

Physics of Waves

A practical review of the workshop on acoustic wavefield imaging

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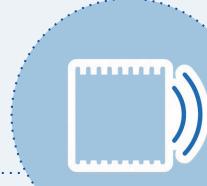


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Introduction to Physics of Waves



Applications Acoustical Imaging

Table 1. Applications Acoustical Imaging [1].

>50 MHz 10 MHz 1 MHz 100 kHz 10 kHz 1 kHz 100 Hz 10 Hz 1 Hz 0.1 Hz

Acoustic Microscopy

Laminated materials (e.g. airplanes) \rightarrow 0.5 cm deep Inspections of welds \rightarrow 0.5-5 cm deep

Medical ultrasound → 0.2-20 cm deep

Ship wreck detection → 1-20 m deep

Near surface inspection → 5-500 m deep

Oil & gas exploration → 0.5-5 km deep

Tectonic imaging → 2-25 km depth

Global seismology → full earth

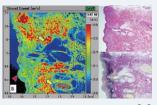


Figure 1. SOS map of skin [2].



Figure 2. Echography [3].



Figure 3. Seismology experiment [4].

^[1] https://kvandongen.tnw.tudelft.nl/wp-content/uploads/2025/02/Koen-Bucaramanga-day-1.pdf

^[2] K. Miura, "Application of Scanning Acoustic Microscopy to Pathological Diagnosis," DOI: 10.5772/63405

^[3] http://langeproductions.com/

^[4] www.iagc.org

Basic Concepts

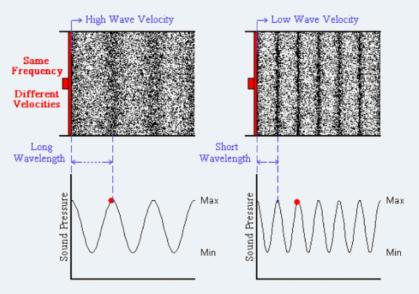


Figure 4. Acoustic Wave Propagation [5].

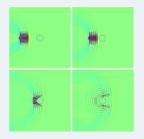


Figure 5. Acoustic Wave Propagation [6].

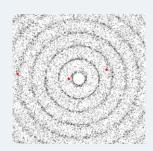
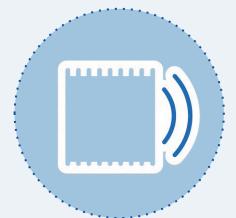


Figure 6. Acoustic Monopole [7].

How a single acoustic source produces waves that propagate in all directions and how they interact in the medium



^[5] http://resource.isvr.soton.ac.uk/spcg/tutorial/tutorial_files/Web-basics-frequency.htm

^[6] http://www.mathworks.com/matlabcentral/fileexchange/screenshots/1347/original.jpg

^[7] http://resource.isvr.soton.ac.uk/spcg/tutorial/tutorial_files/Web-basics-pointsources.htm

Acoustic Field Equations

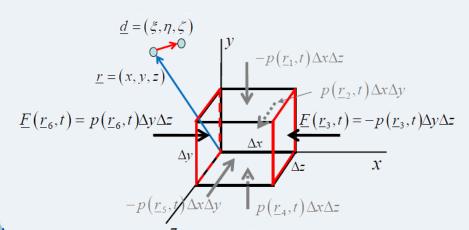
Volumen: $\Delta V = \Delta x \Delta y \Delta z$

Location: $\underline{r} = (x, y, z)$

Displacement field: $\underline{d}(\underline{r},t) = [\xi(\underline{r},t), \eta(\underline{r},t), \zeta(\underline{r},t)]$

Velocity field: $\underline{v}(\underline{r},t) = \frac{d}{dt}\underline{d}(\underline{r},t)$

Pressure field: $p(\underline{r},t) = p_{tot}(\underline{r},t) - p_0$



Hooke's law

A volume source density of injection rate

Equation of deformation

$$\frac{\partial p(\underline{r},t)}{\partial t} = -\frac{1}{\kappa} \nabla \cdot \underline{v}(\underline{r},t) + \frac{1}{\kappa} q(\underline{r},t)$$

$$i\omega \hat{p}(\underline{r},\omega) = -\frac{1}{\kappa} \nabla \cdot \underline{\hat{v}}(\underline{r},\omega)$$

