

Introduction to Dark Matter

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October 27, 2023

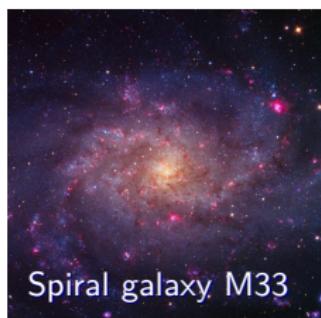


Dark Matter in the Universe

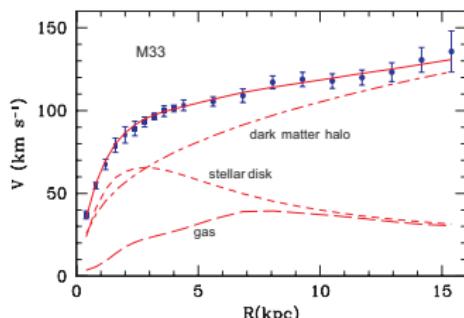
 **Dark matter (DM)** makes up **most** of the **matter** component in the Universe, as suggested by **astrophysical** and **cosmological** observations



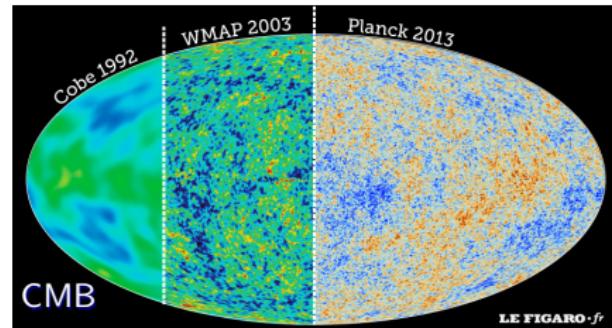
Bullet Cluster



Spiral galaxy M33



Cluster Abell 2218



Coma Cluster (后发座星系团)



Coma Cluster (后发座星系团)



Coma Cluster (后发座星系团)



Coma Cluster



★ In 1933, **Fritz Zwicky** found that the **velocity dispersion** of **galaxies** in the **Coma cluster** was **far too large** to be supported by the **luminous matter**

 **Mass-to-light ratio** $\Upsilon_{\text{Coma}} \sim 240 \Upsilon_{\odot}$

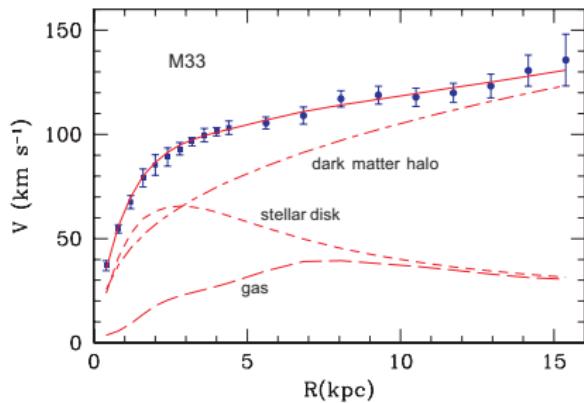
[Kent & Gunn, Astron.J., 87, 945 (1982)]

Typical spiral galaxy: $\mathcal{O}(10)\Upsilon_\odot$

Spiral Galaxies: Rotation Curves



In the 1970s, **Vera Rubin** and her collaborators measured the **rotation curves** of **spiral galaxies** and also found evidence for **non-luminous matter**



[Corbelli & Salucci, astro-ph/9909252, MNRAS]

Spiral Galaxies: Rotation Curves



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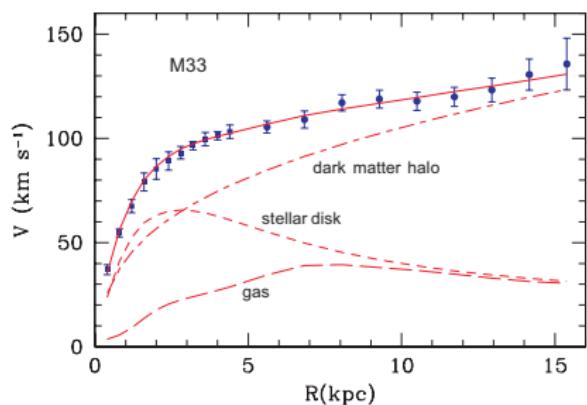


According to **Newton's law of universal gravitation**, the relation between the **rotation velocity** v and the **mass** $M(r)$ within the **radius** r should be

$$\frac{v^2}{r} = \frac{G_N M(r)}{r^2}$$

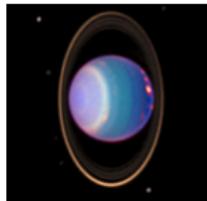
$$M(r) = \text{constant} \quad \text{---} \quad v \propto r^{-1/2}$$

$$M(r) \propto r \quad \text{☞} \quad v = \text{constant}$$



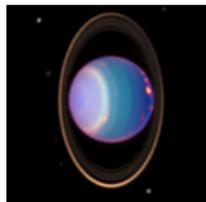
[Corbelli & Salucci, astro-ph/9909252, MNRAS]

How Can We Explain an Anomalous Phenomenon?



Unexpected movement of **Uranus** (after 1821)

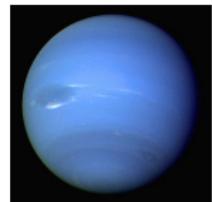
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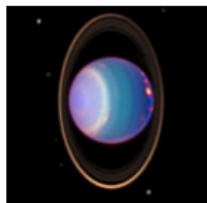


Perturbed by **Neptune** (discovered in 1846)



Calculations independently given by **John Adams** and **Urbain Le Verrier**

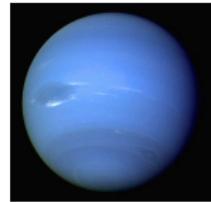
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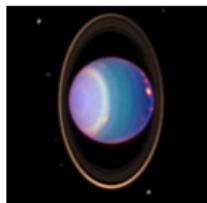
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Search for new objects/substances responsible for it!



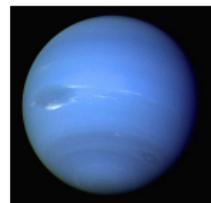
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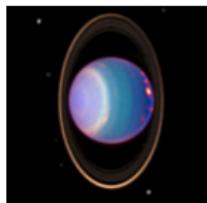
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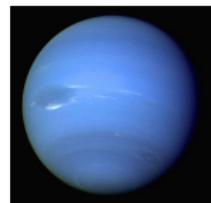
Anomalous perihelion precession of **Mercury** (**Le Verrier**, 1859)



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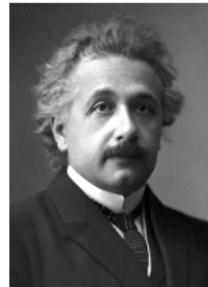
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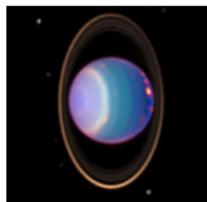
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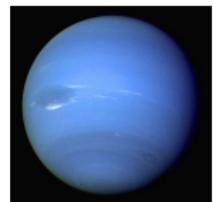
Update Newton's laws to **general relativity** (1915)



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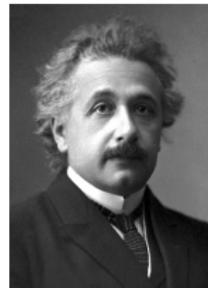
Anomalous perihelion precession of **Mercury** (Le Verrier, 1859)



Update Newton's laws to **general relativity** (1915)



Modify known physical laws!



How about the Anomalous Phenomena Here?



Modify physical laws ➡ MOdified Newtonian Dynamics (MOND)

[Milgrom, ApJ, 270, 365 (1983)]



It is **difficult** to coherently explain data at **all scales** with one model

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Consider new objects ➡ **MAssive Compact Halo Objects (MACHOs)**

They belong to **baryonic dark matter**, including **jupiters, brown dwarfs, white dwarfs, neutron stars, black holes, etc.**



MACHO fraction in the **Galactic dark matter halo** is **< 8% (95% C.L.)**

[EROS-2 Coll., astro-ph/0607207, A&A]

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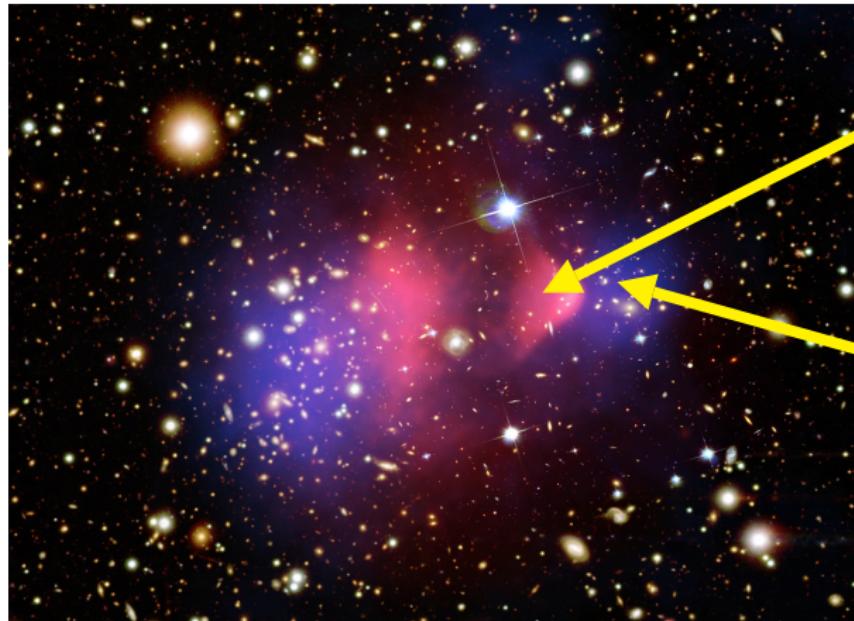
[EROS-2 Coll., astro-ph/0607207, A&A]



Consider new substances **Nonbaryonic Dark Matter**

(not constituted by baryons)

Bullet Cluster: Disfavor MOND



Fluid-like X-ray emitting plasma,
i.e., gas
(luminous matter)

Mass distribution
observed by weak
gravitational lensing
(DM dominated)

😦 An 8σ significance **spatial offset** of the **center** of the **total mass** from the **center** of the **baryonic mass peaks** **cannot be explained** with a **modification** of the **gravitational force law** [Clowe *et al.*, astro-ph/0608407, ApJL]

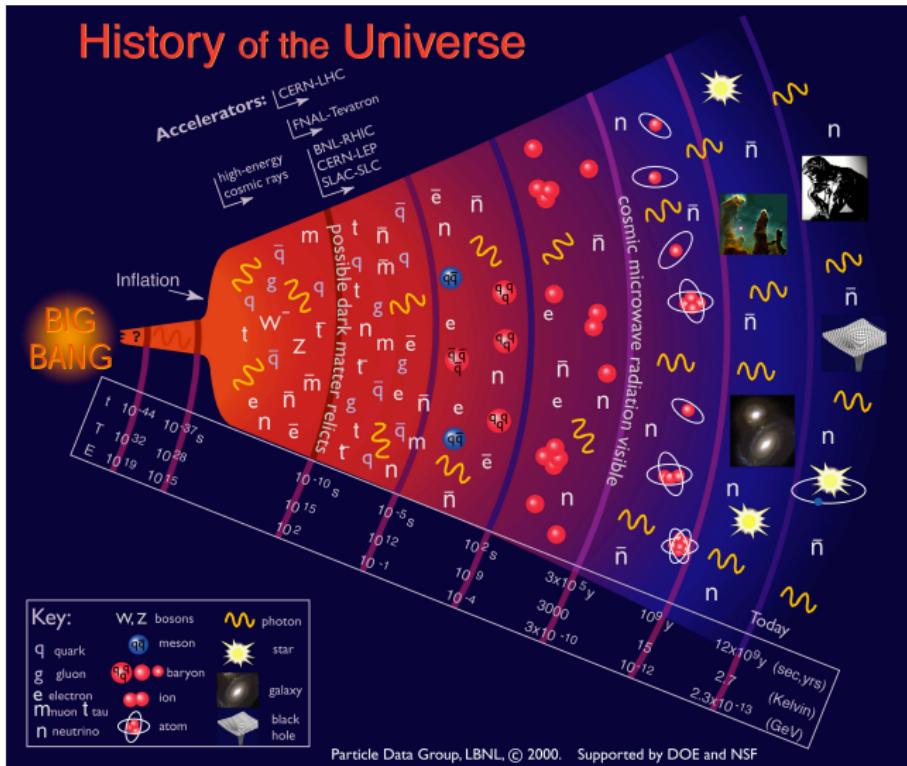
Big Bang Cosmology

~ 13.8 billion years ago, the Universe was extremely **hot, dense, and homogeneous**

Everything was in **thermal equilibrium** and interacted with each other

As the Universe **expanded and cooled down**, its constituents **decoupled** from the thermal bath **one by one**

Then **nuclei, atoms, stars, and galaxies** were **formed**



Structure Formation: Hot, Cold, and Warm Dark Matter



Small initial fluctuations + Gravitational instability



Decoupled matter generates cosmological structures



Baryonic matter decoupled too late



Only baryonic matter  Galaxies would not be formed!



 Needs nonbaryonic dark matter which decoupled much earlier

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Hot dark matter (such as neutrinos): relativistic when it decoupled

 Structures are formed by fragmentation (top-down)



Cold dark matter (CDM): nonrelativistic when it decoupled

 Structures are formed hierarchically (bottom-up)



The observation that galaxies are older than clusters favors cold dark matter

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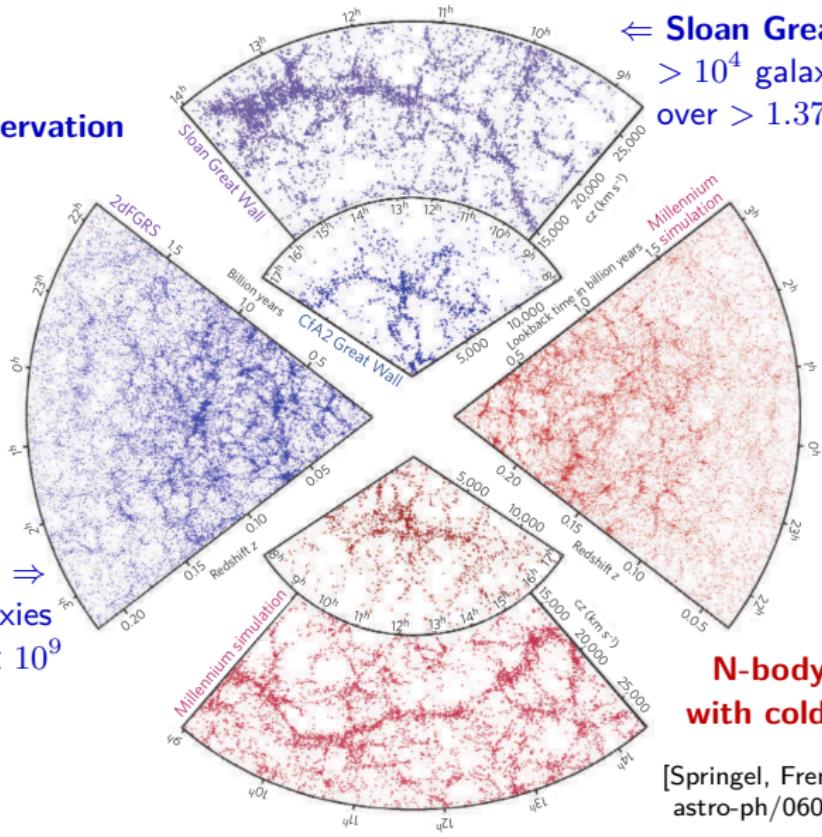
Milky Way dwarf satellites: ~ 60 (observed) vs. ~ 500 (CDM predicted)



“Missing satellites problem” A component of warm dark matter?

Galaxy Distribution: Observation vs Simulation

Observation



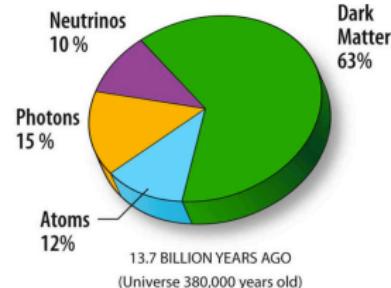
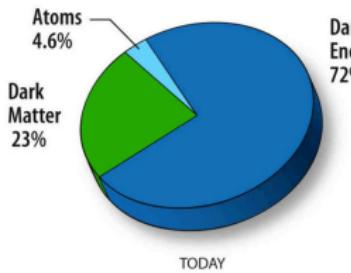
Standard Cosmology: Λ CDM Model



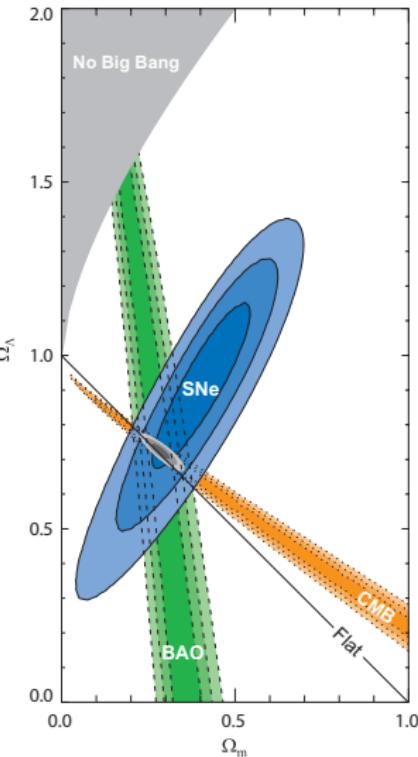
Λ CDM: the **standard cosmological model**

- **Cosmological constant Λ** (dark energy)
- **Cold dark matter (CDM)**
-  The evolution of the Universe is governed by the **Friedmann equation**

$$\frac{k}{H^2 R^2} = \Omega_\Lambda + \Omega_m + \Omega_r - 1$$



[WMAP Science Team]



[Kowalski *et al.*, 0804.4142, ApJ]



Cosmic Microwave Background (CMB)

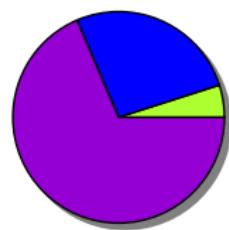
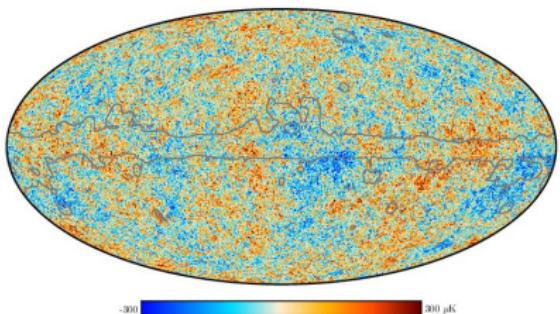
$t \sim 380\,000$ yr, $T \sim 3000$ K

Electrons + Protons \rightarrow Hydrogen Atoms
Photons decoupled

Cools \downarrow down

Today, ~ 2.7 K microwave background

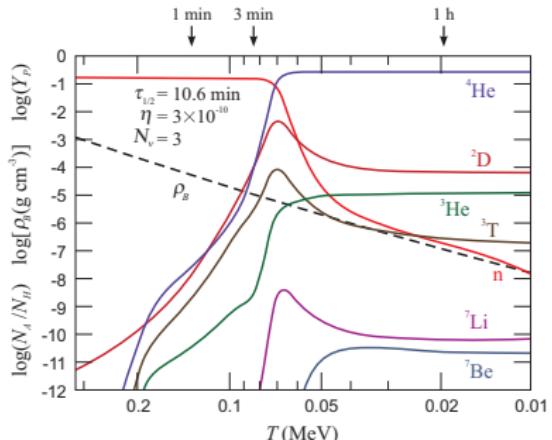
 **Cosmological parameters** Ω_Λ , Ω_c , and Ω_b can be determined by measuring the **CMB anisotropy power spectrum**



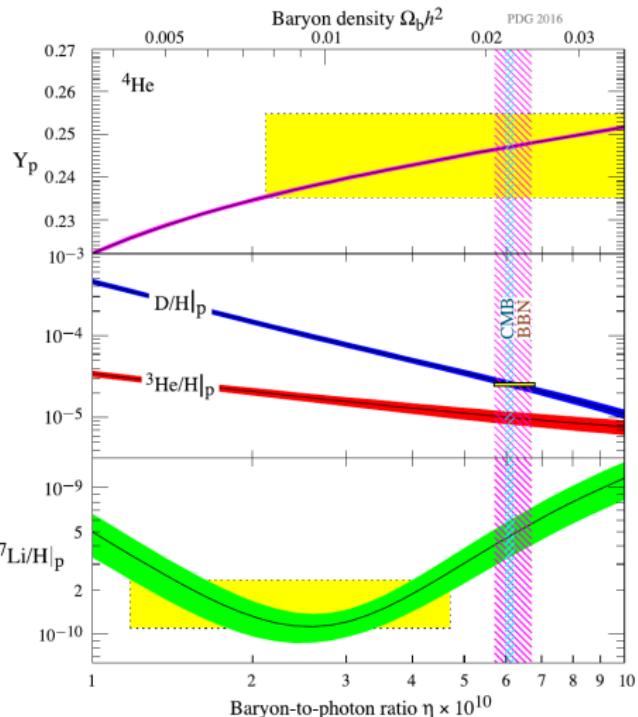
Planck 2018
[1807.06205, 1807.06209]

