# Magnetic Resonance Imaging (MRI)

BME/EECS 516 Lecture #6 – 3<sup>rd</sup> Dim, Noise Notes on MRI, Part 3

#### Announcements

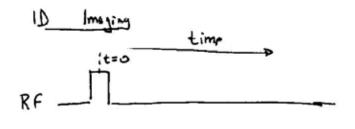
- i US Project Due today.
- i HW #4 due Tuesday, 10/31
- i HW #5 due Tuesday, 10/31
- MRI Project to be assigned next week

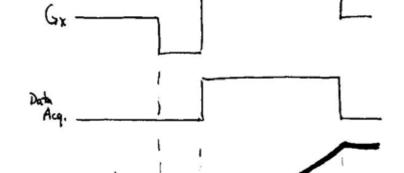
#### Last Lecture

- i 1D Imaging
- 2D Imaging using projections
- i 2D spin-warp imaging
  - Sampling in K-space
  - Point Spread Function
  - Resolution of FFT
  - Resolution, Object, Sampling in MRI

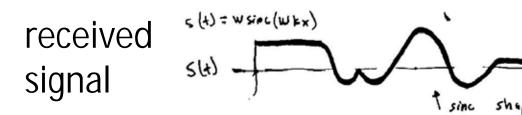
## 1D Imaging

pulse sequence

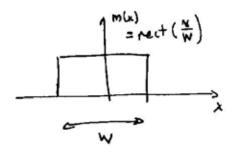




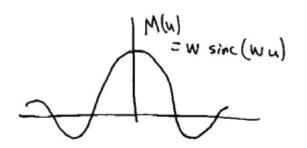




Object



FT of Object

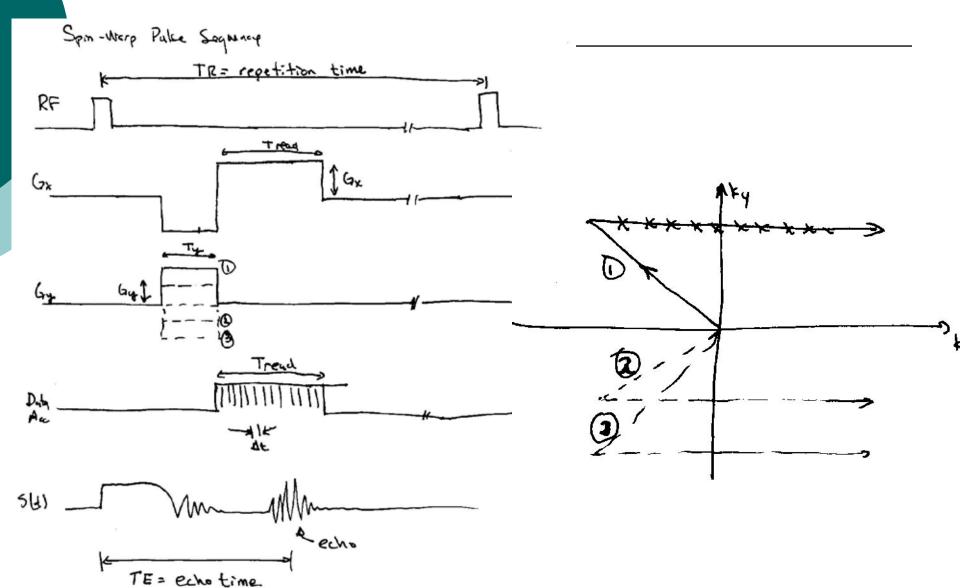


$$k_{x}(t) = \frac{\gamma}{2\pi} \int_{0}^{t} G_{x}(\tau) d\tau$$

## 2D Imaging using Projections

Pulse Sequence - Projection Imaging (2-5ided) 0 & [0, n)

