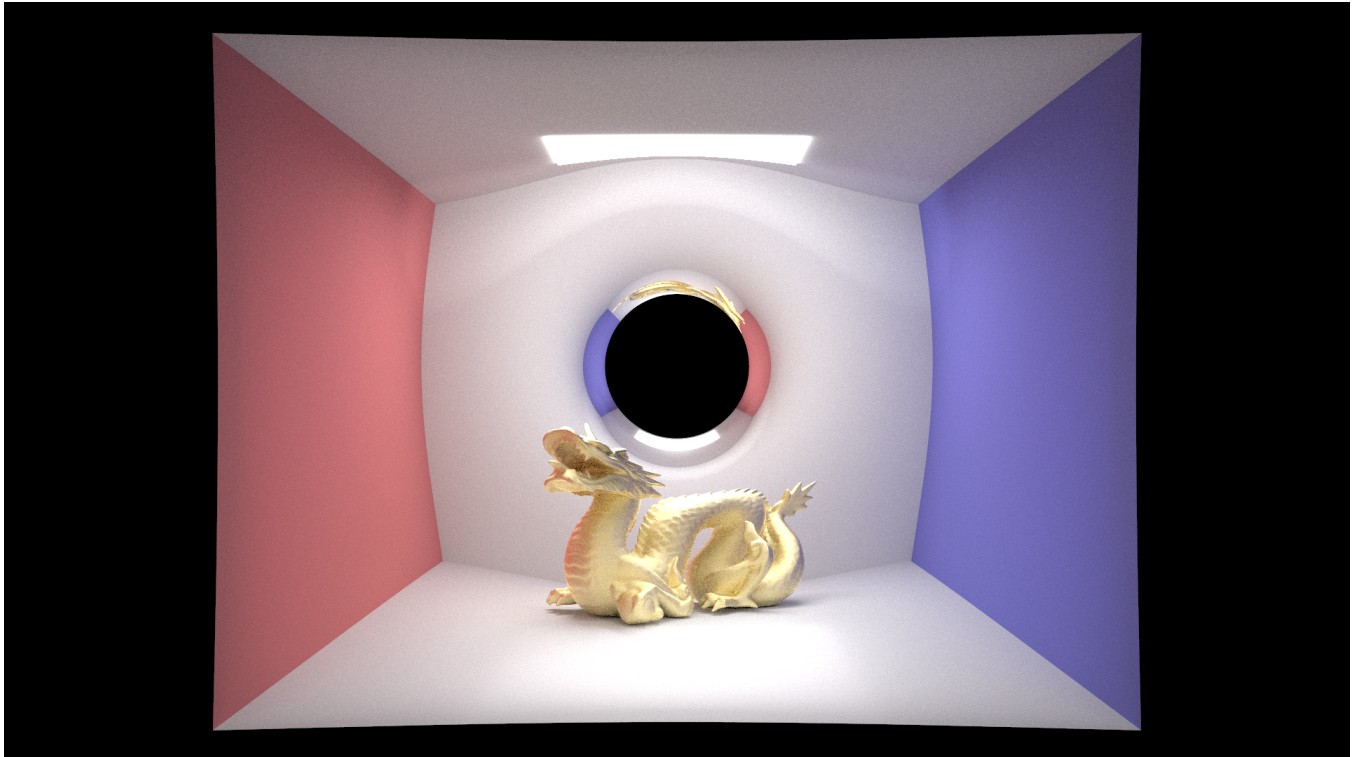


Figure 1: Gold Dragon in Cornell Box Containing a Black Hole

ABSTRACT

We report a ray tracing algorithm in which we can quantitatively generate realistic images in a scene with a black hole. The problem of illustrating what a black hole would look like has been explored previously in works such as the movie *Interstellar*[1]. However, almost all of these assume purely emissive light sources and very distant light sources that would be represented as environment maps. In this paper, we present a ray tracer for a scene with a Schwarzschild black hole in diffuse lighting conditions.

KEYWORDS

ray tracing, black holes, general relativity

ACM Reference Format:

Andrew Aikawa, Yuan Zhou, and Weylan Wang. 2019. Diffuse Relativistic Black Hole in the Cornell Box. In *CS 184 Sp19: Final Project, April 30, 2019, Berkeley, CA*. ACM, New York, NY, USA, ?? pages.