Cold DM (26.5%)
 $\Omega_c h^2 = 0.1200 \pm 0.0012$
Baryons (4.9%)
 $\Omega_b h^2 = 0.02237 \pm 0.00015$
Dark energy (68.6%)
 $\Omega_\Lambda = 0.6847 \pm 0.0073$

Big Bang Nucleosynthesis (BBN): $t \sim 1 \text{ sec} - 1 \text{ hour}$



[Kolb & Turner, *The Early Universe*]



Light element primordial abundances



Infer the **baryon density Ω_b**
(consistent with **CMB observations**)



The majority of matter is **nonbaryonic**

[PDG 2016]



Inferred Properties of Dark Matter

- ① **Dark (electrically neutral): no light emitted** from it
- ② **Nonbaryonic:** BBN & CMB observations
- ③ **Long lived:** survived from early eras of the Universe to now
- ④ **Colorless:** otherwise, it would **bind with nuclei**
- ⑤ **Cold:** structure formation theory
- ⑥ **Abundance:** **more than 80%** of all matter in the Universe

$$\rho_{\text{DM}} \sim 0.3\text{--}0.4 \text{ GeV/cm}^3 \text{ near the earth}$$

Standard Model (SM) of Particle Physics

 $SU(3)_C \times SU(2)_L \times U(1)_Y$ gauge symmetry



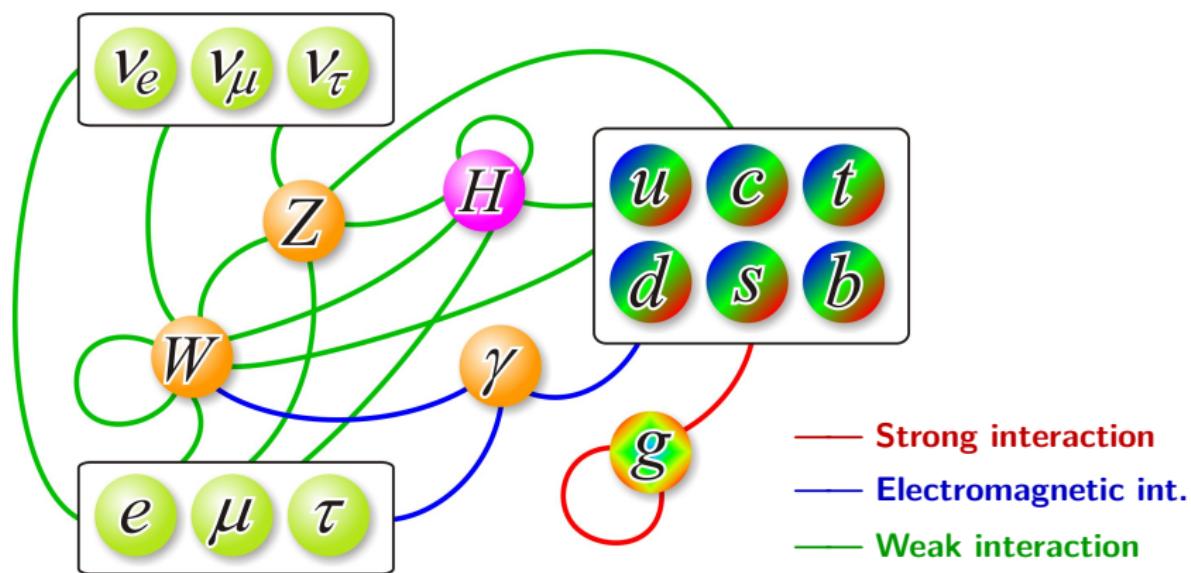
Fermion masses

 Englert-Brout-Higgs mechanism



 $SU(3)_C \times U(1)_{EM}$ gauge symmetry

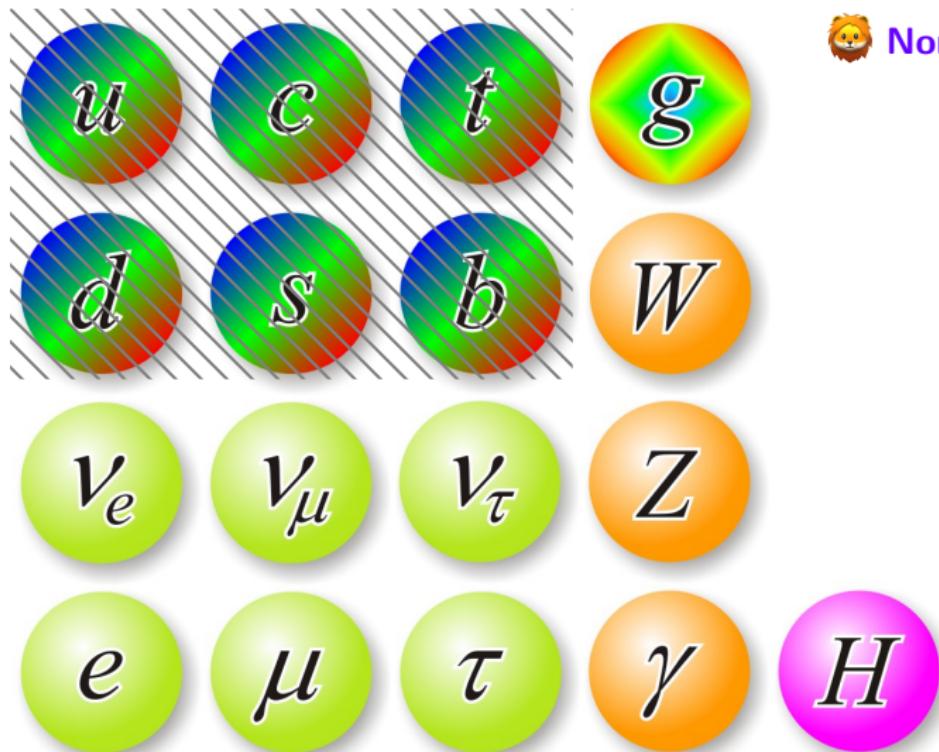
Higgs boson



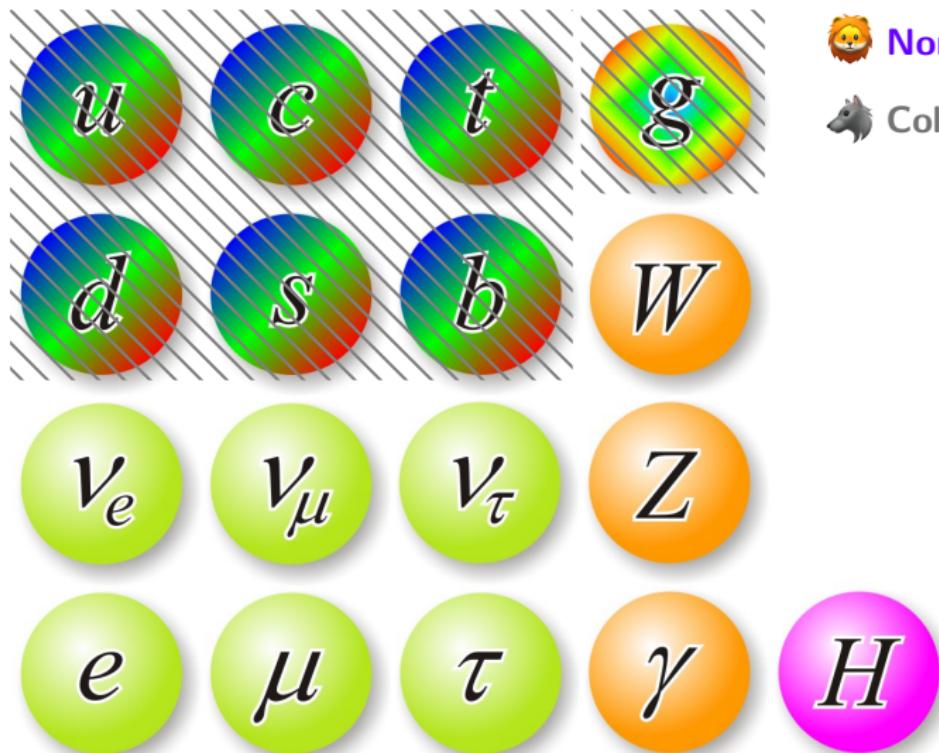
Are There Dark Matter Candidates in the Standard Model?



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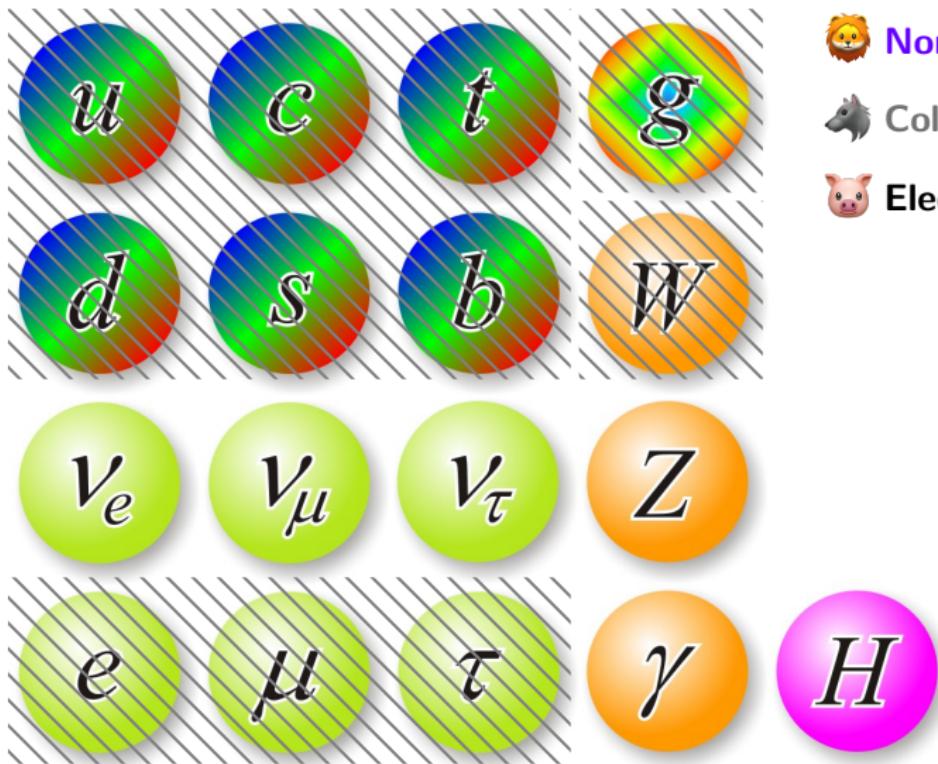


Nonbaryonic



Colorless

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Nonbaryonic

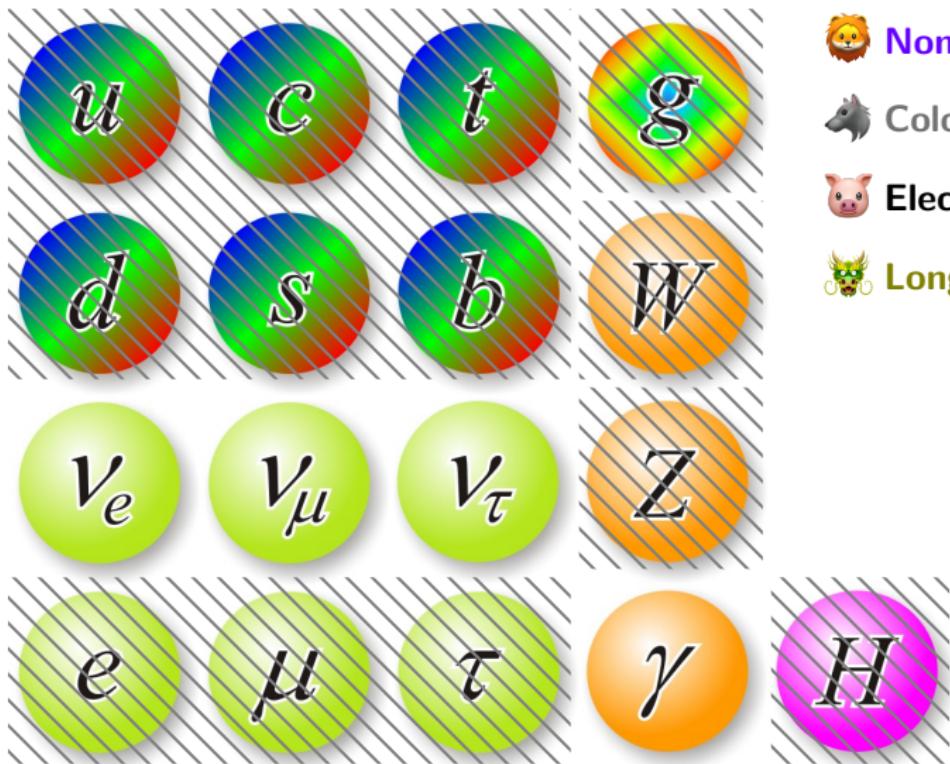


Colorless



Electrically neutral

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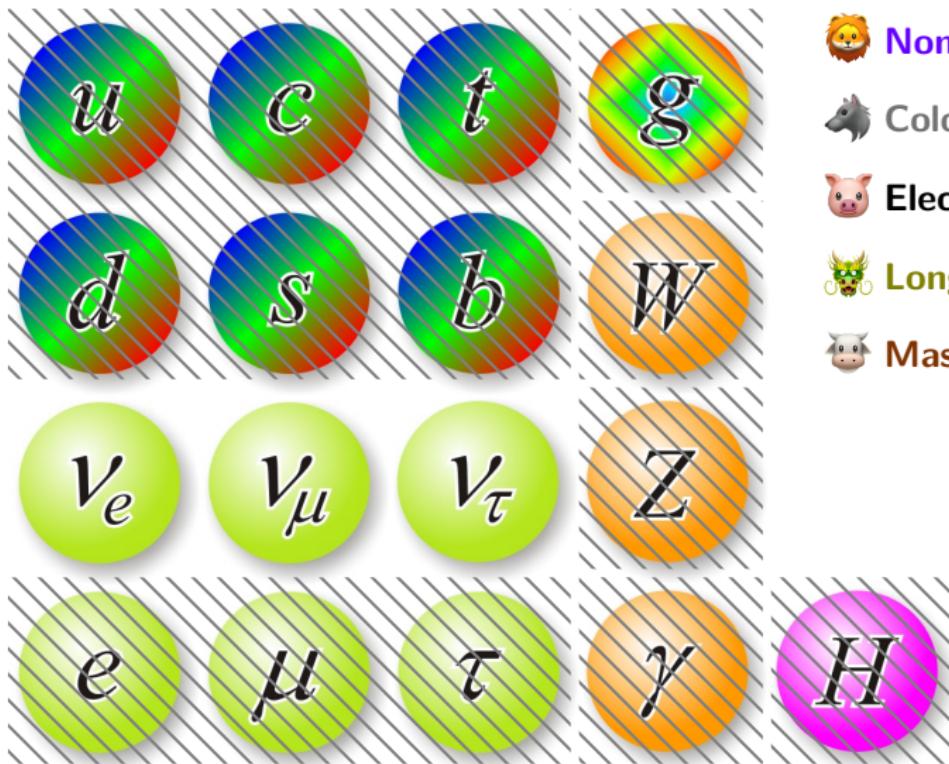


Electrically neutral



Long lived

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Nonbaryonic



Colorless



Electrically neutral

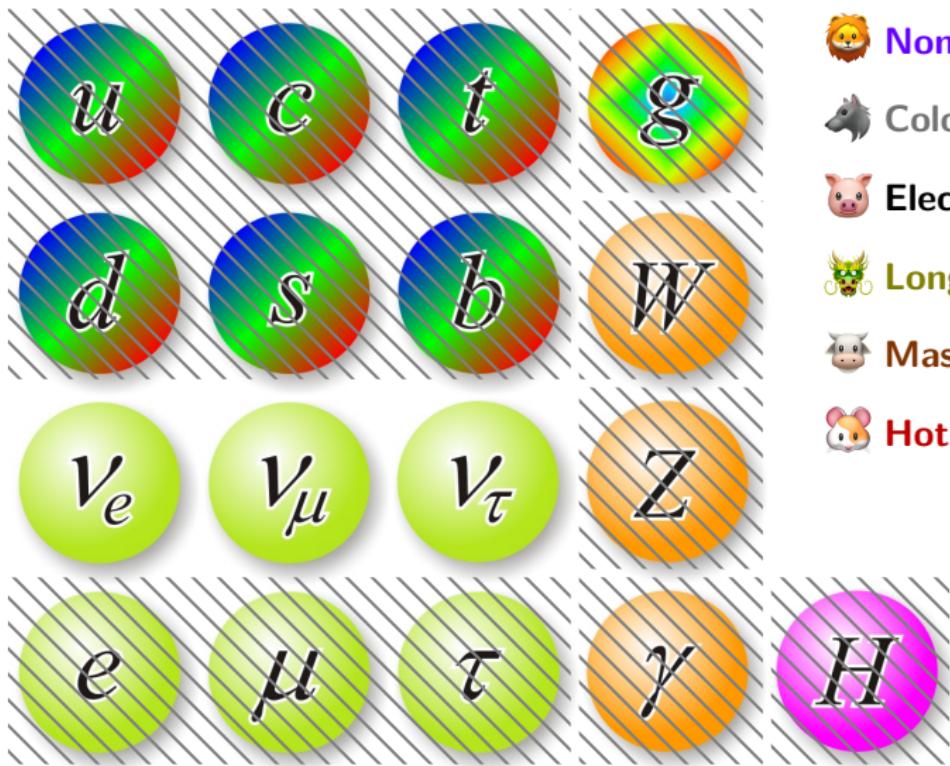


Long lived



Massive

Are There Dark Matter Candidates in the Standard Model?



Nonbaryonic



Colorless



Electrically neutral



Long lived

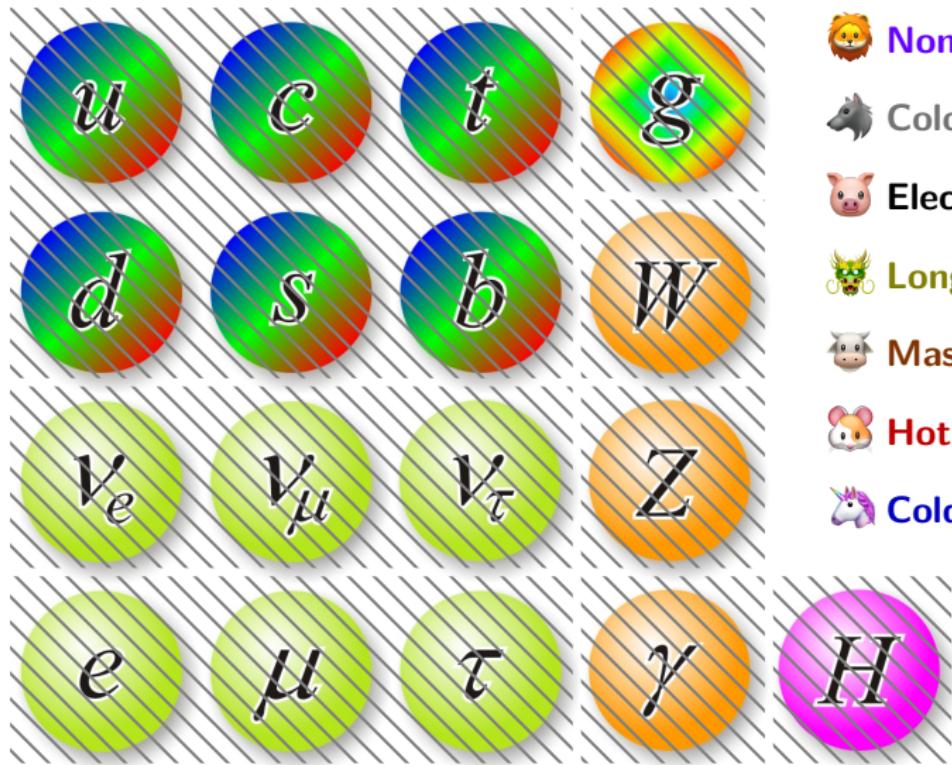


Massive



Hot DM: neutrinos

Are There Dark Matter Candidates in the Standard Model?



