

Introduction to Dark Matter

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Dark Matter in the Universe



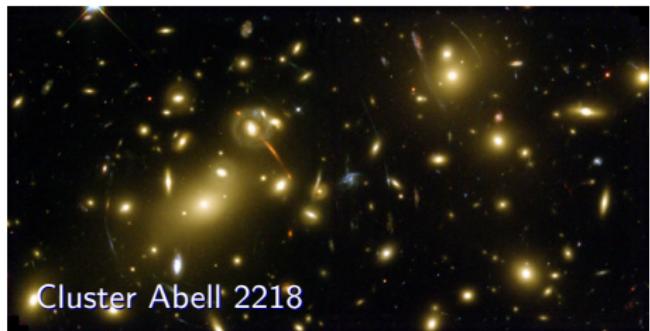
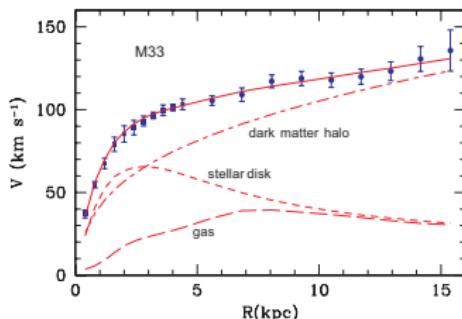
 **Dark matter (DM)** makes up **most** of the **matter** component in the Universe, as suggested by **astrophysical** and **cosmological** observations



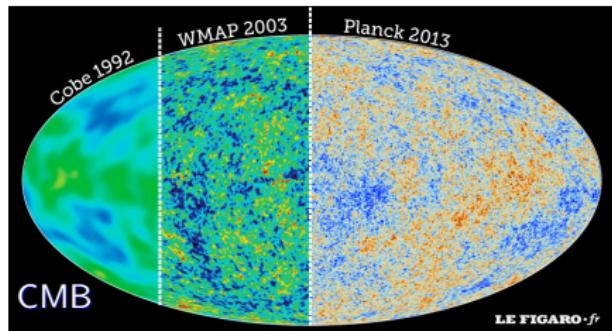
Bullet Cluster



Spiral galaxy M33



Cluster Abell 2218





Coma Cluster (后发座星系团)





Coma Cluster (后发座星系团)





Coma Cluster (后发座星系团)



✨ In 1933, **Fritz Zwicky** found that the **velocity dispersion** of **galaxies** in the **Coma cluster** was **far too large** to be supported by the **luminous matter**

🌧 Mass-to-light ratio $\Upsilon_{\text{Coma}} \sim 240 \Upsilon_{\odot}$

[Kent & Gunn, Astron.J., 87, 945 (1982)]

💧 Typical spiral galaxy: $\mathcal{O}(10) \Upsilon_{\odot}$



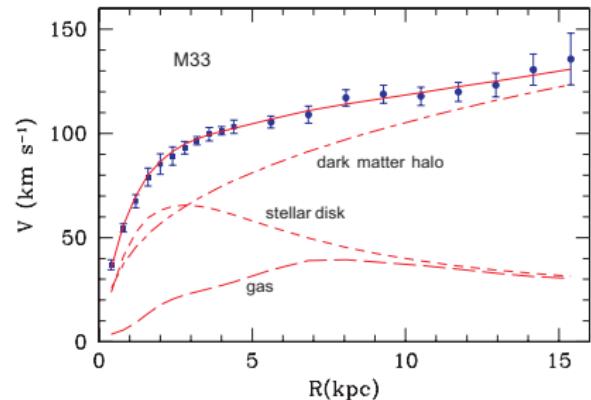
Spiral Galaxies: Rotation Curves



🌀 In the 1970s, **Vera Rubin** and her collaborators measured the **rotation curves** of **spiral galaxies** and also found evidence for **non-luminous matter**



Triangulum galaxy M33



[Corbelli & Salucci, astro-ph/9909252, MNRAS]

Spiral Galaxies: Rotation Curves



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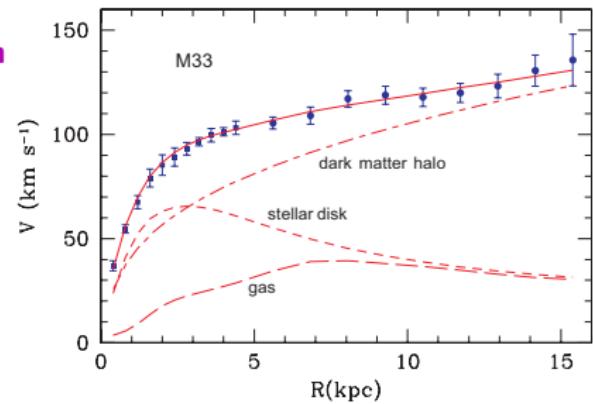


According to **Newton's law of universal gravitation**, the relation between the **rotation velocity** v and the **mass** $M(r)$ within the **radius** r should be

$$\frac{v^2}{r} = \frac{G_N M(r)}{r^2}$$

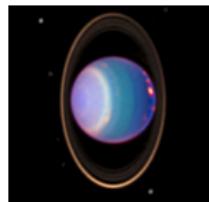
$$M(r) = \text{constant} \quad \text{👉} \quad v \propto r^{-1/2}$$

$$M(r) \propto r \quad \text{→} \quad v = \text{constant}$$



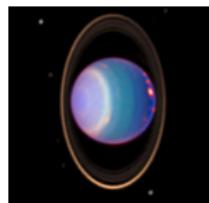
[Corbelli & Salucci, astro-ph/9909252, MNRAS]

How Can We Explain an Anomalous Phenomenon?



Unexpected movement of **Uranus** (after 1821)

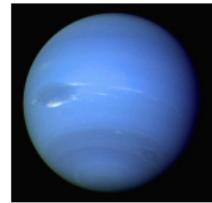
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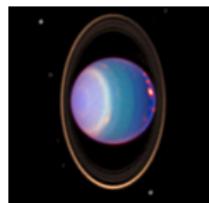


; Perturbed by **Neptune** (discovered in 1846)



Calculations independently given by **John Adams** and **Urbain Le Verrier**

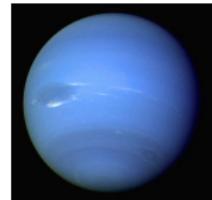
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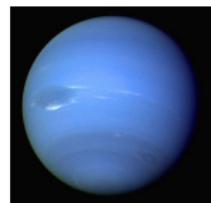
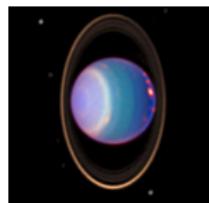


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Search for new objects/substances responsible for it!

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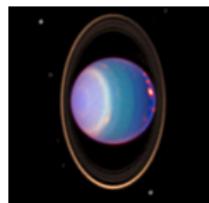


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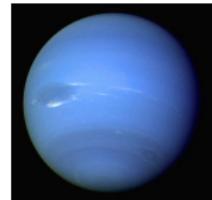


Anomalous perihelion precession of **Mercury** (Le Verrier, 1859)

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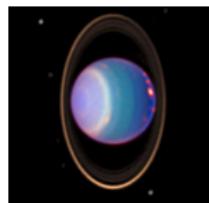
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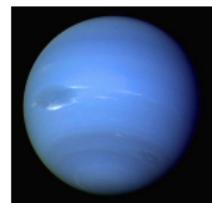
Update Newton's laws to **general relativity** (1915)



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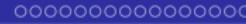
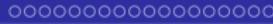


🚀 Update Newton's laws to general relativity (1915)



Modify known physical laws!





How about the Anomalous Phenomena Here?

💣 **Modify physical laws** ➡ MOdified Newtonian Dynamics (MOND)

[Milgrom, ApJ, 270, 365 (1983)]

🤔 It is **difficult** to coherently explain data at **all scales** with one model

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Consider new objects ➡ MAssive Compact Halo Objects (MACHOs)

💡 They belong to **baryonic dark matter**, including **jupiters, brown dwarfs, white dwarfs, neutron stars, black holes, etc.**

😔 **MACHO fraction** in the Galactic dark matter halo is **< 8%** (95% C.L.)

[EROS-2 Coll., astro-ph/0607207, A&A]

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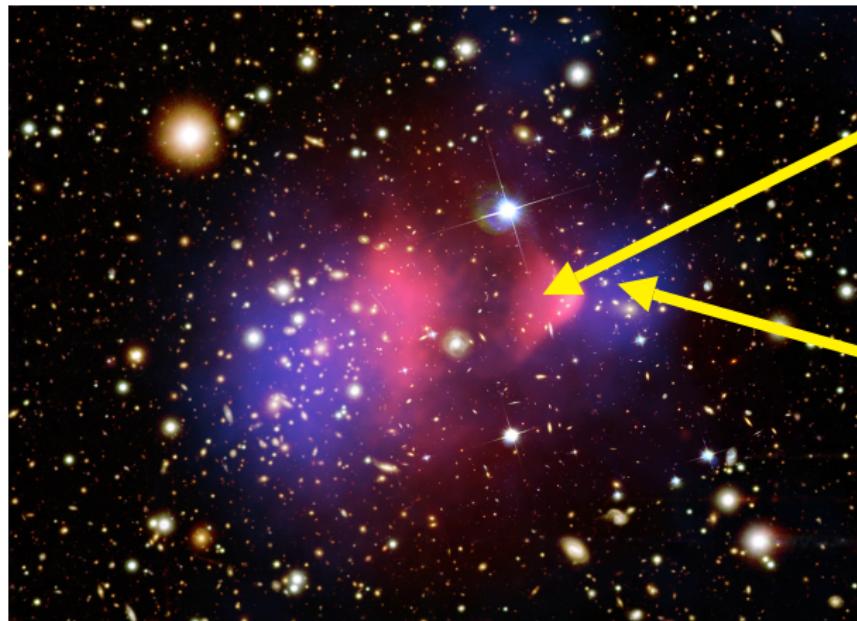
[EROS-2 Coll., astro-ph/0607207, A&A]



Consider new substances ➡ **Nonbaryonic Dark Matter**

(not constituted by baryons)

Bullet Cluster: Disfavor MOND



Fluid-like X-ray emitting plasma,
i.e., gas
(luminous matter)

Mass distribution observed by weak gravitational lensing
(DM dominated)

⛓ An 8σ significance **spatial offset** of the **center** of the **total mass** from the **center** of the **baryonic mass peaks** **cannot be explained** with a **modification** of the **gravitational force law** [Clowe et al., astro-ph/0608407, ApJL]

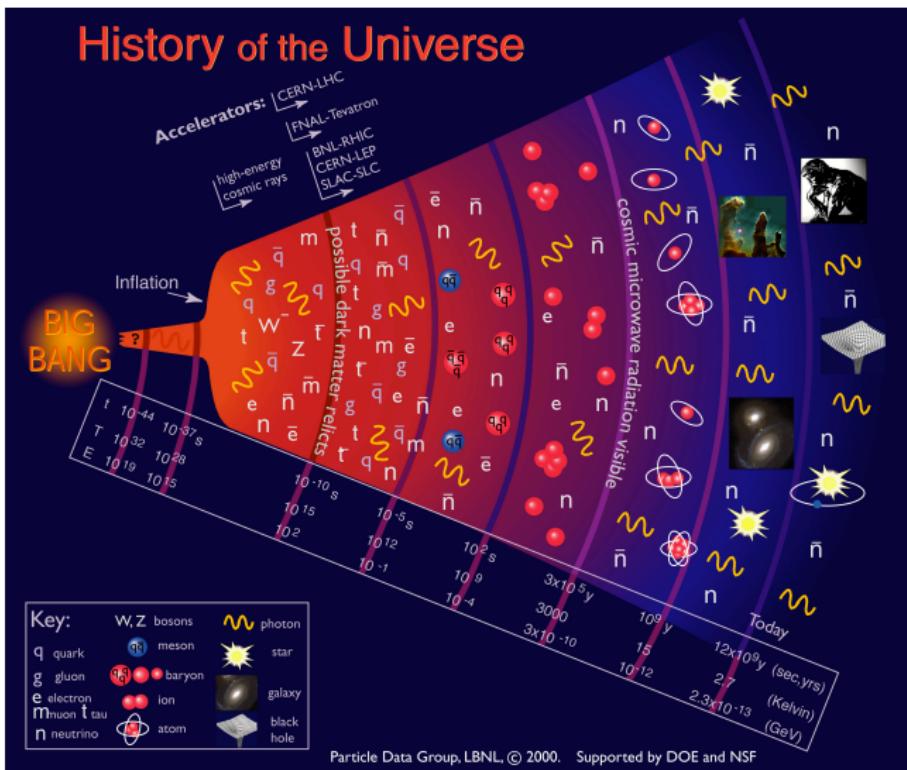
Big Bang Cosmology

~ 13.8 billion years ago, the Universe was extremely **hot, dense, and homogeneous**

Everything was in **thermal equilibrium** and interacted with each other

As the Universe **expanded and cooled down**, its constituents **decoupled** from the thermal bath **one by one**

Then **nuclei, atoms, stars, and galaxies** were **formed**



Structure Formation: Hot, Cold, and Warm Dark Matter



Small initial fluctuations + Gravitational instability



Decoupled matter generates cosmological structures



Baryonic matter decoupled too late



Only baryonic matter  **Galaxies would not be formed!**



 Needs **nonbaryonic dark matter** which **decoupled much earlier**

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Hot dark matter (such as neutrinos): relativistic when it decoupled

Structures are formed by fragmentation (top-down)



Cold dark matter (CDM): nonrelativistic when it decoupled

Structures are formed hierarchically (bottom-up)



The observation that galaxies are older than clusters favors cold dark matter

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Structures are formed **hierarchically (bottom-up)**



The observation that **galaxies** are **older** than **clusters** favors **cold dark matter**



Milky Way dwarf satellites: ~ 60 (**observed**) vs. ~ 500 (**CDM predicted**)

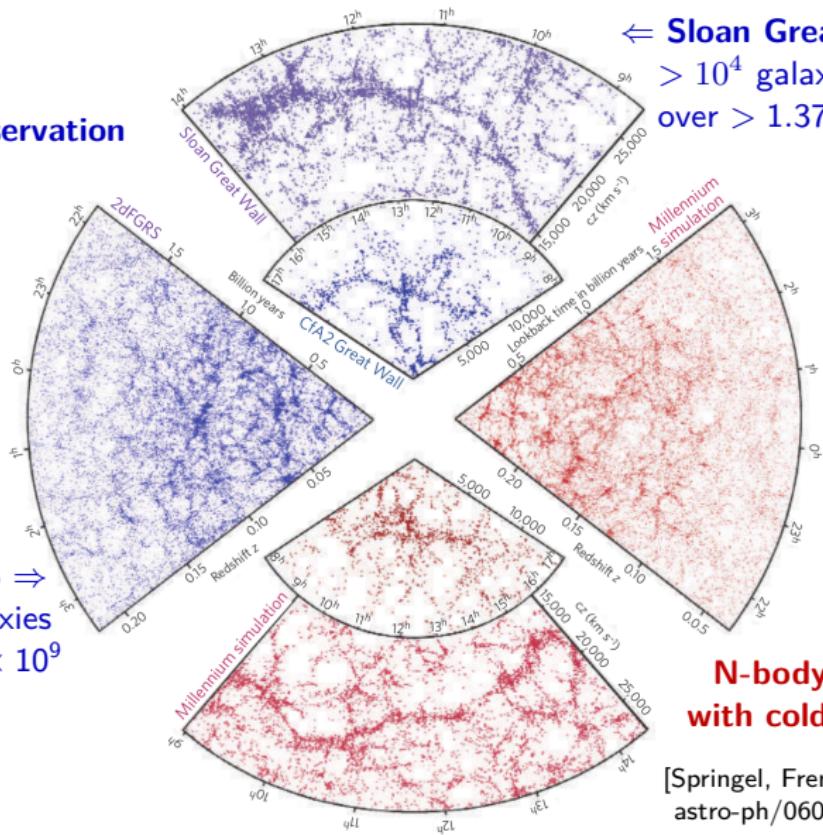


“Missing satellites problem” A component of **warm dark matter?**



Galaxy Distribution: Observation vs Simulation

Observation



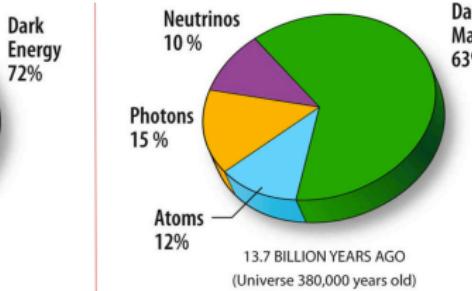
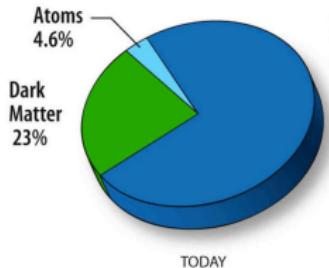
Standard Cosmology: Λ CDM Model



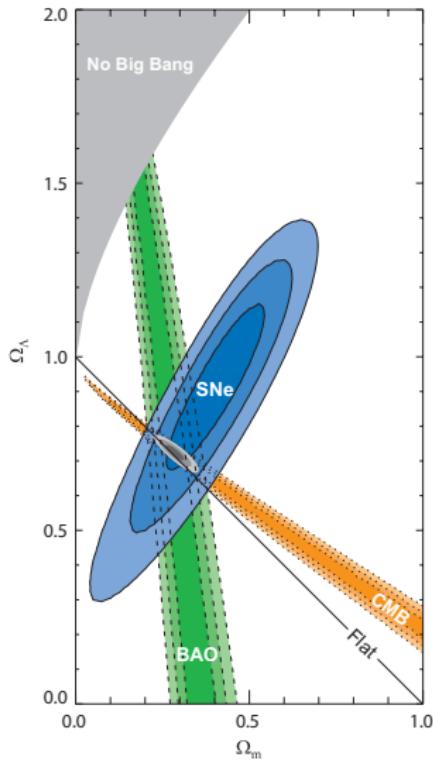
Λ CDM: the **standard cosmological model**

- **Cosmological constant Λ** (dark energy)
- **Cold dark matter (CDM)**
- ⚠ The evolution of the Universe is governed by the **Friedmann equation**

$$\frac{k}{H^2 R^2} = \Omega_\Lambda + \Omega_m + \Omega_r - 1$$



[WMAP Science Team]



[Kowalski *et al.*, 0804.4142, ApJ]



Cosmic Microwave Background (CMB)

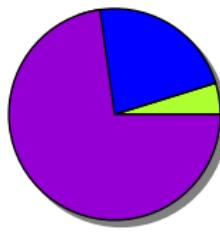
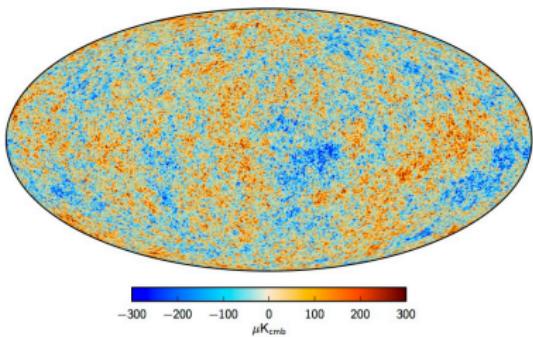
$$t \sim 380\,000 \text{ yr}, T \sim 3000 \text{ K}$$

Electrons + Protons \rightarrow Hydrogen Atoms
Photons decoupled

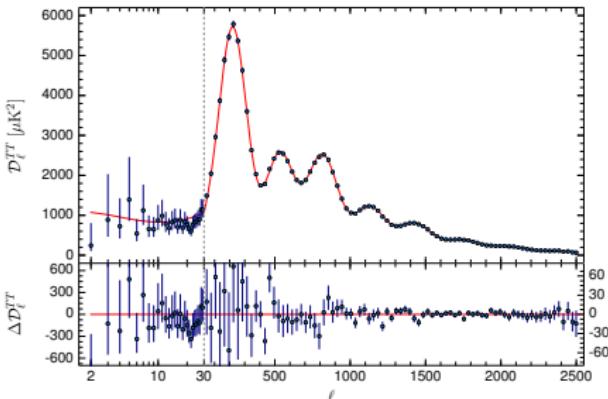
Cools \downarrow down

Today, $\sim 2.7 \text{ K}$ microwave background

 **Cosmological parameters**, e.g., Ω_Λ , Ω_c , and Ω_b , can be determined by measuring the **CMB anisotropy power spectrum**



[1502.01582, 1502.01589]



Cold DM (25.8%)

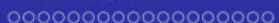
$$\Omega_{ch}^2 = 0.1186 \pm 0.0020$$

Baryons (4.8%)