## 2D Spin-Warp Imaging Pulse Sequence



## 2D Spin-Warp Imaging - Resolution and Object and Sample Spacing in MRI

$$FOV_x = 1/\Delta k_x$$
  $FOV_y = 1/\Delta k_y$   
 $\Delta x = 1/W_{kx}$   $\Delta y = 1/W_{ky}$ 

$$\Delta k_{x} = \frac{\gamma}{2\pi} G_{x} \Delta t \qquad W_{kx} = N_{x} \Delta k_{x} = \frac{\gamma}{2\pi} G_{x} T_{read}$$

$$\Delta k_{y} = \frac{\gamma}{2\pi} \Delta G_{y} T_{y} \qquad W_{ky} = N_{y} \Delta k_{y} = \frac{\gamma}{2\pi} 2G_{y,\text{max}} T_{y}$$

- To achieve a better spatial resolution for a MR image in the x direction, you can
  - A) Increase the duration of Gx (Tread)
  - B) Increase the duration of Gy (Ty)
  - $\Gamma$  C) Increase the number of samples while maintaining time between samples ( $\Delta t$ )
  - D) Increase Gx amplitude

- To increase of the field of view in y direction, you can
- $_{i}$  A) Decrease the time between samples ( $\Delta t$ )
- B) Decrease the number of samples while maintaining time between samples ( $\Delta t$ )
- $_{i}$  C) Decrease the duration of  $G_{y}$  ( $T_{y}$ )
- <sub>i</sub> D) Decrease  $\Delta G_y$

#### This Lecture

- The 3<sup>rd</sup> Dimension Z
  - Extend 2D Spin-Warp Imaging to 3D
- Slice Selective Excitation common approach for dealing with the 3<sup>rd</sup> (z) dimension
  - Small Tip Angle Approximation
  - Putting Slice Selection with the Signal Equation
  - Larger Flip Angles
  - Multi-slice Imaging

#### The 3<sup>rd</sup> Dimension - Z

Recall the MR signal is equal to the Fourier transform of the initial magnetization evaluated at locations defined by  $k_x$  and  $k_y$ .

$$s(t) = \partial m(x,y) \exp(-i2\rho(xk_x(t) + yk_y(t))) dxdy$$
  
=  $F_{2D}\{m(x,y)\}|_{u=k_x(t),v=k_y(t)} = M(k_x(t),k_y(t))$ 

$$k_x(t) = \frac{g}{2\rho} \int_0^t G_x(t) dt, \quad k_y(t) = \frac{g}{2\rho} \int_0^t G_y(t) dt$$

#### The 3<sup>rd</sup> Dimension - Z

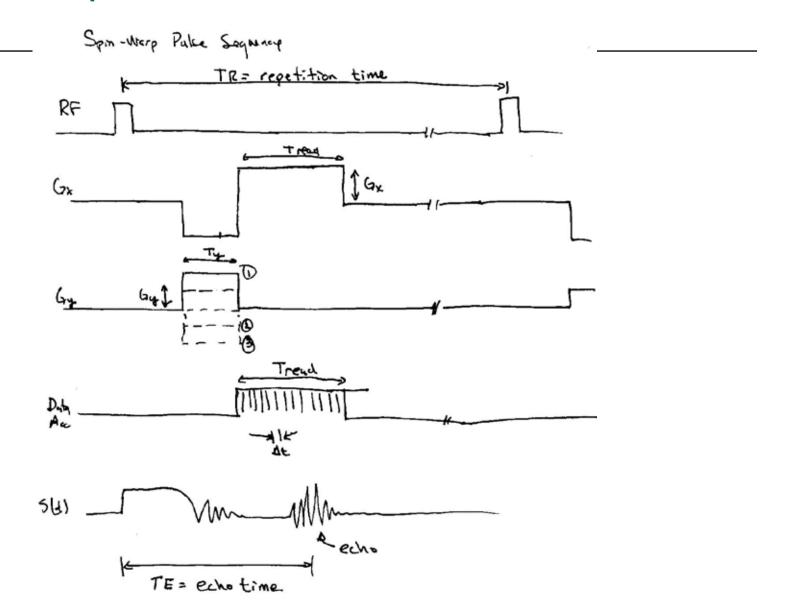
The 3D signal equation can be written as:

$$s(t) = \lim_{x} m(x, y, z) \exp(-i2p(k_x(t)x + k_y(t)y + k_z(t)z) dx dy dz$$

$$= M(u, v, w) \Big|_{u = k_x(t), v = k_y(t), w = k_z(t),}$$

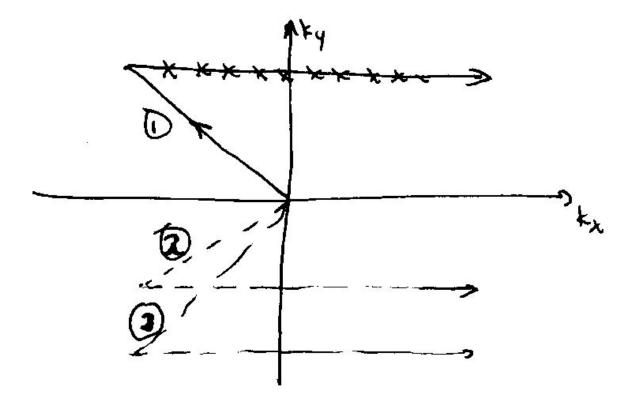
M(u, v, w) is the 3D FT of m(x, y, z)

## Recall 2D Spin-Warp Imaging Pulse Sequence

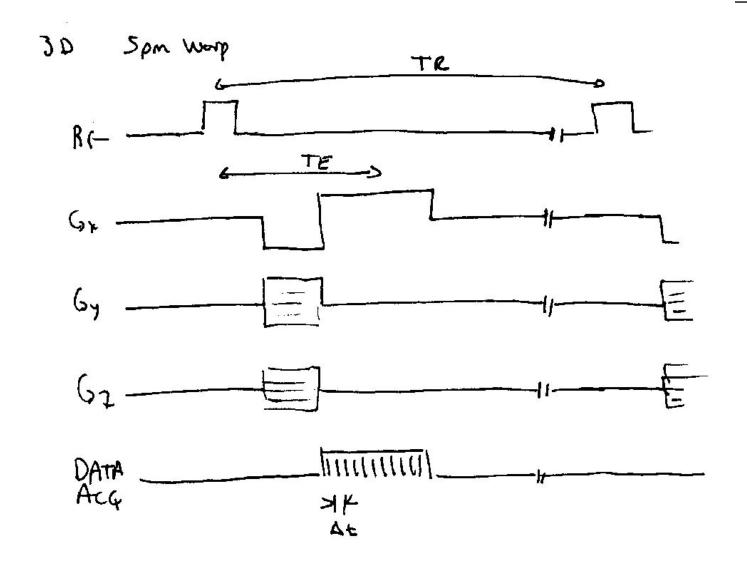


## Recall 2D Spin-Warp Imaging

One line at a time is acquired in the 2D Fourier domain (or k-space).

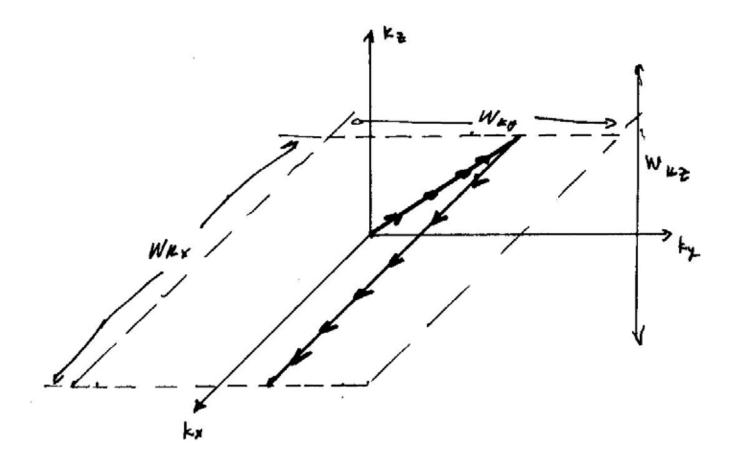


## 3D Spin-Warp Imaging



## 3D Spin-Warp Imaging

Acquire a cubic data set one line at a time:



#### The 3<sup>rd</sup> Dimension - Z

In 2D field:

$$FOV_x = 1/\Delta k_x$$
  $FOV_y = 1/\Delta k_y$   
 $\Delta x = 1/W_{kx}$   $\Delta y = 1/W_{ky}$ 

Add the Z dimension:

$$FOV_z = 1/\Delta k_z$$
  $\Delta z = 1/W_{kz}$ 

#### The 3<sup>rd</sup> Dimension - Z

- Samples in the y direction:  $N_y$ Samples in the z direction:  $N_z$ Total time to acquire 3D volume:  $N_y * N_z * TR$
- For example,  $N_y = N_z = 128$  and TR = 33 ms, Image acquisition time = 128 x 128 x 33 min = 9 min Too long!

reduce the image acquisition time and obtain a high frame rate?

Can we use a short TR (<1ms) to reduce the image acquisition time and obtain a high frame rate? Not easily...

- Too small a TR results in low signal intensity - recall signal intensity
- Limitation on how squeezing too many things in a very short time (RF pulses, gradients, etc.) - bioeffects
- Poor SNR (more on this later) from short readouts

Δt of several μs is typically used in MRI acquisition. Would using a short Δt (sub-μs) improve the imaging?

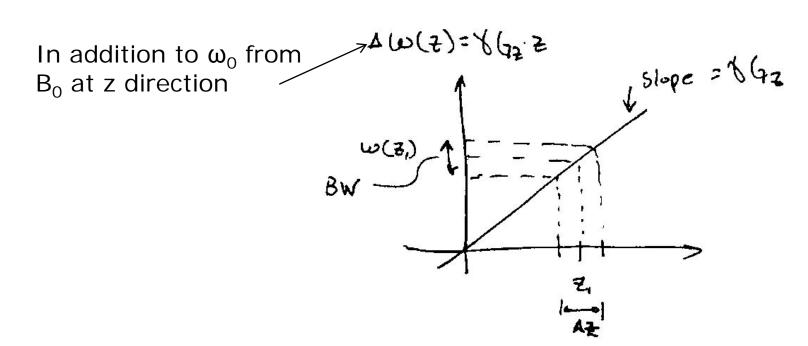
- Δt of several μs is typically used in MRI acquisition. Would using a short Δt (sub-μs) improve the imaging?
- i No.
  - Main issues is reduced signal to noise ratio (more on this later)
  - Field of view is very large, but this isn't necessarily a problem

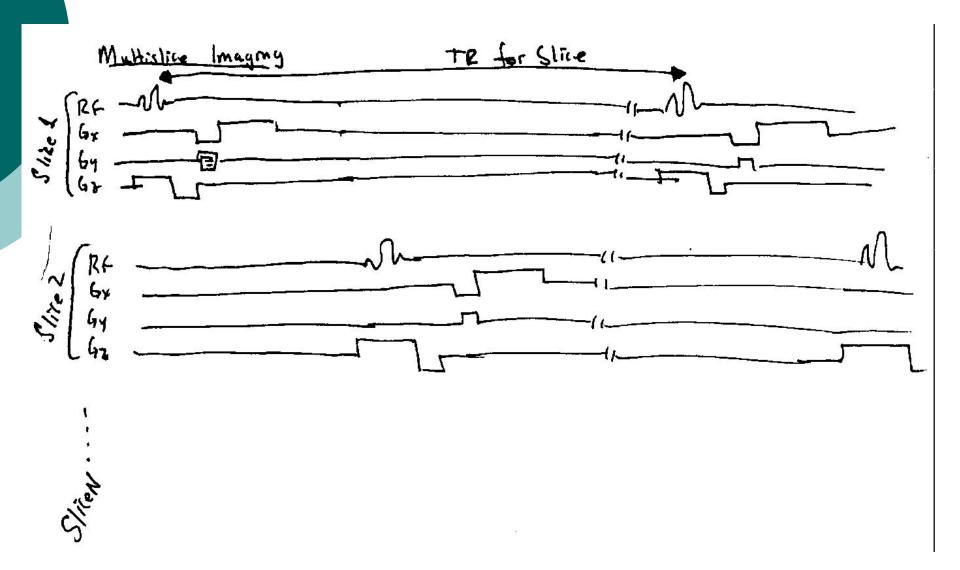
### Questions?

