

ASTR2013 – Foundations of Astrophysics

Week 3: Physics of Stars

Following Dan Maoz – Astrophysics in a Nutshell

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Recap – The Virial theorem and Stellar Timescales

$$E_{
m th}^{
m tot} = -rac{E_{
m gr}}{2}$$
 The Virial Theorem for Stars (written in 3 forms) $E_{
m gr} = -2E_{
m th}^{
m tot}$ $E^{
m tot} = -rac{1}{2}E_{
m gr}$

• The *dynamical timescale* for any body of mass M and radius

R is: $\tau_{\rm ff} \approx \tau_{\rm dyn} = \sqrt{\frac{R^3}{GM}}$ $\propto \bar{\rho}^{-1/2}$

- This is approximately the *free-fall* timescale, $1/2\pi$ times the orbital period of a body almost scraping the surface, or the sound-crossing time.
- We can take the ratio of the *total internal energy* to the *stellar luminosity* to get the Kelvin-Helmholtz timescale:

~10
7
 years $au_{
m KH} pprox rac{1}{2} rac{GM^2}{R} rac{1}{L}$



Week 3 Summary

Textbook: Sections Finishing Chapter 3 (other than quantum mechanics in 3.9 and 3.10).

- 1. Main Sequence Scaling Relations
- 2. Equations of Stellar Structure.
- 3. Equation of State
- 4. Convective instability and convective energy transport.
- 5. Nuclear generation by the proton-proton chain and the CNO cycle.



Solar Parameters

• The sun has an effective temperature of 5772K, and a radius of 6.957 x 10⁸W, resulting in a Luminosity of 3.828 x 10²⁶W. Why not a different radius?

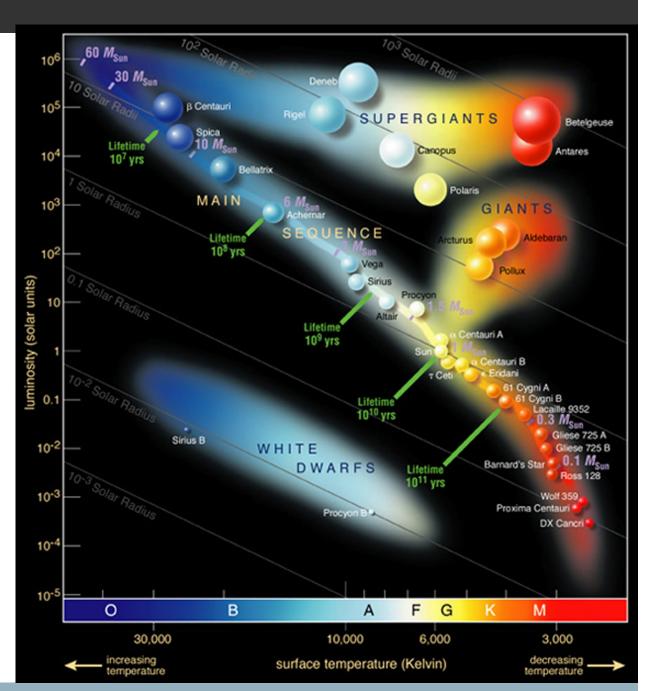
$$L = 4\pi R^2 \sigma_{\rm SB} T_{\rm eff}^4$$

 Radioactive dating of rocks in the solar system given an age of 4.6 Gyr, and the sun appears middle-aged compared to other 1 Msun stars. This means its lifetime is 10Gyr. Why not a different age?



Scaling Relations

- Luminosity is roughly ~M⁴
 for intermediate mass
 stars, ~M⁵ for low mass
 stars and ~M for high mass
 stars. Why?
- Life time goes as M3 for intermediate mass stars – assuming nuclear burning you can figure this out for yourself based on the above...
- Mass appears to relate to radius as roughly M~R. Again... why?