Nonbaryonic



Colorless



Electrically neutral



Long lived



Massive



Hot DM: neutrinos



Cold DM: none

DM Relic Abundance from Thermal Production

If DM particles (χ) were thermally produced in the early Universe, their relic abundance would be determined by the annihilation cross section $\langle\sigma_{\text{ann}}v\rangle$:

$$\Omega_\chi h^2 \simeq \frac{3 \times 10^{-27} \text{ cm}^3 \text{ s}^{-1}}{\langle\sigma_{\text{ann}}v\rangle}$$

Observed value $\Omega_\chi h^2 \simeq 0.1$

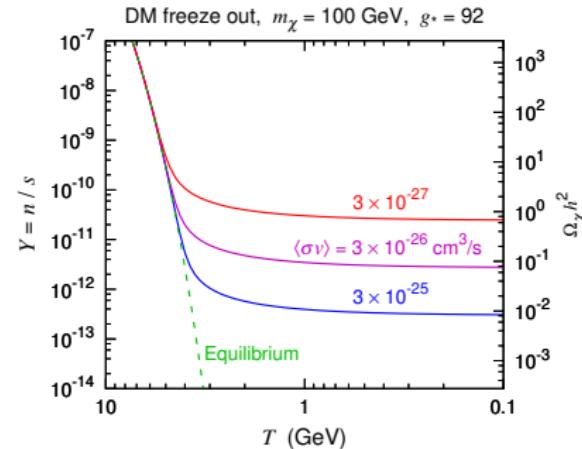
👉 $\langle\sigma_{\text{ann}}v\rangle \simeq 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$

Assuming the annihilation process consists of two weak interaction vertices with the $SU(2)_L$ gauge coupling $g \simeq 0.64$, for $m_\chi \sim \mathcal{O}(\text{TeV})$ we have

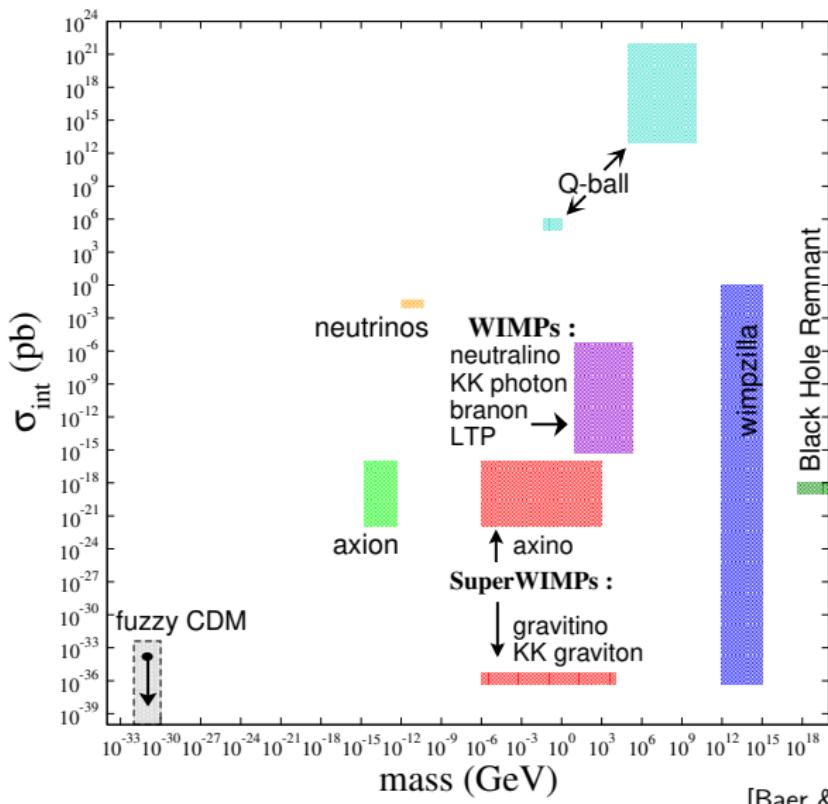
$$\langle\sigma_{\text{ann}}v\rangle \sim \frac{g^4}{16\pi^2 m_\chi^2} \sim \mathcal{O}(10^{-26}) \text{ cm}^3 \text{ s}^{-1}$$

👉 A very attractive class of DM candidates:

Weakly interacting massive particles (WIMPs)

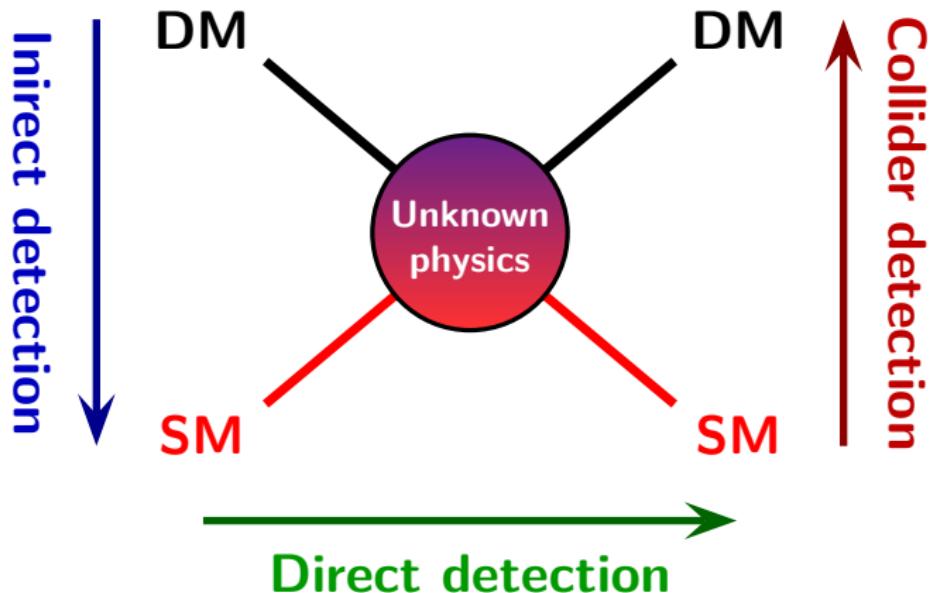


DM Particle Candidates



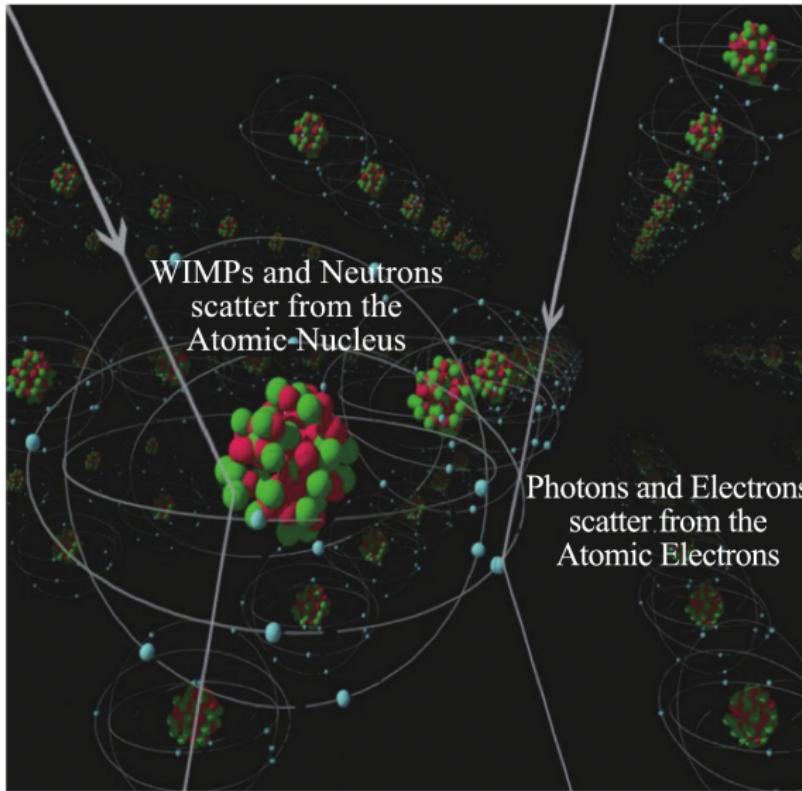
[Baer & Tata, 0805.1905]

Experimental Approaches to Dark Matter



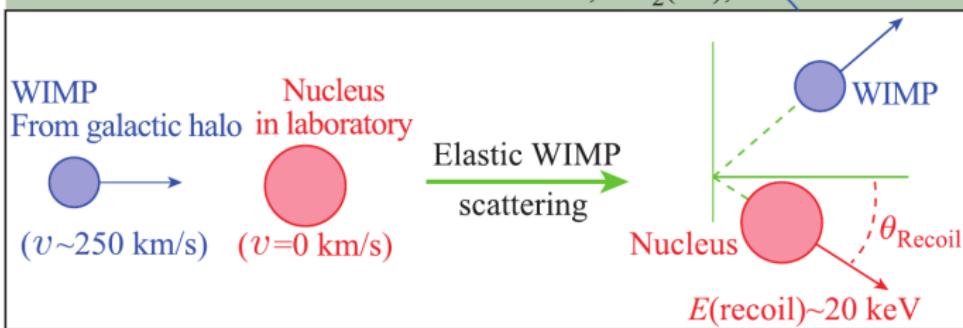
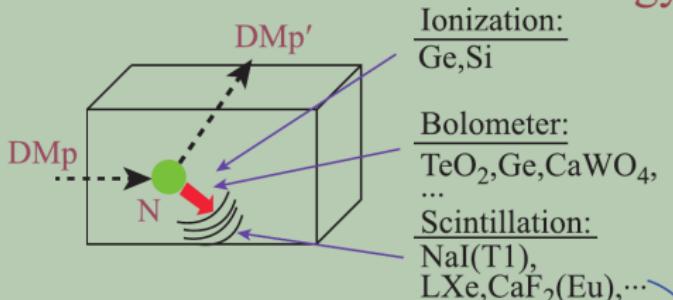


WIMP Scattering off Atomic Nuclei



Direct Detection

- Scatterings on nuclei
→ detection of nuclear recoil energy



[Bing-Lin Young, *Front. Phys.* **12**, 121201 (2017)]

DM Velocity Distribution

 During the collapse process which formed the Galaxy, the velocities of DM particles were “thermalized” by **fluctuations** in the gravitational potential, and DM particles have a **Maxwell-Boltzmann velocity distribution** in the **Galactic rest frame**:

$$\tilde{f}(\tilde{\mathbf{v}}) d^3 \tilde{\mathbf{v}} = \left(\frac{m_\chi}{2\pi k_B T} \right)^{3/2} \exp \left(-\frac{m_\chi \tilde{\mathbf{v}}^2}{2k_B T} \right) d^3 \tilde{\mathbf{v}} = \frac{e^{-\tilde{\mathbf{v}}^2/v_0^2}}{\pi^{3/2} v_0^3} d^3 \tilde{\mathbf{v}}, \quad v_0^2 \equiv \frac{2k_B T}{m_\chi}$$

$$\langle \tilde{\mathbf{v}} \rangle = \int \tilde{\mathbf{v}} \tilde{f}(\tilde{\mathbf{v}}) d^3 \tilde{v} = \mathbf{0}, \quad \langle \tilde{v}^2 \rangle = \int \tilde{v}^2 \tilde{f}(\tilde{\mathbf{v}}) d^3 \tilde{v} = \frac{3}{2} v_0^2$$

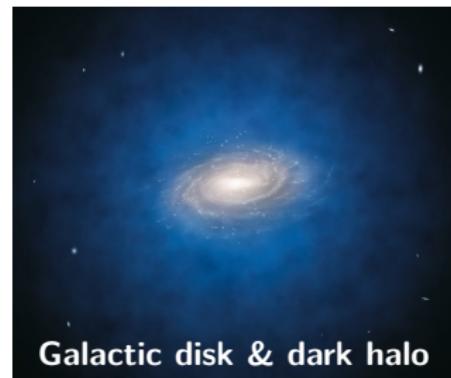
 **Speed distribution:** $\tilde{f}(\tilde{v}) d\tilde{v} = \frac{4\tilde{v}^2}{\sqrt{\pi}v_0^3} e^{-\tilde{v}^2/v_0^2} d\tilde{v}$

For an **isothermal halo**, the **local value** of v_0 equals to the **rotational speed of the Sun**.

$$v_0 = v_\odot \simeq 220 \text{ km/s}$$

[Binney & Tremaine, *Galactic Dynamics*, Chapter 4]

 **Velocity dispersion:** $\sqrt{\langle \tilde{v}^2 \rangle} = \sqrt{\frac{3}{2}} v_0 \simeq 270 \text{ km/s}$



[Credit: ESO/L. Calçada]

Earth Rest Frame

 The DM velocity distribution $f(v)$ seen by an **observer on the Earth** can be derived via **Galilean transformation**

$$\tilde{\mathbf{v}} = \mathbf{v} + \mathbf{v}_{\text{obs}}, \quad \mathbf{v}_{\text{obs}} = \mathbf{v}_{\odot} + \mathbf{v}_{\oplus}$$

 **Velocity distribution:** $f(\mathbf{v}) = \tilde{f}(\mathbf{v} + \mathbf{v}_{\text{obs}})$

Speed distribution:

$$f(v) dv = \frac{4v^2}{\sqrt{\pi}v_0^3} \exp\left(-\frac{v^2 + v_{\text{obs}}^2}{v_0^2}\right) \times \frac{\tilde{v}_0^2}{2vv_{\text{obs}}} \sinh\left(\frac{2vv_{\text{obs}}}{v_0^2}\right) dv$$

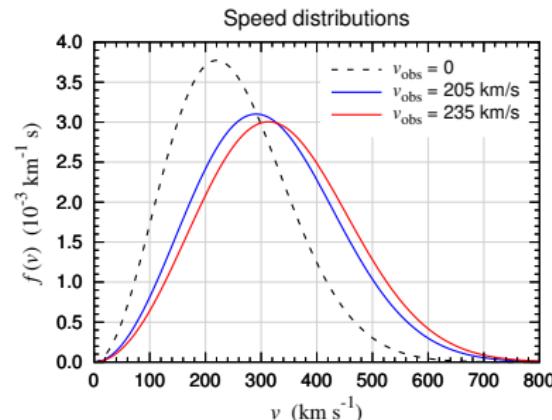
Since $v_{\oplus} \ll v_{\odot}$, we have ($\omega = 2\pi/\text{year}$)

$$v_{\text{obs}}(t) \simeq v_{\odot} + v_{\oplus} \sin \delta \cos[\omega(t - t_0)]$$

$$\simeq 220 \text{ km/s} + 15 \text{ km/s} \cdot \cos[\omega(t - t_0)]$$

👉 Annual modulation signal peaked on June 2

The diagram illustrates the Sun's motion through the WIMP wind. The Sun is shown as a yellow sphere at the center, with a dashed elliptical orbit. The Sun's velocity vector is indicated by a red arrow pointing towards the left, labeled $v_\odot \simeq 220 \text{ km/s}$. The angle between the Sun's velocity vector and the direction of the WIMP wind is labeled $\delta = 30.7^\circ$. The WIMP wind is represented by four purple arrows pointing towards the right. The Earth is shown as a green sphere at the top right of the Sun's orbit. The diagram also shows the Sun's position in the Cygnus constellation, with the label 'Cygnus' and an arrow pointing left. The Sun's motion is also labeled with the months 'June' and 'December' and their corresponding velocity components $v_\odot = 30 \text{ km/s}$.



[Freese *et al.*, PRD 37, 3388 (1988)]

Nuclear Recoil



Energy conservation:

$$\frac{1}{2}m_\chi v^2 = \frac{1}{2}m_\chi v_\chi^2 + \frac{1}{2}m_A v_R^2$$



Momentum conservation:

$$m_\chi v = m_\chi v_\chi \cos \theta_\chi + m_A v_R \cos \theta_R$$

$$m_\chi v_\chi \sin \theta_\chi = m_A v_R \sin \theta_R$$



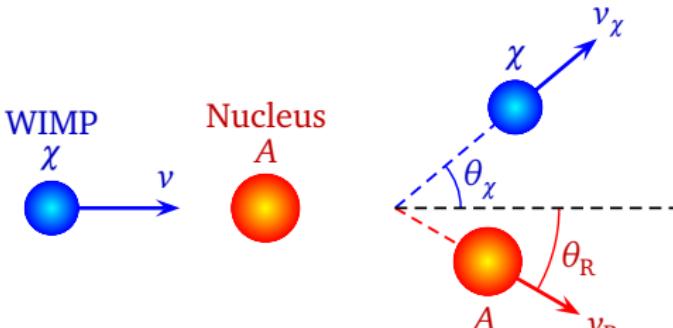
Recoil velocity $v_R = \frac{2m_\chi v \cos \theta_R}{m_\chi + m_A}$



Recoil momentum (momentum transfer) $q_R = m_A v_R = 2\mu_{\chi A} v \cos \theta_R$



Reduced mass of the χA system $\mu_{\chi A} \equiv \frac{m_\chi m_A}{m_\chi + m_A} = \begin{cases} m_A, & \text{for } m_\chi \gg m_A \\ \frac{m_\chi}{2}, & \text{for } m_\chi = m_A \\ m_\chi, & \text{for } m_\chi \ll m_A \end{cases}$



Forward scattering ($\theta_R = 0$) **maximal momentum transfer** $q_R^{\max} = 2\mu_{\chi A} v$

Nuclear Recoil



Energy conservation:

$$\frac{1}{2}m_\chi v^2 = \frac{1}{2}m_\chi v_\chi^2 + \frac{1}{2}m_A v_R^2$$



Momentum conservation:

$$m_\chi v = m_\chi v_\chi \cos \theta_\chi + m_A v_R \cos \theta_R$$

$$m_\chi v_\chi \sin \theta_\chi = m_A v_R \sin \theta_R$$



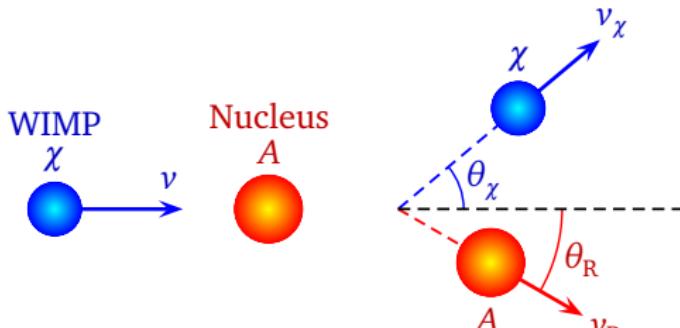
$$\text{Recoil velocity } v_R = \frac{2m_\chi v \cos \theta_R}{m_\chi + m_A}$$



$$\text{Recoil momentum (momentum transfer) } q_R = m_A v_R = 2\mu_{\chi A} v \cos \theta_R$$



$$\text{Kinetic energy of the recoiled nucleus } E_R = \frac{q_R^2}{2m_A} = \frac{2\mu_{\chi A}^2 v^2 \cos^2 \theta_R}{m_A}$$



As $v \sim 10^{-3}c$, for $m_\chi = m_A \simeq 100$ GeV and $\theta_R = 0$,

$$q_R = m_\chi v \sim 100 \text{ MeV}, \quad E_R = \frac{1}{2}m_\chi v^2 \sim 50 \text{ keV}$$

Event Rate



Event rate per unit time per unit energy interval in direct detection experiments:

$$\frac{dR}{dE_R} = N_A \frac{\rho_{\oplus}}{m_X} \int_{v_{\min}}^{v_{\max}} d^3v \ f(v) v \frac{d\sigma_{XA}}{dE_R}$$



Astrophysics factors



Particle physics factors



Detector factors



N_A : target nucleus number



ρ_{\oplus} $\simeq 0.4 \text{ GeV/cm}^3$: DM **mass density** around the Earth



ρ_{\oplus}/m_X is the DM particle **number density** around the Earth



σ_{XA} : DM-nucleus **scattering cross section**



Minimal velocity $v_{\min} = \left(\frac{m_A E_R^{\text{th}}}{2\mu_{XA}^2} \right)^{1/2}$: determined by the **detector threshold** of nuclear recoil energy, E_R^{th}



Maximal velocity v_{\max} : determined by the DM **escape velocity** v_{esc}



$v_{\text{esc}} \simeq 544 \text{ km/s}$ [Smith *et al.*, MNRAS 379, 755]

Cross Section Dependence on Nucleus Spin

There are two kinds of DM-nucleus scattering

Spin-independent (SI) cross section $\sigma_{\chi A}^{\text{SI}} \propto \mu_{\chi A}^2 [ZG_p + (A - Z)G_n]^2$

Spin-dependent (SD) cross section $\sigma_{\chi A}^{\text{SD}} \propto \mu_{\chi A}^2 \frac{J_A + 1}{J_A} (S_p^A G'_p + S_n^A G'_n)^2$

Nucleus properties: mass number A , atomic number Z , spin J_A , expectation value of the proton (neutron) spin content in the nucleus S_p^A (S_n^A)