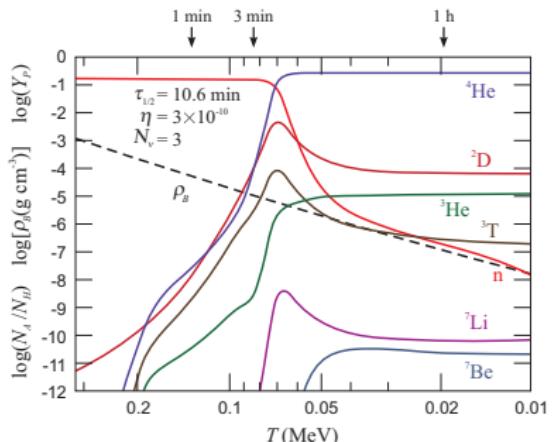
$$\Omega_b h^2 = 0.02226 \pm 0.00023$$

Dark energy (69.3%)

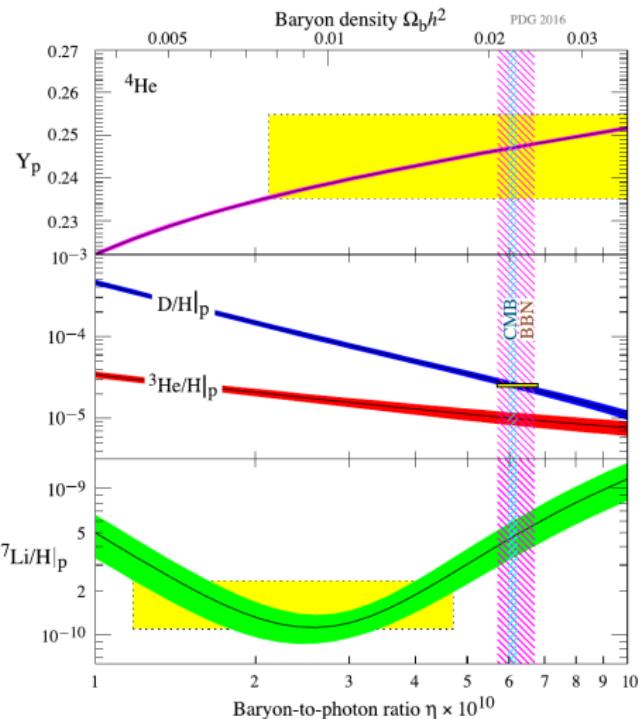
$$\Omega_\Lambda = 0.692 \pm 0.012$$



Big Bang Nucleosynthesis (BBN): $t \sim 1 \text{ sec} - 1 \text{ hour}$



[Kolb & Turner, *The Early Universe*]



Light element primordial abundances



Infer the **baryon density** Ω_b
(consistent with **CMB observations**)



The majority of matter is **nonbaryonic**

[PDG 2016]

Inferred Properties of Dark Matter

- ① **Dark (electrically neutral): no light emitted** from it
- ② **Nonbaryonic:** BBN & CMB observations
- ③ **Long lived:** survived from early eras of the Universe to now
- ④ Colorless: otherwise, it would **bind with nuclei**
- ⑤ **Cold:** structure formation theory
- ⑥ **Abundance:** **more than 80%** of all matter in the Universe

$$\rho_{\text{DM}} \sim 0.3\text{--}0.4 \text{ GeV/cm}^3 \text{ near the earth}$$

Standard Model (SM) of Particle Physics

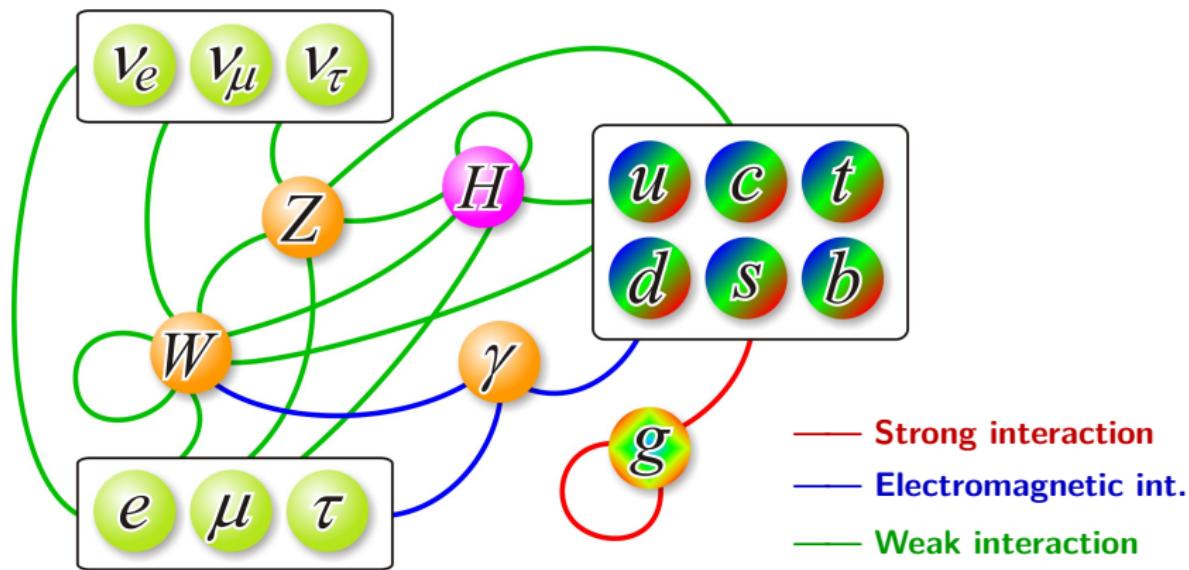
-  $SU(3)_C \times SU(2)_L \times U(1)_Y$ gauge symmetry
-  Englert-Brout-Higgs mechanism
-  $SU(3)_C \times U(1)_{EM}$ gauge symmetry



Fermion masses

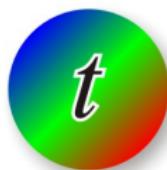


Higgs boson



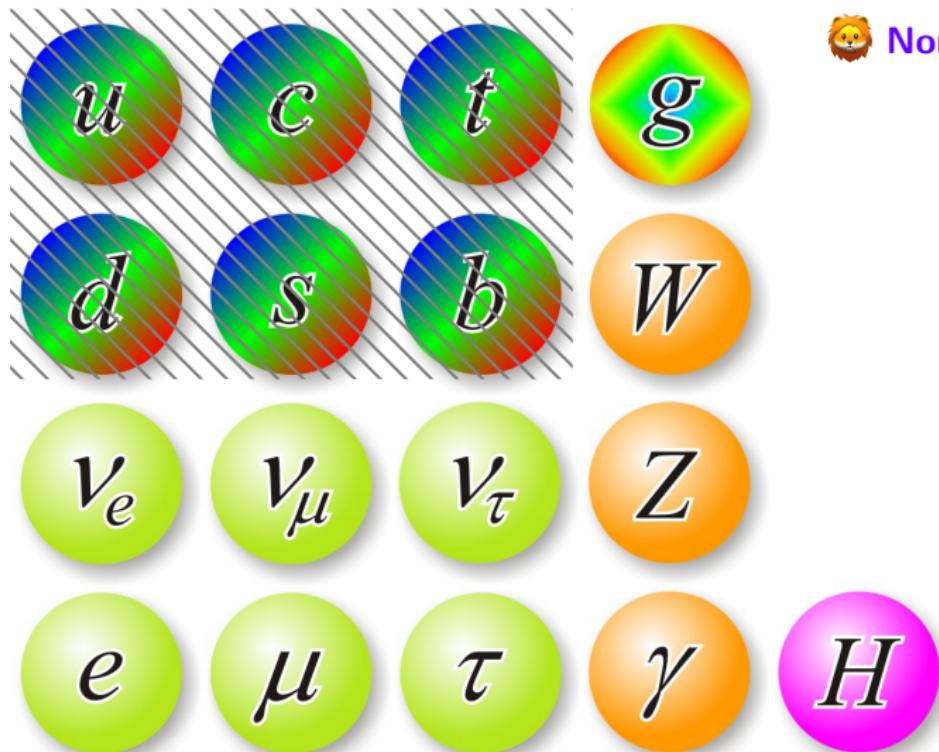


Are There Dark Matter Candidates in the Standard Model?





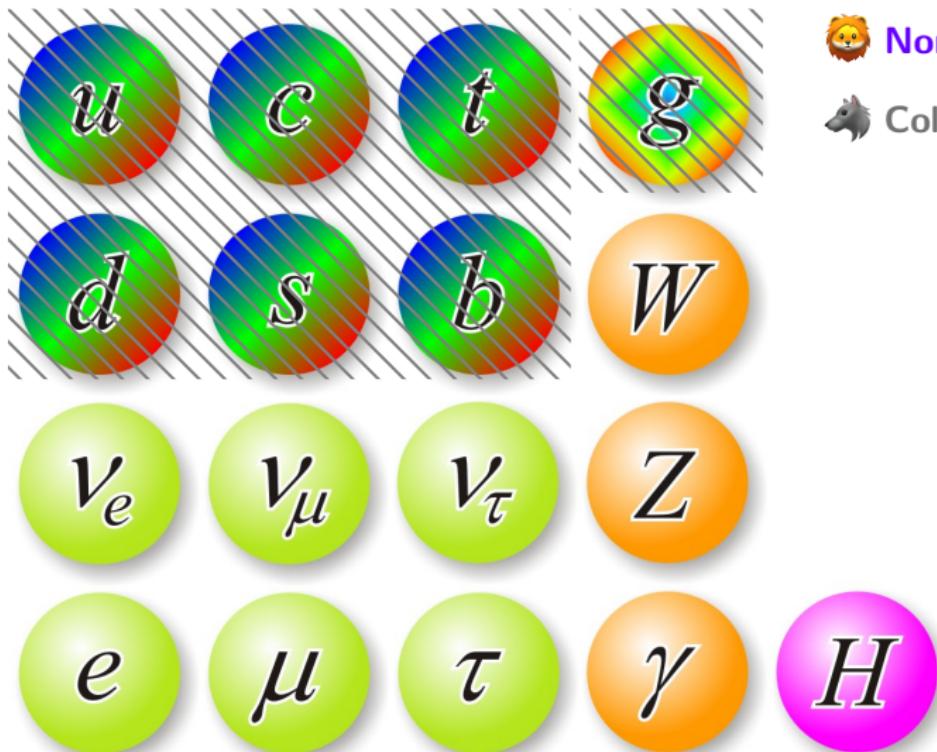
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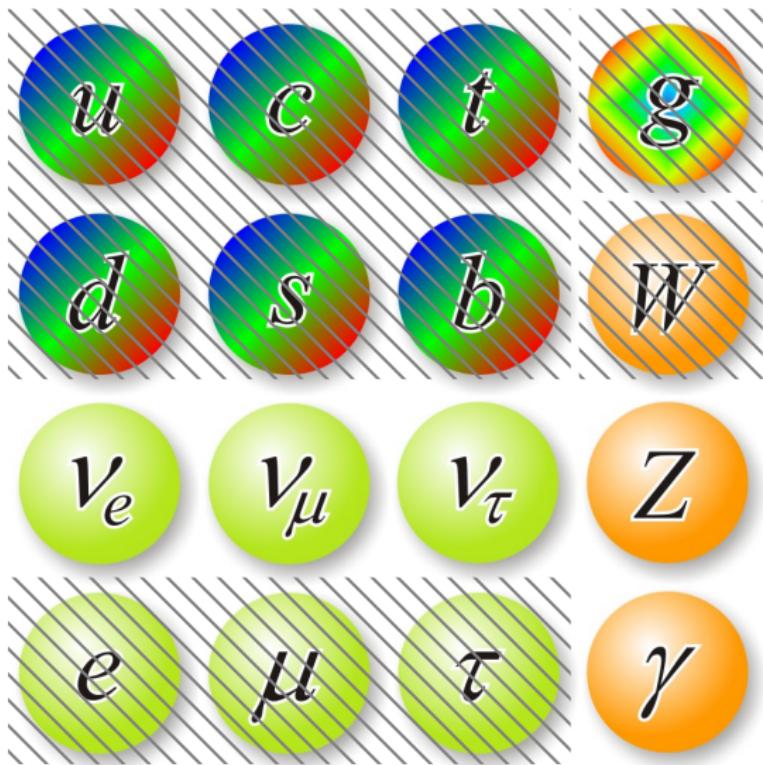
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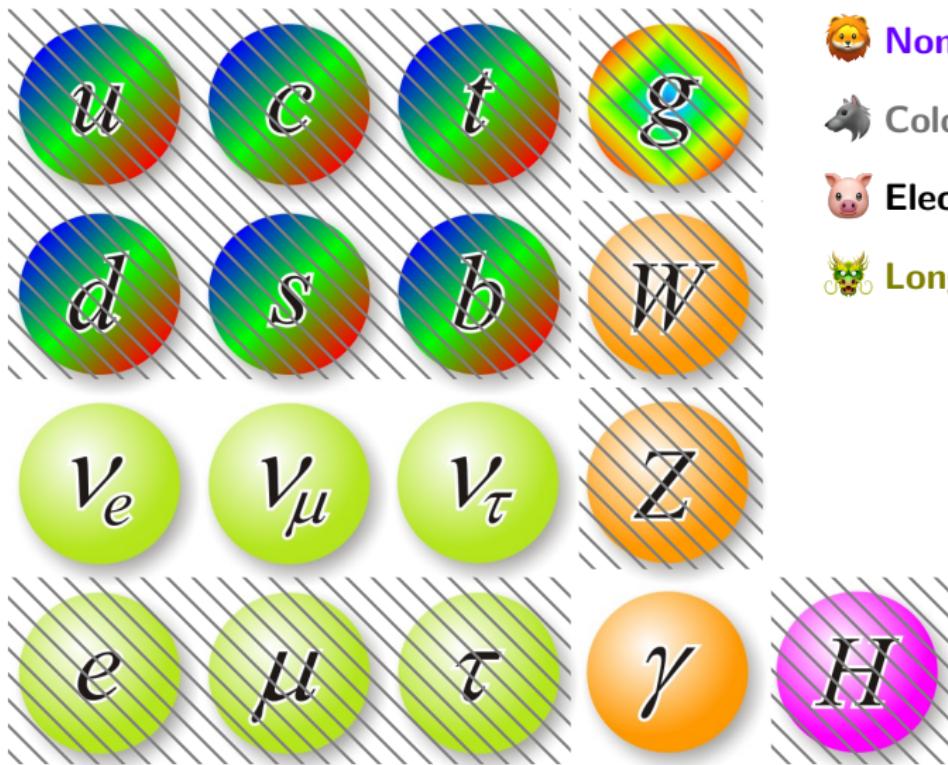
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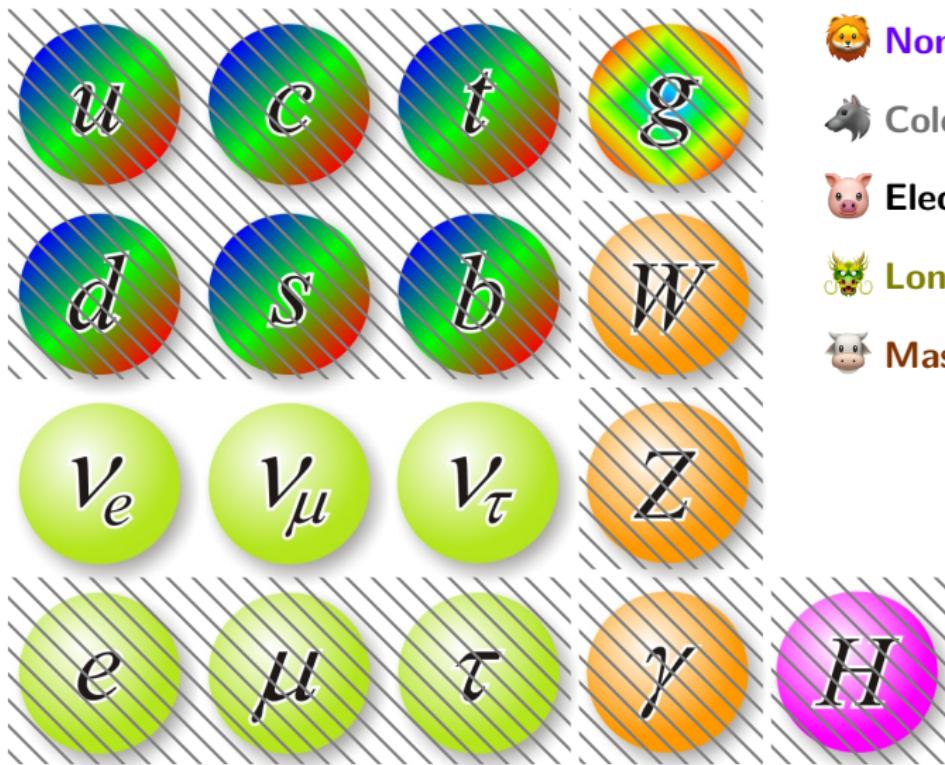
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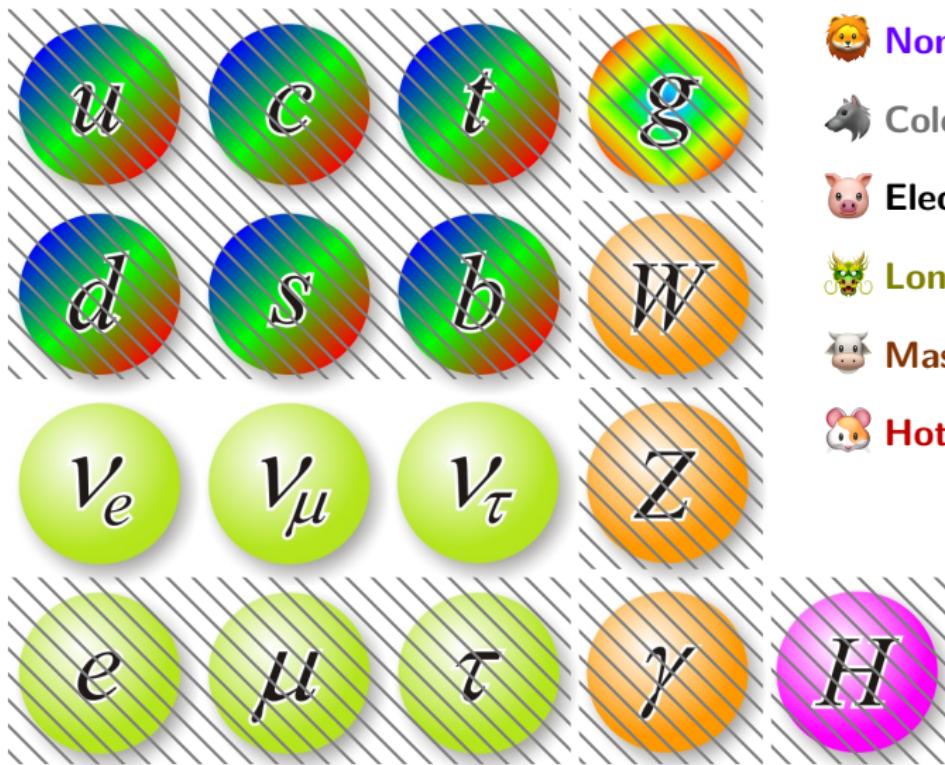
Long lived



Massive



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Colorless



Electrically neutral



Long lived



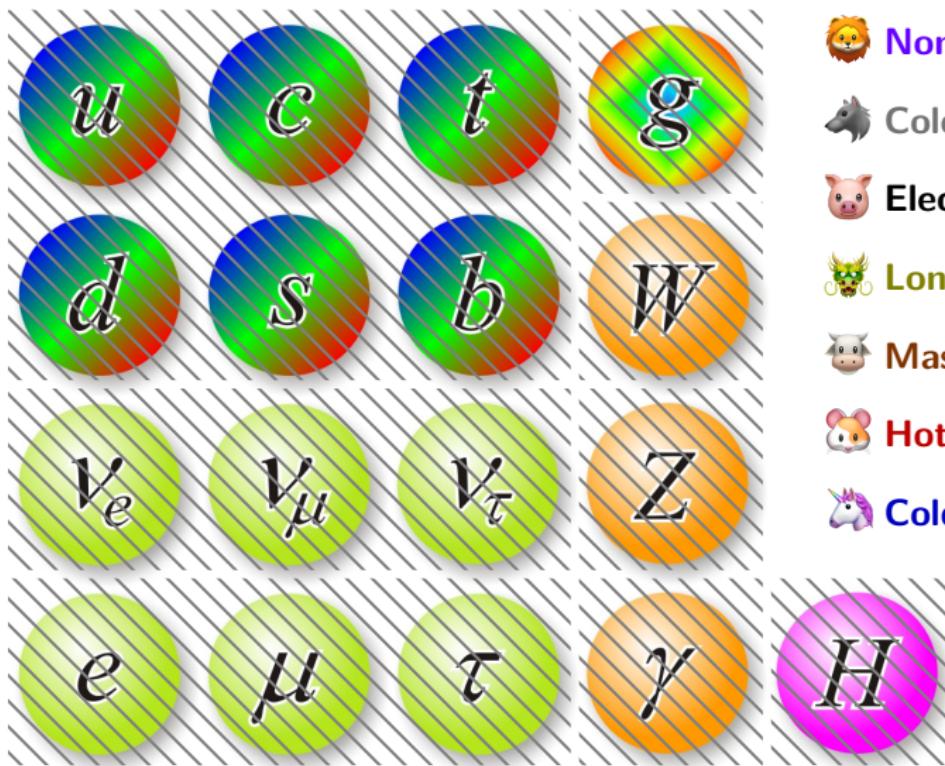
Massive



Hot DM: neutrinos



Are There Dark Matter Candidates in the Standard Model?