3D Spin-Warp Imaging

The most common approach for dealing with the 3<sup>rd</sup> (z) dimension

- Apply a z-gradient so that the resonance frequency varies in the z-direction
- Apply a bandpass RF pulse to excite only the those spins whose resonant frequency lies within the band:





Step 1: Apply a z-gradient so that the resonance frequency varies in the z-direction. Within a slice of thickness Δz:

$$B_{z} = B_{0} + G_{z} \cdot z$$

$$\omega(z) = \gamma(B_{0} + G_{z} \cdot z)$$

$$\omega_{eff}(z) = \gamma G_{z} \cdot z$$

- Step 2: Apply a bandpass RF pulse  $(B_1)$  with bandwidth (BW in Hz)
- RF pulse excite only the those spins whose resonant frequency lies within the band.
  - $[\omega_{RF} BW * 2\pi/2, \omega_{RF} + BW * 2\pi/2]$   $= [\gamma(B_0 + Gzz_1), \gamma(B_0 + Gzz_2)] \text{ to excite spins}$ only within slice  $[z_1, z_2]$

- How to determine  $G_z$  for a slice thickness =  $\Delta z$ ?
  - Bandwidth of the RF pulse BW
  - Bandwidth of the nuclei spin frequency in the magnetic field  $-G_z\Delta z\frac{\gamma}{2\pi}$

$$G_z = \frac{BW}{\left(\frac{\gamma}{2\pi}\right)\Delta z}$$

#### True of False

- When exciting different locations/slices, BW and G<sub>z</sub> remains the same using the same slice thickness (Δz).
- When exciting different locations/slices, different RF excitation pulses with different center frequencies will be used.

When exciting different locations/slices, BW and  $G_z$  remains the same using the same slice thickness ( $\Delta z$ ).

$$G_z = \frac{BW}{\left(\frac{\gamma}{2\pi}\right)\Delta z}$$

 $\omega_{RF}$  is changed for each slice, depending on the center location of each slice.

$$\omega_{RF} = \omega_0 + \gamma G_z z_{slice} = \gamma B_0 + \gamma G_z z_{slice}$$

- How does RF pulse look like in the frequency domain for a rectangular slice profile?
  - $\mid$  A) rect((f-f<sub>RF</sub>)/BW)
  - $\mid$  B) sinc((f-f<sub>RF</sub>)/BW)
  - $\Gamma$  C) exp(- $\pi$ (f-f<sub>RF</sub>)<sup>2</sup>/BW<sup>2</sup>)
  - $\mid$  D)  $\delta(f-f_0)$

How does RF pulse look like in the time domain?

Recall: 
$$F^{-1}\{G(f-f_0)\}=g(t)e^{i2\pi f_0t}$$

Example : 
$$B1(f) = \text{rect}\left(\frac{f - f_{RF}}{BW}\right)$$

$$B_1(t) = BW \operatorname{sinc}(t \cdot BW) e^{i\gamma G_z z_{slice} t}$$

### Typical MRI values

Typical RF pulse – on the order of kW power and 5-10 ms in length.