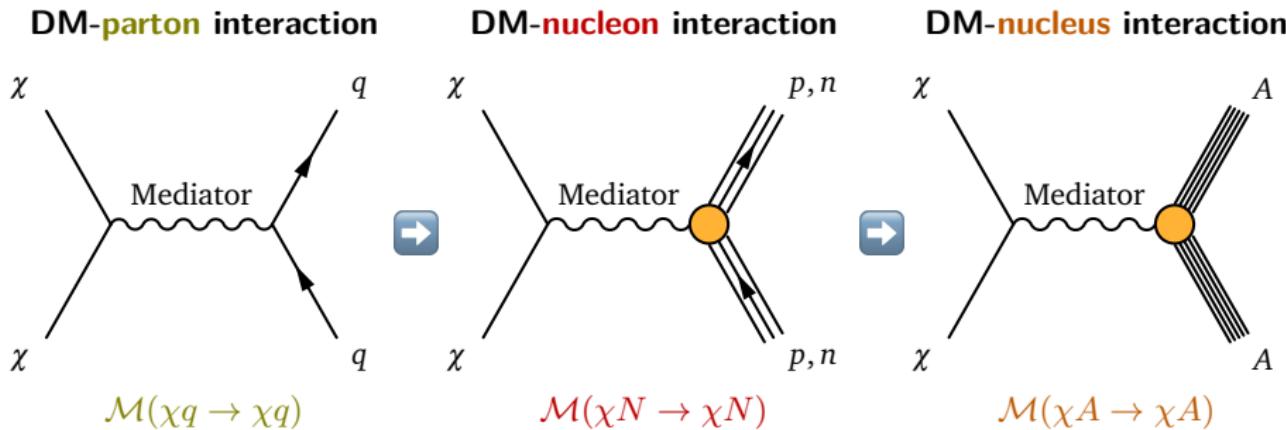
 $G_p^{(\prime)}$ and $G_n^{(\prime)}$: DM effective couplings to the proton and the neutron

- $Z \simeq A/2$  $\sigma_{\chi A}^{\text{SI}} \propto A^2 [(G_p + G_n)/2]^2$
 - ! Strong coherent enhancement for heavy nuclei

- Spins of nucleons tend to cancel out among themselves

- $S_N^A \simeq 1/2$ ($N = p$ or n) for a nucleus with an odd number of N
- $S_N^A \simeq 0$ for a nucleus with an even number of N

Three Levels of Interaction



🚌 As a variety of **target nuclei** are used in **direct detection experiments**, results are usually compared with each other at the **DM-nucleon level**

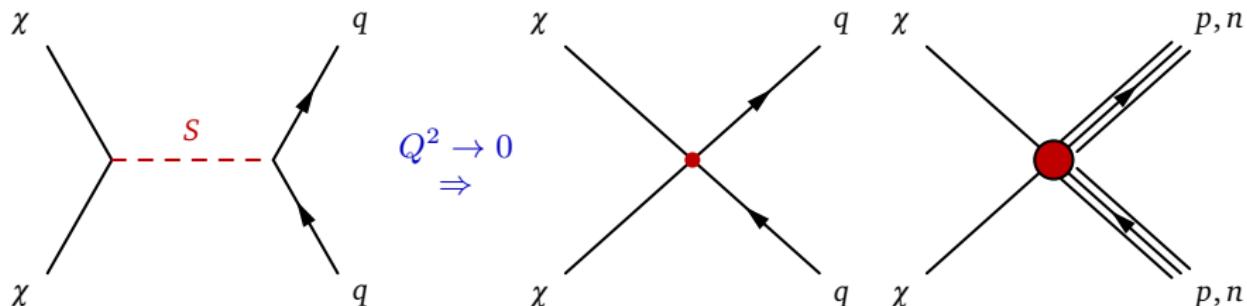
🚌 The **DM-nucleon level** is related to the **DM-parton level** via **form factors**, which describe the **probabilities** of finding **partons** inside **nucleons**

🚐 The **relevant partons** involve not only **valence quarks**, but also **sea quarks** and **gluons**

Zero-Momentum Transfer Limit

⚠️ As the **momentum transfer** is typically **much smaller** than the **underlying energy scale** (e.g., mediator mass), the **zero-momentum transfer limit** is a good approximation for calculation

🚲 In this limit, the mediator field can be integrated out, and the interaction can be described by **effective operators** in **effective field theory**



⚠️ **Scalar mediator propagator:** $\frac{i}{Q^2 - m_S^2} \Rightarrow -\frac{i}{m_S^2}$

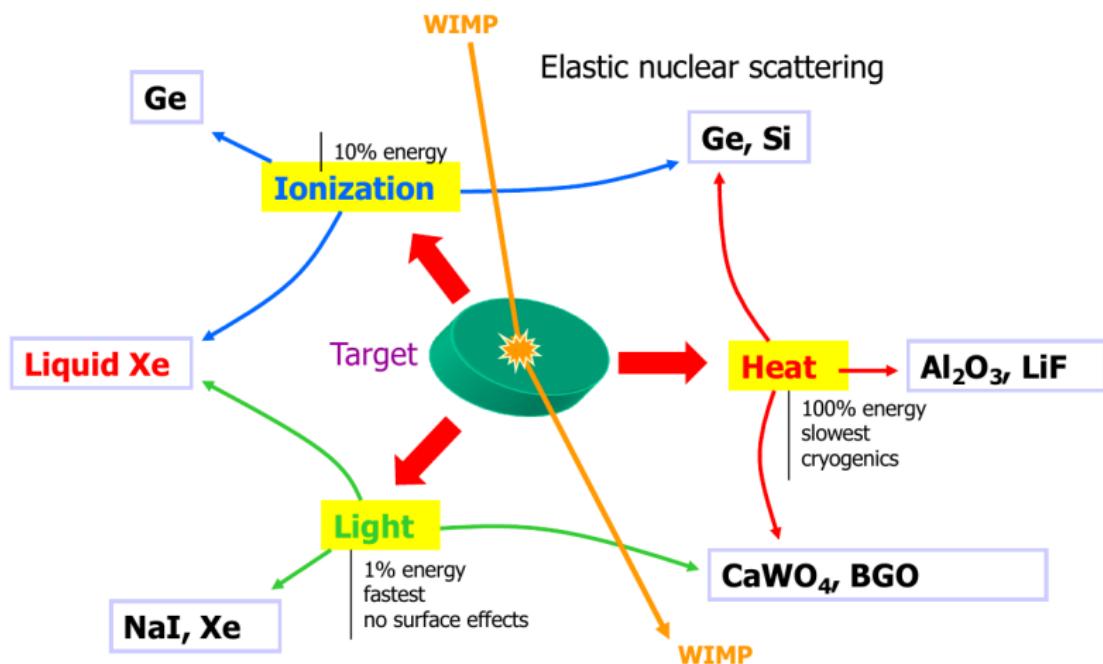
⚠️ **Lagrangian:** $\mathcal{L}_{\text{int}} = g_\chi S \bar{\chi} \chi + g_q S \bar{q} q \Rightarrow \mathcal{L}_{\text{eff}} = G_{\text{eff}} \bar{\chi} \chi \bar{q} q, \quad G_{\text{eff}} = \frac{g_\chi g_q}{m_S^2}$

Effective Operators for DM-quark Interactions

Spin-1/2 DM		Spin-0 DM
SI	$\bar{\chi}\chi\bar{q}q, \bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q$	$\chi^*\chi\bar{q}q, (\chi^*\bar{i}\overleftrightarrow{\partial}^\mu\chi)\bar{q}\gamma_\mu q$
SD	$\bar{\chi}\gamma^\mu\gamma_5\chi\bar{q}\gamma_\mu\gamma_5 q, \bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu} q$	
$\sigma_{\chi N}$ $\propto Q^2 $	$\bar{\chi}\bar{i}\gamma_5\chi\bar{q}\bar{i}\gamma_5 q, \bar{\chi}\chi\bar{q}\bar{i}\gamma_5 q$ $\bar{\chi}\bar{i}\gamma_5\chi\bar{q}q, \bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu\gamma_5 q$ $\bar{\chi}\gamma^\mu\gamma_5\chi\bar{q}\gamma_\mu q, \varepsilon^{\mu\nu\rho\sigma}\bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\rho\sigma} q$	$\chi^*\chi\bar{q}\bar{i}\gamma_5 q$ $(\chi^*\bar{i}\overleftrightarrow{\partial}^\mu\chi)\bar{q}\gamma_\mu\gamma_5 q$
Spin-3/2 DM		Spin-1 DM
SI	$\bar{\chi}^\mu\chi_\mu\bar{q}q, \bar{\chi}^\nu\gamma^\mu\chi_\nu\bar{q}\gamma_\mu q$	$\chi_\mu^*\chi^\mu\bar{q}q, (\chi_\nu^*\bar{i}\overleftrightarrow{\partial}^\mu\chi^\nu)\bar{q}\gamma_\mu q$
SD	$\bar{\chi}^\nu\gamma^\mu\gamma_5\chi_\nu\bar{q}\gamma_\mu\gamma_5 q, \bar{\chi}^\rho\sigma^{\mu\nu}\chi_\rho\bar{q}\sigma_{\mu\nu} q$ $\bar{i}(\bar{\chi}^\mu\chi^\nu - \bar{\chi}^\nu\chi^\mu)\bar{q}\sigma_{\mu\nu} q$	$\bar{i}(\chi_\mu^*\chi_\nu - \chi_\nu^*\chi_\mu)\bar{q}\sigma^{\mu\nu} q$ $\varepsilon^{\mu\nu\rho\sigma}(\chi_\mu^*\bar{\partial}_\nu\chi_\rho)\bar{q}\gamma_\sigma\gamma_5 q$
$\sigma_{\chi N}$ $\propto Q^2 $	$\bar{\chi}^\mu\bar{i}\gamma_5\chi_\mu\bar{q}\bar{i}\gamma_5 q, \bar{\chi}^\mu\chi_\mu\bar{q}\bar{i}\gamma_5 q$ $\bar{\chi}^\mu\bar{i}\gamma_5\chi_\mu\bar{q}q, \bar{\chi}^\nu\gamma^\mu\chi_\nu\bar{q}\gamma_\mu\gamma_5 q$ $\bar{\chi}^\mu\gamma^\mu\gamma_5\chi_\nu\bar{q}\gamma_\mu q, \varepsilon^{\mu\nu\rho\sigma}\bar{i}(\bar{\chi}_\mu\chi_\nu - \bar{\chi}_\nu\chi_\mu)\bar{q}\sigma_{\rho\sigma} q$ $\varepsilon^{\mu\nu\rho\sigma}\bar{\chi}^\alpha\sigma_{\mu\nu}\chi_\alpha\bar{q}\sigma_{\rho\sigma} q, (\bar{\chi}^\mu\gamma_5\chi^\nu - \bar{\chi}^\nu\gamma_5\chi^\mu)\bar{q}\sigma_{\mu\nu} q$ $\varepsilon^{\mu\nu\rho\sigma}(\bar{\chi}_\mu\gamma_5\chi_\nu - \bar{\chi}_\nu\gamma_5\chi_\mu)\bar{q}\sigma_{\rho\sigma} q$	$\chi_\mu^*\chi^\mu\bar{q}\bar{i}\gamma_5 q$ $(\chi_\nu^*\bar{i}\overleftrightarrow{\partial}^\mu\chi^\nu)\bar{q}\gamma_\mu\gamma_5 q$ $\varepsilon^{\mu\nu\rho\sigma}(\chi_\mu^*\bar{\partial}_\nu\chi_\rho)\bar{q}\gamma_\sigma q$ $\varepsilon^{\mu\nu\rho\sigma}\bar{i}(\chi_\mu^*\chi_\nu - \chi_\nu^*\chi_\mu)\bar{q}\sigma_{\rho\sigma} q$

[Zheng, **ZHY**, Shao, Bi, Li, Zhang, arXiv:1012.2022, NPB;
ZHY, Zheng, Bi, Li, Yao, Zhang, arXiv:1112.6052, NPB]

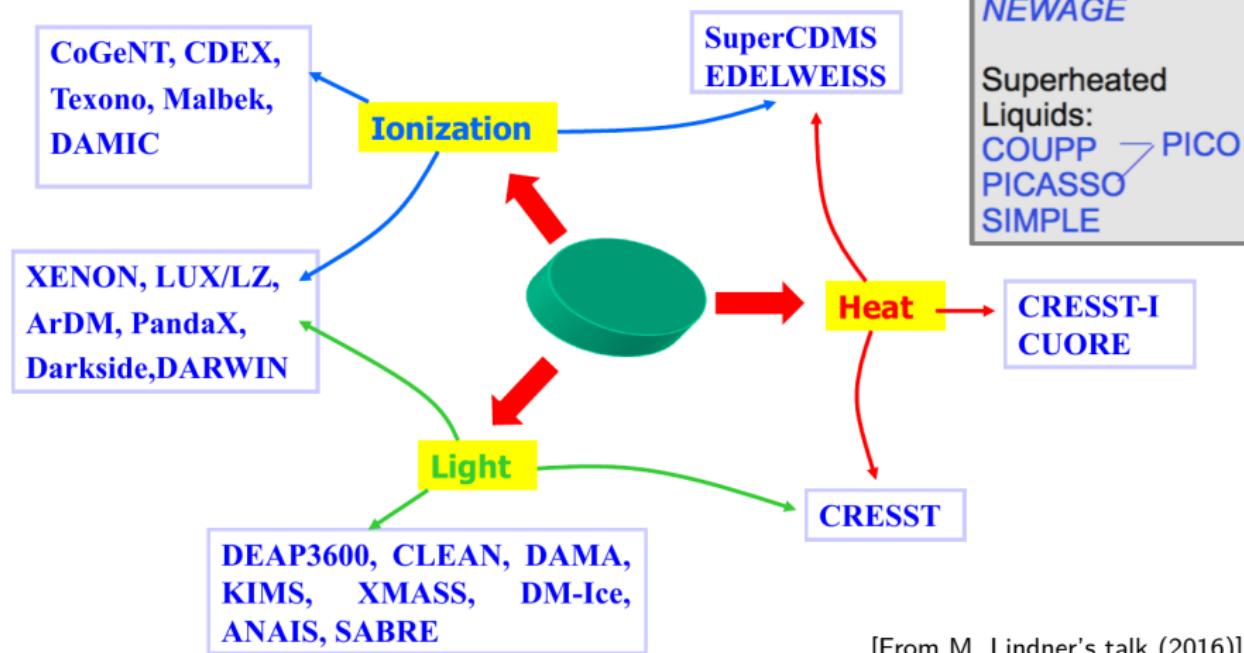
Technologies and Detector Material



[From M. Lindner's talk (2016)]

Technologies and Detector Material

Detection methods: Crystals (NaI, Ge, Si),
Cryogenic Detectors, Liquid Noble Gases



[From M. Lindner's talk (2016)]

Example: Dual-phase Xenon Time Projection Chamber

↑ **Upper: Xenon gas**

↓ **Lower: Liquid Xenon**

UV scintillation photons recorded by photomultiplier tube (PMT) arrays on top and bottom

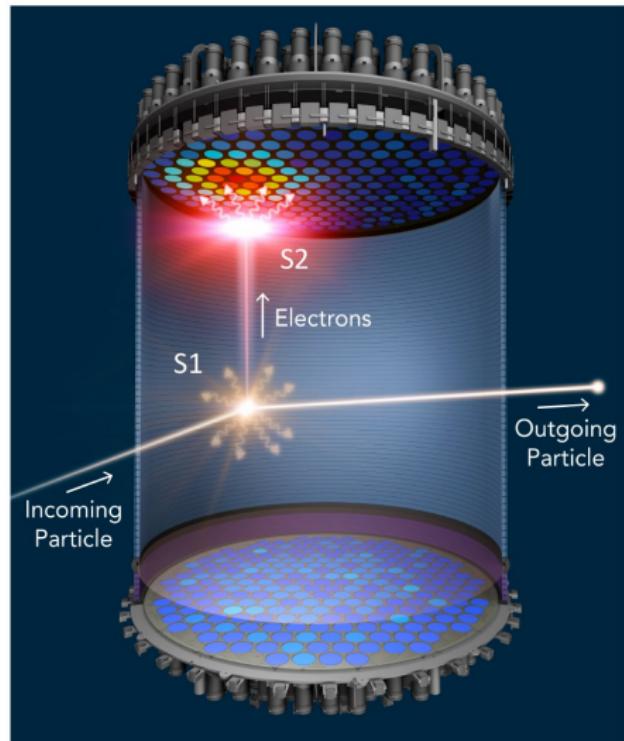
1 Primary scintillation (S1)

Scintillation light **promptly** emitted from the interaction vertex

2 Secondary scintillation (S2)

Ionization electrons emitted from the interaction are **drifted** to the **surface** and into the **gas**, where they emit **proportional scintillation light**

Experiment: **PandaX, XENON, LZ**

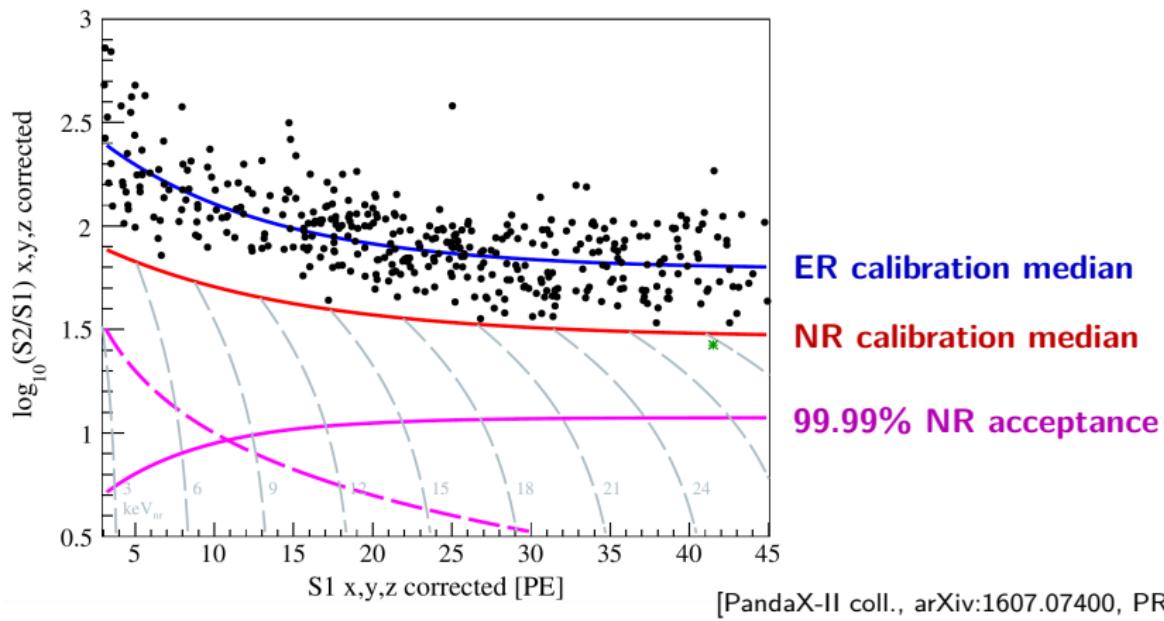


[From A. Cottle's talk (2017)]

PandaX-II Real Data: S1 versus S2

⚠ S1 and S2 are characterized by numbers of **photoelectrons** (PEs) in PMTs

⚠ The γ **background**, which produces **electron recoil (ER)** events, can be distinguished from **nuclear recoil (NR)** events using the **S2-to-S1 ratio**



Backgrounds



Background suppression:

- ## Deep underground Shielded environments



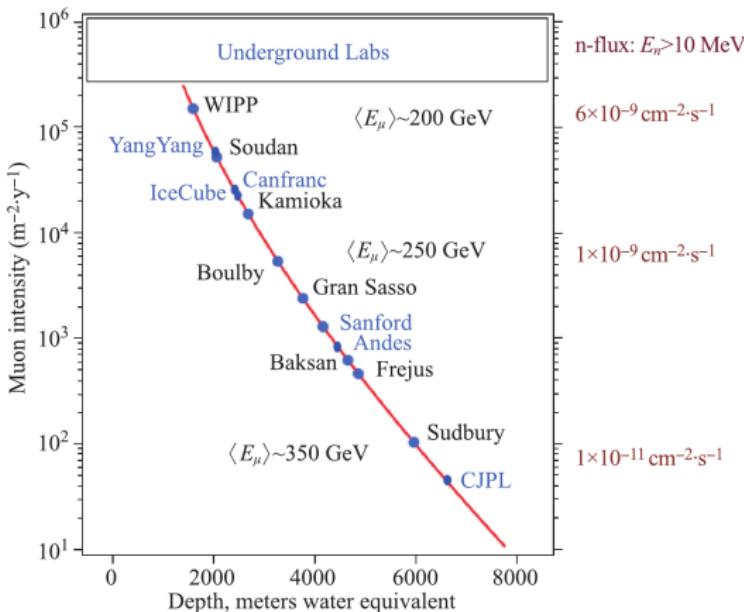
Cosmogenic backgrounds

- **Cosmic rays** and secondary reactions
 - **Activation products** in shields and detectors