Stellar Structure Equations

- Conservation of mass, momentum and energy are the underlying principles, resulting in 4 differential equations as a function of radius coordinate r or in other text and codes, mass coordinate M_r.
- Mass Continuity:

$$\frac{dM(r)}{dr} = 4\pi r^2 \rho(r)$$

Hydrostatic Equilibrium:

$$\frac{dP(r)}{dr} = -\frac{GM(r)\rho(r)}{r^2}$$



Pressure in a little more detail

- Pressure actually as 3 components:
 - Ideal gas pressure:

$$P_{\rm gas} = nk_B T = \frac{\rho k_B}{\mu u} T$$

– Radiation pressure:

$$P_{\rm rad} = \frac{u}{3} = \frac{4\sigma_{\rm SB}}{3c}T^4$$

 Degeneracy pressure... important as the electron density approaches 1 electron per cube of the de Broglie wavelength:

$$\lambda_e = \frac{n}{mv_e}$$



Energy Transport

 If radiation transports energy, then the derivative of the radiative energy density is proportional to luminosity.

$$\frac{du(r)}{dr} \propto \frac{dT(r)^4}{dr} \propto F(r) = \frac{L(r)}{4\pi R^2}$$

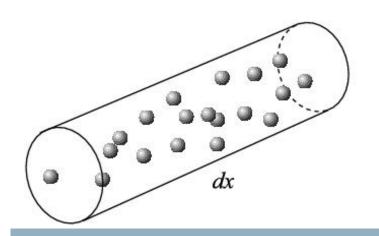
 If internal kinetic energy of gas transports energy perfectly, then the temperature gradient with pressure is the adiabatic gradient:

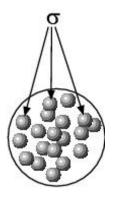
$$\frac{dT}{dr} = \frac{\gamma - 1}{\gamma} \frac{T}{P} \frac{dP}{dr} \qquad \text{from } PV^{\gamma} = \text{const}$$

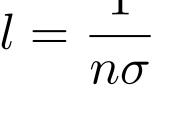
more on this later...

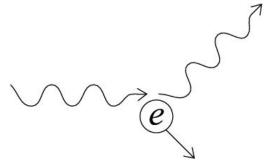


- Electrons, atoms and ions absorb or scatter photons inside a star.
- Each particle has a cross-section $\sigma(\lambda)$ to absorbing radiation of wavelength λ .
- Given a particle number density n, we can define a mean-free path:











• A mean free path is inversely proportional to density, so we instead typically talk about opacity: $\Sigma n_i \sigma_i$

 $\kappa = \frac{\Sigma n_i \sigma_i}{
ho}$

 For mostly ionized gas, Thompson scattering provides a lower limit to opacity:

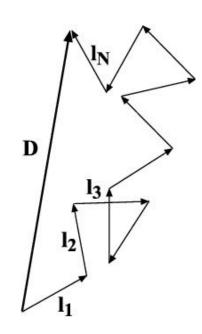
$$\sigma_{\rm T} = \frac{8\pi}{3} \left(\frac{e^2}{m_e c^2} \right)^2 = 6.7 \times 10^{-25} \, {\rm cm}^2.$$

 Free-Free, and bound-free opacities are proportional to density. Bound-bound opacities matter much less in stellar interiors.

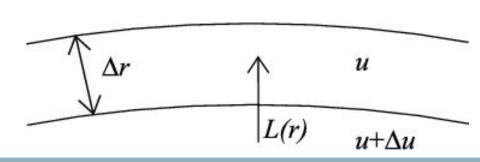


- Photons diffuse through the star, with an average mean free path of 2cm.
- Considering neighboring shells with radiative energy density u, changing with r, we can write:

 $\frac{L(r)}{4\pi r^2} = -\frac{cl}{3} \frac{du}{dr}$



• In a more rigorous derivation, the factor of 3 is called the *Eddington Approximation*, and is accurate deep in the star.



The energy density derivative is:

$$\frac{du}{dr} = \frac{du}{dT} \frac{dT}{dr}$$
$$= \frac{4\sigma_{\rm SB}}{c} T^3 \frac{dT}{dr}$$

• We can write the mean free path in terms of the Rosseland Mean Opacity as: $l = (\kappa_R \rho)^{-1}$

This gives the radiative energy flow equation:

$$\frac{dT(r)}{dr} = -\frac{3L(r)\kappa_R(\rho,T)\rho(r)}{64\pi r^2\sigma_{\rm SB}T^3(r)} \quad \text{(textbook uses a instead of } \sigma_{\rm SB})$$



