

Solving the Heat Equation

Background

We consider a thin, perfectly insulated, rod of length L [m] (see Figure 1). Time-varying temperatures are imposed on each end. At the left end the temperature is given by $\text{left}(t)$, while the at the right end the temperature is given by $\text{right}(t)$.

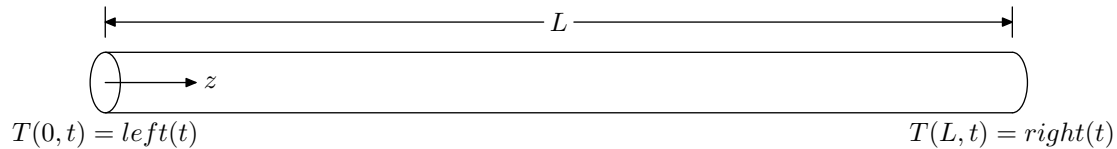


Figure 1: A thin, perfectly-insulated, rod with specified, time-varying, temperatures imposed on both ends.

We are interested in modeling the temperature along the entire rod as a function of location and time: $T(x, t)$. The initial temperature of the rod is uniformly 0. That is,

$$T(x, 0) = \text{left}(0) = \text{right}(0) = 0 \quad (1)$$

for all $0 \leq x \leq L$.

Mathematics

Temperature as a function of time and locations is described by a well-known partial differential equation. Not surprisingly, this partial differential equation is known as the *Heat Equation*¹.

$$\frac{\partial T}{\partial t} = \frac{k}{c_p \rho} \frac{\partial^2 T}{\partial x^2} \quad (2)$$

where

- T is the temperature,
- t is the time,
- k is the *thermal conductivity*²,
- c_p is the *specific heat capacity*³,
- ρ is the *mass density* of the material, and
- x is the space coordinate.

¹See, for example, http://en.wikipedia.org/wiki/Heat_equation.

²See, for example, http://en.wikipedia.org/wiki/Thermal_conductivity. According to this page, the thermal conductivity for Portland Cement is approximately 0.29 [W/(mKelvin)], which that of concrete and stone is approximately 1.7 [W/mKelvin].

³See, for example, http://en.wikipedia.org/wiki/Specific_heat_capacity.

The coefficient $\frac{k}{c_p \rho}$ is called the *thermal diffusivity*⁴ and is often denoted using α . This substitution visually simplifies the partial differential equation to

$$\frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial x^2} \quad (3)$$

This is the form that we will be working with in this problem.

MATLAB

The heat equation is the a *parabolic partial differential equation*⁵, and is amenable to a numerical solution using various methods. In fact, MATLAB has a built-in function for doing just that: `pdepe`.

`pdepe` Solve initial-boundary value problems for parabolic-elliptic PDEs in 1-D. `SOL = pdepe(M,PDEFUN,ICFUN,BCFUN,XMESH,TSPAN)` solves initial-boundary value problems for small systems of parabolic and elliptic PDEs in one space variable x and time t to modest accuracy. There are `npde` unknown solution components that satisfy a system of `npde` equations of the form

$$c(x,t,u,Du/Dx) * Du/Dt = x^{(-m)} * D(x^m * f(x,t,u,Du/Dx))/Dx + s(x,t,u,Du/Dx)$$

Here $f(x,t,u,Du/Dx)$ is a flux and $s(x,t,u,Du/Dx)$ is a source term. m must be 0, 1, or 2, corresponding to slab, cylindrical, or spherical symmetry, respectively. The coupling of the partial derivatives with respect to time is restricted to multiplication by a diagonal matrix $c(x,t,u,Du/Dx)$. The diagonal elements of c are either identically zero or positive.

An entry that is identically zero corresponds to an elliptic equation and otherwise to a parabolic equation. There must be at least one parabolic equation. An entry of c corresponding to a parabolic equation is permitted to vanish at isolated values of x provided they are included in the mesh `XMESH`, and in particular, is always allowed to vanish at the ends of the interval. The PDEs hold for $t_0 \leq t \leq t_f$ and $a \leq x \leq b$. The interval $[a,b]$ must be finite. If $m > 0$, it is required that $0 \leq a$. The solution components are to have known values at the initial time $t = t_0$, the initial conditions. The solution components are to satisfy boundary conditions at $x=a$ and $x=b$ for all t of the form

$$p(x,t,u) + q(x,t) * f(x,t,u,Du/Dx) = 0$$

$q(x,t)$ is a diagonal matrix. The diagonal elements of q must be either identically zero or never zero. Note that the boundary conditions are expressed in terms of the flux rather than Du/Dx . Also, of the two coefficients, only p can depend on u .