Newton's law

Equation of motion

A volume source density of volume source

 $-\nabla p(\underline{r},t) = \rho \frac{\partial \underline{v}(\underline{r},t)}{\partial t} - \underline{f}(\underline{r},t)$

$$-\nabla \hat{p}(\underline{r},\omega) = i\omega \rho \hat{\underline{v}}(\underline{r},\omega)$$

Acoustic Field Equations

Wave Equation

In the absence of sources:

A escalar wave equation for the pressure field

$$\rho \kappa \frac{\partial}{\partial t} \left[\frac{\partial p(\underline{r}, t)}{\partial t} \right] = \rho \kappa \frac{\partial}{\partial t} \left[-\frac{1}{\kappa} \nabla \cdot \underline{v}(\underline{r}, t) \right]$$

$$\nabla \cdot \left[-\nabla p(\underline{r}, t) \right] = \nabla \cdot \left[\rho \frac{\partial \underline{v}(\underline{r}, t)}{\partial t} \right]$$

$$\Rightarrow \boxed{\nabla^2 p(\underline{r}, t) - \rho \kappa \frac{\partial^2 p(\underline{r}, t)}{\partial t^2} = 0}$$

$$\kappa \nabla \left[\frac{\partial p(\underline{r},t)}{\partial t} \right] = \kappa \nabla \left[-\frac{1}{\kappa} \nabla \cdot \underline{v}(\underline{r},t) \right]$$

$$\kappa \frac{\partial}{\partial t} \left[-\nabla p(\underline{r},t) \right] = \kappa \frac{\partial}{\partial t} \left[\rho \frac{\partial \underline{v}(\underline{r},t)}{\partial t} \right]$$

$$\Rightarrow \nabla \left[\nabla \nabla \cdot \underline{v}(\underline{r},t) \right] - \rho \kappa \frac{\partial^2 \underline{v}(\underline{r},t)}{\partial t^2} = 0$$

Helmholtz Equation

$$\nabla^2 \hat{p}\left(\underline{r},\omega\right) + \frac{\omega^2}{c^2} \hat{p}\left(\underline{r},\omega\right) = 0$$

$$\nabla^{2} \hat{p}(\underline{r}, \omega) + \frac{\omega^{2}}{c^{2}} \hat{p}(\underline{r}, \omega) = -\hat{S}(\underline{r}, \omega)$$

Spherical Waves

A vectorial wave equation for the velocity wavefield
$$=\frac{1}{c^2} \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \hat{p}(r,\omega) \right) + \frac{\omega^2}{c^2} \hat{p}(r,\omega) = -\delta(\underline{r}) \hat{S}(\omega)$$

$$\hat{p}(r,\omega) = \hat{A}(\omega) \frac{e^{-i\omega r/c_0}}{r} + \frac{\hat{B}(\omega)}{r} \frac{e^{i\omega r/c_0}}{r}, \quad \text{for } r \neq 0.$$

Green's function or impulse response of the medium

$$\hat{G}(\underline{r},\omega) = \frac{e^{-i\omega|\underline{r}|/c_0}}{4\pi|\underline{r}|}$$

Acoustic Field Equations

Green's functions

In the absence of sources:

the absence of sources:
1D:
$$\left(\frac{d^2}{dx^2} + k^2\right) \hat{G}(x) = -\delta(x-a) \rightarrow \hat{G}(x) = \frac{i}{2k} e^{-ikr}$$
 with $r = |x-a|$

$$\frac{i}{2ka} \left[H_1^{(1)}(ka) + \frac{2i}{\pi ka} \right] \qquad r = 0$$
2D: $(\nabla^2 + k^2) \hat{G}(\vec{x}) = -\delta(\vec{x} - \vec{a}) \rightarrow \hat{G}(\vec{x}) = -\frac{i}{2k} H_2^{(2)}(kr)$ with $r = |\vec{x} - \vec{a}|$

2D:
$$(\nabla^2 + k^2) \hat{G}(\vec{x}) = -\delta(\vec{x} - \vec{a}) \rightarrow \hat{G}(\vec{x}) = -\frac{i}{4} H_0^{(2)}(kr) \text{ with } r = |\vec{x} - \vec{a}|$$