Nonbaryonic



Colorless



Electrically neutral



Long lived



Massive



Hot DM: neutrinos



Cold DM: none



DM Relic Abundance from Thermal Production

If DM particles (χ) were thermally produced in the early Universe, their relic abundance would be determined by the annihilation cross section $\langle \sigma_{\text{ann}} v \rangle$:

$$\Omega_\chi h^2 \simeq \frac{3 \times 10^{-27} \text{ cm}^3 \text{ s}^{-1}}{\langle \sigma_{\text{ann}} v \rangle}$$

Observed value $\Omega_\chi h^2 \simeq 0.1$

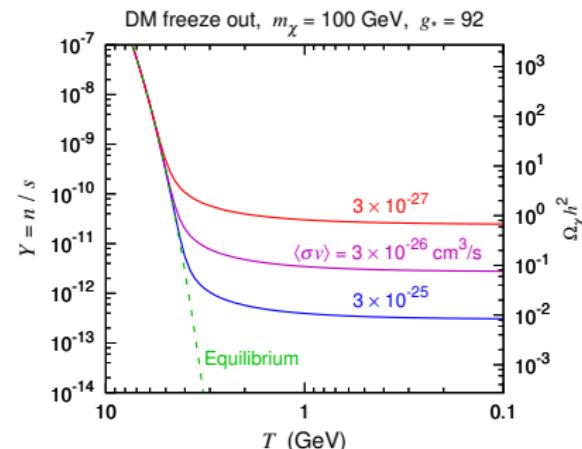
$\langle \sigma_{\text{ann}} v \rangle \simeq 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$

Assuming the annihilation process consists of two weak interaction vertices with the $SU(2)_L$ gauge coupling $g \simeq 0.64$, for $m_\chi \sim \mathcal{O}(\text{TeV})$ we have

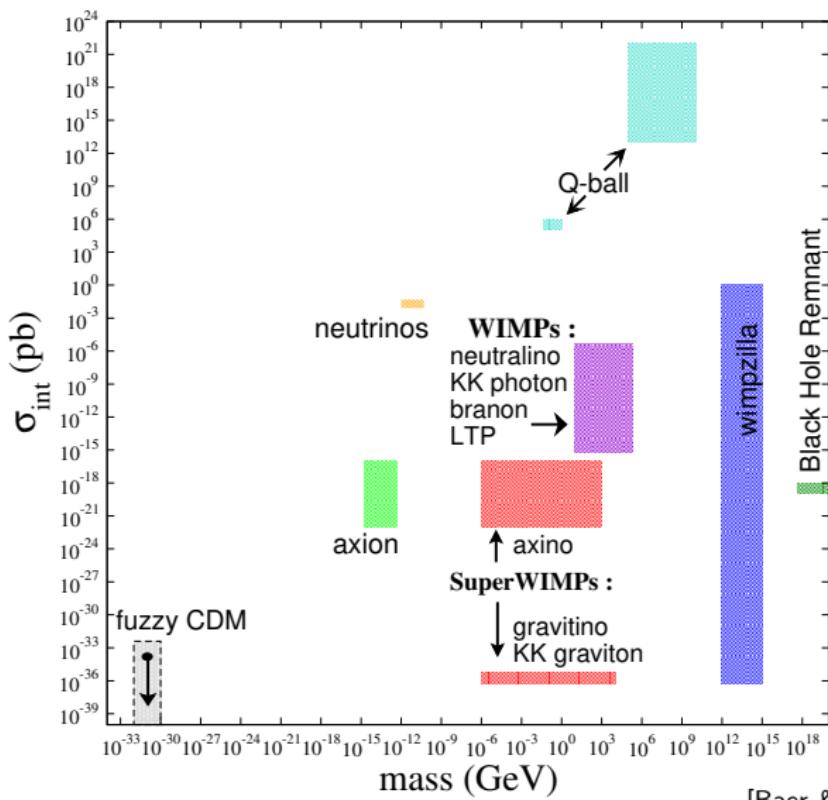
$$\langle \sigma_{\text{ann}} v \rangle \sim \frac{g^4}{16\pi^2 m_\chi^2} \sim \mathcal{O}(10^{-26}) \text{ cm}^3 \text{ s}^{-1}$$

A very attractive class of DM candidates:

Weakly interacting massive particles (WIMPs)

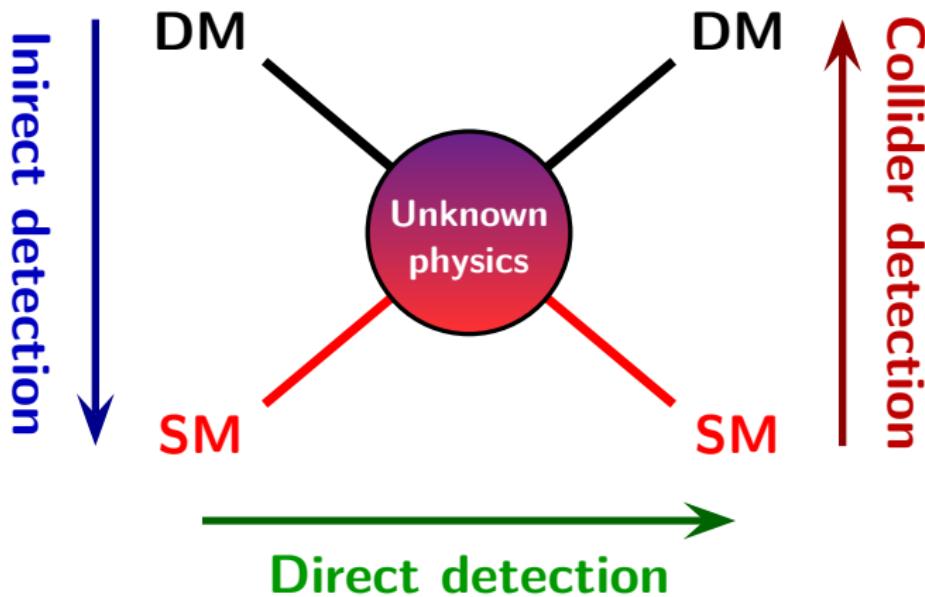


DM Particle Candidates



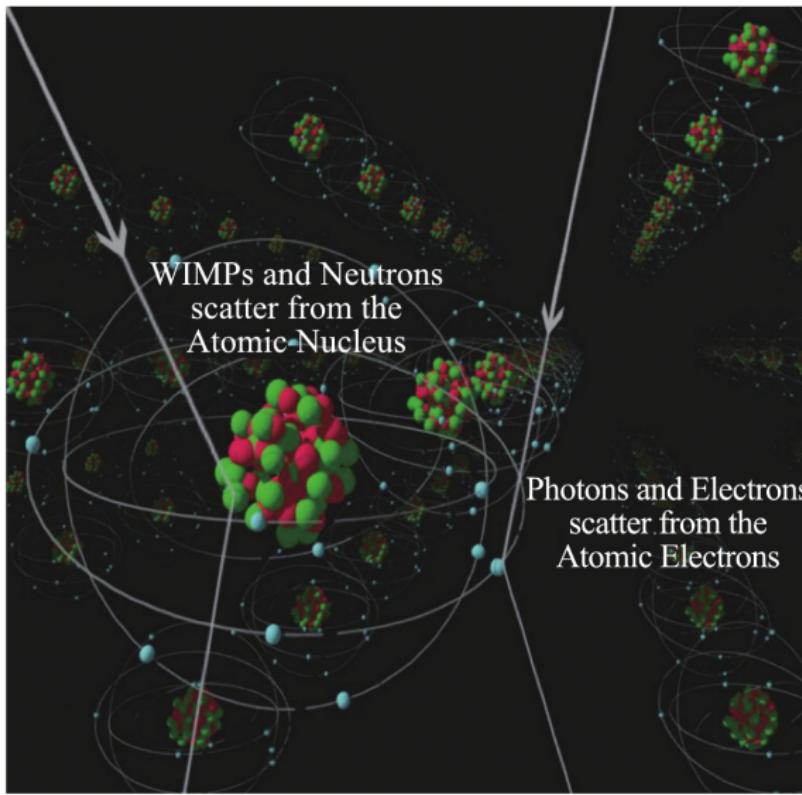
[Baer & Tata, 0805.1905]

Experimental Approaches to Dark Matter



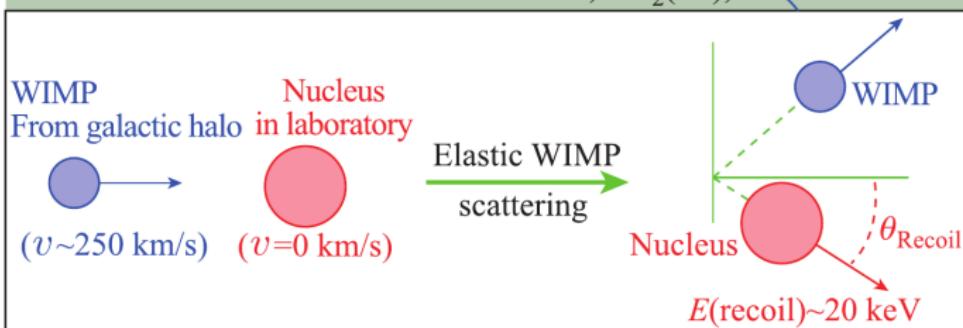
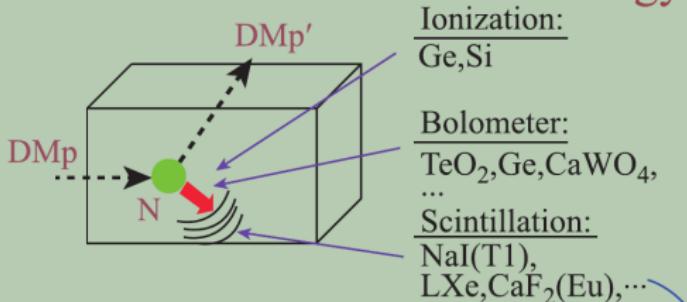


WIMP Scattering off Atomic Nuclei



Direct Detection

- Scatterings on nuclei
 - detection of nuclear recoil energy



[Bing-Lin Young, Front.Phys. 12, 121201 (2017)]

DM Velocity Distribution

 During the collapse process which formed the Galaxy, the velocities of DM particles were “thermalized” by **fluctuations** in the gravitational potential, and DM particles have a **Maxwell-Boltzmann velocity distribution** in the **Galactic rest frame**:

$$\tilde{f}(\tilde{\mathbf{v}}) d^3 \tilde{v} = \left(\frac{m_\chi}{2\pi k_B T} \right)^{3/2} \exp \left(-\frac{m_\chi \tilde{v}^2}{2k_B T} \right) d^3 \tilde{v} = \frac{e^{-\tilde{v}^2/v_0^2}}{\pi^{3/2} v_0^3} d^3 \tilde{v}, \quad v_0^2 \equiv \frac{2k_B T}{m_\chi}$$

$$\langle \tilde{\mathbf{v}} \rangle = \int \tilde{\mathbf{v}} \tilde{f}(\tilde{\mathbf{v}}) d^3 \tilde{v} = \mathbf{0}, \quad \langle \tilde{v}^2 \rangle = \int \tilde{v}^2 \tilde{f}(\tilde{\mathbf{v}}) d^3 \tilde{v} = \frac{3}{2} v_0^2$$

 **Speed distribution:** $\tilde{f}(\tilde{v}) \, d\tilde{v} = \frac{4\tilde{v}^2}{\sqrt{\pi}v_0^3} e^{-\tilde{v}^2/v_0^2} \, d\tilde{v}$

 For an **isothermal halo**, the **local value** of v_0 equals to the **rotational speed of the Sun**:

$$v_0 = v_\odot \simeq 220 \text{ km/s}$$

[Binney & Tremaine, *Galactic Dynamics*, Chapter 4]



Galactic disk & dark halo

[Credit: ESO/L. Calçada]

 **Velocity dispersion:** $\sqrt{\langle \tilde{v}^2 \rangle} = \sqrt{\frac{3}{2}} v_0 \simeq 270 \text{ km/s}$

Earth Rest Frame

The DM velocity distribution $f(v)$ seen by an **observer on the Earth** can be derived via **Galilean transformation**.

$$\tilde{\mathbf{v}} \equiv \mathbf{v} + \mathbf{v}_{\text{obs}}, \quad \mathbf{v}_{\text{obs}} \equiv \mathbf{v}_{\odot} + \mathbf{v}_{\oplus}$$

 **Velocity distribution:** $f(\mathbf{v}) = \tilde{f}(\mathbf{v} + \mathbf{v}_{\text{obs}})$

Speed distribution:

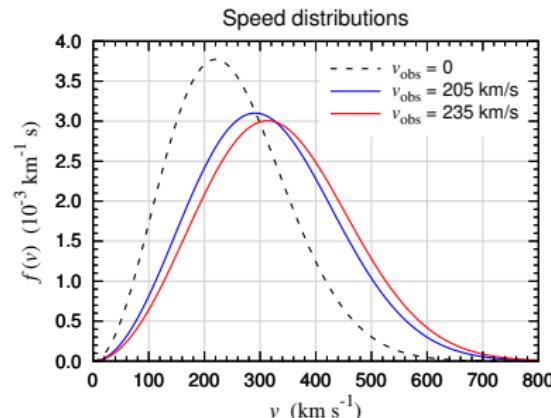
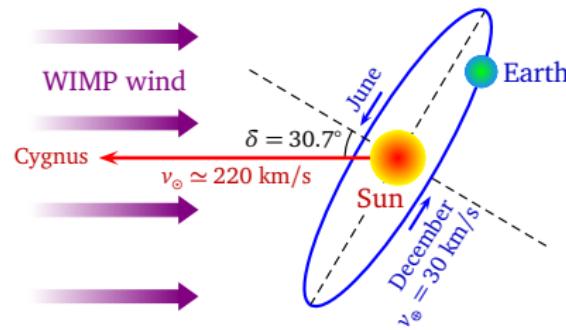
$$f(v) dv = \frac{4v^2}{\sqrt{\pi} v_0^3} \exp\left(-\frac{v^2 + \textcolor{blue}{v}_{\text{obs}}^2}{v_0^2}\right) \times \frac{\tilde{v}_0^2}{2vv_{\text{obs}}} \sinh\left(\frac{2vv_{\text{obs}}}{\tilde{v}_0^2}\right) dv$$

Since $v_{\oplus} \ll v_{\odot}$, we have ($\omega = 2\pi/\text{year}$)

$$v_{\text{obs}}(t) \simeq v_{\odot} + v_{\oplus} \sin \delta \cos[\omega(t - t_0)]$$

$$\simeq 220 \text{ km/s} + 15 \text{ km/s} \cdot \cos[\omega(t - t_0)]$$

👉 Annual modulation signal peaked on June 2



Nuclear Recoil



Energy conservation:

$$\frac{1}{2}m_\chi v^2 = \frac{1}{2}m_\chi v_\chi^2 + \frac{1}{2}m_A v_R^2$$



Momentum conservation:

$$m_\chi v = m_\chi v_\chi \cos \theta_\chi + m_A v_R \cos \theta_R$$

$$m_\chi v_\chi \sin \theta_\chi = m_A v_R \sin \theta_R$$



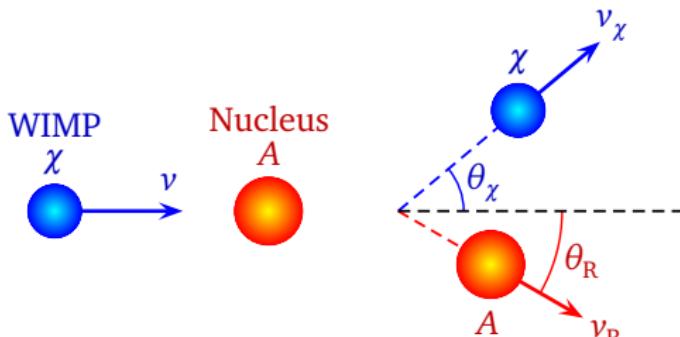
Recoil velocity $v_R = \frac{2m_\chi v \cos \theta_R}{m_\chi + m_A}$



Recoil momentum (momentum transfer) $q_R = m_A v_R = 2\mu_{\chi A} v \cos \theta_R$



Reduced mass of the χA system $\mu_{\chi A} \equiv \frac{m_\chi m_A}{m_\chi + m_A} = \begin{cases} m_A, & \text{for } m_\chi \gg m_A \\ \frac{m_\chi}{2}, & \text{for } m_\chi = m_A \\ m_\chi, & \text{for } m_\chi \ll m_A \end{cases}$



Forward scattering ($\theta_R = 0$) ➡ maximal momentum transfer $q_R^{\max} = 2\mu_{\chi A} v$



Nuclear Recoil



Energy conservation:

$$\frac{1}{2}m_\chi v^2 = \frac{1}{2}m_\chi v_\chi^2 + \frac{1}{2}m_A v_R^2$$



Momentum conservation:

$$m_\chi v = m_\chi v_\chi \cos \theta_\chi + m_A v_R \cos \theta_R$$

$$m_\chi v_\chi \sin \theta_\chi = m_A v_R \sin \theta_R$$



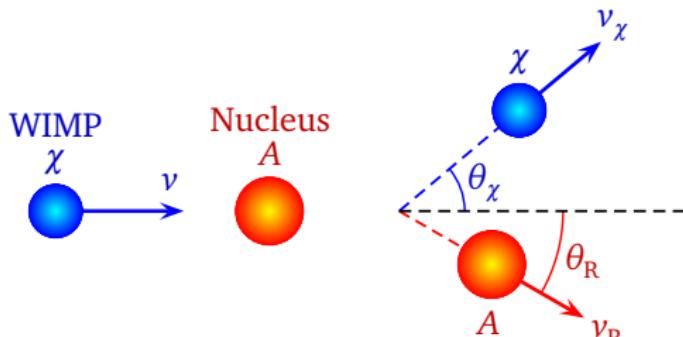
$$\text{Recoil velocity } v_R = \frac{2m_\chi v \cos \theta_R}{m_\chi + m_A}$$



Recoil momentum (momentum transfer) $q_R = m_A v_R = 2\mu_{\chi A} v \cos \theta_R$



$$\text{Kinetic energy of the recoiled nucleus } E_R = \frac{q_R^2}{2m_A} = \frac{2\mu_{\chi A}^2}{m_A} v^2 \cos^2 \theta_R$$



As $v \sim 10^{-3}c$, for $m_\chi = m_A \simeq 100$ GeV and $\theta_R = 0$,

$$q_R = m_\chi v \sim 100 \text{ MeV}, \quad E_R = \frac{1}{2}m_\chi v^2 \sim 50 \text{ keV}$$

Event Rate



Event rate per unit time per unit energy interval in direct detection experiments:

$$\frac{dR}{dE_R} = N_A \frac{\rho_{\oplus}}{m_{\chi}} \int_{v_{\min}}^{v_{\max}} d^3 v f(\mathbf{v}) v \frac{d\sigma_{\chi A}}{dE_R}$$



Astrophysics factors



Particle physics factors



Detector factors



N_A : target nucleus number



$\rho_{\oplus} \simeq 0.4 \text{ GeV/cm}^3$: DM **mass density** around the Earth



ρ_{\oplus}/m_{χ} is the DM particle **number density** around the Earth



$\sigma_{\chi A}$: DM-nucleus **scattering cross section**



Minimal velocity $v_{\min} = \left(\frac{m_A E_R^{\text{th}}}{2\mu_{\chi A}^2} \right)^{1/2}$: determined by the **detector threshold** of nuclear recoil energy, E_R^{th}



Maximal velocity v_{\max} : determined by the DM **escape velocity** v_{esc}



$v_{\text{esc}} \simeq 544 \text{ km/s}$ [Smith *et al.*, MNRAS 379, 755]

Cross Section Dependence on Nucleus Spin

 There are two kinds of DM-nucleus scattering

 **Spin-independent (SI)** cross section $\sigma_{\chi A}^{\text{SI}} \propto \mu_{\chi A}^2 [ZG_p + (A - Z)G_n]^2$

 **Spin-dependent (SD)** cross section $\sigma_{\chi A}^{\text{SD}} \propto \mu_{\chi A}^2 \frac{J_A + 1}{J_A} (S_p^A G'_p + S_n^A G'_n)^2$

 **Nucleus properties:** mass number A , atomic number Z , spin J_A , expectation value of the proton (neutron) spin content in the nucleus S_p^A (S_n^A)

 $G_p^{(t)}$ and $G_n^{(t)}$: DM effective couplings to the proton and the neutron