Radiogenic backgrounds

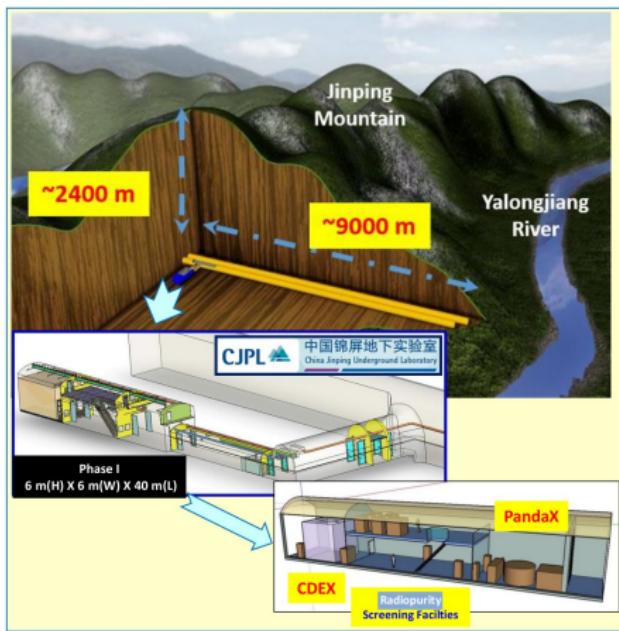
- **External natural** radioactivity:
walls, structures of site, radon
 - **Internal** radioactivity:
shield and construction materials
naturally occurring radio-isotopes



[From P. Cushman's talk (2014)]



China JinPing Underground Laboratory (CJPL)



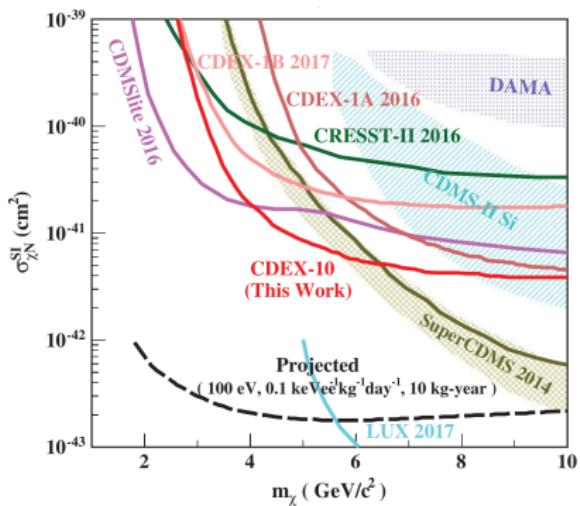
[Yue et al., arXiv:1602.02462]



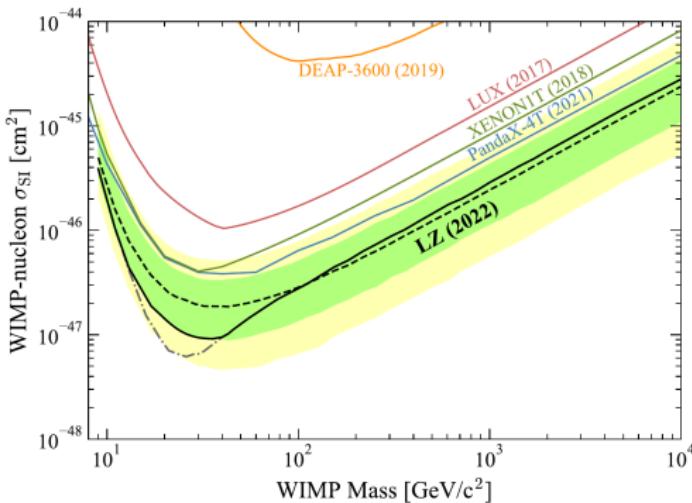
Experiments: CDEX, PandaX

Exclusion Limits for SI Scattering

- For isospin-conserving SI scattering, protons and neutrons can be treated as the same species, i.e., “nucleons”
 - Lower thresholds and lighter targets are needed to probe the low mass regime
 - It requires more exposure, heavier targets, and fewer backgrounds to explore the high mass regime



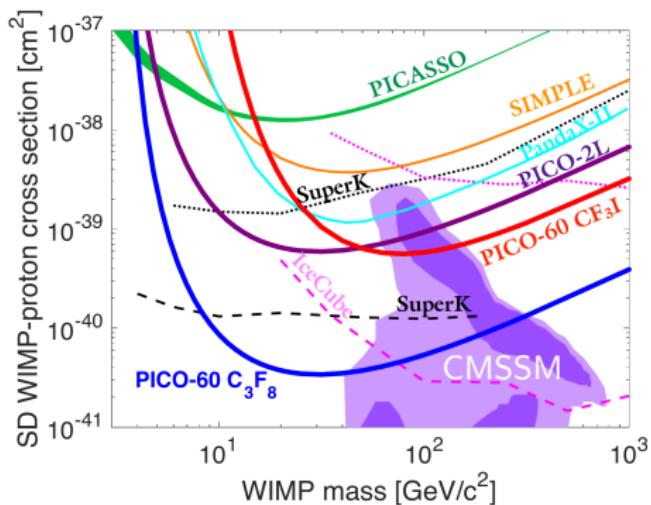
[CDEX Coll., arXiv:1802.09016, PRL]



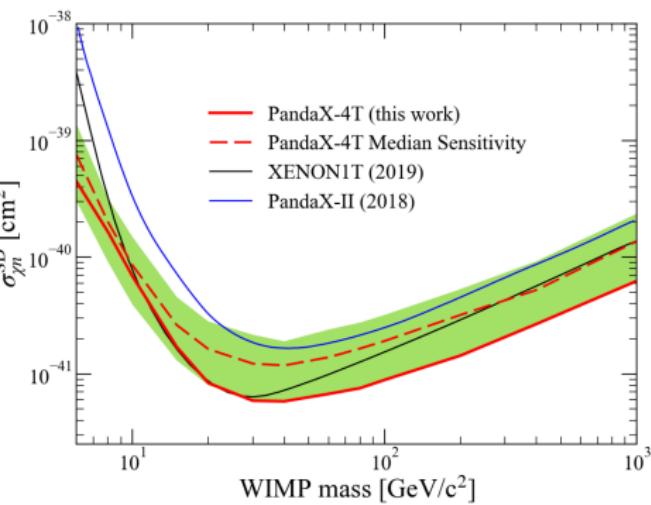
[LZ Coll., arXiv:2207.03764, PRL]

Exclusion Limits for SD Scattering

- For **SD scattering**, specific detection material usually has **very different** sensitivities to **DM-proton** and **DM-neutron** cross sections
 - As there is no coherent enhancement for **SD scattering**, the sensitivity is **lower** than the **SI case** by **several orders of magnitude**



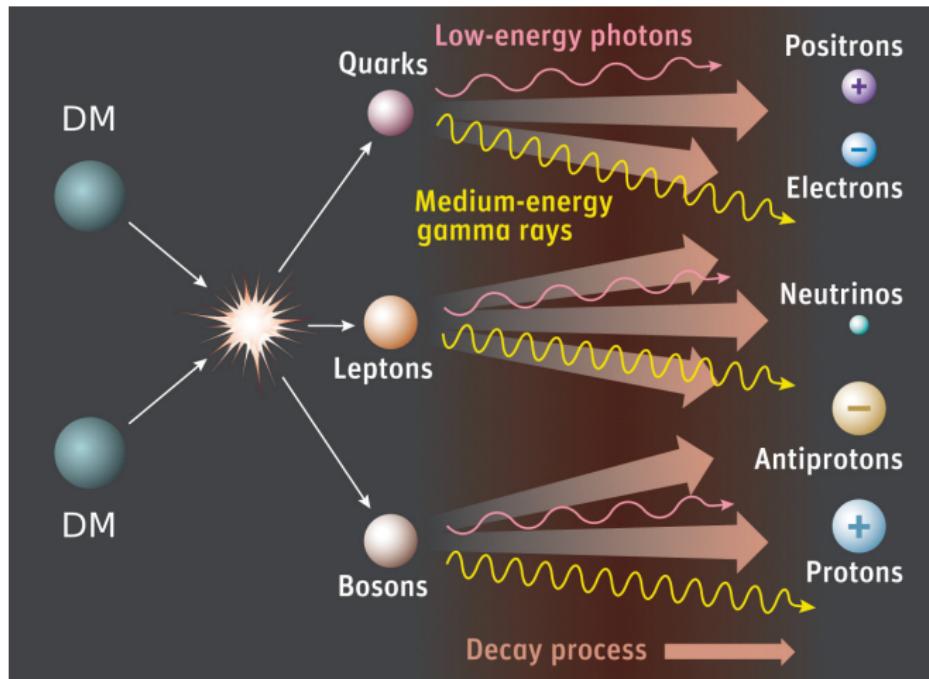
[PICO coll., arXiv:1702.07666, PRL]



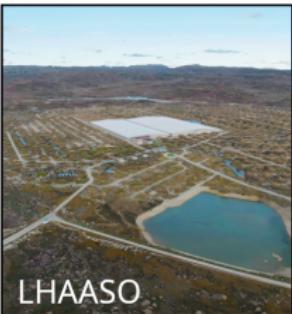
[PandaX-4T coll., arXiv:2208.03626, PLB]

Indirect Detection

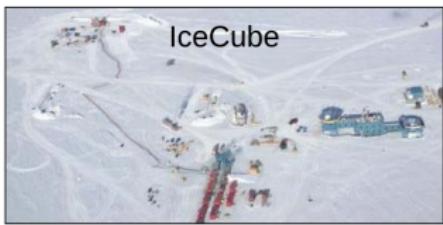
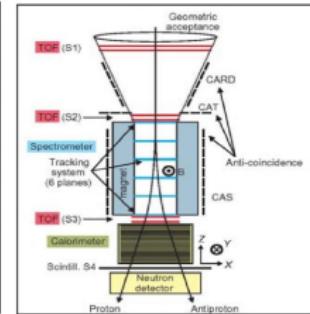
 **Indirect detection** looks for **stable products**, e.g., γ rays, cosmic rays, and neutrinos, from **DM annihilation** or **decay** (if DM is not totally stable) **in space**



Indirect Detection Experiments



PAMELA



Dark Matter Source Function



 **Particle number** per unit time per unit volume per unit energy interval of a **stable species** (γ , e^\pm , ν , p , \bar{p} , \dots) produced from **DM annihilation** or **decay**:

$$Q_{\text{ann}}(\mathbf{x}, E) = \frac{\langle \sigma_{\text{ann}} v \rangle_{\text{tot}}}{2m_\chi^2} \rho^2(\mathbf{x}) \sum_i F_i \left(\frac{dN}{dE} \right)_i \text{(annihilation)}$$

Astrophysics factors

$$Q_{\text{dec}}(\mathbf{x}, E) = \frac{1}{\tau_\chi m_\chi} \rho(\mathbf{x}) \sum_i B_i \left(\frac{dN}{dE} \right)_i \text{ (decay)}$$

Particle physics factors



$\rho(x)$: DM mass density at the source position x



$(dN/dE)_i$: **number per unit energy interval** from a single event in the **channel *i***



$\langle \sigma_{\text{ann}} v \rangle_{\text{tot}}$: thermal average of the total **annihilation cross section** multiplied by the relative velocity between the two incoming DM particles



$F_i \equiv \langle \sigma_{\text{ann}} v \rangle_i / \langle \sigma_{\text{ann}} v \rangle_{\text{tot}}$: **branching fraction** of the annihilation channel i



$\tau_x \equiv 1/\Gamma_x$: mean **lifetime** of the DM particle



$B_i \equiv \Gamma_i / \Gamma_\chi$: **branching ratio** of the decay channel i

γ rays from DM: Continuous Spectrum

DM annihilation or decay into e^+e^- , $\mu^+\mu^-$, $\tau^+\tau^-$, $q\bar{q}$, W^+W^- , Z^0Z^0 , h^0h^0



γ -ray emissions from final state radiation or particle decays



Cutoff energy:

- m_χ for DM annihilation
- $m_\chi/2$ for DM decay

More promising to look at

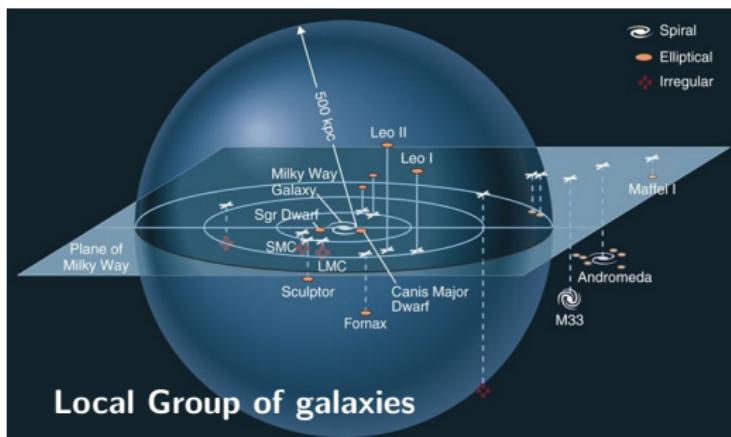
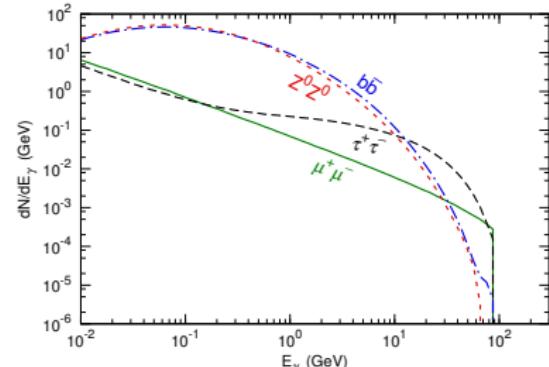
DM-dominated regions:

Galactic Center

Galactic halo

dwarf galaxies

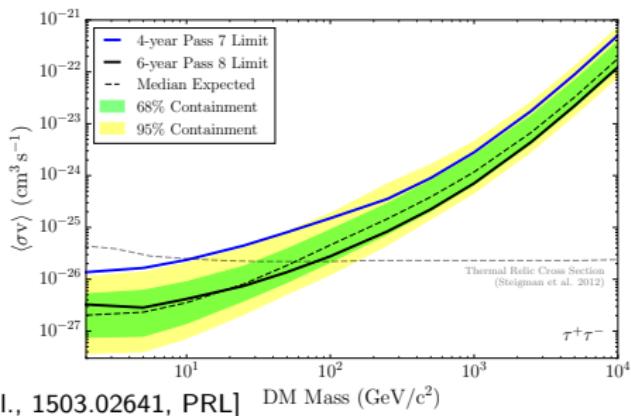
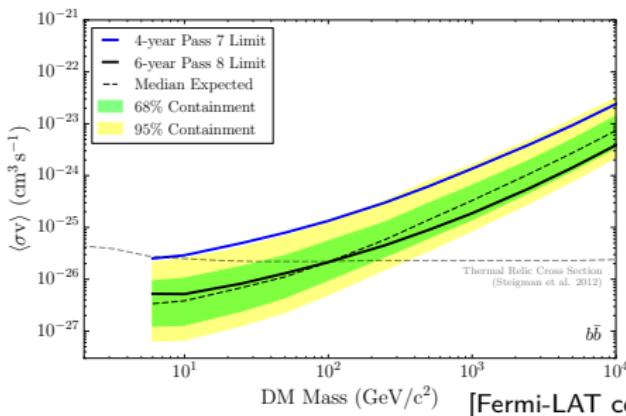
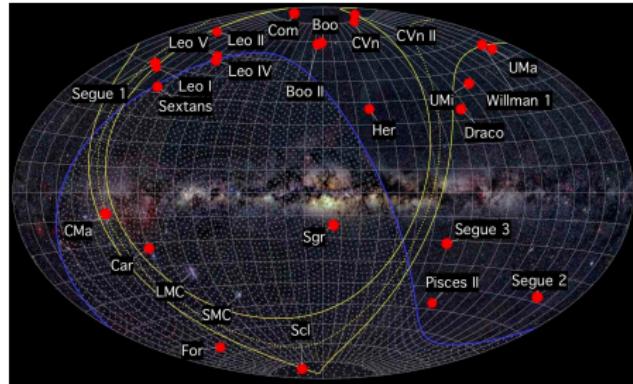
clusters of galaxies



γ -ray Observation of Dwarf Galaxies

The space experiment **Fermi-LAT** searched for **γ -ray emissions** from **dwarf spheroidal satellite galaxies** of the **Milky Way** and found **no significant signal**

Based on the 6-year data, **upper limits** on the **DM annihilation cross section** are given



GeV Excess at the Galactic Center?

💡 Since 2009, several research groups reported an **excess** of **continuous spectrum γ -rays** in the **Fermi-LAT** data after **subtracting** well-known **astrophysical backgrounds**,

🕯 It locates in the **Galactic Center (GC)** region and peaks at a few GeV

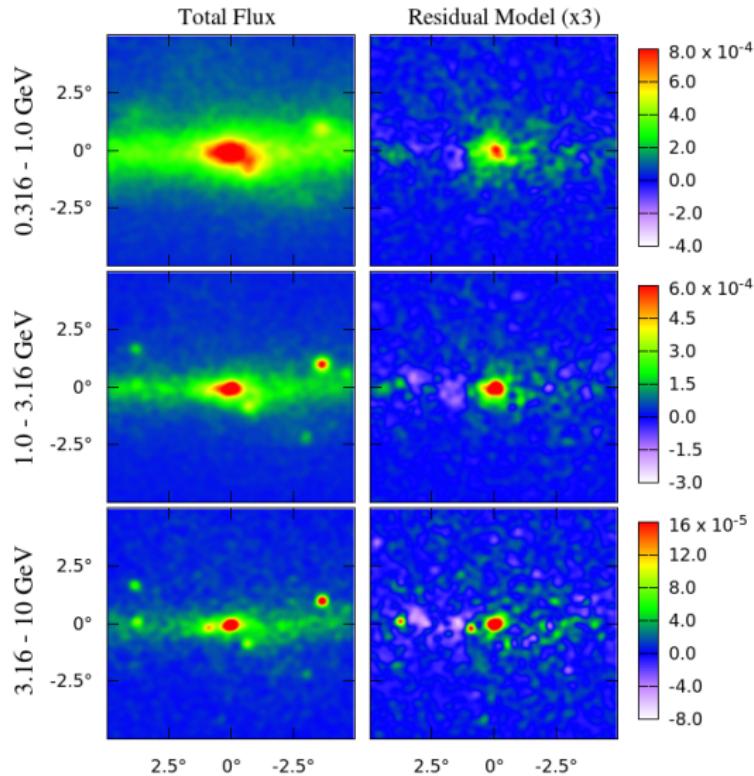


Left: raw γ -ray maps



Right: residual maps

subtracting the Galactic diffuse model, the 20 cm template, point sources, and the isotropic template



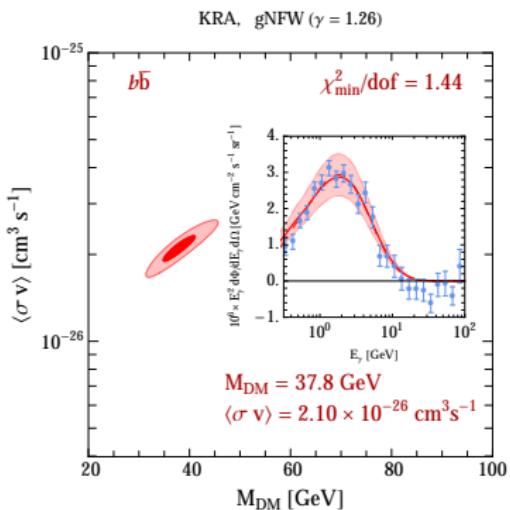
[Daylan et al., 1402.6703, PDU]

Interpretation with Dark Matter Annihilation

DM annihilation into $b\bar{b}$

$$m_\chi \simeq 30 - 40 \text{ GeV}$$

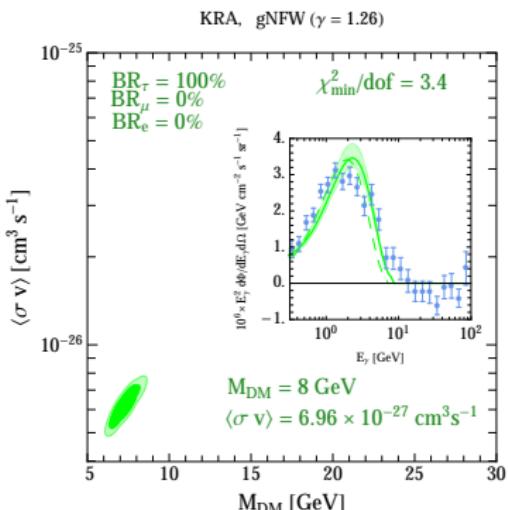
$$\langle \sigma_{\text{ann}} v \rangle \sim 10^{-26} \text{ cm}^3 \text{ s}^{-1}$$



DM annihilation into $\tau^+\tau^-$

$$m_\chi \sim 9 \text{ GeV}$$

$$\langle \sigma_{\text{ann}} v \rangle \sim 5 \times 10^{-27} \text{ cm}^3 \text{ s}^{-1}$$



[Cirelli *et al.*, arXiv:1407.2173, JCAP]