3D:
$$(\nabla^2 + k^2)\hat{G}(\vec{x}) = -\delta(\vec{x} - \vec{a}) \rightarrow \hat{G}(\vec{x}) = \frac{1}{4\pi} \frac{e^{-ikr}}{r}$$
 with $r = |\vec{x} - \vec{a}|$
$$\begin{cases} \frac{3e^{-ikr}}{4k^3\pi a^3r} \left[\sin(ka) - ka\cos(ka)\right] & r \ge a \\ \frac{3}{4k^2\pi a^3} \left[(1 + ika)e^{-ika} - 1\right] & r = 0 \end{cases}$$

$$\begin{cases} \frac{i}{2ka}J_1(ka)H_0^{(1)}(kr) & r \ge a \end{cases}$$

$$\left[\frac{i}{2ka}\left[H_1^{(1)}(ka) + \frac{2i}{\pi ka}\right]\right] \qquad r = 0$$

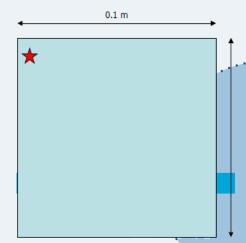
$$\left[\frac{3e^{-ikr}}{4k^3\pi a^3r}\left[\sin(ka) - ka\cos(ka)\right] \qquad r \ge$$

$$\frac{3}{4k^2\pi a^3} \Big[(1+ika)e^{-ika} - 1 \Big] \qquad r = 0$$



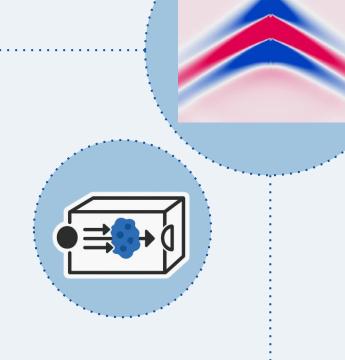
Exercise 1

- Compute the pressure field generated by a point source (volume source density of injection rate)
 - Centre frequency Gaussian pulse f₀ = 1 MHz;
 - Medium water (c = 1500 m/s, ρ = 1000 kg/m³);
 - Locate the point source in the centre of a volume (e.g. 0.1 m x 0.1 m x 0.0015 m).
- Compute the pressure field in the frequency domain and transform the resulting field to the time domain using FFT.
- $\Delta x = ?$
- $\Delta t = ?$
- Nt = ?

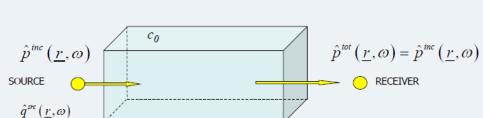


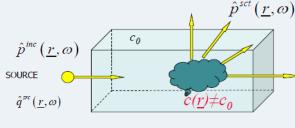
02

Heterogeneous Media



Equations





Spatial variations in the volume density of mass are neglected

 $\hat{p}^{tot}(r,\omega) = \hat{p}^{inc}(r,\omega) + \hat{p}^{sct}(r,\omega)$

RECEIVER

The acoustic media parameters become spatially varying

Hooke's law

$$\nabla \bullet \underline{v}(\underline{r}, t) + \kappa (\underline{r}) \partial_{t} p(\underline{r}, t) = q(\underline{r}, t)$$

$$\downarrow \qquad \qquad \qquad \downarrow$$

$$\nabla \bullet \underline{v}(\underline{r}, t) + \kappa_{0} \partial_{t} p(\underline{r}, t)$$

$$= q(\underline{r}, t) + \{\kappa_{0} - \kappa (\underline{r})\} \partial_{t} p(\underline{r}, t)$$

Newton's law

$$\nabla^{2} p\left(\underline{r},t\right) - \frac{1}{c_{0}^{2}} \partial_{t}^{2} p\left(\underline{r},t\right) = -\underbrace{\left\{\rho_{0} \partial_{t} q\left(\underline{r},t\right) - \nabla \bullet \underline{f}(\underline{r},t)\right\}}_{S_{pr}(\underline{r},t)} - \underbrace{\left\{\rho_{0} \left(\kappa_{0} - \kappa\left(\underline{r}\right)\right) \partial_{t}^{2} p(\underline{r},t) - \nabla \bullet \underbrace{\left\{\rho_{0} - \rho\left(\underline{r}\right)\right\} \partial_{t} v\left(\underline{r},t\right)\right\}}_{S_{cs}(\underline{r},t)}$$

$$\nabla^{2} \hat{p}(\underline{r}, \omega) + \frac{\omega^{2}}{c_{0}^{2}} \hat{p}(\underline{r}, \omega) = -\left\{\rho_{0} i \omega \hat{q}(\underline{r}, \omega) - \nabla \cdot \underline{\hat{f}}(\underline{r}, \omega)\right\} - \left\{-\rho_{0}(\kappa_{0} - \kappa(\underline{r})) \omega^{2} \hat{p}(\underline{r}, \omega) - \nabla \cdot \left[\left\{\rho_{0} - \rho(\underline{r})\right\} i \omega \underline{\hat{v}}(\underline{r}, \omega)\right]\right\}$$