1 RAY TRACING IN GENERAL RELATIVITY

Einstein's field equations predict changes in the spacetime curvature in the vicinity of a black hole. The metric tensor, $g_{\mu\nu}$, the solution to the field equations, determines the metric, how 'distance' is measured in spacetime. The path taken by a ray of light is determined by Fermat's principle, dictating that the ray will take whatever path takes the ray from point A to point B with the least action, a quantity describing how much metric distance the ray has traveled. In a typical Euclidean metric space, where the metric is uniform over all space and in a medium with a uniform index of refraction, rays will trace lines, a central assumption in traditional ray tracing algorithms. However, in the presence of a massive body, a spatially varying metric will cause light rays to curve in order to

Permission to make digital or hard copies of all or part of this work for personal or classroom use is granted without fee provided that copies are not made or distributed for profit or commercial advantage and that copies bear this notice and the full citation on the first page. Copyrights for components of this work owned by others than the author(s) must be honored. Abstracting with credit is permitted. To copy otherwise, or republish, to post on servers or to redistribute to lists, requires prior specific permission and/or a fee. Request permissions from permissions@acm.org.

CS 184 Sp19, April 30, 2019, Berkeley, CA

© 2019 Copyright held by the owner/author(s). Publication rights licensed to ACM.
ACM ISBN 978-x-xxxx-xxxx-x/YY/MM...\$15.00

minimize the action, how much spacetime 'distance' the ray travels. We examine the particular case of a Schwarzschild black hole, a black hole with no spin for which the metric is analytic.

2 PHOTON ORBITS IN SCHWARZCHILD SPACETIME

Due to the spherical symmetry of the Schwarzschild spacetime, photon orbitals are confined to a plane. The equations of motion are found by minimizing the action by means of Euler-Lagrange equations which yield[2]:

$$u'' + u - \frac{3r_s u^2}{2} = 0$$

r_s being the Schwarzschild radius, a characteristic quantity of a black hole which is solely determined by its mass, and $u(\theta) = 1/r(\theta)$, where r is the polar coordinate of the orbit and θ the angular coordinate of the orbit. This equation has no closed form solutions and in general can only be solved numerically. For our ray tracer, we used a 4th order Runge-Kutta method.

3 RAY CASTING ALGORITHM

Our ray casting algorithm works by treating the geodesic paths as a piece wise connection of small linear rays. By numerically integrating the photon orbital equation, we get the next point on the geodesic, which gives us a ray defined by the line segment connecting the original point to the next point. Since we are building on top of our project 3 [3] and photon orbit equation is in terms of polar coordinates, we needed to convert between black hole polar coordinates and the world space. The plane of the orbit is entirely determined by the black hole position, ray origin, and ray initial direction. To form the black hole centered coordinates, define the unit vector pointing from the black hole to the ray origin as the x-axis and the orthogonal projection of the initial ray direction as the y-direction. The initial conditions are given by:

$$u(\theta = 0) = \frac{1}{r(\theta = 0)}, u'(\theta = 0) = -\frac{1}{y'r(\theta = 0)}$$

where y' is the slope of the ray direction in the black hole centered coordinates. The second equation was obtained by inverting the equation for the slope of a polar curve for r' and differentiating $u = 1/r$. At each step of the numerical integration, we check if the micro-ray formed by each step intersects with the scene and if so, to return the outgoing spectrum at that point through uniform hemisphere sampling.

4 EMERGENT BEHAVIOR

In our images, we find some notable features that have been reported in other black hole renderers. Perhaps the most iconic feature of the black hole that we reproduced in our images is the circular distortion around the black hole due to the bending of light. In addition, we also find that there is a circular region immediately around the black hole with the view of the room inverted, a result of intense bending near the black hole. Another very important feature which is of importance in astronomy, is the lensing effect of the black hole. Light from the lamp in the scene becomes focused by the black hole and converges onto the surfaces behind the black

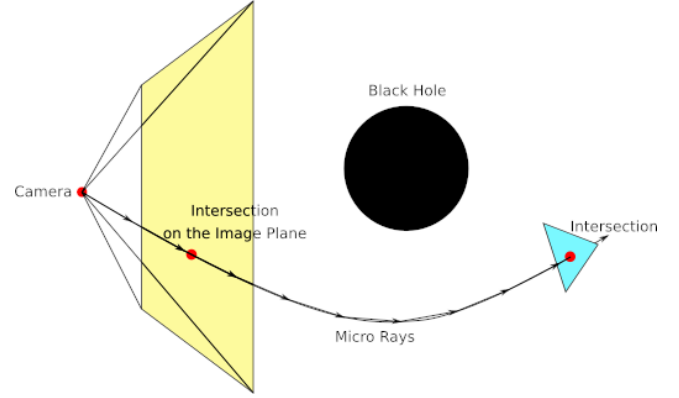


Figure 2: Micro Ray Intersection

hole with respect to the light source, manifesting as an increase in the intensity of the lighting.