Typical gradient amplitude – 4 Gauss/cm = 40 mT/m.

Typical Resolution – 1 mm.

#### Questions?

Slice Selective Excitation – Bloch equation

## Small Tip Angle Approximation

Solve Bloch Equation for Slice Selective Excitation

- Solve the Bloch equations for this specific case, Let
  - $B_1$  time-varying magnetic field rotating at  $\omega_0$
  - Rotating frame be at  $\omega_{frame} = \omega_0$

$$\mathbf{B}_{1}(t) = B_{1}(t)(\cos \mathbf{W}_{0}t\mathbf{i} + \sin \mathbf{W}_{0}t\mathbf{j})$$
$$\mathbf{B}_{1,eff}(t) = B_{1}(t)\mathbf{i}'$$

Recall  $B_1$  is in x-y plane

A z-gradient is applied, so the component in the z-direction is:

$$\mathbf{B}_{\mathbf{z}}(z) = (B_0 + G_z \times z)\mathbf{k}$$

$$\mathbf{B}_{\mathbf{z},eff}(z) = (G_z \times z)\mathbf{k}$$

The net effective applied field is:

$$\mathbf{B}_{eff} = B_1(t)\mathbf{i}' + (G_z \times z)\mathbf{k}$$

Bloch equation for this case reduces to

$$\frac{d\mathbf{M}_{rot}}{dt} = \mathbf{M}_{rot} ' \mathcal{B}_{eff} = \hat{\mathbf{e}} 0 \mathcal{G}_z z 0 \hat{\mathbf{u}}$$

$$\hat{\mathbf{e}} 0 \mathcal{G}_z z 0 \mathcal{B}_1(t) \hat{\mathbf{u}} \mathbf{M}_{rot}$$

$$\hat{\mathbf{e}} 0 - \mathcal{B}_1(t) 0 \hat{\mathbf{u}}$$

What we would like to know is how  $M_{rot}$ , varies vs. z following the application of  $B_1$ . This is a very difficult equation to solve because it is non-linear.

To solve the above Bloch equation, use small tip angle approximation. Assume the  $B_1$  produces a small net rotation angle:

$$\grave{O}B_1(t)dt < \frac{p}{6}$$
 (30°)

In this case, we can assume  $m_z \approx m_0$  during the RF pulse. Essentially, we are saying that:  $\cos \mathcal{E}_0 \mathcal{B}_1(t) dt \stackrel{\circ}{\dot{\Xi}} 1$ 

Under this assumption,  $\frac{dm_{z(t)}}{dt} = 0$  à  $m_z(t) = m_0$ .

The above Bloch equation can then be rewritten as:

$$\frac{d\mathbf{M_{rot}}}{dt} = \mathbf{\hat{e}} \quad 0 \qquad \mathbf{\mathcal{G}}_z z \qquad 0 \quad \mathbf{\hat{e}} \quad \mathbf{\hat{e}} \quad \mathbf{\mathcal{G}}_z z \qquad 0 \qquad \mathbf{\hat{e}} \quad \mathbf{\hat{e}} \quad \mathbf{\mathcal{G}}_z z \qquad 0 \qquad \mathbf{\mathcal{G}}_1(t) \mathbf{\hat{e}} \quad 0 \qquad \mathbf{\hat{e}} \quad \mathbf{\hat{e}}$$

Now solve for

$$m_{xy,rot}(z,t) = m_{x,rot}(z,t) + im_{y,rot}(z,t).$$

Write a differential equation using for the transverse component:

$$\frac{dm_{xy,rot}}{dt} = \frac{dm_{x,rot}}{dt} + i\frac{dm_{y,rot}}{dt}$$

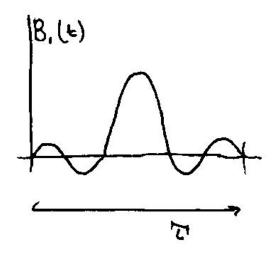
$$= \mathcal{G}G_z z m_{y,rot} - i\mathcal{G}G_z z m_{x,rot} + i\mathcal{G}B_1(t) m_0$$