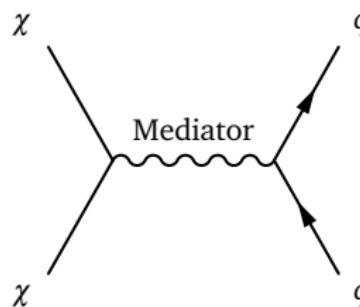
- $Z \simeq A/2$  $\sigma_{\chi A}^{\text{SI}} \propto A^2 [(G_p + G_n)/2]^2$
! Strong coherent enhancement for heavy nuclei

- Spins of nucleons tend to cancel out among themselves

- $S_N^A \simeq 1/2$ ($N = p$ or n) for a nucleus with an odd number of N
- $S_N^A \simeq 0$ for a nucleus with an even number of N

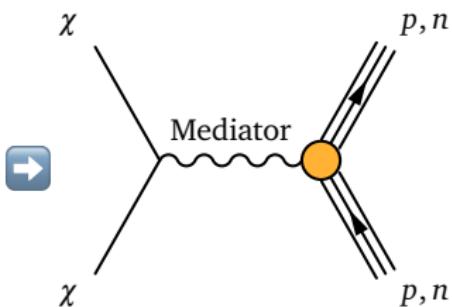
Three Levels of Interaction

DM-parton interaction



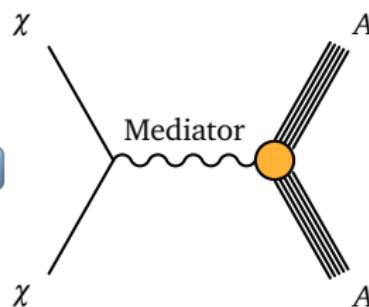
$$\mathcal{M}(\chi q \rightarrow \chi q)$$

DM-nucleon interaction



$$\mathcal{M}(\chi N \rightarrow \chi N)$$

DM-nucleus interaction



$$\mathcal{M}(\chi A \rightarrow \chi A)$$



 As a variety of **target nuclei** are used in **direct detection experiments**, results are usually compared with each other at the **DM-nucleon level**



 The **DM-nucleon level** is related to the **DM-parton level** via **form factors**, which describe the **probabilities** of finding **partons** inside **nucleons**

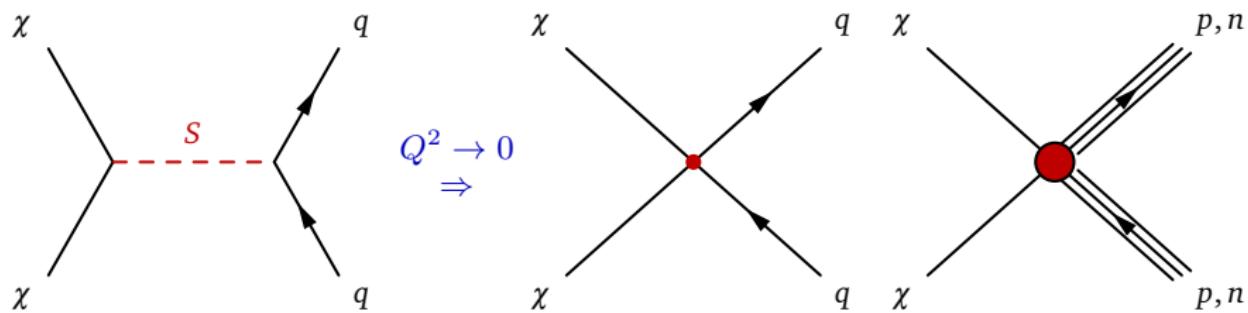


 The relevant partons involve not only valence quarks, but also sea quarks and gluons

Zero-Momentum Transfer Limit

🛵 As the **momentum transfer** is typically **much smaller** than the **underlying energy scale** (e.g., **mediator mass**), the **zero-momentum transfer limit** is a good approximation for calculation

🚲 In this limit, the mediator field can be integrated out, and the interaction can be described by **effective operators** in **effective field theory**



🛵 **Scalar mediator propagator:** $\frac{i}{Q^2 - m_S^2} \Rightarrow -\frac{i}{m_S^2}$

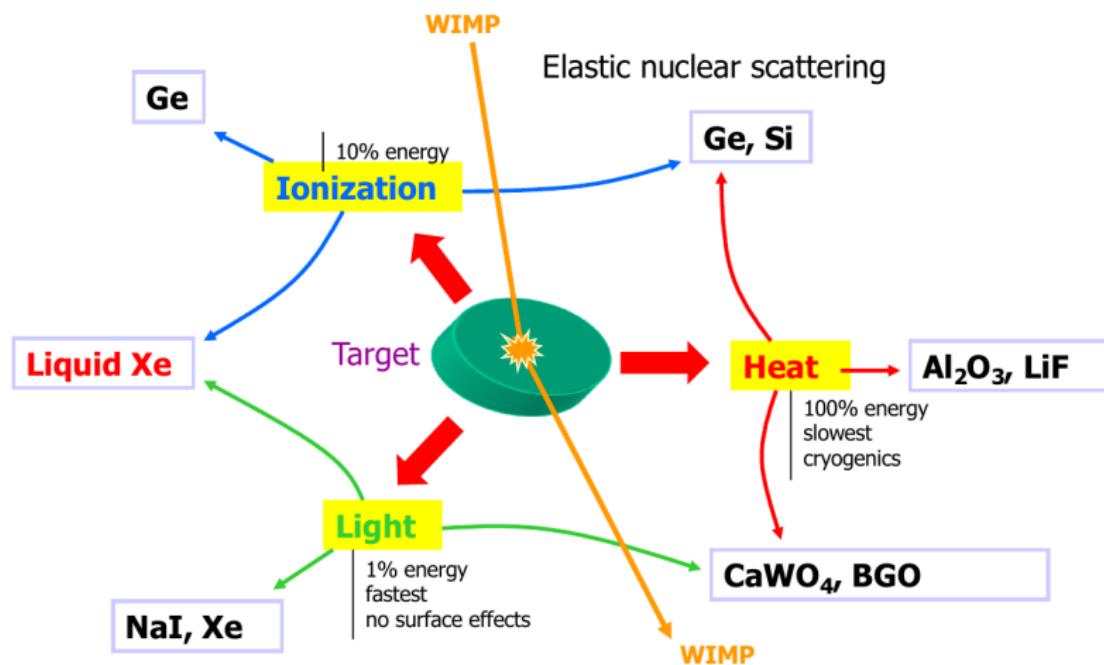
🛴 **Lagrangian:** $\mathcal{L}_{\text{int}} = g_\chi S \bar{\chi} \chi + g_q S \bar{q} q \Rightarrow \mathcal{L}_{\text{eff}} = G_{\text{eff}} \bar{\chi} \chi \bar{q} q, \quad G_{\text{eff}} = \frac{g_\chi g_q}{m_S^2}$

Effective Operators for DM-quark Interactions

	Spin-1/2 DM	Spin-0 DM
SI	$\bar{\chi}\chi\bar{q}q, \bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q$	$\chi^*\chi\bar{q}q, (\chi^*i\overleftrightarrow{\partial}^\mu\chi)\bar{q}\gamma_\mu q$
SD	$\bar{\chi}\gamma^\mu\gamma_5\chi\bar{q}\gamma_\mu\gamma_5 q, \bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu} q$	
$\sigma_{\chi N}$ $\propto Q^2 $	$\bar{\chi}i\gamma_5\chi\bar{q}i\gamma_5 q, \bar{\chi}\chi\bar{q}i\gamma_5 q$ $\bar{\chi}i\gamma_5\chi\bar{q}q, \bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu\gamma_5 q$ $\bar{\chi}\gamma^\mu\gamma_5\chi\bar{q}\gamma_\mu q, \epsilon^{\mu\nu\rho\sigma}\bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\rho\sigma} q$	$\chi^*\chi\bar{q}i\gamma_5 q$ $(\chi^*i\overleftrightarrow{\partial}^\mu\chi)\bar{q}\gamma_\mu\gamma_5 q$
	Spin-3/2 DM	Spin-1 DM
SI	$\bar{\chi}^\mu\chi_\mu\bar{q}q, \bar{\chi}^\nu\gamma^\mu\chi_\nu\bar{q}\gamma_\mu q$	$\chi_\mu^*\chi^\mu\bar{q}q, (\chi_\nu^*i\overleftrightarrow{\partial}^\mu\chi^\nu)\bar{q}\gamma_\mu q$
SD	$\bar{\chi}^\nu\gamma^\mu\gamma_5\chi_\nu\bar{q}\gamma_\mu\gamma_5 q, \bar{\chi}^\rho\sigma^{\mu\nu}\chi_\rho\bar{q}\sigma_{\mu\nu} q$ $i(\bar{\chi}^\mu\chi^\nu - \bar{\chi}^\nu\chi^\mu)\bar{q}\sigma_{\mu\nu} q$	$i(\chi_\mu^*\chi_\nu - \chi_\nu^*\chi_\mu)\bar{q}\sigma^{\mu\nu} q$ $\epsilon^{\mu\nu\rho\sigma}(\chi_\mu^*\overleftrightarrow{\partial}_\nu\chi_\rho)\bar{q}\gamma_\sigma\gamma_5 q$
$\sigma_{\chi N}$ $\propto Q^2 $	$\bar{\chi}^\mu i\gamma_5\chi_\mu\bar{q}i\gamma_5 q, \bar{\chi}^\mu\chi_\mu\bar{q}i\gamma_5 q$ $\bar{\chi}^\mu i\gamma_5\chi_\mu\bar{q}q, \bar{\chi}^\nu\gamma^\mu\chi_\nu\bar{q}\gamma_\mu\gamma_5 q$ $\bar{\chi}^\mu\gamma^\mu\gamma_5\chi_\nu\bar{q}\gamma_\mu q, \epsilon^{\mu\nu\rho\sigma}i(\bar{\chi}_\mu\chi_\nu - \bar{\chi}_\nu\chi_\mu)\bar{q}\sigma_{\rho\sigma} q$ $\epsilon^{\mu\nu\rho\sigma}\bar{\chi}^\alpha\sigma_{\mu\nu}\chi_\alpha\bar{q}\sigma_{\rho\sigma} q, (\bar{\chi}^\mu\gamma_5\chi^\nu - \bar{\chi}^\nu\gamma_5\chi^\mu)\bar{q}\sigma_{\mu\nu} q$ $\epsilon^{\mu\nu\rho\sigma}(\bar{\chi}_\mu\gamma_5\chi_\nu - \bar{\chi}_\nu\gamma_5\chi_\mu)\bar{q}\sigma_{\rho\sigma} q$	$\chi_\mu^*\chi^\mu\bar{q}i\gamma_5 q$ $(\chi_\nu^*i\overleftrightarrow{\partial}^\mu\chi^\nu)\bar{q}\gamma_\mu\gamma_5 q$ $\epsilon^{\mu\nu\rho\sigma}(\chi_\mu^*\overleftrightarrow{\partial}_\nu\chi_\rho)\bar{q}\gamma_\sigma q$ $\epsilon^{\mu\nu\rho\sigma}i(\chi_\mu^*\chi_\nu - \chi_\nu^*\chi_\mu)\bar{q}\sigma_{\rho\sigma} q$

[Zheng, **ZHY**, Shao, Bi, Li, Zhang, arXiv:1012.2022, NPB;
ZHY, Zheng, Bi, Li, Yao, Zhang, arXiv:1112.6052, NPB]

Technologies and Detector Material

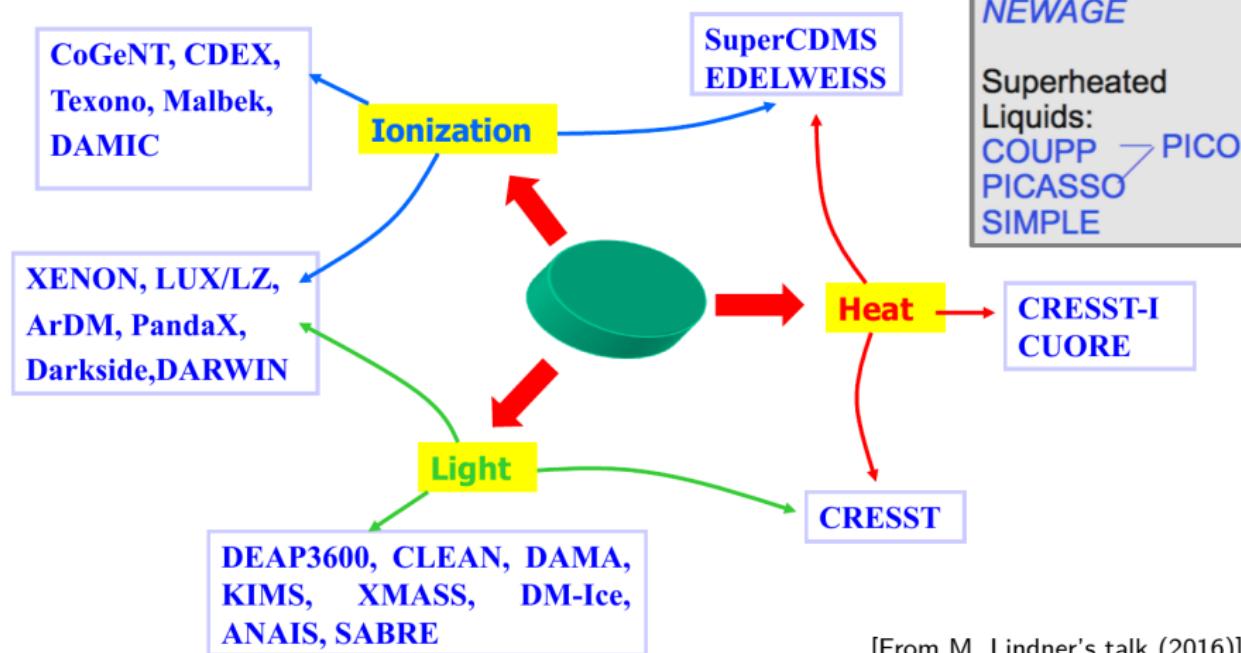


[From M. Lindner's talk (2016)]



Technologies and Detector Material

Detection methods: Crystals (NaI, Ge, Si),
Cryogenic Detectors, Liquid Noble Gases



[From M. Lindner's talk (2016)]



Example: Dual-phase Xenon Time Projection Chamber

↑ Upper: Xenon gas

↓ Lower: Liquid Xenon

UV scintillation photons recorded by photomultiplier tube (PMT) arrays on top and bottom

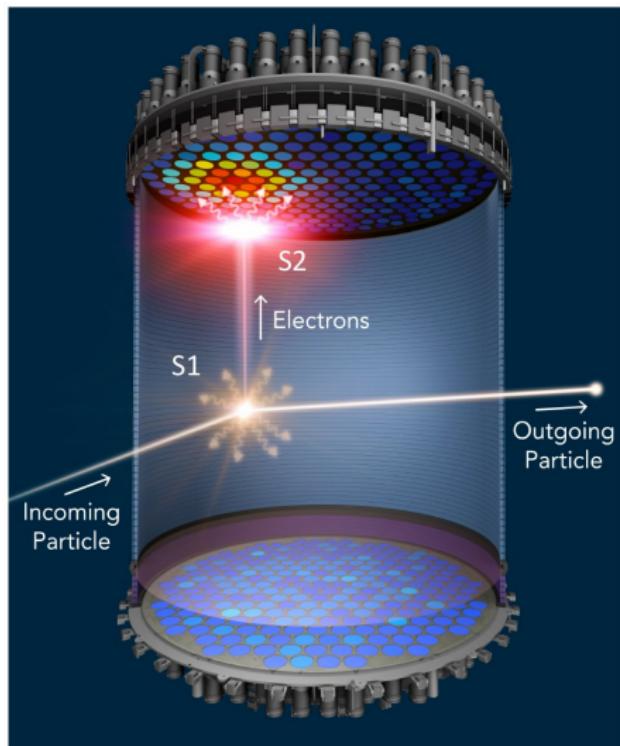
① Primary scintillation (S1)

Scintillation light promptly emitted from the interaction vertex

② Secondary scintillation (S2)

Ionization electrons emitted from the interaction are drifted to the surface and into the gas, where they emit proportional scintillation light

Experiment: PandaX, XENON, LZ

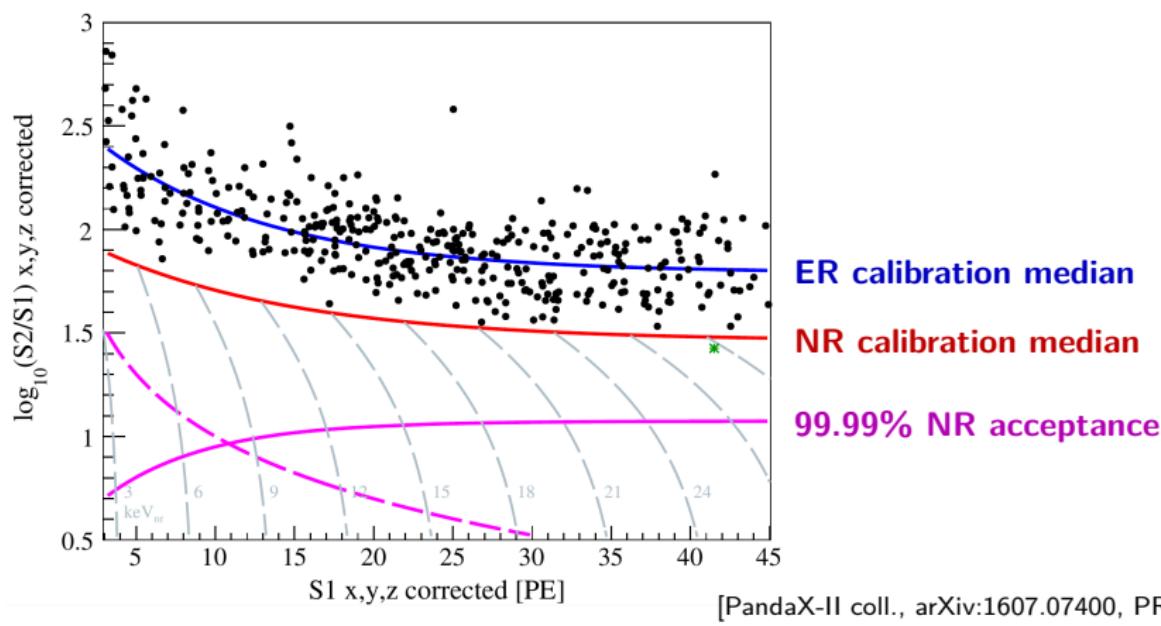


[From A. Cottle's talk (2017)]

PandaX-II Real Data: S1 versus S2

S1 and **S2** are characterized by numbers of **photoelectrons** (PEs) in **PMTs**

The γ background, which produces **electron recoil (ER)** events, can be distinguished from **nuclear recoil (NR)** events using the **S2-to-S1 ratio**





Backgrounds



Background suppression:

- Deep underground

- Shielded environments



Cosmogenic backgrounds

- Cosmic rays and secondary reactions

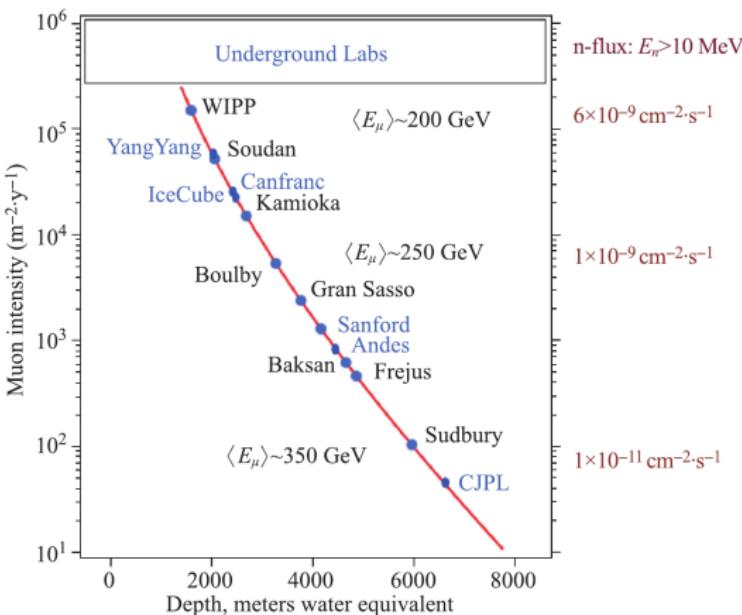
- Activation products in shields and detectors



Radiogenic backgrounds

- External natural radioactivity: walls, structures of site, radon

- Internal radioactivity: shield and construction materials, detector contamination in manufacture, naturally occurring radio-isotopes in target material



[From P. Cushman's talk (2014)]



China JinPing Underground Laboratory (CJPL)



[Yue et al., arXiv:1602.02462]