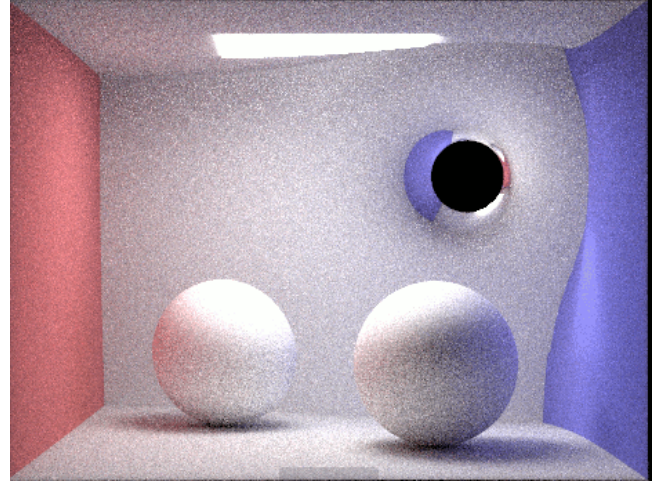


Figure 3: Black hole in Cornell Box. Note the bright spot on the blue wall as a result of the gravitational lensing

5 GRAVITATIONAL REDSHIFT

The distortion of spacetime induces changes in the local speed of light. Due to the wave nature of light, light spectra undergo a wavelength shift in a phenomenon known as gravitational redshift. Emission near the black hole becomes longer in wavelength as it moves away from the black hole and is described by the equation:

$$\frac{\lambda_1}{\lambda_2} = [(1 - r_s/R_1)(1 - r_s/R_2)]^{1/2}$$

where R_1, R_2 are the distances of the emitter and observer respectively relative to the black hole and λ_1, λ_2 are the emitted and observed wavelength respectively. In order to implement redshift, we took spectral data for the Cornell box and used spectral intensity values ranging from 400nm to 700nm with 4nm resolution[5].

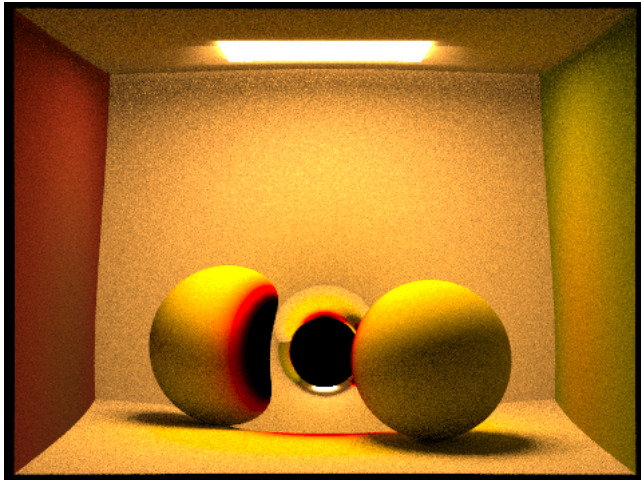


Figure 4: Black hole with redshift. Note that the spheres appear black due to visible wavelengths being pushed out the visible range. In addition, parts of the green wall appear more yellow-orange and the red wall more off red-grayish as a result of the redshift

6 PROGRESS

For our future goals, we would like to be able to implement a black hole with spin. This would be very hard however as this requires solving 5 differential equations simultaneously numerically. We can also see about optimizing our code to run on a GPU through CUDA or OptiX.

7 ACKNOWLEDGMENTS

Special thanks to the Open Computing Facility (ocf.io) for letting us run our CPU intensive ray tracer. And Ren Ng for answering our questions regarding color spaces and where to find spectral data.

8 REFERENCES

- [1] Oliver James, Eugenie von Tunzelmann, Paul Franklin, Kip S Thorne. 2015. Gravitational lensing by spinning black holes in astrophysics, and in the movie Interstellar. Classical and Quantum Gravity, Vol. 32, Number 6. IOP Science.
- [2] https://hepweb.ucsd.edu/ph110b/110b_notes//cs184.eecs.berkeley.edu/sp19/article/18/assignment-3-1-pathtracer
- [4] https://en.wikipedia.org/wiki/Gravitational_redshift
- [5] <https://www.graphics.cornell.edu/online/box/>