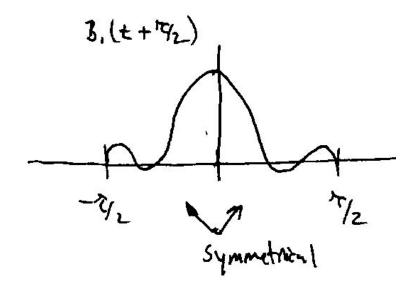
$$= -i\mathcal{G}G_z z m_{xy,rot} + i\mathcal{G}B_1(t) m_0$$

- It is a first order differential equation
  - $i\gamma G_z z$  is a constant with respect to time
  - I Driving function  $i\gamma B_1(t)m_0$
  - Initial condition,  $m_{xy,rot}(z,t) = 0$
- Solution to this differential equation

$$m_{xy,rot}(z,t) = im_0 e^{-igG_z zt} \sum_{0}^{\infty} \exp(igG_z z X)gB_1(X)dX$$

We want the solution to this differential equation at the time of the end of the RF pulse -  $\tau$ ,  $m_{xy,rot}(z,\tau)$ :





- Substitute s = x t/2
- Assume RF pulse is symmetrical (even) around  $\tau/2$  and is zero outside of the interval  $[0,\tau]$ .

$$\begin{split} &m_{xy,rot}(z,t) = im_0 e^{-igG_z zt} \underbrace{\stackrel{t/2}{\overset{t}{\bigcirc}}}_{-t/2} \exp(igG_z z(s+t/2)) gB_1(s+t/2) ds \\ &= ign_0 e^{-igG_z zt/2} \underbrace{\stackrel{t/2}{\overset{t}{\bigcirc}}}_{-t/2} \exp(igG_z zs) B_1(s+t/2) ds \\ &= ign_0 e^{-igG_z zt/2} \underbrace{\stackrel{t}{\overset{t}{\bigcirc}}}_{-t/2} B_1(s+t/2) \exp \underbrace{\stackrel{\text{def}}{\overset{t}{\bigcirc}}}_{\overset{\text{def}}}{\overset{\text{def}}}{\overset{\text{def}}{\overset{\text{def}}{\overset{\text{def}}{\overset{\text{def}}{\overset{\text{def}}{\overset{\text{def}}}{\overset{\text{def}}{\overset{\text{def}}{\overset{\text{def}}}{\overset{\text{def}}}{\overset{\text{def}}}{\overset{\text{def}}}{\overset{\text{def}}}{\overset{\text{def}}}}$$

For symmetrical RF pulses, the forward and inverse FT are the same.

Thus, under the small tip angle approximation, the slice profile (variation of  $m_{xy,rot}$  with z) is related to the spectrum of the RF pulse:

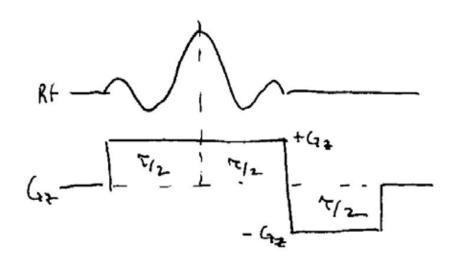
$$\left| m_{xy,rot}(z,t) \right| \mu \left| F \left\{ B_1(t) \right\} \right|_{f=\frac{g}{2\rho}G_zz}$$

$$m_{xy,rot}(z,\tau) = i\gamma m_0 e^{-i\gamma G_z z \tau/2} F\left\{ B_1(t+\tau/2) \right\} \Big|_{f=\frac{\gamma}{2\pi}G_z z}$$

- We're almost there, but still have some undesired phase variation in the z-direction  $\exp(-igG_zzt/2)$ 
  - can lead to undesired phase destruction when integrated by the RF coil.

$$m_{xy,rot}(z,\tau) = i\gamma m_0 e^{-i\gamma G_z z \tau/2} F\{B_1(t+\tau/2)\}\Big|_{f=\frac{\gamma}{2\pi}G_z z}$$

How is this fixed? Simply apply a negative  $G_z$  for a period  $\tau/2$  – often called a slice rephasing pulse.



