

Data Review

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Intermolecular Energies

Supermolecular approach

$$E_{int} = E_{AB} - E_A - E_B$$

- ► Straightforward, but cannot separate different types of interactions
- ► Can adopt to different electronic structure methods
- ► DFT-D3 with proper functional can be both cheap and accurate
- Symmetry-Adapted Perturbation Theory
 - ► Can give details about different types of interactions; important in understanding their nature
 - ▶ Not as cheap as DFT-D3
 - SAPT0 is somewhat cheap, but does not include intramonomer correlation

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SAPT(DFT)

- Attempt to inlude intramonomer correlation in a cheap way
- Replaces HF orbitals with KS orbitals
- Needs to consider orbital response for dispersion terms
- Exchange-dispersion term needs to be estimated from scaling
- Investigate the accuracy and efficiency of SAPT(DFT)

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Three-Body Interaction

- Crucial in computing lattice energies
- DFT-D3 does not perform well for three-body interaction
- ▶ MP2.5 scales as $O(N^6)$, MP2 is $O(N^5)$ but lacks three-body dispersion
- Three-body dispersion can be implemented with SAPT(DFT) in ${\cal O}(N^5)$
- Combine MP2 with SAPT(DFT) dispersion to model three-body interaction

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Idea of SAPT(DFT)

SAPT energy in orders of interaction and fluctuation potentials; n denotes order in V and k, l for W_A, W_B

$$H = F_A + F_B + V + W_A + W_B$$

$$E_{int} = \sum_{n=1}^{\infty} \sum_{k=0}^{\infty} \sum_{l=0}^{\infty} \left(E_{pol}^{(nkl)} + E_{exch}^{(nkl)} \right)$$

- SAPT0: n = 2, k = l = 0, no intramonomer correlation, $O(N^5)$ cost
- ▶ Many-body SAPT: $k, l \ge 2$, $O(N^7)$ or higher cost
- SAPT(DFT): Use Kohn-Sham operator $K_{A,B}$ instead of Fock operator $F_{A,B},\,O(N^5)$ cost
- Primitive SAPT(DFT) works well on 1st-order terms, but not 2nd-order terms (especially dispersion). Needs orbital response for them

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Dispersion Term

- ALDA kernel good for pure GGA functional, but not for hybrid functional
- ▶ Exact exchange in v_{xc} → increased ϵ^{ab}_{ij} → decreased $E^{(2)}_{disp}$

$$E_{disp,u}^{(2)} = -4 \sum_{ia \in A, jb \in B} \frac{\left| (ia|jb) \right|^2}{\epsilon_{ij}^{ab}}$$

▶ Hybrid ALDA kernel to compensate, or localized HF (LHF) exchange to avoid increase in ϵ^{ab}_{ij}

$$f_{xc} = \alpha f_{xc}^{HF} + (1 - \alpha) f_{xc}^{ALDA}$$

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Equation for $E_{disp}^{(2)}$

Computing coupled FDDS:

$$\chi = \chi'_0 + (\chi'_0 \mathbf{S}^{-1} \mathbf{W} + \mathbf{K}) \left[\mathbf{S} - (\chi'_0 \mathbf{S}^{-1} \mathbf{W} + \mathbf{K}) \right]^{-1} \chi'_0$$
$$\chi'_0 = \chi_0 - \alpha \mathbf{K}_2 (\lambda)$$
$$\mathbf{K} = \left[-\alpha \mathbf{K}_1 (\lambda d) - \alpha \mathbf{K}_2 (\lambda d) + \alpha^2 \mathbf{K}_{21} (\lambda) \right] (\mathbf{R}^t)^{-1} \mathbf{S}$$