⁴See, for example, http://en.wikipedia.org/wiki/Thermal_diffusivity. This page gives various values for numerous materials, but nothing for concrete. You may also find http://www.wbdg.org/ccb/ARMYCOE/COESTDS/crd_c37.pdf interesting, if a bit dated.

⁵See, for example, http://en.wikipedia.org/wiki/Parabolic_partial_differential_equation.

The input argument `M` defines the symmetry of the problem. `PDEFUN`, `ICFUN`, and `BCFUN` are function handles.

`[C,F,S] = PDEFUN(X,T,U,DUDX)` evaluates the quantities defining the differential equation. The input arguments are scalars `X` and `T` and vectors `U` and `DUDX` that approximate the solution and its partial derivative with respect to `x`, respectively. `PDEFUN` returns column vectors: `C` (containing the diagonal of the matrix $c(x,t,u,Dx/Du)$), `F`, and `S` (representing the flux and source term, respectively).

`U = ICFUN(X)` evaluates the initial conditions. For a scalar `X`, `ICFUN` must return a column vector, corresponding to the initial values of the solution components at `X`.

`[PL,QL,PR,QR] = BCFUN(XL,UL,XR,UR,T)` evaluates the components of the boundary conditions at time `T`. `XL` and `XR` are scalars representing the left and right boundary points. `UL` and `UR` are column vectors with the solution at the left and right boundary, respectively. `PL` and `QL` are column vectors corresponding to `p` and the diagonal of `q`, evaluated at the left boundary, similarly `PR` and `QR` correspond to the right boundary. When $m > 0$ and $a = 0$, boundedness of the solution near $x = 0$ requires that the flux f vanish at $a = 0$. `pdepe` imposes this boundary condition automatically.

Making some sense of pdepe

Our problem at hand is

$$\frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial x^2} \quad (4)$$

As given in the MATLAB documentation, `pdepe` solves a partial differential equation of the following form:

$$c \left(x, t, u, \frac{\partial u}{\partial x} \right) \frac{\partial u}{\partial t} = x^{-m} \frac{\partial}{\partial x} \left(x^m f \left(x, t, u, \frac{\partial u}{\partial x} \right) \right) + s \left(x, t, u, \frac{\partial u}{\partial x} \right) \quad (5)$$

We are not working with an axis symmetric, or spherically symmetric, problem so $m = 0$, and (5) simplifies to

$$c \left(x, t, u, \frac{\partial u}{\partial x} \right) \frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left(f \left(x, t, u, \frac{\partial u}{\partial x} \right) \right) + s \left(x, t, u, \frac{\partial u}{\partial x} \right) \quad (6)$$

Our “`c`” function is simply a constant, $c = 1$, and we do not have an “`s`” function at all, $s = 0$. Thus, (6) simplifies to

$$\frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left(f \left(x, t, u, \frac{\partial u}{\partial x} \right) \right) \quad (7)$$

Looking at (7) we see that $f = \alpha \frac{\partial u}{\partial x}$, so we have

$$\frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left(\alpha \frac{\partial u}{\partial x} \right) = \alpha \frac{\partial^2 u}{\partial x^2} \quad (8)$$

Which is right where we wanted to be.

In summary,

- $m = 0$
- $u = T$
- $c = 1$
- $s = 0$
- $f = \alpha \frac{\partial u}{\partial x}$

An example

```
function T = Example( alpha, lambda1, lambda2 )
    m = 0;
    z = linspace(0,1,50);
    t = linspace(0,5,100);

    sol = pdepe(0, @pdefun, @icfun, @bcfun, z, t);
    % Extract the first solution component as T.
    T = sol(:,:,1);

    % A surface plot is often a good way to study a solution.
    surf(z,t,T)
    title('Numerical solution.')
    xlabel('Depth z')
    ylabel('Time t')

    % A solution profile can also be illuminating.
    figure
    plot(z,T(end,:))
    title('Solution at t = 2')
    xlabel('Depth z')
    ylabel('T(z,2)')

    % -----
    function [c,f,s] = pdefun(x, t, T, dTdx)
        c = 1;
        f = alpha * dTdx;
        s = 0;
    end

    % -----
    function T0 = icfun(z)
        T0 = 0;
    end

    % -----
```

```
function [pl,ql,pr,qr] = bcfun(zl, Tl, zr, Tr, t)
    pl = Tl - 2*sin(2*pi*t/lambda1);
    ql = 0;
    pr = Tr - sin(2*pi*t/lambda2);
    qr = 0;
end
end
```