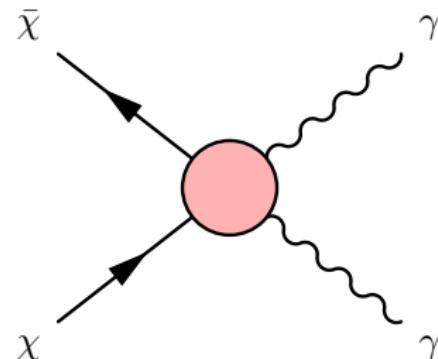
γ rays from DM: Line Spectrum

⌚ DM particles should **not have electric charge** and thus **not directly couple to photons**

👓 Nonetheless, DM particles **may couple to photons** via **high-order loop diagrams**



🚧 **Highly suppressed:** the $\chi\bar{\chi} \rightarrow \gamma\gamma$ branching fraction may be only $\sim 10^{-4} - 10^{-1}$



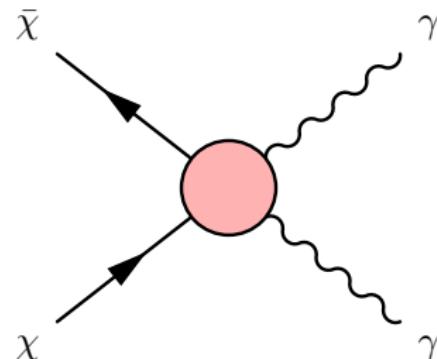
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👢 Highly suppressed: the $\chi\bar{\chi} \rightarrow \gamma\gamma$ branching fraction may be only $\sim 10^{-4} - 10^{-1}$



👢 For **nonrelativistic** DM particles in space, the photons produced in $\chi\bar{\chi} \rightarrow \gamma\gamma$ would be **mono-energetic**



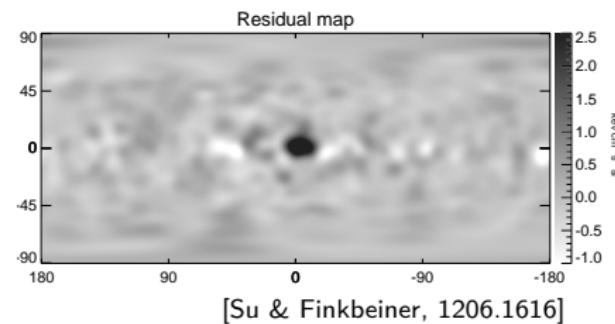
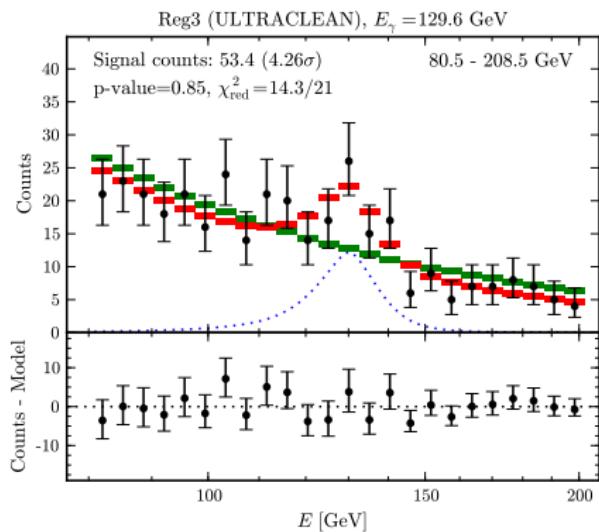
👠 A γ -ray line at energy $\sim m_\chi$
(“smoking gun” for DM particles)



A γ -ray Line Signal at the Galactic Center?

Using the **3.7-year Fermi-LAT γ -ray data**, several analyses showed that there might be evidence of a **monochromatic γ -ray line** at **energy ~ 130 GeV**, originating from the **Galactic Center region** (about $3-4\sigma$)

It may be explained by **DM annihilation** with $\langle \sigma_{\text{ann}} v \rangle \sim 10^{-27} \text{ cm}^3 \text{ s}^{-1}$



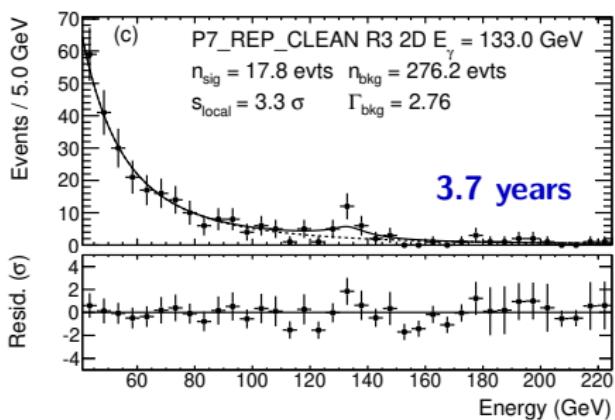
[Weniger, 1204.2797, JCAP]

Fermi-LAT Official Results: Not Confirmed with More Data



3.7-year data

⚠ The most significant fit occurred at $E_\gamma = 133$ GeV and had a **local significance** of 3.3σ , translating to a global significance of 1.6σ



[Fermi-LAT Coll., 1305.5597, PRD]

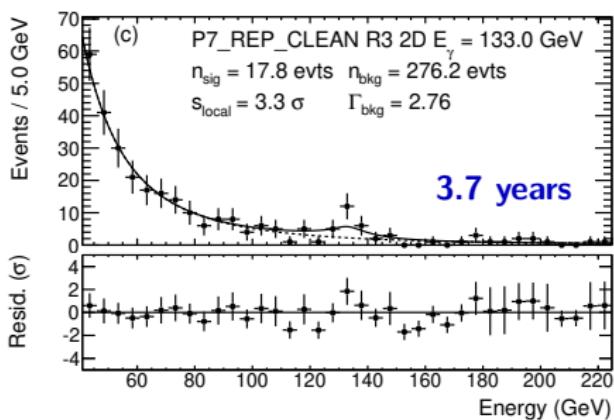
Fermi-LAT Official Results: Not Confirmed with More Data

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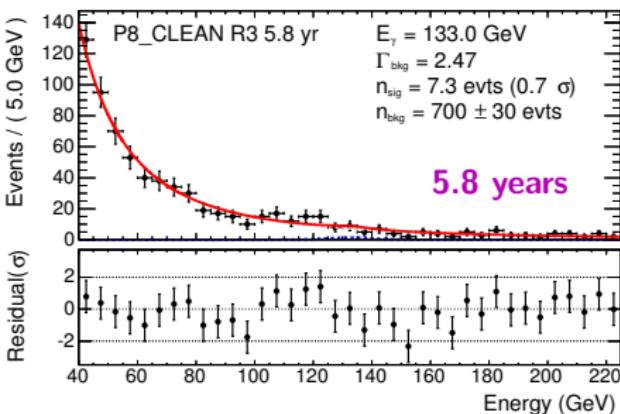
⚠ The most significant fit occurred at $E_\gamma = 133$ GeV and had a **local significance** of 3.3σ , translating to a global significance of 1.6σ

● 5.8-year data

🚫 The **local significance** has dropped to 0.72σ



[Fermi-LAT Coll., 1305.5597, PRD]



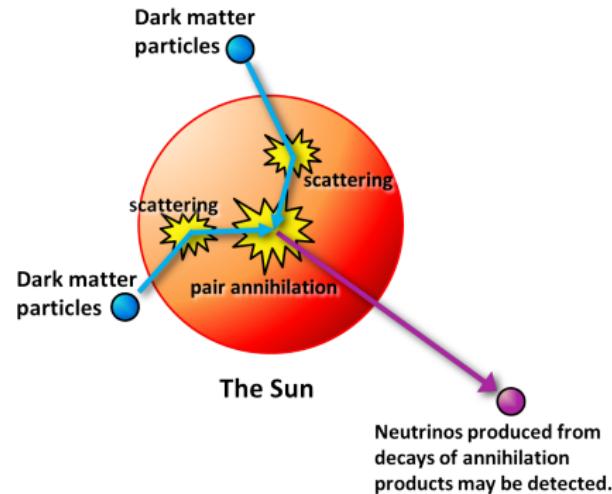
[Fermi-LAT Coll., 1506.00013, PRD]

Neutrinos from DM

🕸 Dark matter may be **captured** and **accumulated** at the **core** of the **Sun** ☀ (or the **Earth** 🌎), producing **high energy neutrinos** that could **freely go out**

🐜 Change Rate of the **number** of **DM particles** in the **Sun**:

$$\frac{dN_\chi}{dt} = C_\odot(\sigma_{\chi H}, \sigma_{\chi He}) - A_\odot(\sigma_{\text{ann}})N_\chi^2$$

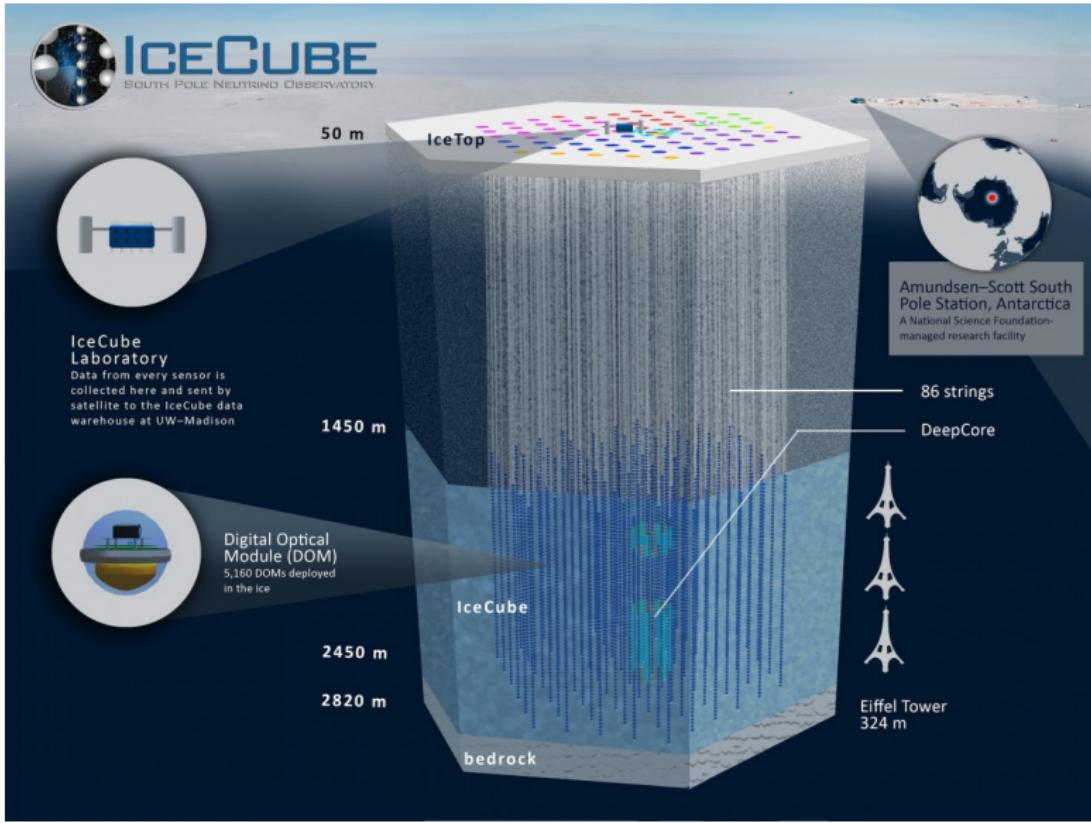


🐛 Capture rate C_\odot depends on **DM scattering** on Hydrogen and Helium

🐞 Annihilation rate $A_\odot = \langle \sigma_{\text{ann}} v \rangle / V_{\text{eff}}$ depends on **DM annihilation** as well as the effective volume of the solar core

🐌 The **age** of the **Sun** is **long enough** (~ 4.6 billion years) to make the **capture** and **annihilation** processes reach **equilibrium**: $dN_\chi/dt = 0$

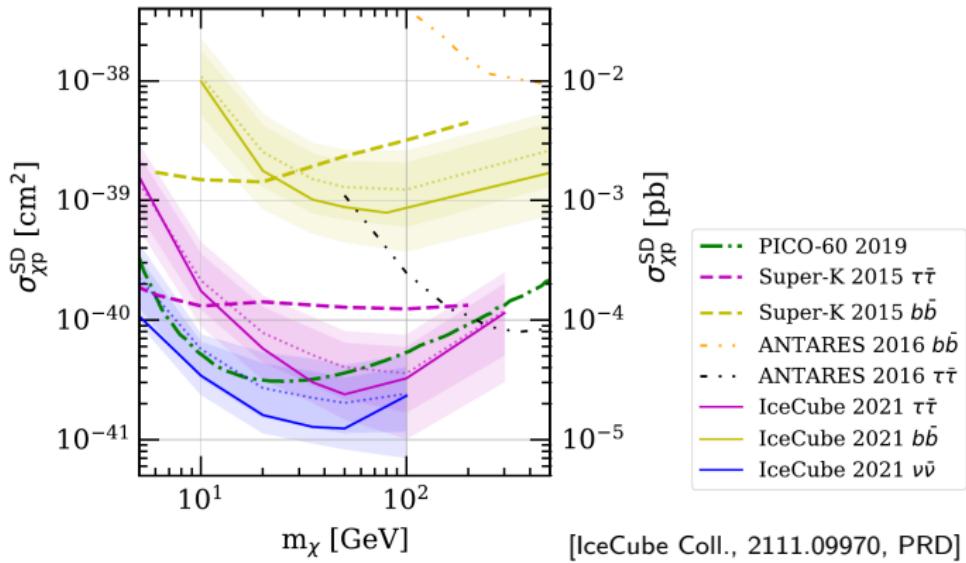
IceCube: South Pole Neutrino Observatory



Searches for Neutrinos from DM Annihilation within the Sun

Using the **7-year IceCube data**, **no significant detection** of **neutrinos** with energies < 500 GeV from **DM annihilation** at the **solar core** is found

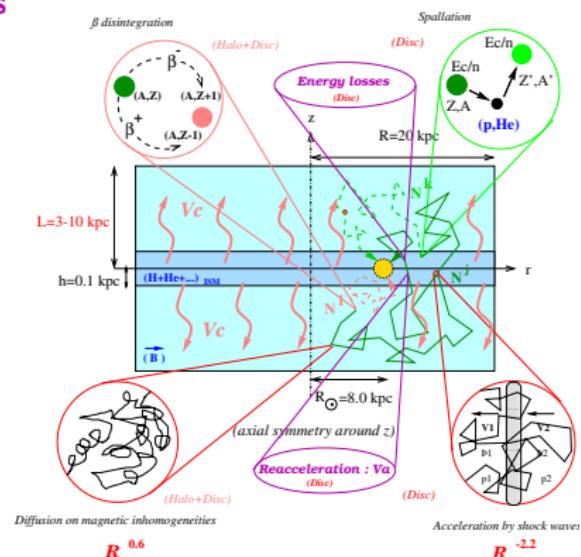
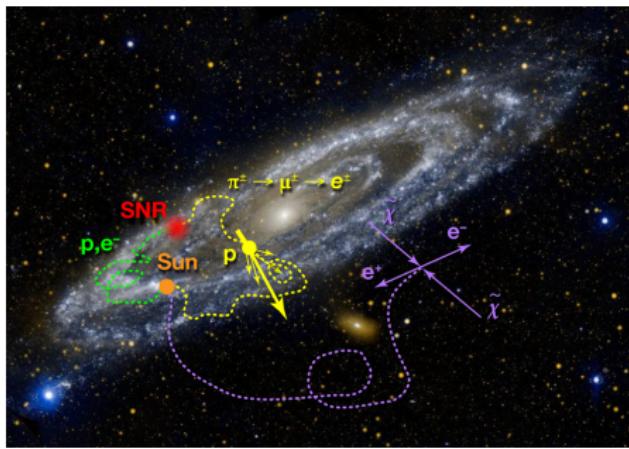
Assuming **equilibrium** in the **capture** and **annihilation** processes, the constraints can be converted to those on the **DM scattering cross section**



Cosmic Rays from DM

After produced in sources, **Galactic cosmic rays** diffuse in the interstellar space, suffering from several **propagation effects** before they arrive at the **Earth**: **diffusion**, energy losses, **convection**, **reacceleration**, **spallation**, ...

Unlike γ rays and neutrinos, cosmic rays typically **do not contain direction information** of their **sources**

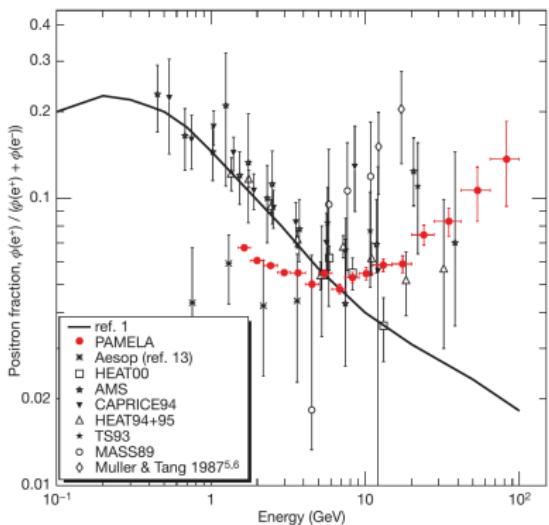


[Maurin et al., astro-ph/0212111]

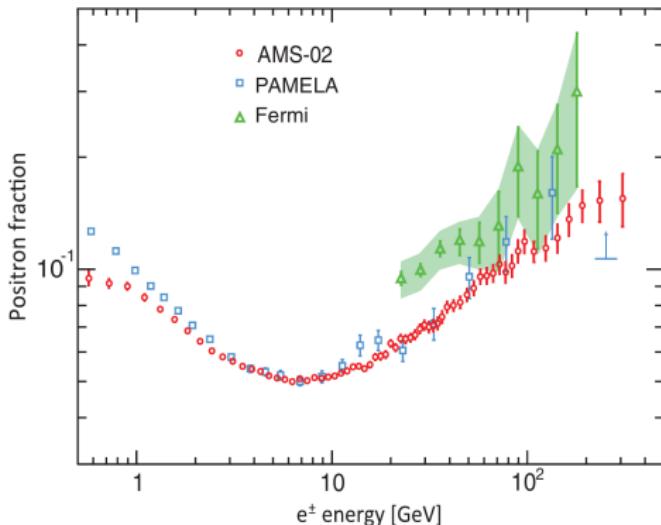
Cosmic-ray Positron Excess

 In 2008, the **PAMELA** experiment found an **unexpected increase** in the **cosmic-ray positron fraction** with $E \gtrsim 10$ GeV

In 2013, the **AMS-02** experiment **confirmed** such a **positron excess**



[PAMELA Coll., 0810.4995, Nature]

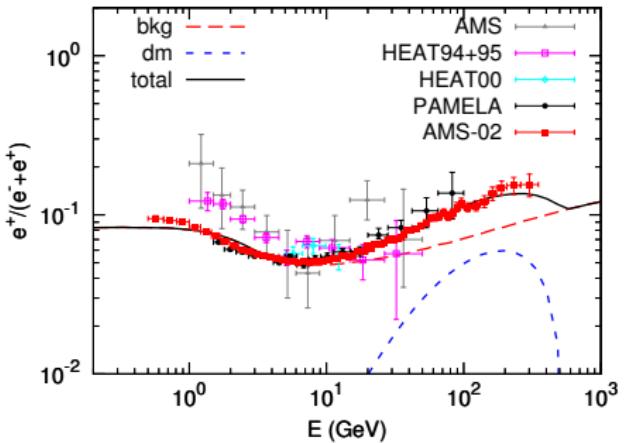


[AMS Coll., PRL 110, 141102 (2013)]

Interpretation: Dark Matter vs Pulsar



Interpretation with Galactic
DM annihilation into $\tau^+\tau^-$

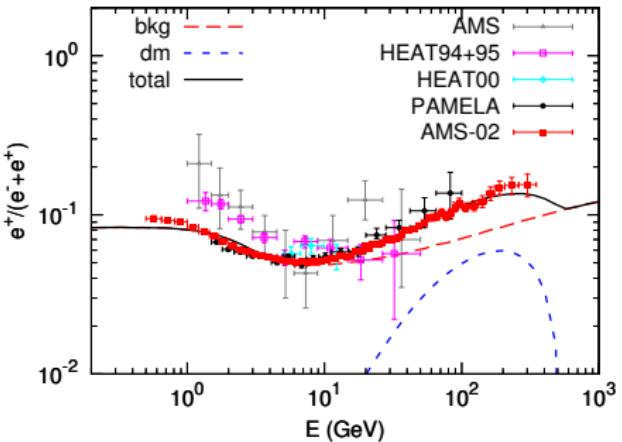


[Yuan, Bi, et al., 1304.1482, APP]