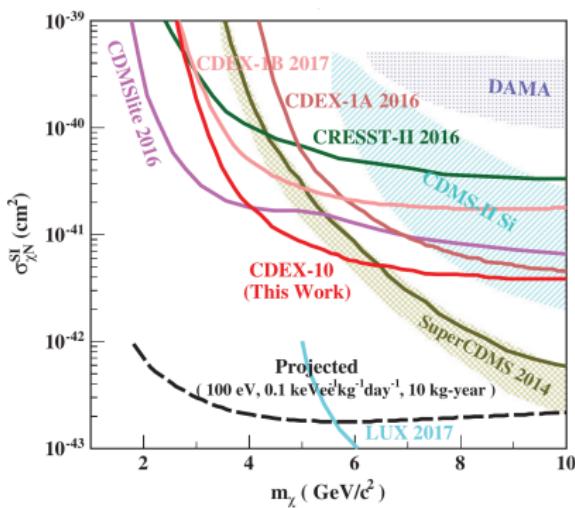


Experiments: CDEX, PandaX

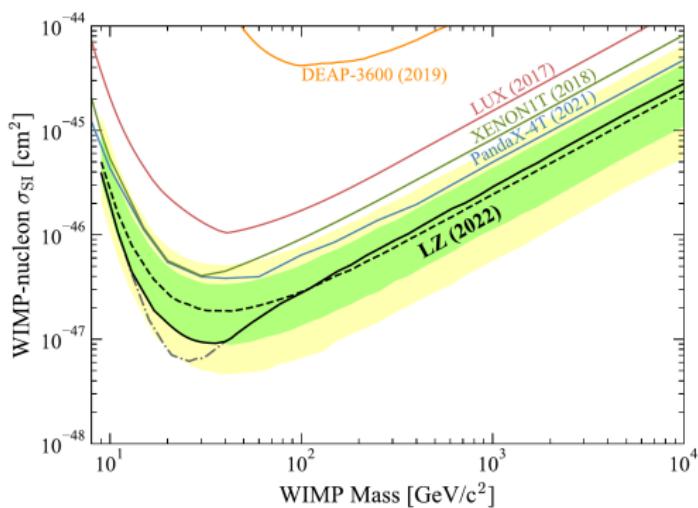


Exclusion Limits for SI Scattering

- For **isospin-conserving SI scattering**, protons and **neutrons** can be treated as the **same species**, i.e., "**nucleons**"
- Lower thresholds** and **lighter targets** are needed to probe the **low mass regime**
- It requires **more exposure**, **heavier targets**, and **fewer backgrounds** to explore the **high mass regime**



[CDEX Coll., arXiv:1802.09016, PRL]



[LZ Coll., arXiv:2207.03764, PRL]

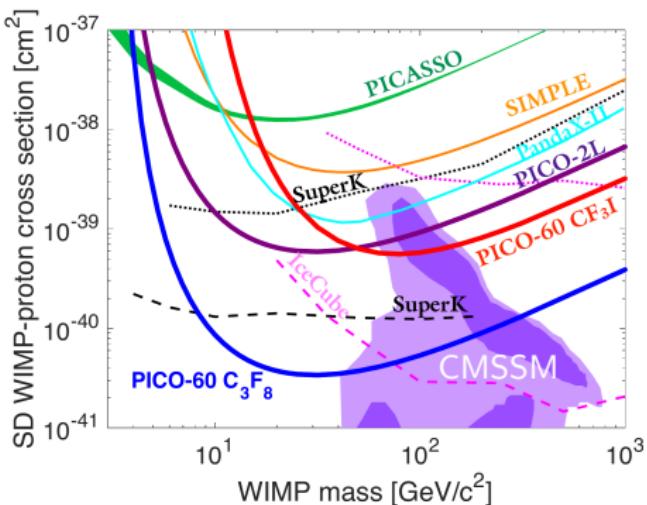
Exclusion Limits for SD Scattering



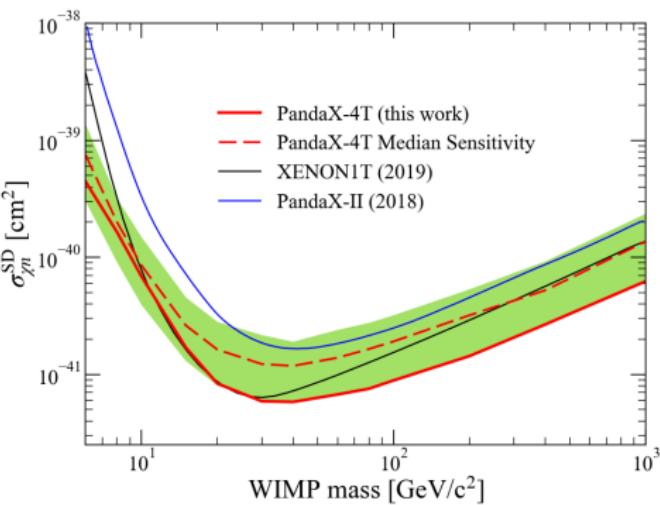
 For **SD scattering**, specific detection material usually has **very different** sensitivities to **DM-proton** and **DM-neutron** cross sections



As there is no coherent enhancement for **SD scattering**, the sensitivity is **lower** than the **SI case** by **several orders of magnitude**



[PICO coll., arXiv:1702.07666, PRL]

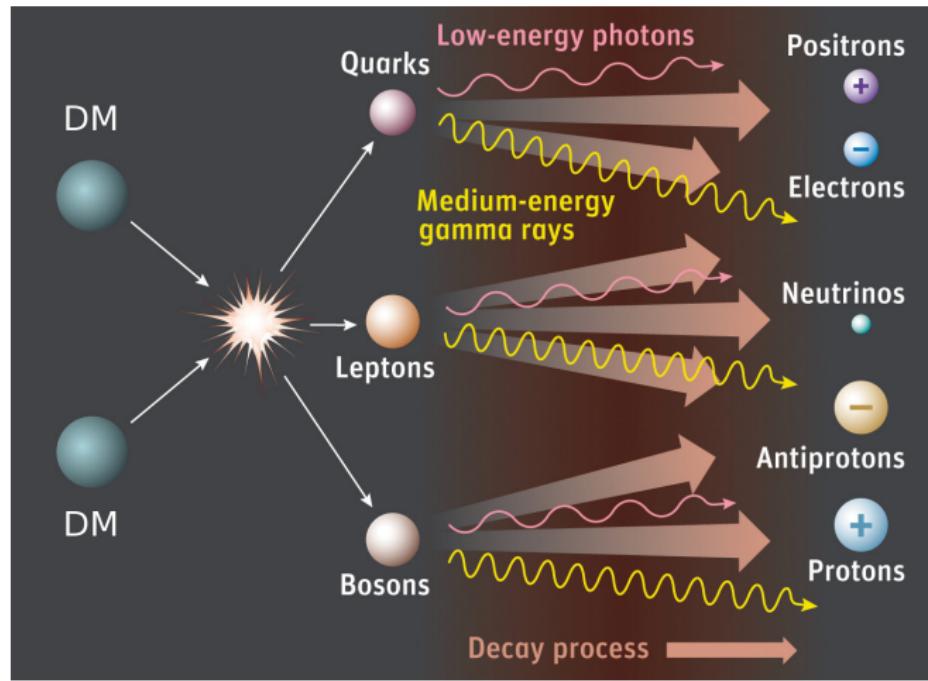


[PandaX-4T coll., arXiv:2208.03626, PLB]



Indirect Detection

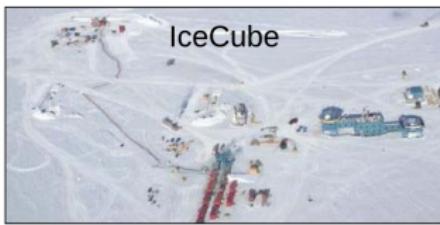
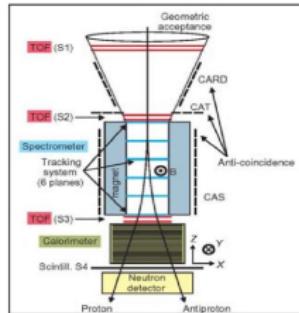
🔭 Indirect detection looks for **stable products**, e.g., γ rays, cosmic rays, and neutrinos, from **DM annihilation** or **decay** (if DM is not totally stable) **in space**



Indirect Detection Experiments



PAMELA



Dark Matter Source Function

 **Particle number per unit time per unit volume per unit energy interval** of a **stable species** ($\gamma, e^\pm, \nu, p, \bar{p}, \dots$) produced from **DM annihilation** or **decay**:

$$Q_{\text{ann}}(\mathbf{x}, E) = \frac{\langle \sigma_{\text{ann}} v \rangle_{\text{tot}}}{2m_\chi^2} \rho^2(\mathbf{x}) \sum_i F_i \left(\frac{dN}{dE} \right)_i \quad (\text{annihilation})$$

$$Q_{\text{dec}}(\mathbf{x}, E) = \frac{1}{\tau_\chi m_\chi} \rho(\mathbf{x}) \sum_i B_i \left(\frac{dN}{dE} \right)_i \quad (\text{decay})$$

Astrophysics factors

Particle physics factors

 $\rho(\mathbf{x})$: **DM mass density** at the source position \mathbf{x}

 $(dN/dE)_i$: **number per unit energy interval** from a single event in the **channel i**

 $\langle \sigma_{\text{ann}} v \rangle_{\text{tot}}$: thermal average of the total **annihilation cross section** multiplied by the relative velocity between the two incoming DM particles

 $F_i \equiv \langle \sigma_{\text{ann}} v \rangle_i / \langle \sigma_{\text{ann}} v \rangle_{\text{tot}}$: **branching fraction** of the annihilation channel i

 $\tau_\chi \equiv 1/\Gamma_\chi$: mean **lifetime** of the DM particle

 $B_i \equiv \Gamma_i / \Gamma_\chi$: **branching ratio** of the decay channel i

γ rays from DM: Continuous Spectrum

DM annihilation or decay into e^+e^- , $\mu^+\mu^-$, $\tau^+\tau^-$, $q\bar{q}$, W^+W^- , Z^0Z^0 , h^0h^0



γ -ray emissions from final state radiation or particle decays

Cutoff energy:

- m_χ for DM annihilation
- $m_\chi/2$ for DM decay

More promising to look at

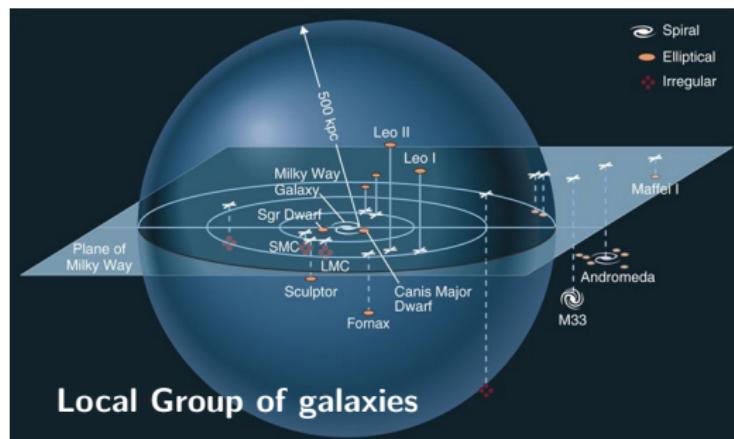
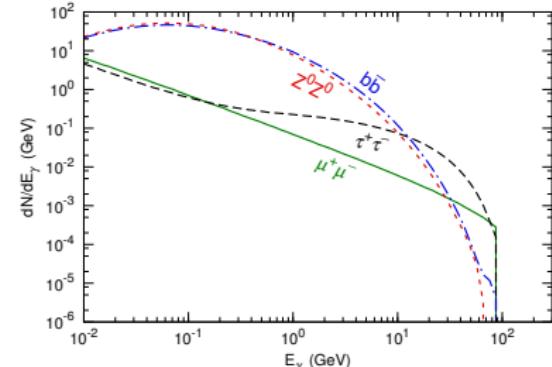
DM-dominated regions:

Galactic Center

Galactic halo

dwarf galaxies

clusters of galaxies

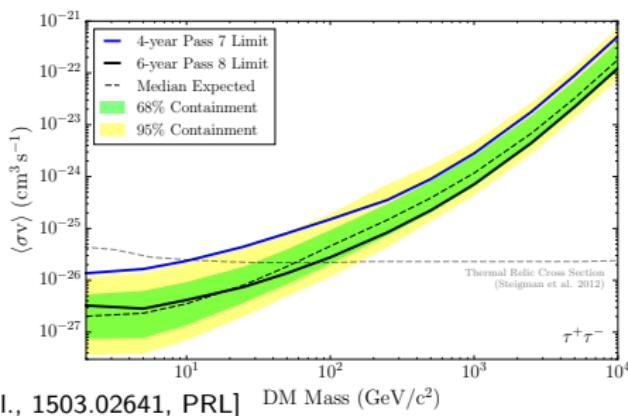
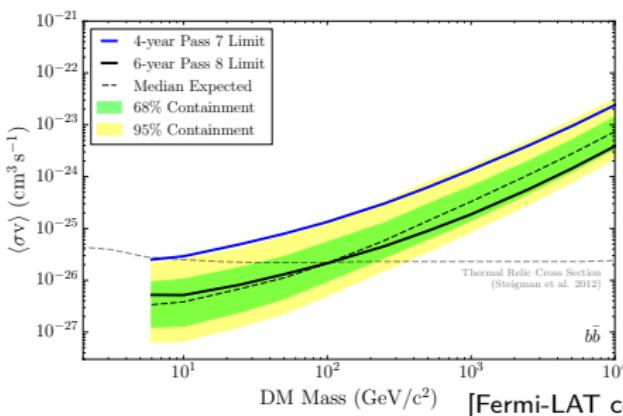
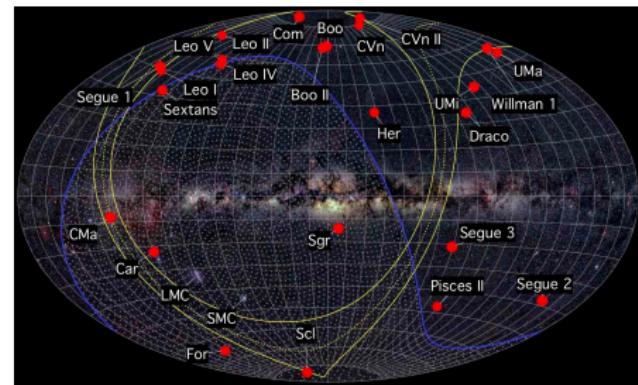




γ -ray Observation of Dwarf Galaxies

The space experiment **Fermi-LAT** searched for γ -ray emissions from **dwarf spheroidal satellite galaxies** of the **Milky Way** and found **no significant signal**

Based on the 6-year data, **upper limits** on the **DM annihilation cross section** are given

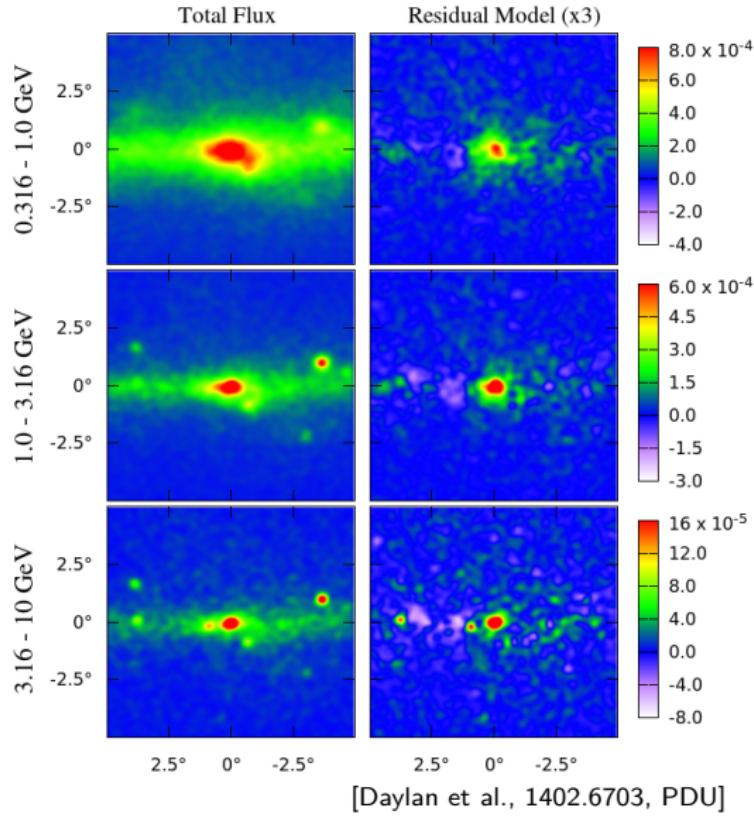


GeV Excess at the Galactic Center?

💡 Since 2009, several research groups reported an **excess** of **continuous spectrum γ -rays** in the **Fermi-LAT data** after **subtracting** well-known **astrophysical backgrounds**,

🕯 It locates in the **Galactic Center (GC)** region and peaks at a few GeV

📘 Left: raw γ -ray maps
📘 Right: residual maps
subtracting the Galactic diffuse model, the 20 cm template, point sources, and the isotropic template



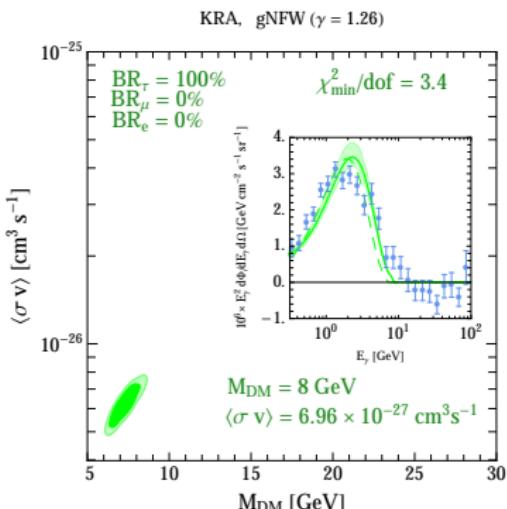
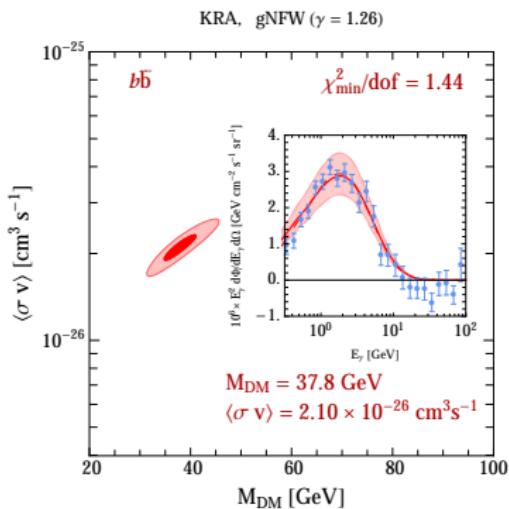
Interpretation with Dark Matter Annihilation

DM annihilation into $b\bar{b}$

$$\langle \sigma_{\text{ann}} v \rangle \sim 10^{-26} \text{ cm}^3 \text{ s}^{-1}$$

DM annihilation into $\tau^+\tau^-$

$$m_\chi \sim 9 \text{ GeV}$$



[Cirelli et al., arXiv:1407.2173, JCAP]



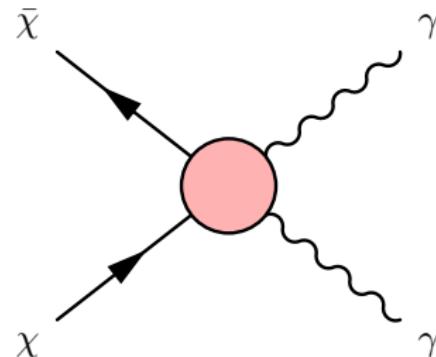
γ rays from DM: Line Spectrum

⌚ DM particles should **not have electric charge** and thus **not directly couple to photons**

👓 Nonetheless, DM particles **may couple to photons** via **high-order loop diagrams**



🚧 Highly suppressed: the $\chi\bar{\chi} \rightarrow \gamma\gamma$ branching fraction may be only $\sim 10^{-4} - 10^{-1}$





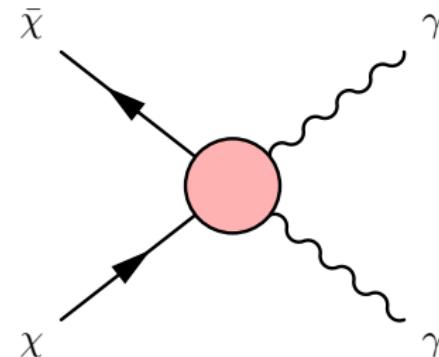
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👢 For **nonrelativistic** DM particles in space, the photons produced in $\chi\bar{\chi} \rightarrow \gamma\gamma$ would be **mono-energetic**



👠 A γ -ray line at energy $\sim m_\chi$
("smoking gun" for DM particles)

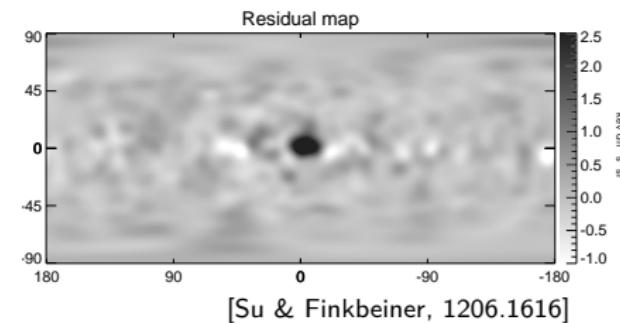
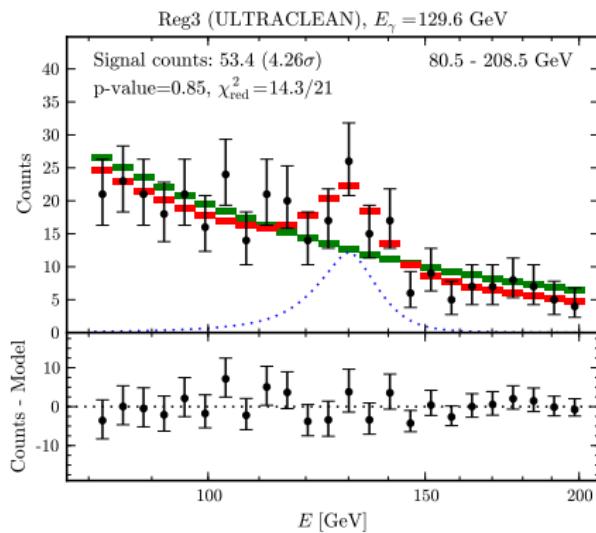




A γ -ray Line Signal at the Galactic Center?