Equations

Primary sources

$$S_{pr}(\underline{r}',t)$$

Green's function

$$\hat{G}(\underline{r}, \omega) = \frac{e^{-i\omega|\underline{r}|/c_0}}{4\pi|\underline{r}|}$$

Field generated

$$\hat{G}(\underline{r},\omega) = \frac{e^{-i\omega|\underline{r}|/c_0}}{4\pi|\underline{r}|} \qquad \hat{p}^{inc}(\underline{r},\omega) = \int_{r'\in D} \hat{G}(\underline{r}-\underline{r}',\omega) S_{pr}(\underline{r}',\omega) dV(\underline{r}').$$

Principle of superposition

$$\hat{p}(r,\omega) = \hat{p}^{inc}(\underline{r},\omega) + \int_{\underline{r}' \in D} \hat{G}(\underline{r} - \underline{r}',\omega) \chi(\underline{r}') \omega^2 \hat{p}(\underline{r}',\omega) dV(\underline{r}')$$

with
$$\chi(\underline{r}') = \frac{1}{c^2(\underline{r}')} - \frac{1}{c_0^2}$$

Born approximation

$$\hat{p}(r,\omega) = \hat{p}^{\textit{inc}}(\underline{r},\omega) + \int_{r' \in D} \hat{G}(\underline{r} - \underline{r}',\omega) \; \chi(\underline{r}') \; \omega^2 \; \hat{p}^{\textit{inc}}(\underline{r}',\omega) \; dV(\underline{r}')$$

Forward and Inverse Problem

Forward problem

- Source and contrast are known
- Total/actual field is unknown
- → Linear Problem

Inverse problem

- Sources are known
- Total/actual field is known at the boundary
- Contrast and field in ROI is unknown
- → Non-linear problem

Solutions

Neumann Series

$$\hat{p}^{\left(1\right)}$$
 Born approximation
$$\hat{p}^{\left(2\right)}(r,\omega) = \hat{p}^{inc}(\underline{r},\omega) + \int_{\underline{r}'\in D} \hat{G}(\underline{r}-\underline{r}',\omega) \; \chi\left(\underline{r}'\right)\omega^2 \; \hat{p}^{\left(1\right)}(\underline{r}',\omega) \; dV(\underline{r}')$$

 $\hat{p}^{(1)}, \, \hat{p}^{(2)}, ..., \, \hat{p}^{(n)}$: Neumann Series

Warning: For strong contrast, the iteration may not converge to the true solution.

Conjugate gradient method

$$\hat{p}(r,\omega) = \hat{p}^{inc}(\underline{r},\omega) + \int_{r'\in D} \hat{G}(\underline{r} - \underline{r}',\omega) \ \chi(\underline{r}') \ \omega^2 \ \hat{p}(\underline{r}',\omega) \ dV(\underline{r}')$$

$$\mathsf{E}_\mathsf{n} = \| \ \mathsf{r}_\mathsf{n} \ \|^2 = \| \ \mathsf{f} - \mathsf{L}\mathsf{u}_\mathsf{n} \ \|^2$$

$$\mathsf{u}_\mathsf{n} = \mathsf{u}_\mathsf{n-1} + \alpha_\mathsf{n} \mathsf{d}_\mathsf{n}$$

$$\mathsf{d}_\mathsf{n} = \mathsf{L}^\dagger \mathsf{r}_\mathsf{n}$$

Known incident field
$$\mathbf{f}=\hat{p}^{inc}(\underline{r},\omega)$$
 Unknown total field $\mathbf{u}=\hat{p}(\underline{r},\omega)$ Remaining integral operator

Finite Difference Time Domain (FDTD)

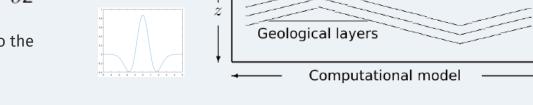
$$\frac{1}{\mathbf{c}^2(x,z)} \frac{\partial^2 p(x,z)}{\partial t^2} = \frac{\partial^2 p(x,z)}{\partial x^2} + \frac{\partial^2 p(x,z)}{\partial z^2} + \operatorname{src}(x,z,t) \Big|_{z}$$