Interpretation: Dark Matter vs Pulsar



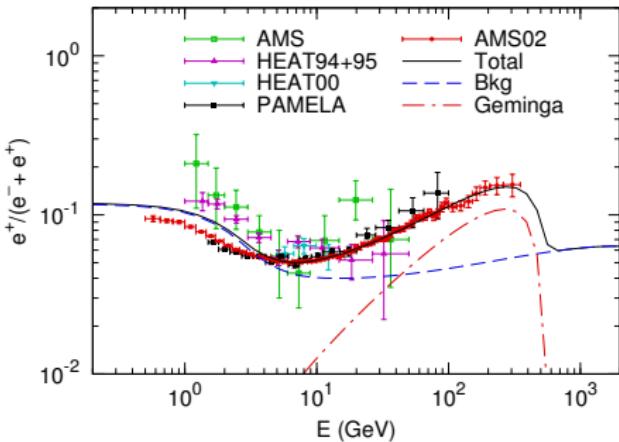
Interpretation with Galactic
DM annihilation into $\tau^+\tau^-$



[Yuan, Bi, et al., 1304.1482, APP]



Interpretation with the
nearby pulsar Geminga



[Yin, ZHY, Yuan, Bi, 1304.4128, PRD]

Past and Current High Energy Colliders

● **TEVATRON**: $p\bar{p}$ collider, 1987–2011

Circumference: 6.28 km

Collision energy: $\sqrt{s} = 1.96$ TeV

Luminosity: $\mathcal{L} \sim 4.3 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$

Detectors: CDF, DØ

● **LEP**: e^+e^- collider, 1989–2000

Circumference: 26.66 km

Collision energy: $\sqrt{s} = 91\text{--}209$ GeV

Luminosity: $\mathcal{L} \sim (2\text{--}10) \times 10^{31} \text{ cm}^{-2} \text{ s}^{-1}$

Detectors: ALEPH, DELPHI, OPAL, L3

● **LHC**: pp ($p\text{Pb}$, PbPb) collider, 2009–Now

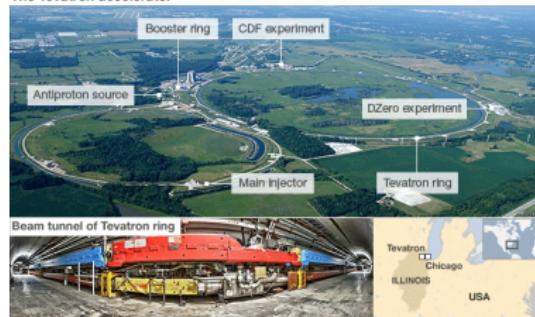
Circumference: 26.66 km

Collision energy: $\sqrt{s} = 7, 8, 13, 14$ TeV

Luminosity: $\mathcal{L} \sim (1\text{--}5) \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$

Detectors: ATLAS, CMS, ALICE, LHCb

The Tevatron accelerator



Source: Fermilab

Large Hadron Collider



Future Projects

- ## ● **ILC**: International Linear Collider

e^+e^- collider, $\sqrt{s} = 0.25-1$ TeV

$$\mathcal{L} \sim 1.5 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$

Detectors: SiD, ILD

- ## ● **CEPC**: Circular Electron-Positron Collider (China)

e^+e^- collider, $\sqrt{s} \sim 91\text{--}350$ GeV

$$\mathcal{L} \sim (0.5\text{--}100) \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$

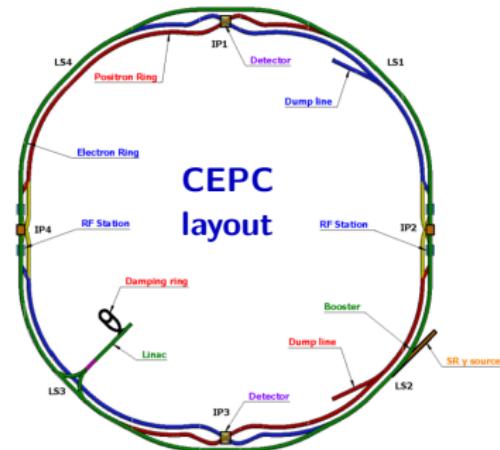
- **SPPC**: Super Proton-Proton Collider (China)

pp collider, $\sqrt{s} \sim 50\text{--}70$ TeV, $\mathcal{L} \sim 2.15 \times 10^{35}$ cm $^{-2}$ s $^{-1}$

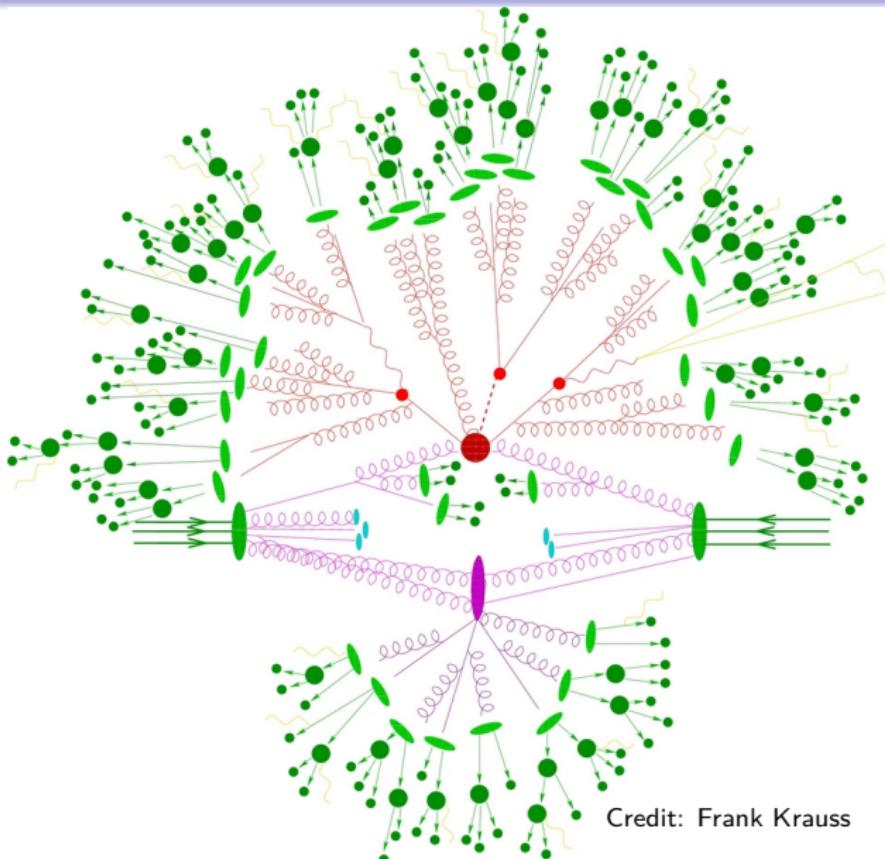
- ## ● FCC: Future Circular Collider (CERN)

- **FCC-ee**: e^+e^- collider, $\sqrt{s} \sim 91\text{--}350$ GeV, $\mathcal{L} \sim (1\text{--}200) \times 10^{34}$ cm $^{-2}$ s $^{-1}$
 - **FCC-hh**: pp collider, $\sqrt{s} \sim 100$ TeV, $\mathcal{L} \sim 5 \times 10^{34}$ cm $^{-2}$ s $^{-1}$

- **CLIC**: Compact Linear Collider, $\sqrt{s} \sim 1\text{--}3$ TeV, $\mathcal{L} \sim 6 \times 10^{34}$ cm $^{-2}$ s $^{-1}$