Using the **3.7-year Fermi-LAT γ -ray data**, several analyses showed that there might be evidence of a **monochromatic γ -ray line** at **energy ~ 130 GeV**, originating from the **Galactic Center region** (about $3\text{--}4\sigma$)

It may be explained by **DM annihilation** with $\langle \sigma_{\text{ann}} v \rangle \sim 10^{-27} \text{ cm}^3 \text{ s}^{-1}$

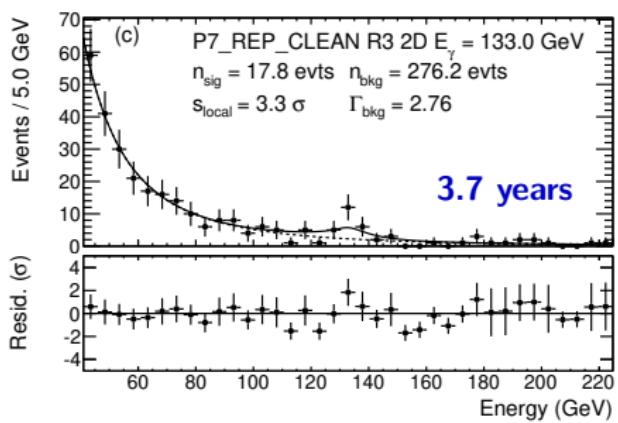


[Weniger, 1204.2797, JCAP]

Fermi-LAT Official Results: Not Confirmed with More Data

3.7-year data

⚠ The most significant fit occurred at $E_\gamma = 133$ GeV and had a **local significance** of 3.3σ , translating to a global significance of 1.6σ



[Fermi-LAT Coll., 1305.5597, PRD]

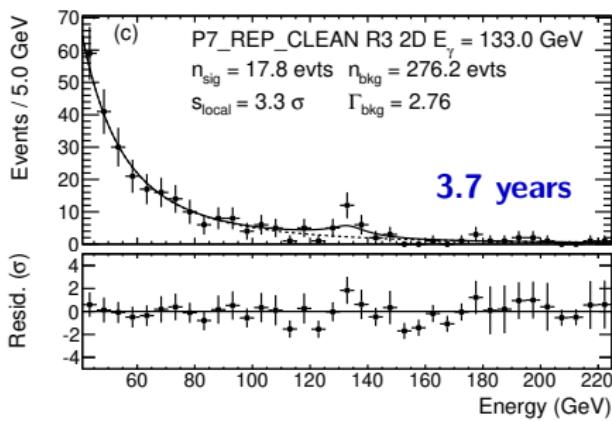
Fermi-LAT Official Results: Not Confirmed with More Data

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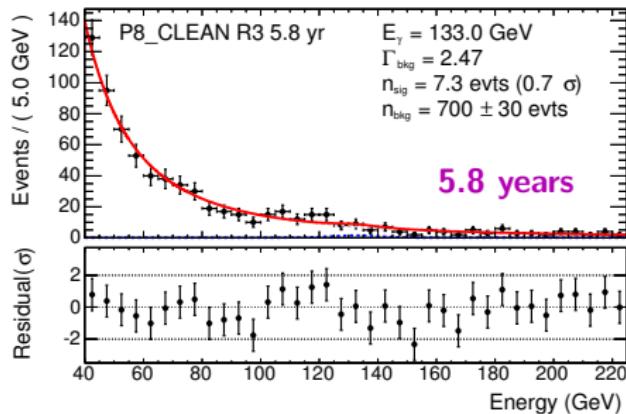
⚠ The most significant fit occurred at $E_\gamma = 133$ GeV and had a **local significance** of 3.3σ , translating to a global significance of 1.6σ

● 5.8-year data

 The local significance has dropped to 0.72σ



[Fermi-LAT Coll., 1305.5597, PRD]



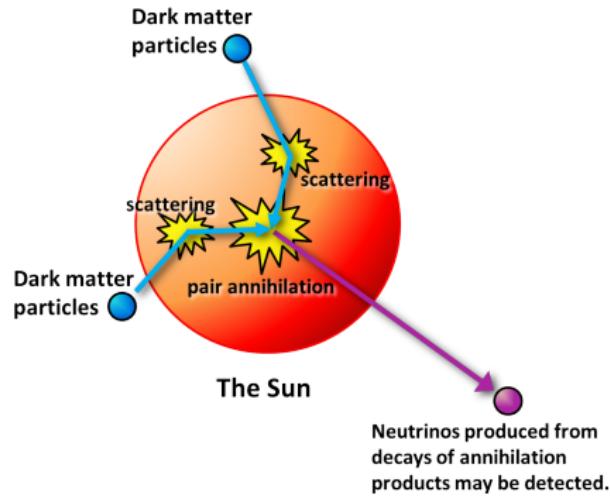
[Fermi-LAT Coll., 1506.00013, PRD]

Neutrinos from DM

🕸 Dark matter may be **captured** and **accumulated** at the **core** of the **Sun** ☀ (or the **Earth** 🌎), producing **high energy neutrinos** that could **freely go out**

🐜 Change Rate of the **number** of **DM particles** in the **Sun**:

$$\frac{dN_\chi}{dt} = C_\odot(\sigma_{\chi H}, \sigma_{\chi He}) - A_\odot(\sigma_{\text{ann}}) N_\chi^2$$

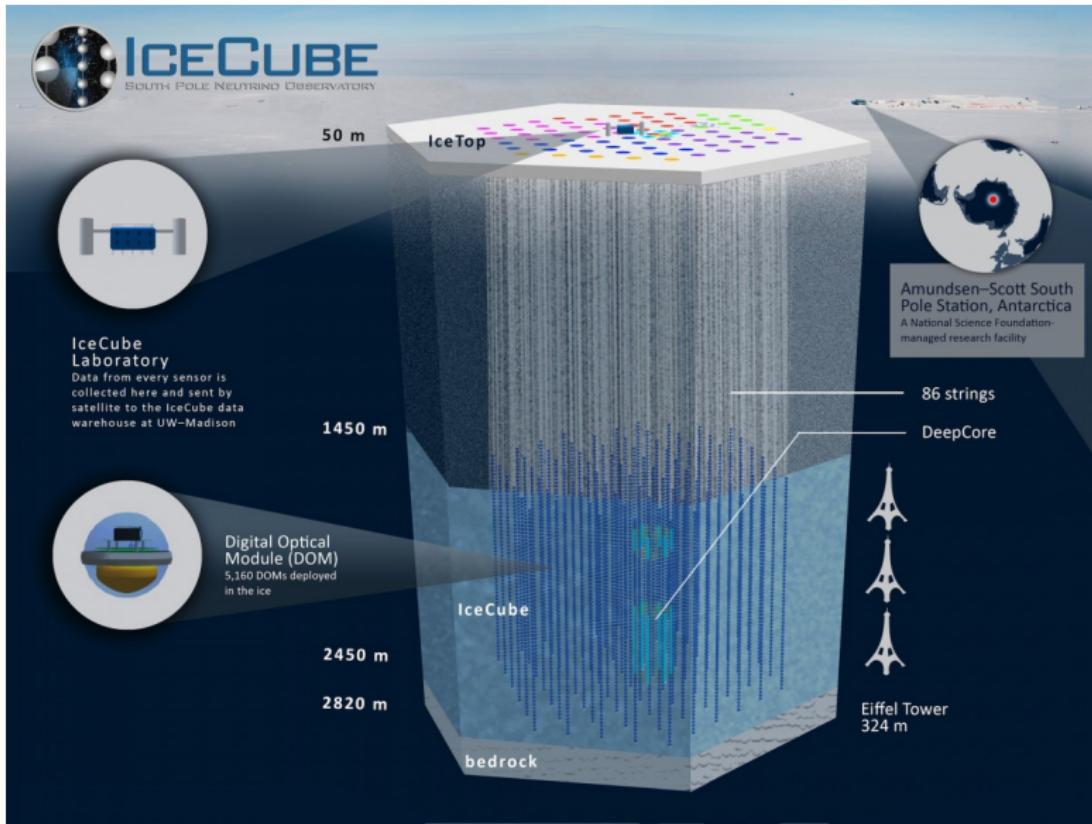


🐞 Capture rate C_\odot depends on **DM scattering** on Hydrogen and Helium

🐞 Annihilation rate $A_\odot = \langle \sigma_{\text{ann}} v \rangle / V_{\text{eff}}$ depends on **DM annihilation** as well as the effective volume of the solar core

🐌 The **age** of the **Sun** is **long enough** (~ 4.6 billion years) to make the **capture** and **annihilation** processes reach **equilibrium**: $dN_\chi/dt = 0$

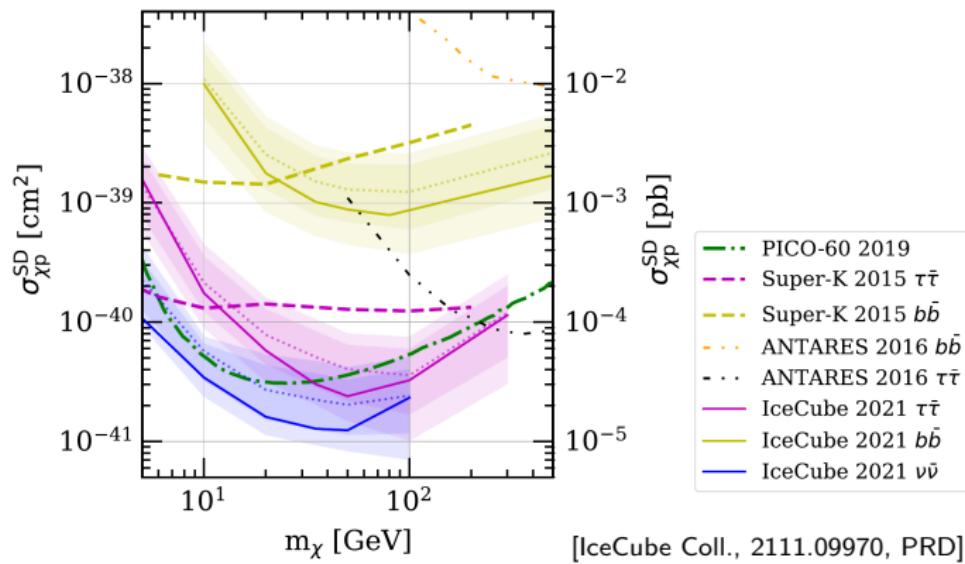
IceCube: South Pole Neutrino Observatory



Searches for Neutrinos from DM Annihilation within the Sun

 Using the **7-year IceCube data**, no significant detection of **neutrinos** with energies $< 500 \text{ GeV}$ from **DM annihilation** at the **solar core** is found

Assuming **equilibrium** in the **capture** and **annihilation** processes, the constraints can be converted to those on the **DM scattering cross section**

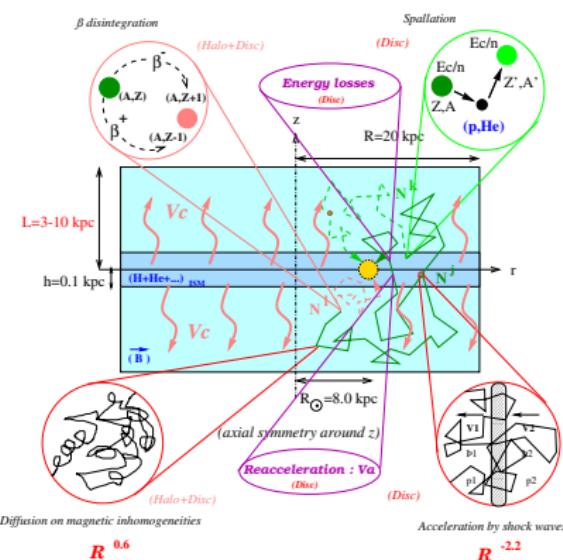
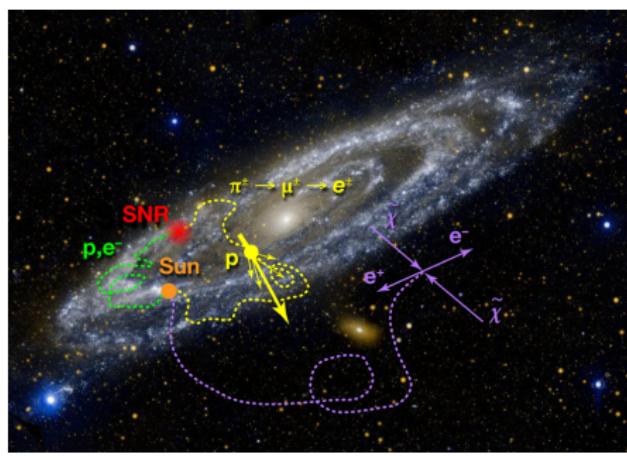




Cosmic Rays from DM

After produced in sources, **Galactic cosmic rays** diffuse in the interstellar space, suffering from several **propagation effects** before they arrive at the **Earth**: **diffusion**, energy losses, **convection**, **reacceleration**, **spallation**, ...

Unlike γ rays and neutrinos, cosmic rays typically **do not contain direction information** of their **sources**

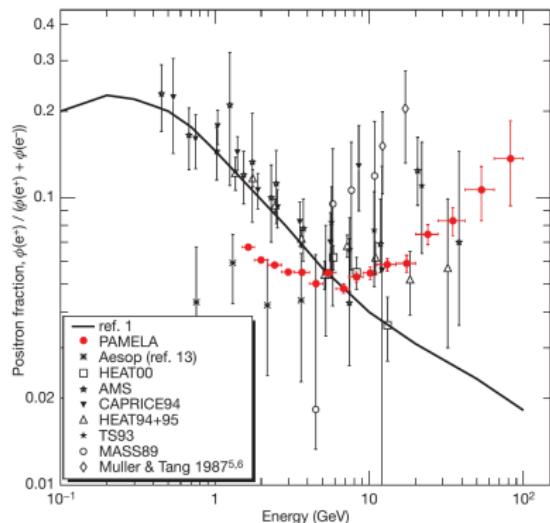


[Maurin et al., astro-ph/0212111]

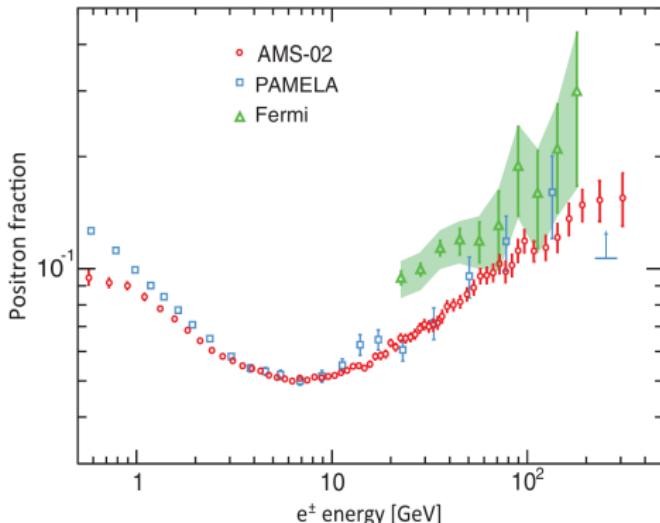
Cosmic-ray Positron Excess

 In 2008, the **PAMELA** experiment found an **unexpected increase** in the **cosmic-ray positron fraction** with $E \geq 10$ GeV

 In 2013, the **AMS-02** experiment **confirmed** such a **positron excess**



[PAMELA Coll., 0810.4995, Nature]



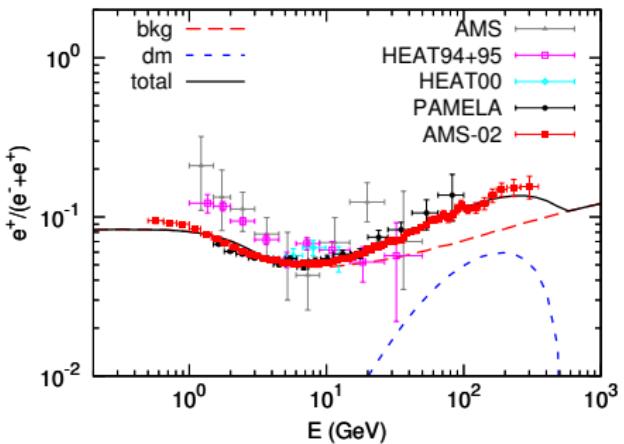
[AMS Coll., PRL 110, 141102 (2013)]



Interpretation: Dark Matter vs Pulsar



Interpretation with Galactic
DM annihilation into $\tau^+\tau^-$

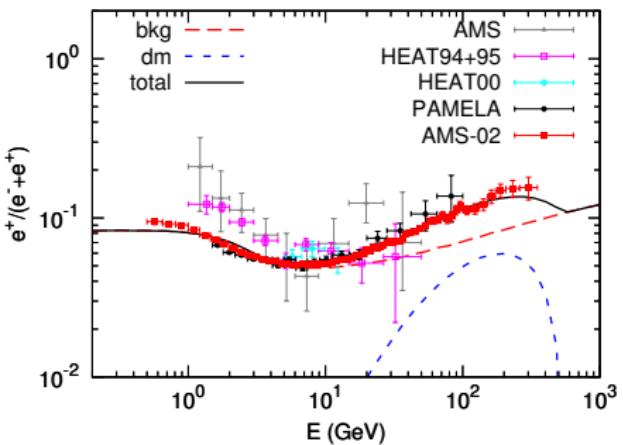


[Yuan, Bi, et al., 1304.1482, APP]

Interpretation: Dark Matter vs Pulsar



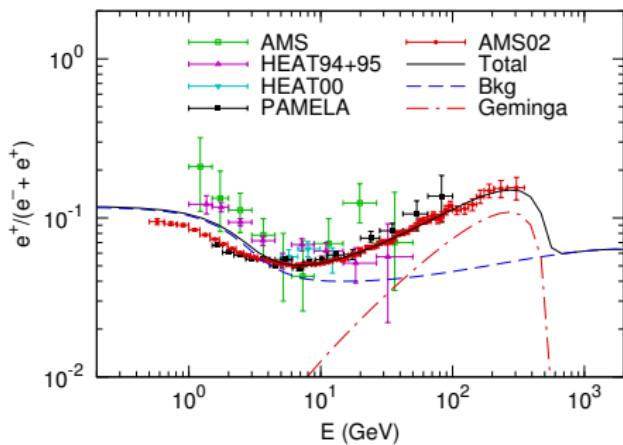
Interpretation with Galactic
DM annihilation into $\tau^+\tau^-$



[Yuan, Bi, et al., 1304.1482, APP]



Interpretation with the
nearby pulsar Geminga



[Yin, ZHY, Yuan, Bi, 1304.4128, PRD]



Past and Current High Energy Colliders

● TEVATRON: $p\bar{p}$ collider, 1987–2011

Circumference: 6.28 km

Collision energy: $\sqrt{s} = 1.96$ TeV

Luminosity: $\mathcal{L} \sim 4.3 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$

Detectors: CDF, DØ

● LEP: e^+e^- collider, 1989–2000

Circumference: 26.66 km

Collision energy: $\sqrt{s} = 91\text{--}209$ GeV

Luminosity: $\mathcal{L} \sim (2\text{--}10) \times 10^{31} \text{ cm}^{-2} \text{ s}^{-1}$