Energy Generation

- On the *nuclear timescale*, we can approximate a star as being in steady state, with constant energy generation from nuclear reactions.
- Energy generation is conventionally parameterized by ε, which is the power generated per unit mass.
- Energy conservation gives:

$$\frac{dL(r)}{dr} = 4\pi r^2 \rho(r) \epsilon(\rho, T)$$



Equations of Stellar Structure

$$\frac{dP(r)}{dr} = -\frac{GM(r)\rho(r)}{r^2},\tag{3.56}$$

$$\frac{dM(r)}{dr} = 4\pi r^2 \rho(r), \tag{3.57}$$

$$\frac{dT(r)}{dr} = -\frac{3L(r)\kappa(r)\rho(r)}{4\pi r^2 4acT(r)^3},$$
(3.58)

$$\frac{dL(r)}{dr} = 4\pi r^2 \rho(r) \epsilon(r). \tag{3.59}$$

The opacity, nuclear generation rate and pressure are really a complex function of density and temperature.



Equations of Stellar Structure

 Pressure as a function of density and temperature is and equation of state, e.g.

$$P_{\rm gas} = nk_B T = \frac{\rho k_B}{\mu u} T$$

- [noting that $\mu = \mu(\rho, T)$ in general]
- We also need boundary conditions including the obvious ones and something like:

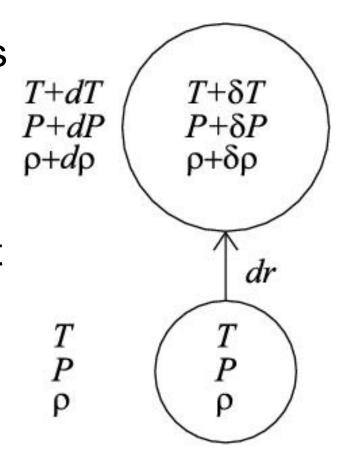
$$L(r_*) = 4\pi R^2 \sigma_{\rm SB} T^4(r \text{ near surface})$$

 Here "near surface" is often 2/3 of a photon mean-free path from r*, but needs a multiwavelength photospheric model to be accurate.



Convective Instability

- A parcel of gas that moves upwards without exchanging energy with its surroundings is expanding adiabatically.
- The gradient dP/dT, or the gradient of logarithms for this parcel, is called the *adiabatic gradient*.
- If the parcel expands and cools more than the surrounding gas, it becomes buoyant and accelerates upwards, i.e. it is *unstable*.





Convective Instability

• The adiabatic gradient is therefore the steepest possible temperature gradient in a star. If γ is too small, or κ too large (or both) the star has a convective region.

 γ =5/3 for a monatomic gas γ ~1.1 when hydrogen is 50% ionized

Maximum luminosity carried by convection for a gas with energy density u_{gas} = is of order:

$$L_{\rm conv.} < 4\pi r^2 u_{\rm gas} v_s$$

$$PV^{\gamma} = \text{const.}$$
 $P \propto \rho^{\gamma}$
 $\frac{d \log(P)}{d \log(\rho)} = \gamma$

$$\nabla_{\text{ad}} = \frac{d \log(T)}{d \log(P)} = \frac{\gamma - 1}{\gamma}$$



Nuclear Generation

 Nuclear energy generation in the main sequence comes from two key reaction chains. The protonproton chain:

$$p + p \to d \text{ (1.4 MeV)}$$

 $p + d \to {}^{3}_{2}\text{He (5.5 MeV)}$
 ${}^{3}_{2}\text{He} + {}^{3}_{2}\text{He} \to {}^{4}_{2}\text{He} + 2p \text{ (12.9MeV)}$

• ... and the CNO cycle:

More on this with numbers in the labs!

$$p + {}^{12}C \rightarrow {}^{13}N + \gamma$$

 ${}^{13}N \rightarrow {}^{13}C + e^{+} + \nu_{e}$,
 $p + {}^{13}C \rightarrow {}^{14}N + \gamma$,
 $p + {}^{14}N \rightarrow {}^{15}O + \gamma$,
 $p + {}^{15}O \rightarrow {}^{15}N + e^{+} + \nu_{e}$,
 $p + {}^{15}N \rightarrow {}^{12}C + {}^{4}He$.



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