The finite difference method can be applied to the acoustical scalar wave equation

$$\left(\frac{\partial^2 p}{\partial t^2}\right)_{i,j}^n = \frac{p_{i,j}^{n+1} - 2p_{i,j}^n + p_{i,j}^{n-1}}{(\Delta t)^2}$$

$$\left(\frac{\partial^2 p}{\partial x^2}\right)_{i,j}^n = \frac{p_{i+1,j}^n - 2p_{i,j}^n + p_{i-1,j}^n}{(\Delta x)^2}$$

$$\left(\frac{\partial^2 \rho}{\partial z^2}\right)_{i,j}^n = \frac{\rho_{i,j+1}^n - 2\rho_{i,j}^n + \rho_{i,j-1}^n}{(\Delta z)^2}$$



Receivers

$$p_{i,j}^{n+1} = 2(1-2G^2)p_{i,j}^n - p_{i,j}^{n-1} + G^2(p_{i+1,j}^n + p_{i-1,j}^n + p_{i,j+1}^n + p_{i,j-1}^n)$$

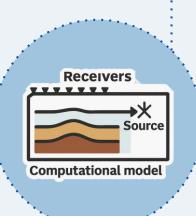
$$G=rac{c\Delta t}{\Delta h}$$
 Courant parameter $\Delta h=\Delta x=\Delta z$

Stability is archieved when
$$\alpha \leq \frac{1}{\sqrt{2}}$$
 $\alpha = \frac{\mathbf{c}\Delta t}{\Delta h}$

Numerical dispersion is avoided when $\frac{v_p}{c}=1$ $v_p=\frac{\omega}{k}$

$$\frac{v_p}{c} = 1$$
 $v_p = \frac{\omega}{k}$

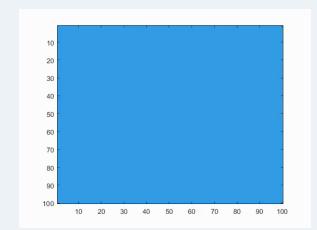
X Source

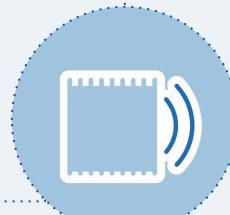


Exercise 2

Simulate wave propagation using FDTD using the Ricket function as source. Define your parameters properly and check the stability of your simulation.

$$\frac{1}{\mathbf{c}^2(x,z)}\frac{\partial^2 p(x,z)}{\partial t^2} = \frac{\partial^2 p(x,z)}{\partial x^2} + \frac{\partial^2 p(x,z)}{\partial z^2} + \operatorname{src}(x,z,t),$$





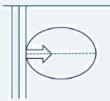
03

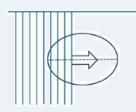
Ultrasound Imaging

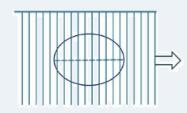




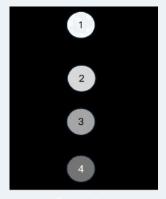
B-mode ultrasound



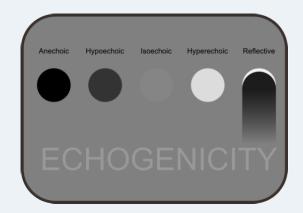




An 2D B-mode image is formed line-by-line as the beam moves along the transducer array.

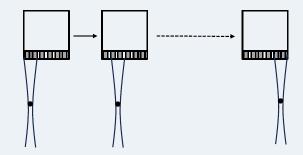


B-mode



Linear Array Imaging

- In linear array imaging, a few elements are grouped and fired to get an A-line
- Delays can be used to create a focused beam



Number of A lines to make an image

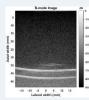
- = number of elements
- number of grouped elements + 1

Exercise 3



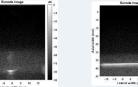
Cross-sectional image

Longitudinal image

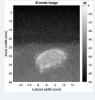


Needle

Tube



Finger





A series of ultrasound data acquisitions were conducted using the L11-5V linear array transducer, which has a center frequency of 7.6 MHz. The experiment was performed in a controlled water tank environment to ensure consistent acoustic propagation conditions. Obtain the B-mode image from the channel data.