Detectors: ALEPH, DELPHI, OPAL, L3

● LHC: pp ($p\text{Pb}$, PbPb) collider, 2009–Now

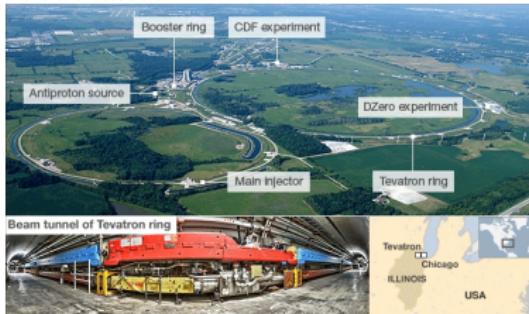
Circumference: 26.66 km

Collision energy: $\sqrt{s} = 7, 8, 13, 14$ TeV

Luminosity: $\mathcal{L} \sim (1\text{--}5) \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$

Detectors: ATLAS, CMS, ALICE, LHCb

The Tevatron accelerator



Source: Fermilab

Large Hadron Collider



Future Projects

- **ILC**: International Linear Collider

e^+e^- collider, $\sqrt{s} = 0.25-1$ TeV

$$\mathcal{L} \sim 1.5 \times 10^{34} \text{ cm}^{-2} \text{s}^{-1}$$

Detectors: SiD, ILD

- **CEPC**: Circular Electron-Positron Collider (China)

e^+e^- collider, $\sqrt{s} \approx 91\text{--}350$ GeV

$$f \approx (0.5-100) \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$

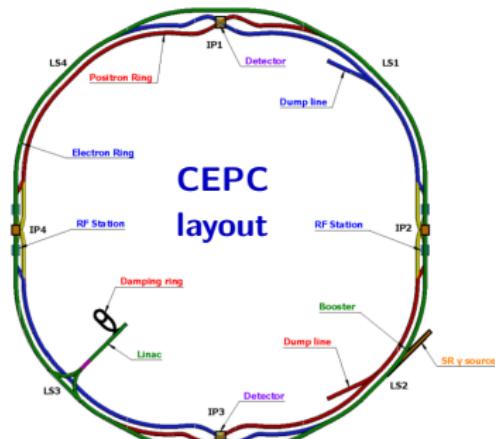
- **SPPC**: Super Proton-Proton Collider (China)

pp collider, $\sqrt{s} \sim 50\text{--}70$ TeV, $\mathcal{L} \sim 2.15 \times 10^{35}$ cm $^{-2}$ s $^{-1}$

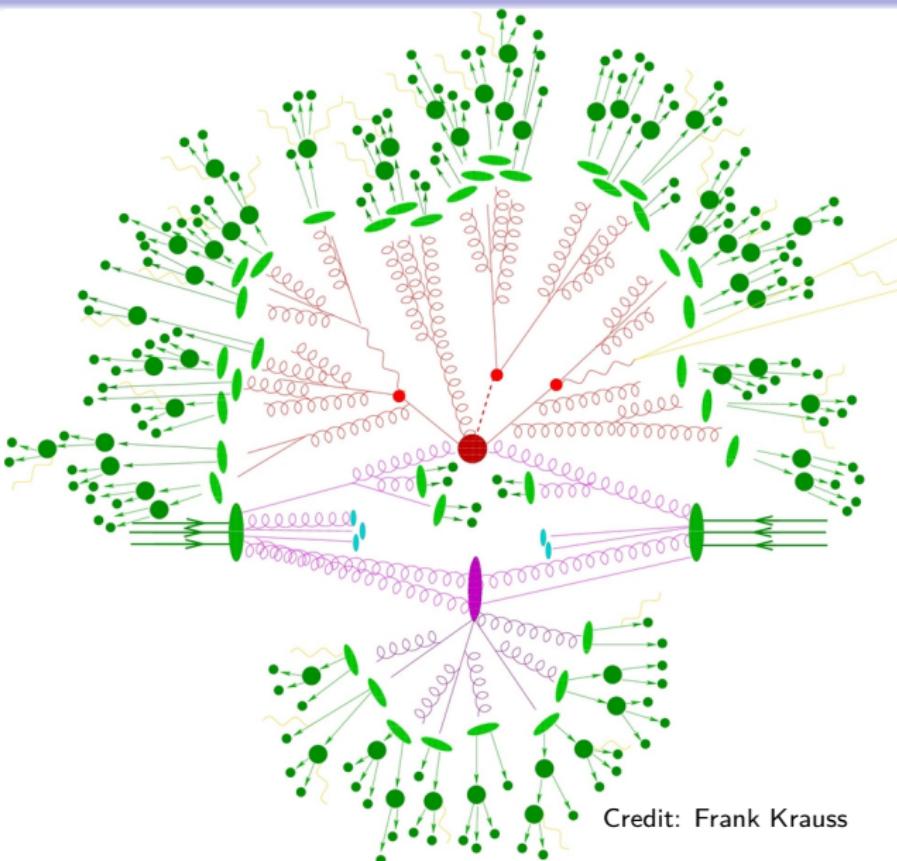
- ## ● FCC: Future Circular Collider (CERN)

- **FCC-ee**: e^+e^- collider, $\sqrt{s} \sim 91\text{--}350$ GeV, $\mathcal{L} \sim (1\text{--}200) \times 10^{34}$ cm $^{-2}$ s $^{-1}$
 - **FCC-hh**: pp collider, $\sqrt{s} \sim 100$ TeV, $\mathcal{L} \sim 5 \times 10^{34}$ cm $^{-2}$ s $^{-1}$

- **CLIC**: Compact Linear Collider, $\sqrt{s} \sim 1\text{--}3$ TeV, $\mathcal{L} \sim 6 \times 10^{34}$ cm $^{-2}$ s $^{-1}$