- $lackbox{O}(N^5)$ scaling is limited to forming ${f K}_1$, ${f K}_2$ and ${f K}_{21}$
- Separates nontrivial and trivial parts of the code, also highly reduces need of disk I/O operations
- $\triangleright E_{disp}^{(2)}$ from coupled FDDS:

$$E_{\text{disp}}^{(2)} = -\frac{1}{2\pi} \int_0^\infty d\omega \int d\mathbf{r}_A d\mathbf{r}_A' d\mathbf{r}_B d\mathbf{r}_B'$$
$$\frac{1}{|\mathbf{r}_A - \mathbf{r}_B|} \frac{1}{|\mathbf{r}_A' - \mathbf{r}_B''|} \chi^A (\mathbf{r}_A, \mathbf{r}_A' | i\omega) \chi^B (\mathbf{r}_B, \mathbf{r}_B' | i\omega)$$



Exchange-Dispersion Term

- Coupled exchange-dispersion requires storing 4-index tensors on disk
- Estimate from scaling uncoupled exchange-dispersion; Scale with ratio in dispersion term or with pre-fitted (with S22×5) fixed factor

$$\tilde{E}_{\text{exch-disp,r}}^{(2)} = E_{\text{exch-disp,u}}^{(2)} \cdot \frac{E_{\text{disp,r}}^{(2)}}{E_{\text{disp,u}}^{(2)}}$$

$$\tilde{E}_{\text{exch-disp,r}}^{(2)} = \alpha \cdot E_{\text{exch-disp,u}}^{(2)}(\alpha = 0.686361)$$

- lacktriangle The value above is fitted from $E_{
 m exch-disp,u}^{(2)}$ with LHF orbitals
- Non-LHF orbitals have greater o-v gaps and smaller $E_{
 m disp,u}^{(2)}$, needs to re-fit with non-LHF results

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LHF vs non-LHF orbitals

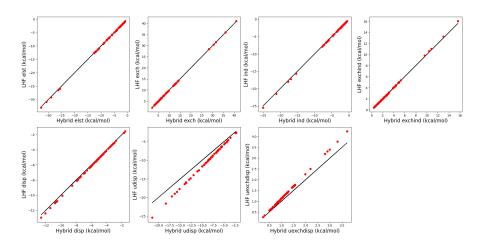


Figure: *

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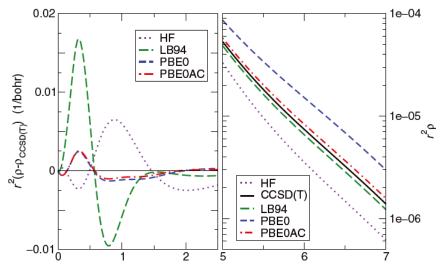
GRAC

- ► (Hybrid-)GGA functionals does not have correct long-range behavior $v_{xc}(r) \rightarrow -1/r + (I_p + \epsilon_{\text{HOMO}})$
- Underestimates o-v gap as a consequence
- Functionals like LB94 have correct asymptotic behavior, but poor in bulk region
- ▶ Using gradient-regulated asymptotic correction (GRAC) scheme to connect PBE0 and LB94 with the switching function f[g(r)]:

$$\begin{aligned} v_{xc}^{\text{GRAC}} &= \left\{1 - f[g(\boldsymbol{r})]\right\} v_{xc}^{\text{PBE0}} + f[g(\boldsymbol{r})] v_{xc}^{\text{LB94}} \\ &f[g(\boldsymbol{r})] = \left(1 + e^{-\alpha[g(\boldsymbol{r}) - \beta]}\right)^{-1} \\ &g(\boldsymbol{r}) = \frac{|\nabla \rho(\boldsymbol{r})|}{\rho^{4/3}(\boldsymbol{r})} \end{aligned}$$



GRAC & Long-Range Behavior



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Test 5

Test0	Test1	Test2
Test1	Test2	Test3 LATEX
Test4	Test5	Test6

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Test 6

Example block

Test0	Test1	Test2
Test1	Test2	Test3 LATEX
Test4	Test5	Test6

Block

Test0	Test1	Test2
Test1	Test2	Test3 LATEX
Test4	Test5	Test6

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Test 7

Block A

short Short stuff

long Longer stuff

longest label Longest stuff (insert cat)

- item1
- ▶ item2
- item3



Frame with Columns

Block 1

Text here

Block 2

More text here



Frame without Columns

Block

Even more text here

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