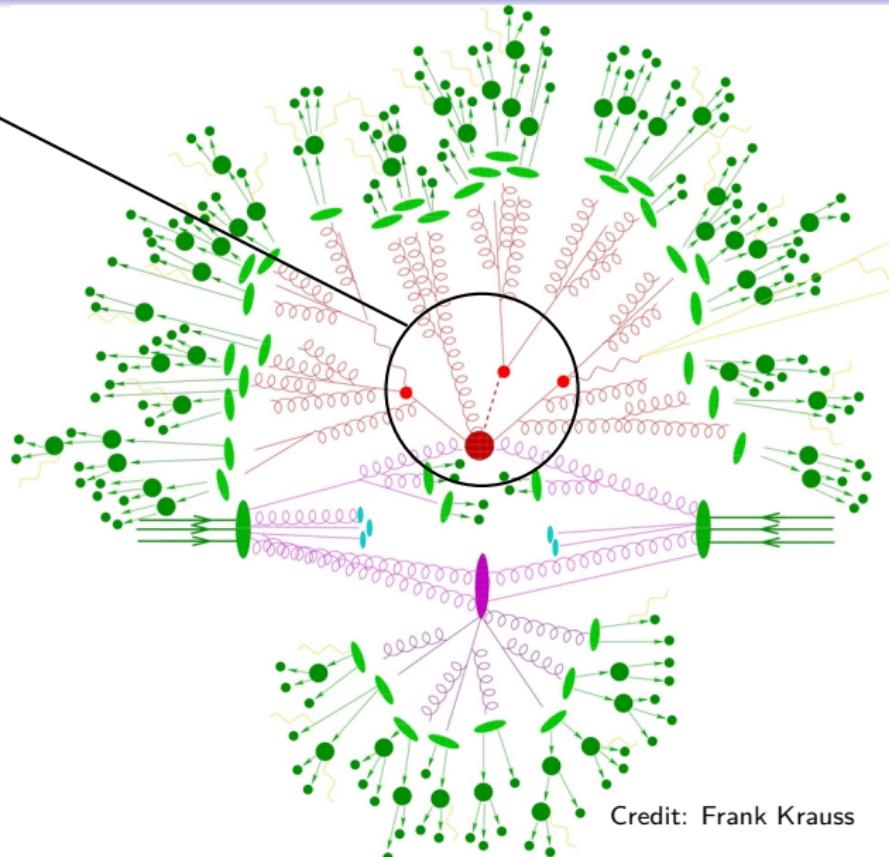


Typical Event in High Energy pp Collisions



Typical Event in High Energy pp Collisions

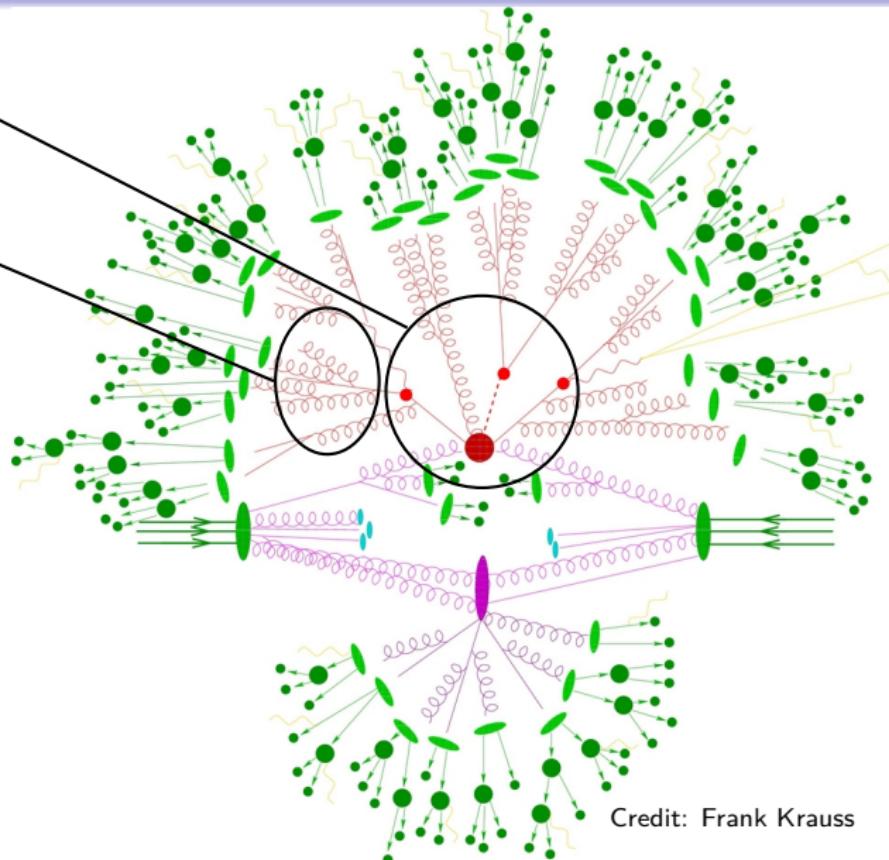
Hard scattering



Typical Event in High Energy pp Collisions

Hard scattering

Parton shower



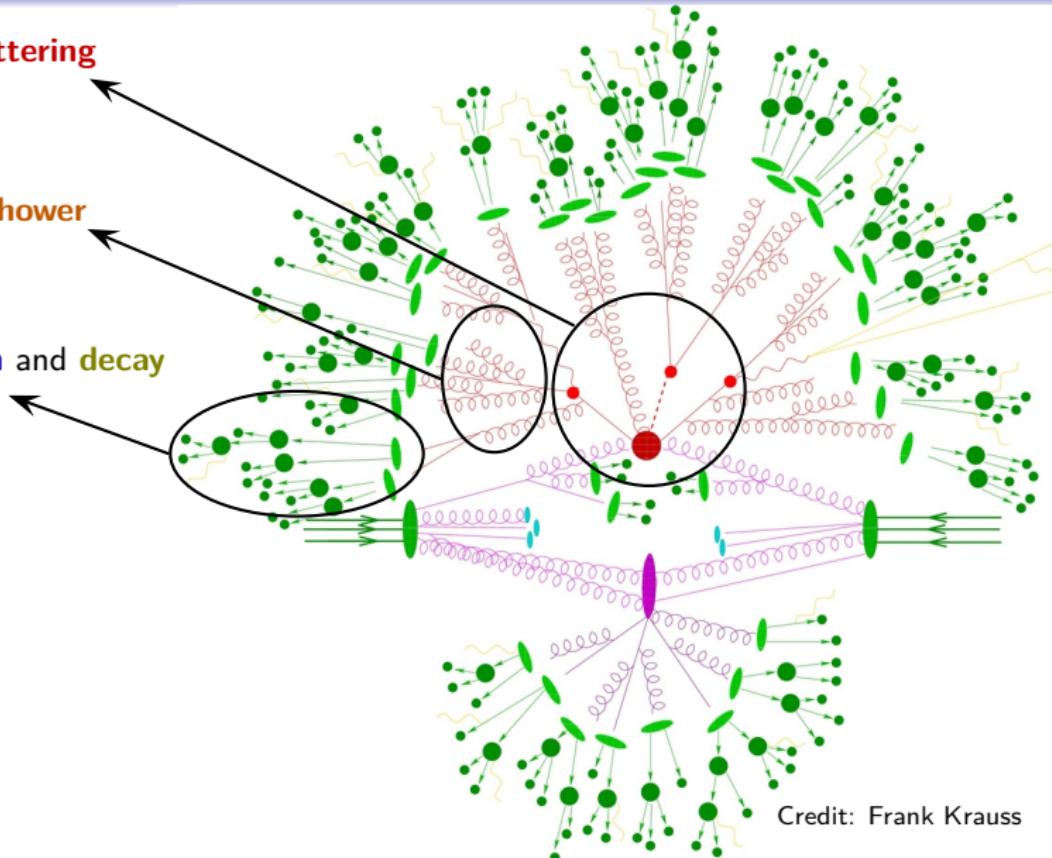
Credit: Frank Krauss

Typical Event in High Energy pp Collisions

Hard scattering

Parton shower

Hadronization and decay



Credit: Frank Krauss

Typical Event in High Energy pp Collisions

Hard scattering

Parton shower

Hadronization and decay

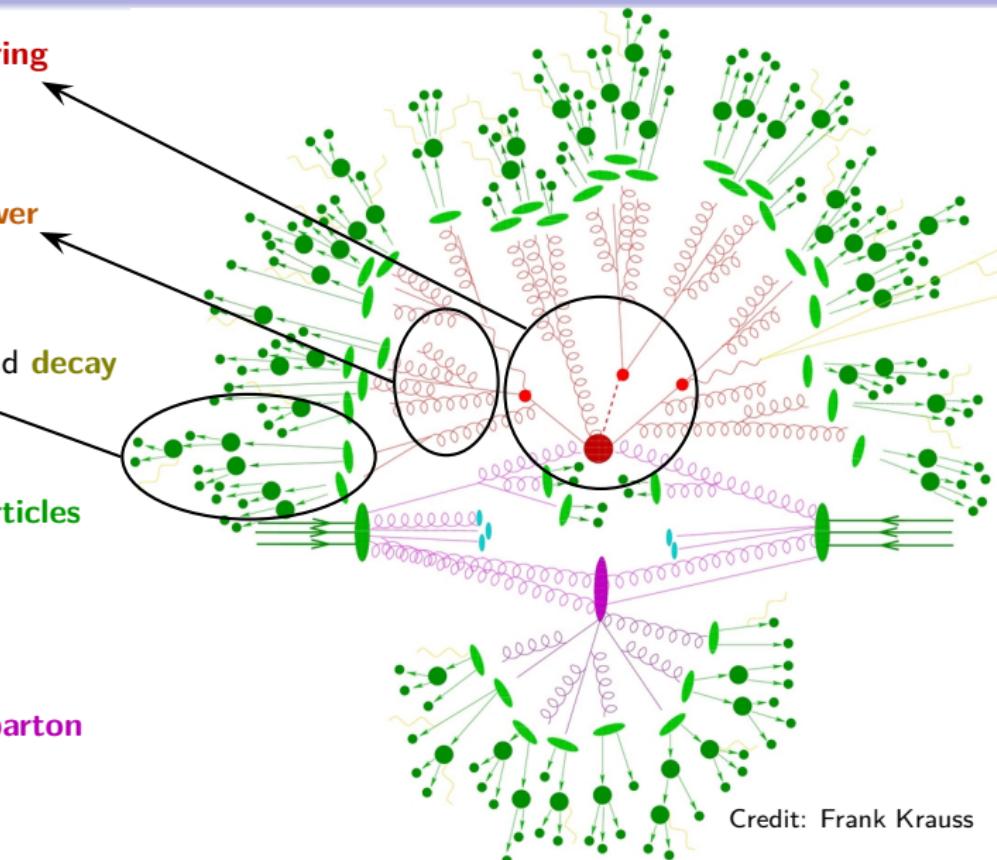
Collimated particles



Jet

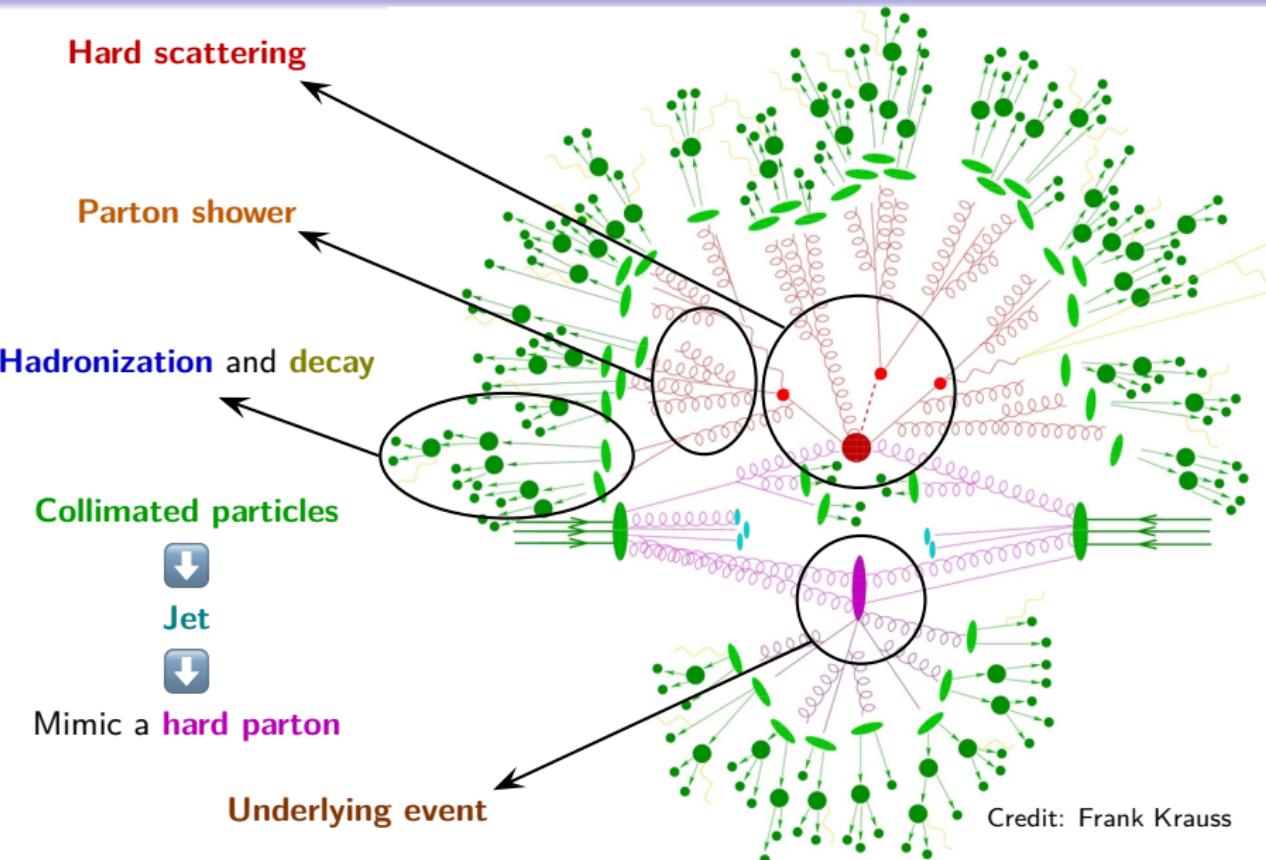


Mimic a hard parton

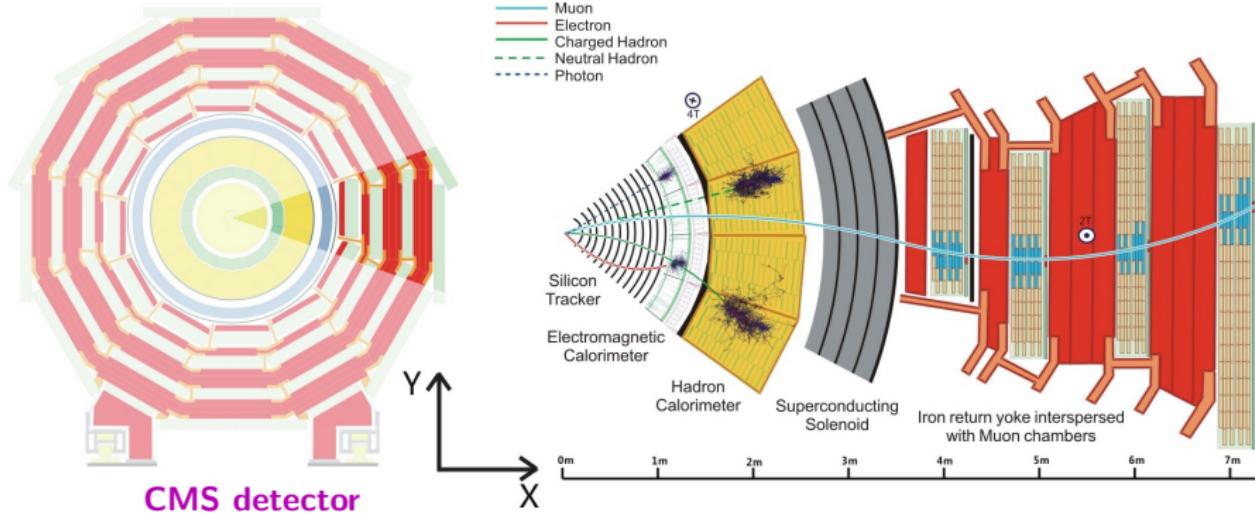


Credit: Frank Krauss

Typical Event in High Energy pp Collisions

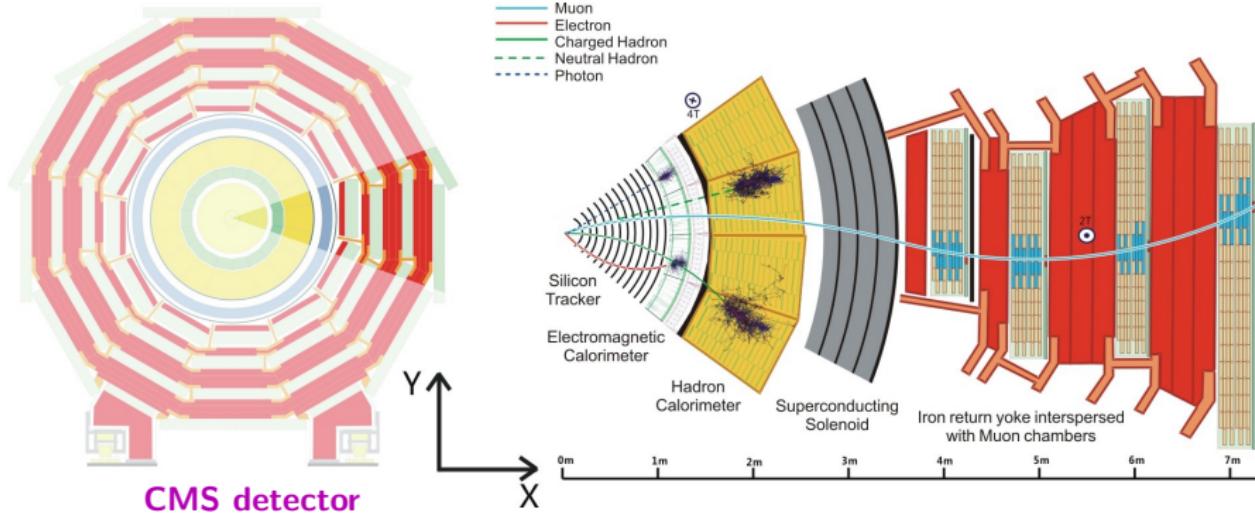


Particle Detectors at Colliders



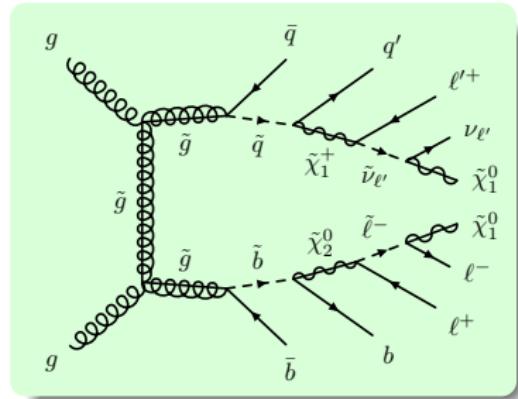
Sub-detectors	γ	e^\pm	μ^\pm	Charged hadrons	Neutral hadrons	ν , DM
Tracker, $ \eta \lesssim 2.5$	×	✓	✓	✓	×	×
ECAL, $ \eta \lesssim 3$	✖	✖	✓	✓	×	×
HCAL, $ \eta \lesssim 5$	×	✗	✗	✖	✖	✗
Muon detectors, $ \eta \lesssim 2.4$	×	✗	✓	✗	✗	✗

Particle Detectors at Colliders



Sub-detectors	γ	e^\pm	μ^\pm	Charged hadrons	Neutral hadrons	ν , DM
Tracker, $ \eta \lesssim 2.5$	✗	✓	✓	✓	✗	✗
ECAL, $ \eta \lesssim 3$	✗	✗	✓	✓		✗
HCAL, $ \eta \lesssim 5$	✗	✗	✗	✗		✗
Muon detectors, $ \eta \lesssim 2.4$	✗	✗	✓	✗	✗	✗

DM Production



Social dark matter

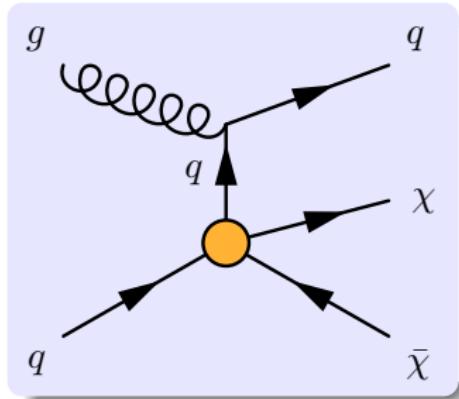
Accompanied by **other new particles**

Complicated decay chains

Decay products of other particles

Various final states

(jets + leptons + \cancel{E}_T , ...)



Maverick dark matter

DM particle is the **only new particle**

reachable at the collision energy

Direct production

Mono-X + \cancel{E}_T final states

(monojet, mono- γ , mono- W/Z , ...)

τ -portal Simplified DM Models

>We studied four **τ -portal simplified models** involving a mediator with additive quantum numbers identical to the right-handed τ^-

We interpreted the **GC GeV excess signal** as **DM annihilation into $\tau^+\tau^-$** , and discussed **how to test this interpretation at the LHC**

Spin-1/2 fermion χ , spin-0 mediator ϕ :

$$\mathcal{L}_\phi = \lambda \phi \bar{\tau}_R \chi_L + \text{H.c.}$$

● **DFDM model:** χ is a Dirac fermion

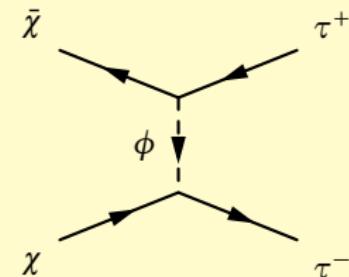
● **MFDM model:** χ is a Majorana fermion

Spin-0 scalar χ , spin-1/2 mediator ψ :

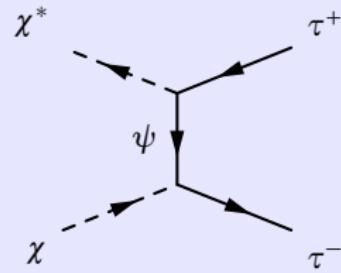
$$\mathcal{L}_\psi = \kappa \chi \bar{\tau}_R \psi_L + \text{H.c.}$$

● **CSDM model:** χ is a complex scalar

● **RSDM model:** χ is a real scalar



DFDM: annihilation

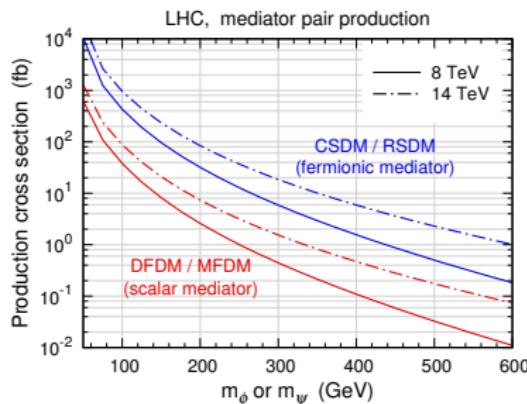


CSDM: annihilation

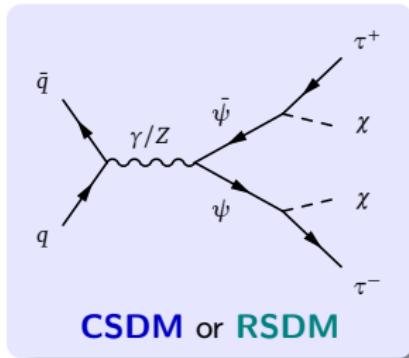
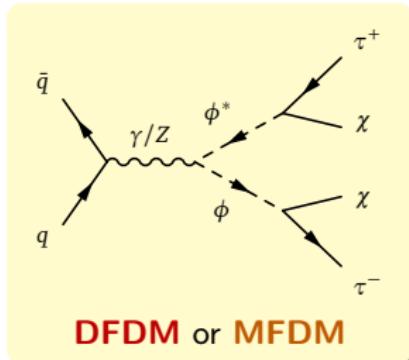
Mediator Pair Production at the LHC

 The mediators ϕ and ψ could be produced at the LHC through **Drell-Yan processes** exchanging s -channel γ or Z , and then decay into τ^\pm and χ

 We found that the 8 TeV LHC data cannot explore the interesting regions in these models, and went further to investigate the LHC sensitivity at $\sqrt{s} = 14$ TeV with **tight τ_h -tagging** techniques



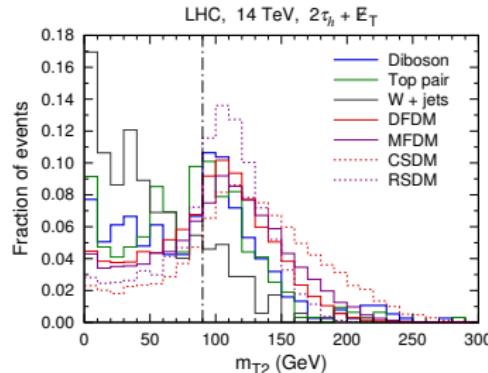
[ZHY, Bi, Yan, Yin, 1410.3347, PRD]





14 TeV LHC Searches for $pp \rightarrow \phi\phi^*/\psi\bar{\psi} \rightarrow \tau^+\tau^-\chi\chi$

🎃 **$2\tau_h + \cancel{E}_T$ channel:** two opposite-sign tau-jet (τ_h);
without any other particle; $m_{T2} > 90$ GeV



✳️ **Signals:**

DFDM model

$m_\phi = 225$ GeV

MFDM model

$m_\phi = 250$ GeV

CSDM model

$m_\psi = 300$ GeV

RSDM model

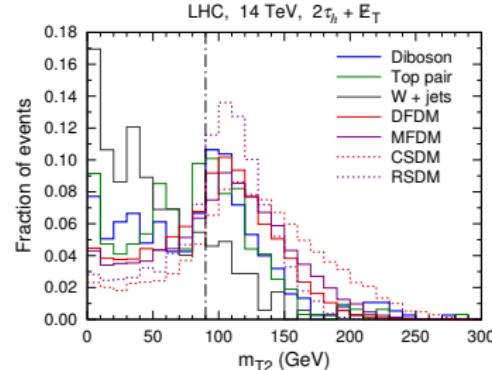
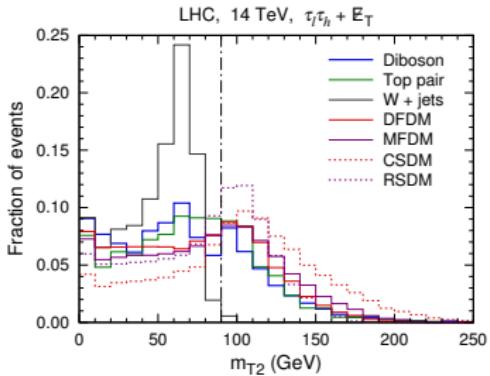
$m_\psi = 200$ GeV



14 TeV LHC Searches for $pp \rightarrow \phi\phi^*/\psi\bar{\psi} \rightarrow \tau^+\tau^-\chi\chi$

🎃 **$2\tau_h + \cancel{E}_T$ channel:** two opposite-sign tau-jet (τ_h); without any other particle; $m_{T2} > 90$ GeV

🧩 **$\tau_\ell\tau_h + \cancel{E}_T$ channel:** one τ_h and one light lepton ($\ell = \mu, e$) with opposite signs; without any other particle; $m_{T2} > 90$ GeV



🧩 **Signals:**

DFDM model

$m_\phi = 225$ GeV

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$m_\psi = 300$ GeV

RSDM model

$m_\psi = 200$ GeV

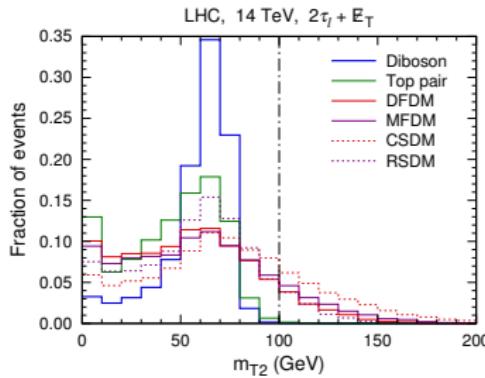
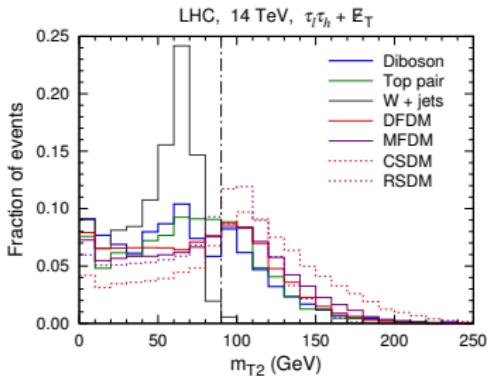
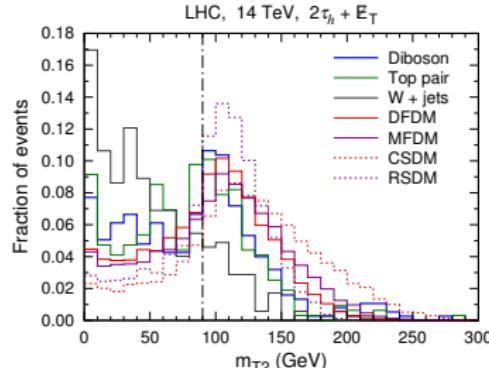


14 TeV LHC Searches for $pp \rightarrow \phi\phi^*/\psi\bar{\psi} \rightarrow \tau^+\tau^-\chi\chi$

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风筝 **$\tau_\ell\tau_h + \cancel{E}_T$ channel:** one τ_h and one light lepton ($\ell = \mu, e$) with opposite signs; without any other particle; $m_{T2} > 90$ GeV

紫色气球 **$2\tau_\ell + \cancel{E}_T$ channel:** two opposite-sign light leptons; $|m_{\ell\ell} - m_Z| > 10$ GeV for the same-favor case; without any other particle; $m_{T2} > 100$ GeV



Signals:

DFDM model

$$m_\phi = 225 \text{ GeV}$$

MFDM model

$$m_\phi = 250 \text{ GeV}$$

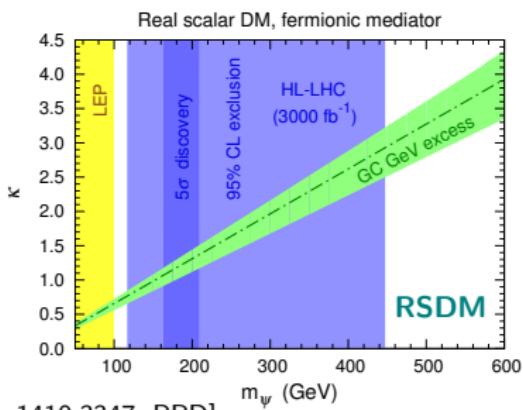
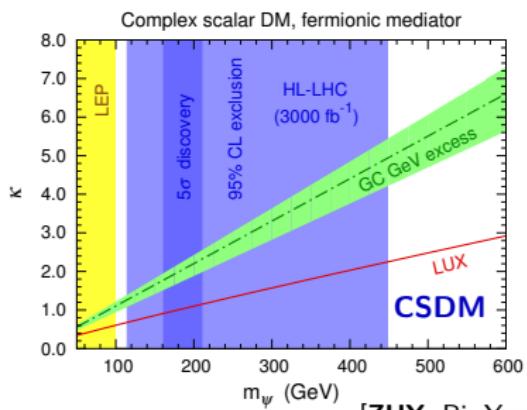
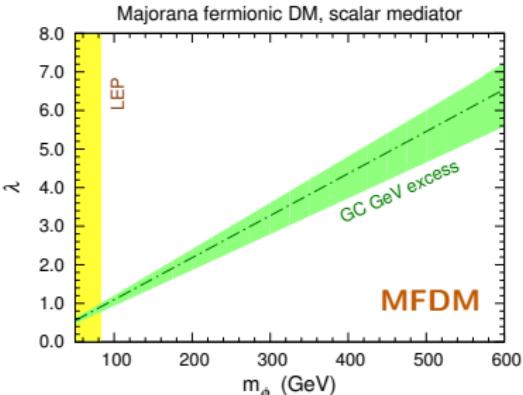
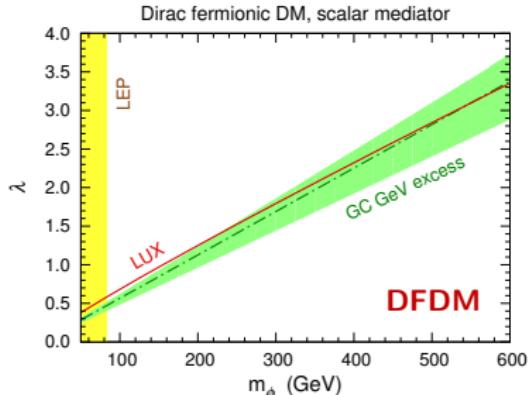
CSDM model

$$m_\psi = 300 \text{ GeV}$$

RSDM model

$$m_\psi = 200 \text{ GeV}$$

Sensitivity of the 14 TeV High-Luminosity LHC

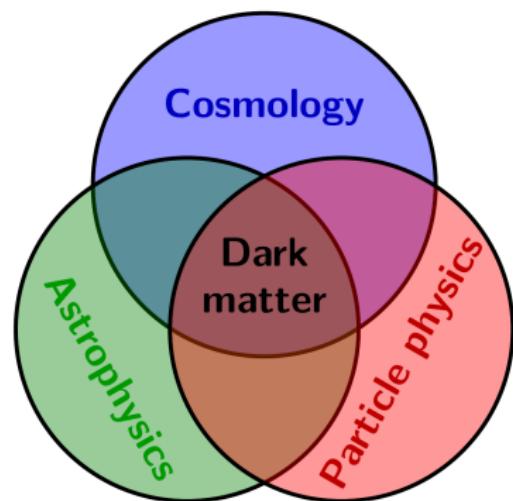


[ZHY, Bi, Yan, Yin, 1410.3347, PRD]



Summary

- **Dark matter** connects our knowledge of the Universe from the **largest** to the **smallest** scales
- Although several anomalous observations have been found in direct and indirect searches, there is **no absolutely solid DM detection signal so far**
- **DM detection sensitivities are being improved quickly**, and it is promising to detect robust DM signals in the future



Summary

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- Although several anomalous observations have been found in direct and indirect searches, there is **no absolutely solid DM detection signal so far**
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Thank you!