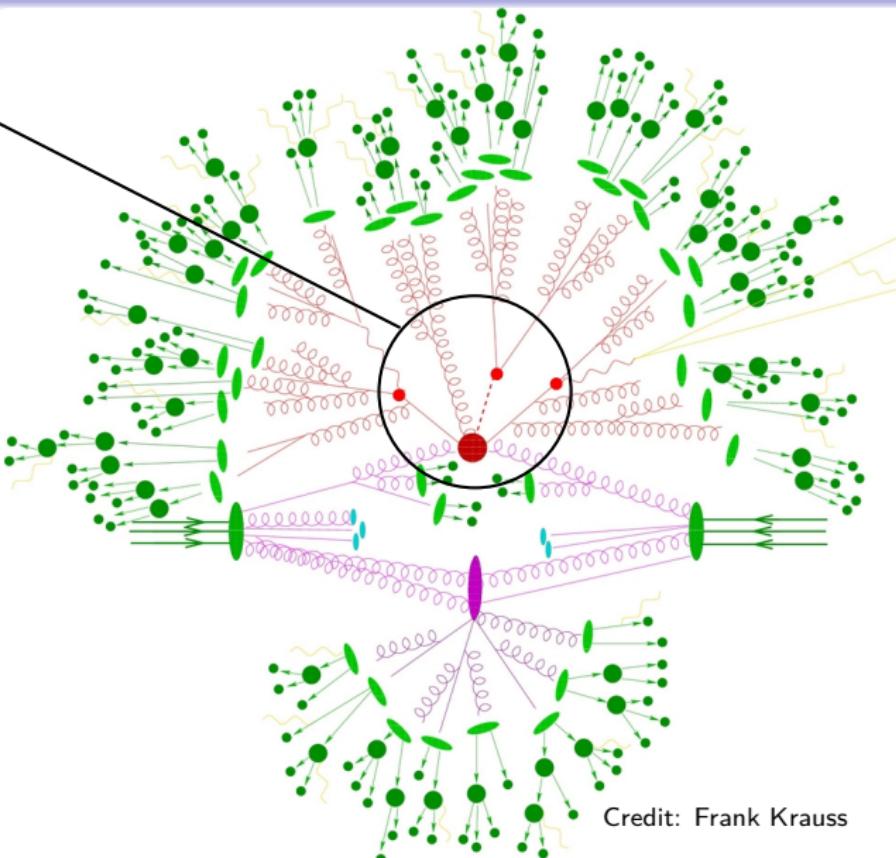


Typical Event in High Energy pp Collisions



Typical Event in High Energy pp Collisions

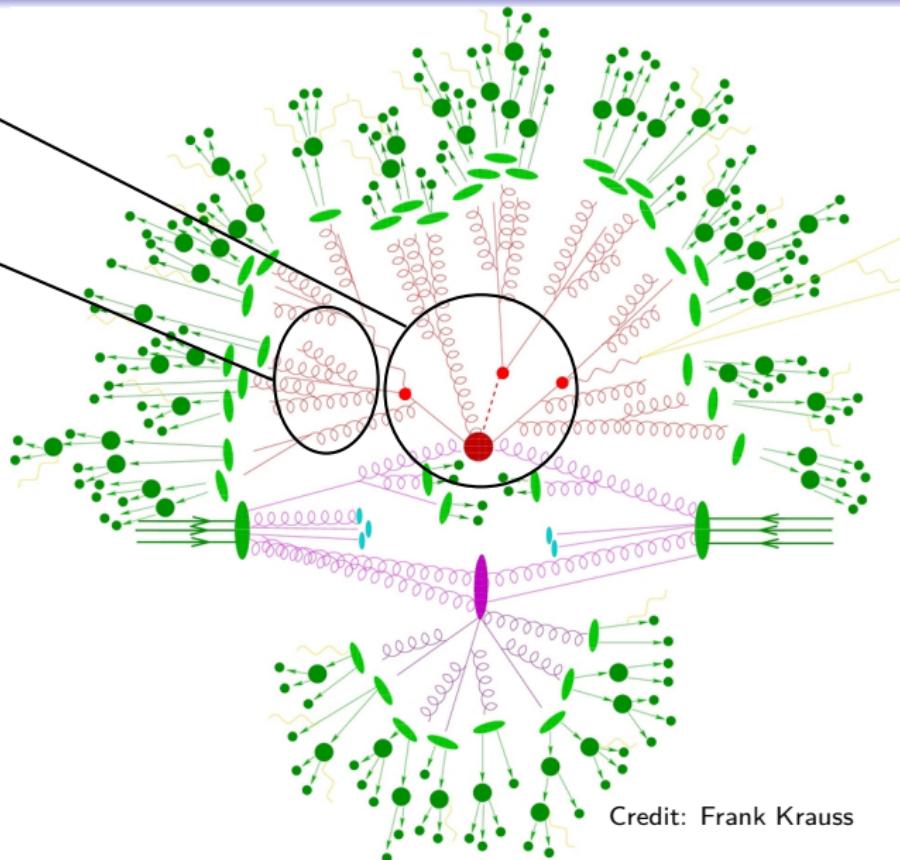
Hard scattering



Typical Event in High Energy pp Collisions

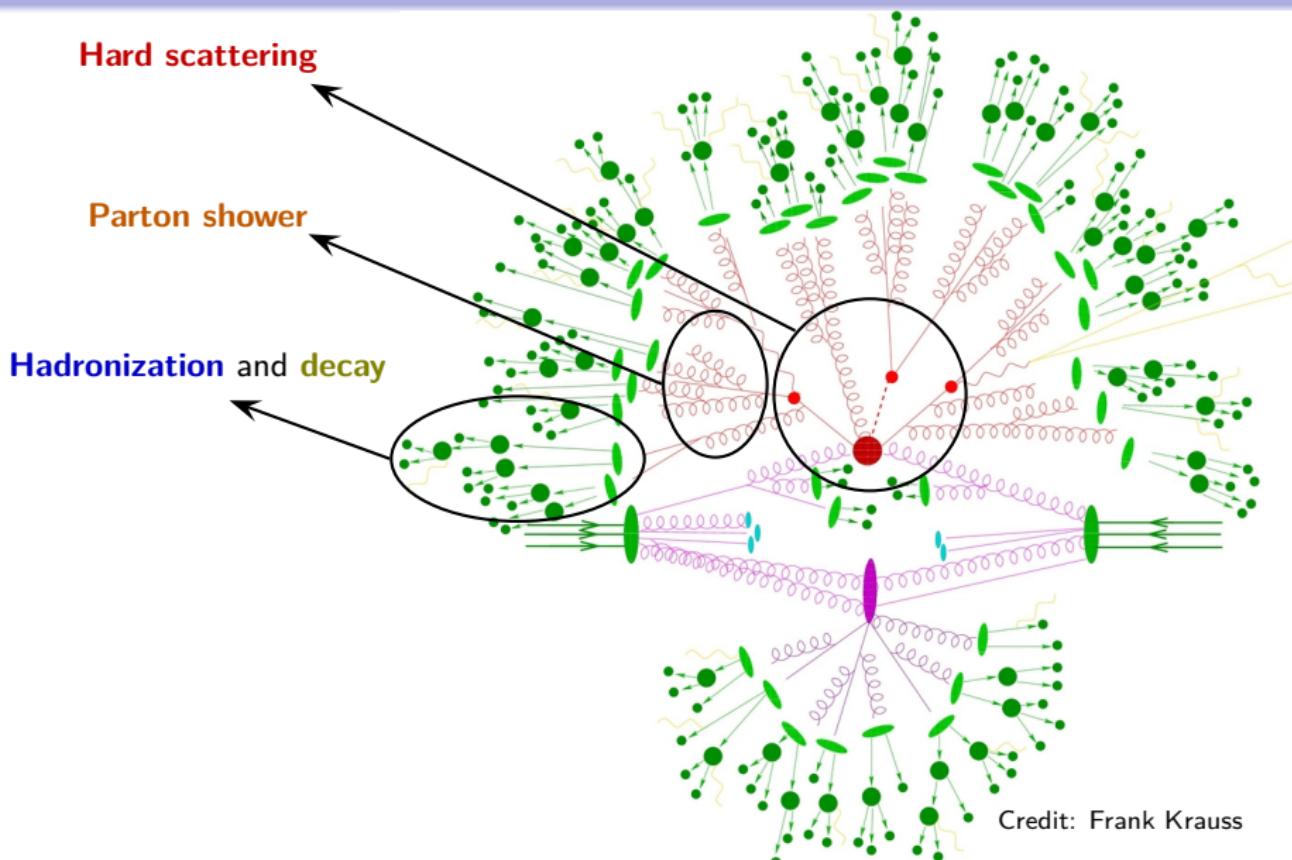
Hard scattering

Parton shower



Credit: Frank Krauss

Typical Event in High Energy pp Collisions



Credit: Frank Krauss

Typical Event in High Energy pp Collisions

Hard scattering

Parton shower

Hadronization and decay

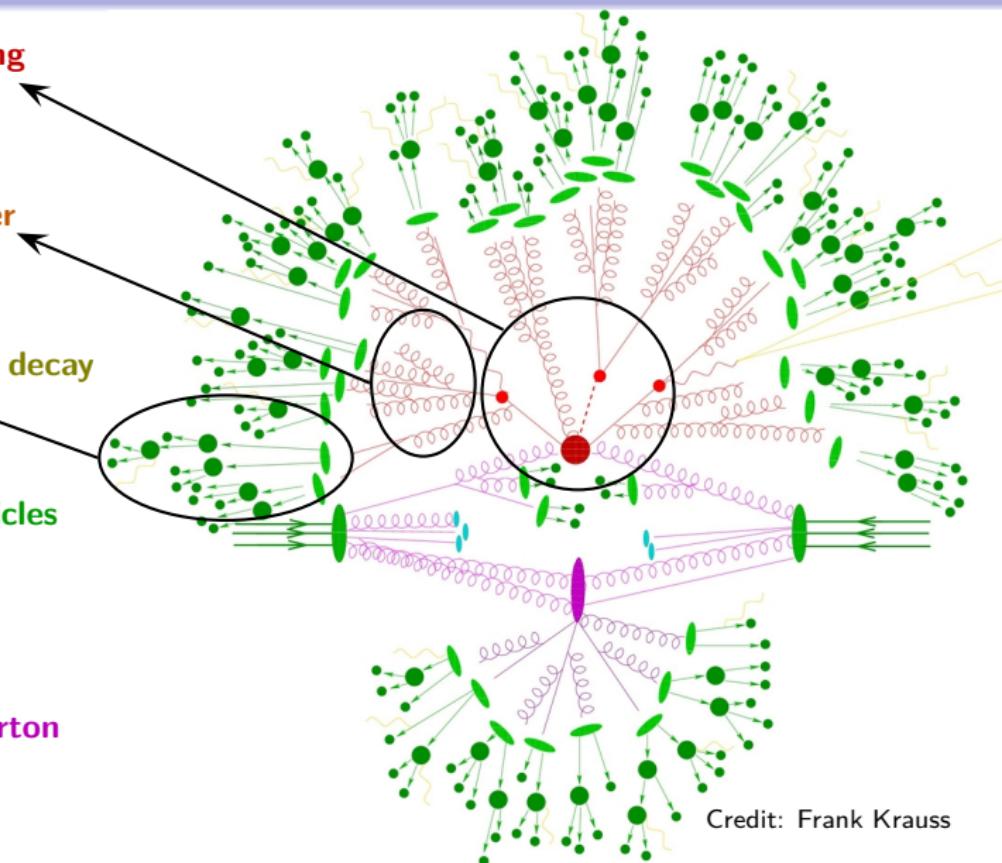
Collimated particles



Jet

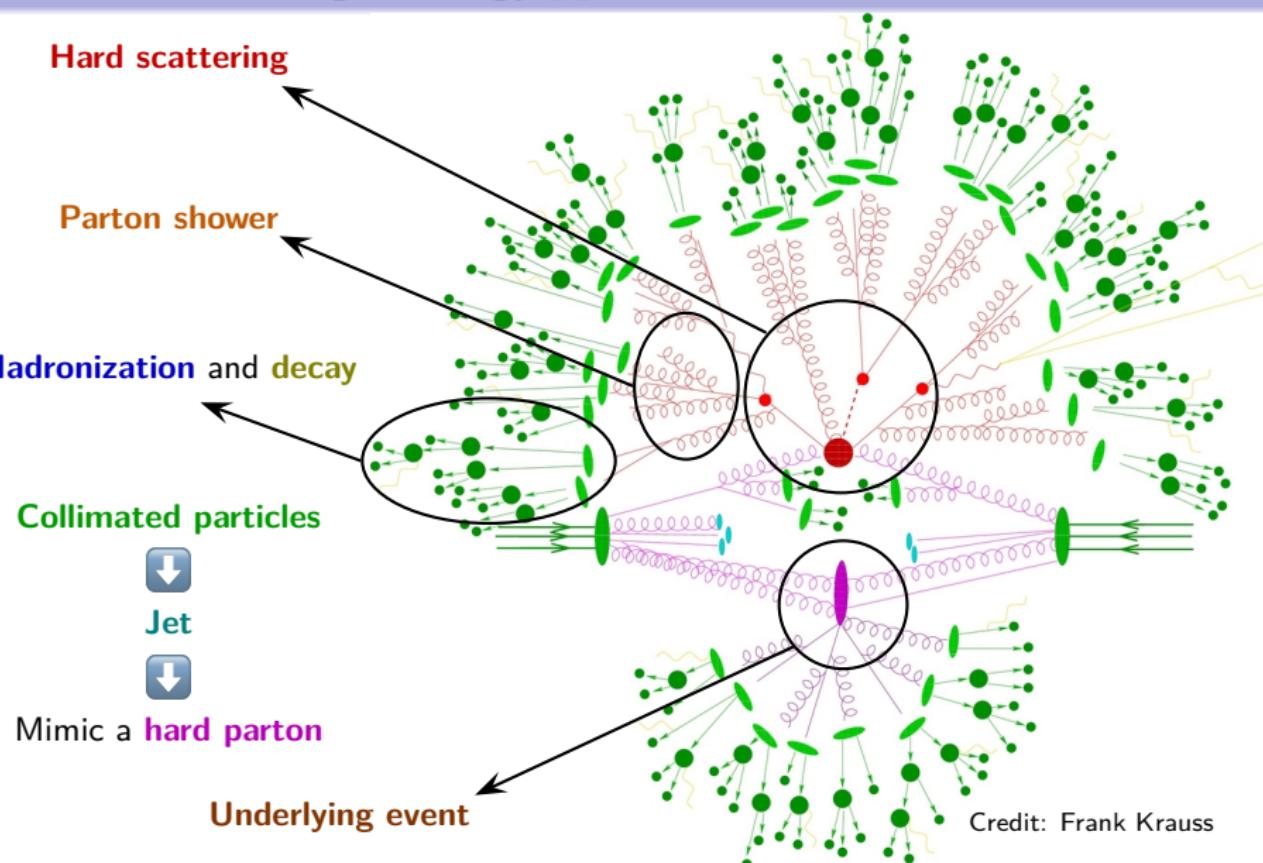


Mimic a hard parton

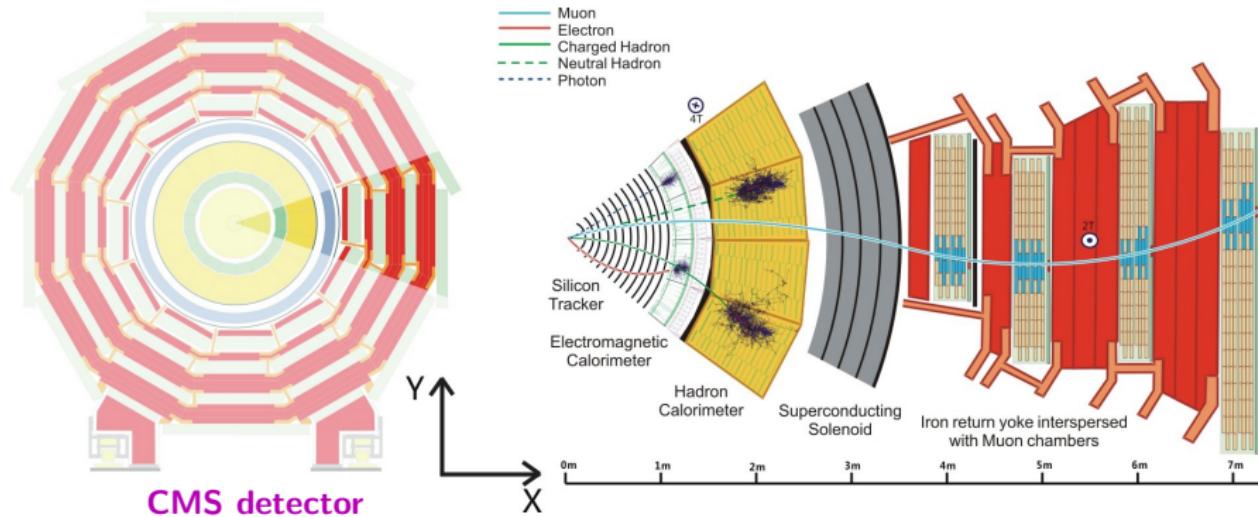


Credit: Frank Krauss

Typical Event in High Energy pp Collisions

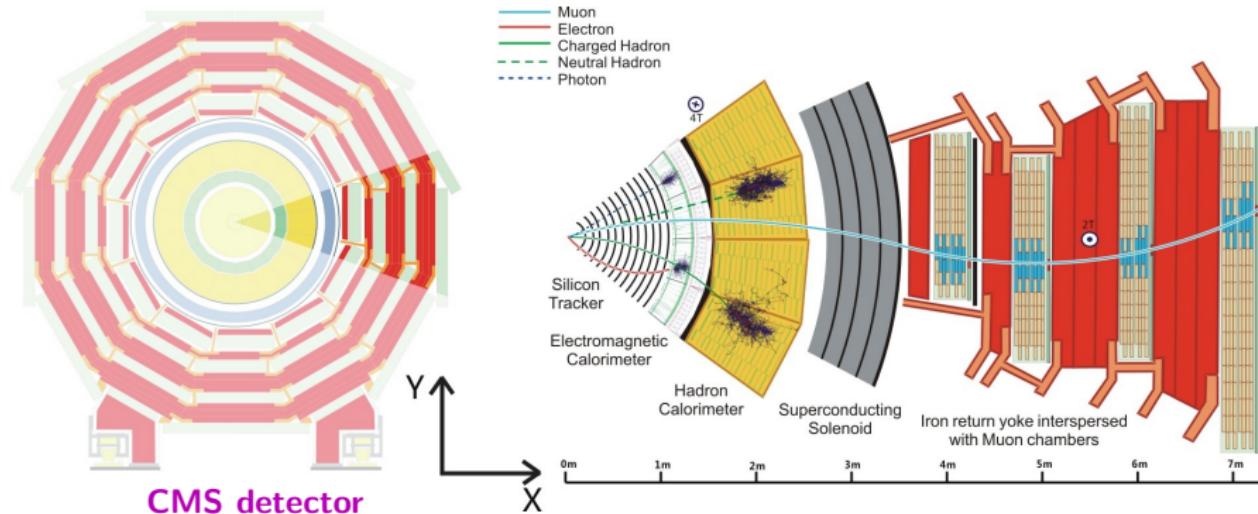


Particle Detectors at Colliders



Sub-detectors	γ	e^\pm	μ^\pm	Charged hadrons	Neutral hadrons	ν , DM
Tracker, $ \eta \lesssim 2.5$	✗	✓	✓	✓	✗	✗
ECAL, $ \eta \lesssim 3$	✗	✗	✓	✓	✗	✗
HCAL, $ \eta \lesssim 5$	✗	✗	✗	✗	✗	✗
Muon detectors, $ \eta \lesssim 2.4$	✗	✗	✓	✗	✗	✗

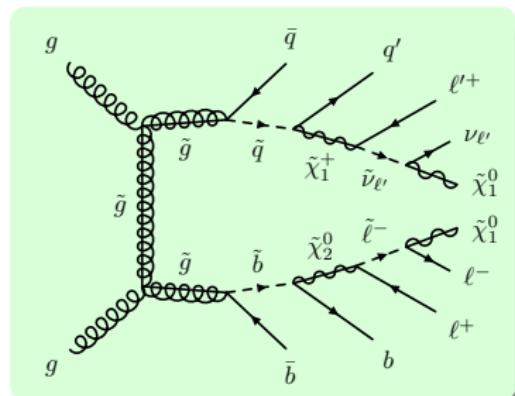
Particle Detectors at Colliders



Sub-detectors	γ	e^\pm	μ^\pm	Charged hadrons	Neutral hadrons	ν , DM
Tracker, $ \eta \lesssim 2.5$	✗	✓	✓	✓	✗	✗
ECAL, $ \eta \lesssim 3$	✗	✗	✓	✓	✗	✗
HCAL, $ \eta \lesssim 5$	✗	✗	✗	✗	✗	✗
Muon detectors, $ \eta \lesssim 2.4$	✗	✗	✓	✗	✗	✗

Missing energy
 E_T

DM Production



Social dark matter

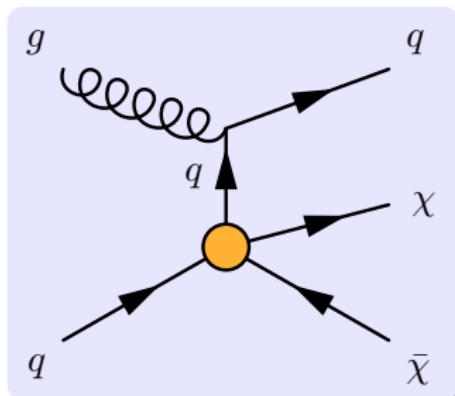
Accompanied by other new particles

Complicated decay chains

Decay products of other particles

Various final states

(jets + leptons + \cancel{E}_T , ...)



Maverick dark matter

DM particle is the **only new particle**
reachable at the collision energy

Direct production

Mono-X + E_T final states

(monojet, mono- γ , mono- W/Z , ...)



τ -portal Simplified DM Models

We studied four **τ -portal simplified models** involving a mediator with additive quantum numbers identical to the right-handed τ^-

We interpreted the **GC GeV excess signal** as **DM annihilation into $\tau^+\tau^-$** , and discussed **how to test this interpretation at the LHC**

Spin-1/2 fermion χ , spin-0 mediator ϕ :

$$\mathcal{L}_\phi = \lambda \phi \bar{\tau}_R \chi_L + \text{H.c.}$$

- **DFDM model:** χ is a Dirac fermion

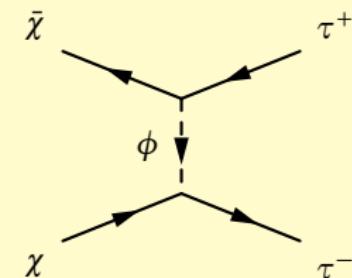
- **MFDM model:** χ is a Majorana fermion

Spin-0 scalar χ , spin-1/2 mediator ψ :

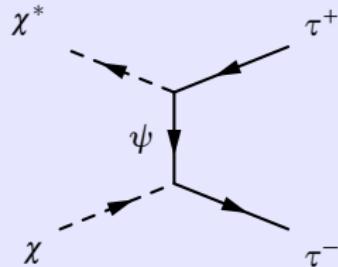
$$\mathcal{L}_\psi = \kappa \chi \bar{\tau}_R \psi_L + \text{H.c.}$$

- **CSDM model:** χ is a complex scalar

- **RSDM model:** χ is a real scalar



DFDM: annihilation

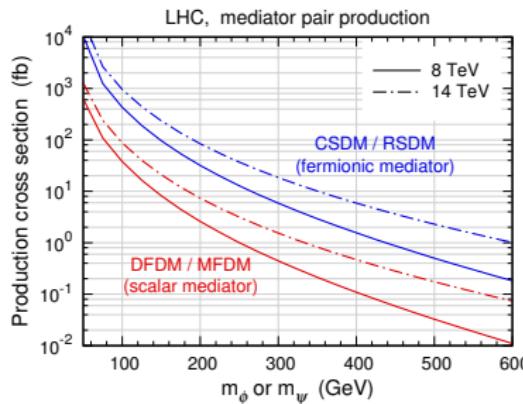


CSDM: annihilation

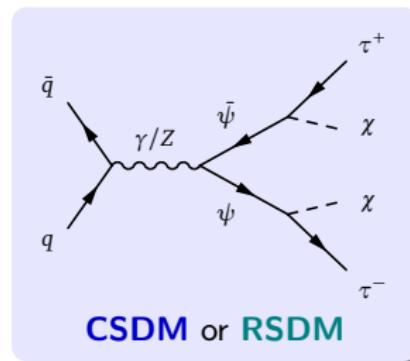
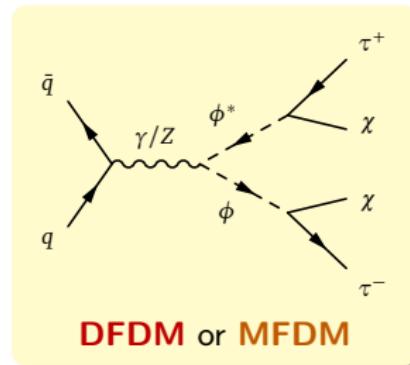
Mediator Pair Production at the LHC

The mediators ϕ and ψ could be produced at the LHC through **Drell-Yan processes** exchanging s -channel γ or Z , and then decay into τ^\pm and χ

 We found that the 8 TeV LHC data cannot explore the interesting regions in these models, and went further to investigate the LHC sensitivity at $\sqrt{s} = 14$ TeV with **tight τ_h -tagging** techniques



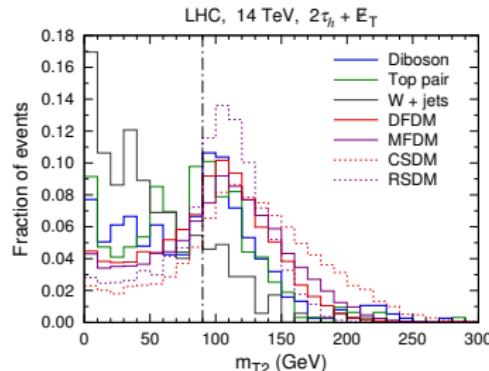
[ZHY, Bi, Yan, Yin, 1410.3347, PRD]





14 TeV LHC Searches for $pp \rightarrow \phi\phi^*/\psi\bar{\psi} \rightarrow \tau^+\tau^-\chi\chi$

2 $\tau_h + \cancel{E}_T$ channel: two opposite-sign tau-jet (τ_h);
without any other particle; $m_{T2} > 90$ GeV



Signals:

DFDM model

$m_\phi = 225$ GeV

MFDM model

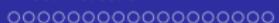
$m_\phi = 250$ GeV

CSDM model

$m_\psi = 300$ GeV

RSDM model

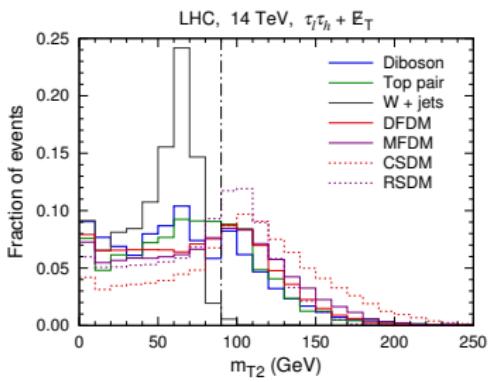
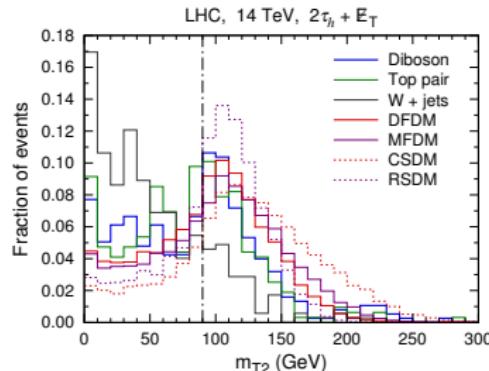
$m_\psi = 200$ GeV



14 TeV LHC Searches for $pp \rightarrow \phi\phi^*/\psi\bar{\psi} \rightarrow \tau^+\tau^-\chi\chi$

🎃 **$2\tau_h + \cancel{E}_T$ channel:** two opposite-sign tau-jet (τ_h); without any other particle; $m_{T2} > 90$ GeV

风筝 **$\tau_\ell\tau_h + \cancel{E}_T$ channel:** one τ_h and one light lepton ($\ell = \mu, e$) with opposite signs; without any other particle; $m_{T2} > 90$ GeV



拼图 **Signals:**

DFDM model

$$m_\phi = 225 \text{ GeV}$$

MFDM model

$$m_\phi = 250 \text{ GeV}$$

CSDM model

$$m_\psi = 300 \text{ GeV}$$

RSDM model

$$m_\psi = 200 \text{ GeV}$$

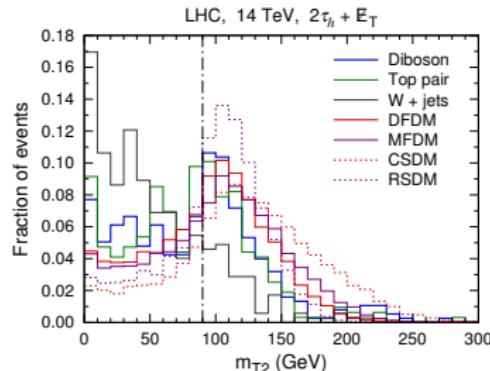
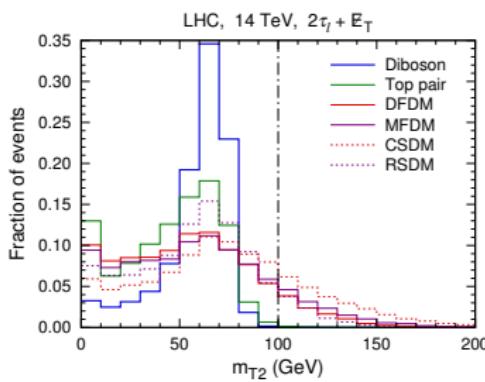
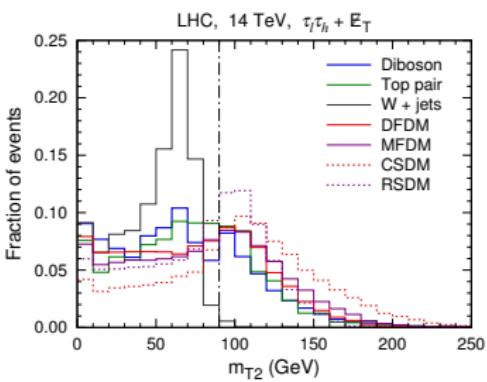


14 TeV LHC Searches for $pp \rightarrow \phi\phi^*/\psi\bar{\psi} \rightarrow \tau^+\tau^-\chi\chi$

$2\tau_h + \cancel{E}_T$ channel: two opposite-sign tau-jet (τ_h); without any other particle; $m_{T2} > 90$ GeV

$\tau_\ell\tau_h + \cancel{E}_T$ channel: one τ_h and one light lepton ($\ell = \mu, e$) with opposite signs; without any other particle; $m_{T2} > 90$ GeV

$2\tau_\ell + \cancel{E}_T$ channel: two opposite-sign light leptons; $|m_{\ell\ell} - m_Z| > 10$ GeV for the same-favor case; without any other particle; $m_{T2} > 100$ GeV



Signals:

DFDM model

$$m_\phi = 225 \text{ GeV}$$

MFDM model

$$m_\phi = 250 \text{ GeV}$$

CSDM model

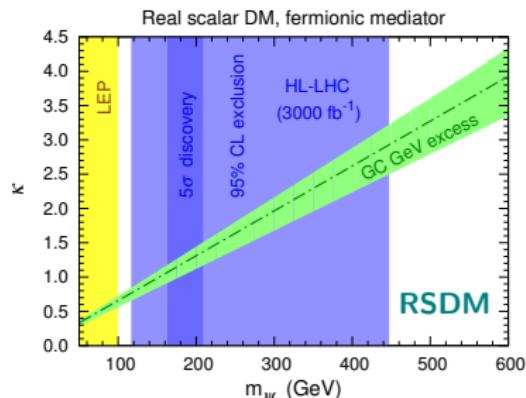
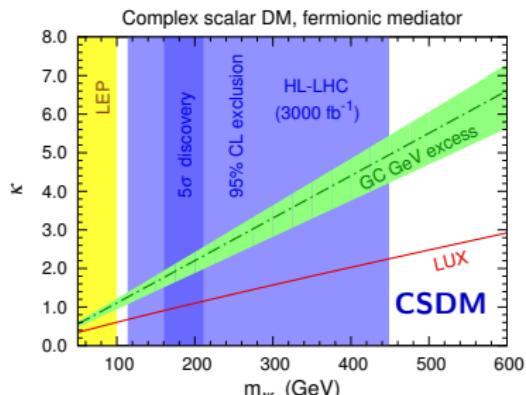
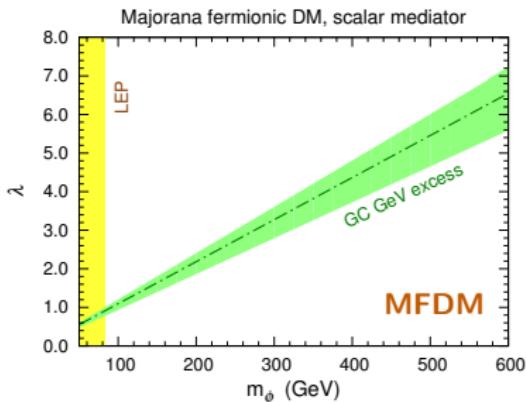
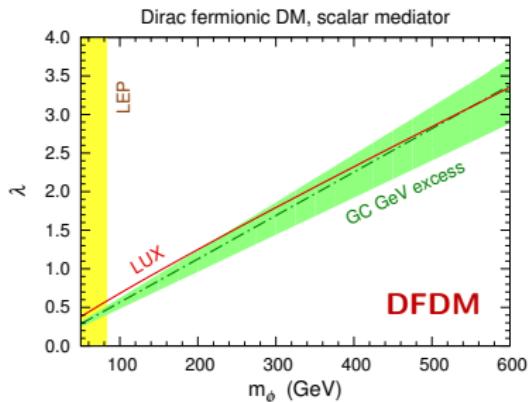
$$m_\psi = 300 \text{ GeV}$$

RSDM model

$$m_\psi = 200 \text{ GeV}$$



Sensitivity of the 14 TeV High-Luminosity LHC

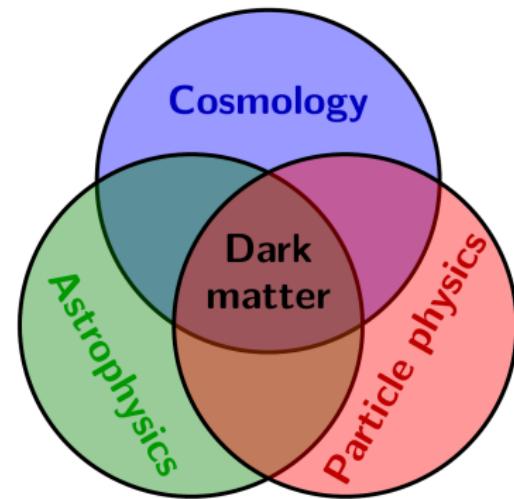


[ZHY, Bi, Yan, Yin, 1410.3347, PRD]



Summary

- **Dark matter** connects our knowledge of the Universe from the **largest** to the **smallest** scales
- Although several anomalous observations have been found in direct and indirect searches, there is **no absolutely solid DM detection signal so far**
- **DM detection sensitivities are being improved quickly**, and it is promising to detect robust DM signals in the future



Summary

- **Dark matter** connects our knowledge of the Universe from the **largest** to the **smallest** scales
- Although several anomalous observations have been found in direct and indirect searches, there is **no absolutely solid DM detection signal so far**
- **DM detection sensitivities are being improved quickly**, and it is promising to detect robust DM signals in the future

Thank you!