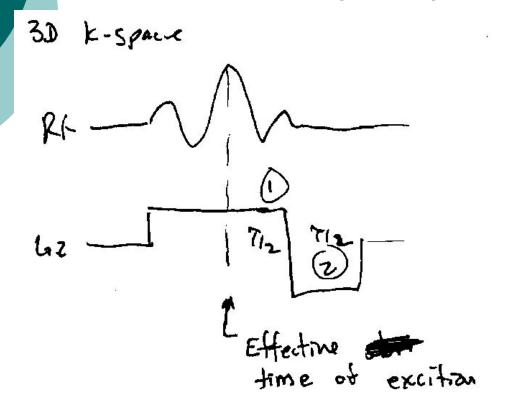
Slice rephasing pulse results in phase accumulation of:

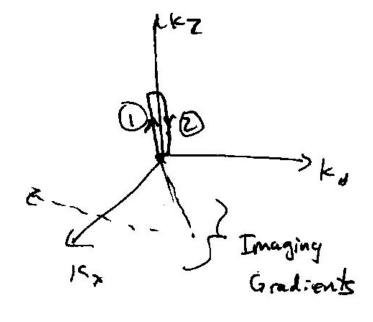
$$\exp \sum_{t=0}^{3t/2} \frac{\partial^{3t/2}}{\partial t} DB(z,t) dt = \exp(igG_z zt/2)$$

thus:

$$m_{xy,rot}(z,\tau) = i\gamma m_0 F \left\{ B_1(t+\tau/2) \right\} \Big|_{f=\frac{\gamma}{2\pi}G_z z}$$

#### k-space picture





RF pulse applied along real (i') axis à magnetization end up along imaginary (j')axis.

$$m_{xy,rot}(z,3\tau/2) = i\gamma m_0 F \left\{ B_1(t+\tau/2) \right\} \Big|_{f=\frac{\gamma}{2\pi}G_z z}$$

The flip angle at the center of the slice is:

$$\alpha = \int_{-\tau/2}^{\tau/2} \gamma B_1(s + \tau/2) ds = \gamma F \{ B_1(s + \tau/2) \} \Big|_{z=0}$$

After an  $\alpha$  pulse,  $|m_{xy,rot}|=m_0 \sin\alpha \sim m_0\alpha$  (for small  $\alpha$ ,  $\sin\alpha \sim \alpha$ ).

## RF pulse - MRI project

- $_{i}$  B<sub>1</sub>(t) = A\*RF\_shape(t)
  - A amplitude of RF pulse
  - RF\_shape(t) shape of RF pulse

For a 90° pulse (flip angle = 90°), if RF\_shape(t) is known, how to calculate A?

## RF pulse - MRI project

- $_{i}$  B<sub>1</sub>(t) = A\*RF\_shape(t)
  - A amplitude of RF pulse
  - RF\_shape(t) shape of RF pulse
- ¡ For a 90° pulse

$$\gamma \int_{0}^{t} B_{1}(\tau) d\tau = \frac{\pi}{2}$$

$$B_1(t) = A \cdot RF \_shape(t)$$

$$\gamma \int_{0}^{L} A \cdot RF \_shape(\tau) d\tau = \frac{\pi}{2}$$

#### MRI project

- For a slice thickness  $\Delta z$ , calculate  $G_{7}$ :
  - $BW * 2\pi = \gamma G_z \Delta z ? G_z = BW / (\gamma / 2\pi) \Delta z$
  - In time domain, the RF pulse is:

$$RF_{shif}(t) = RF(t)e^{i\omega_{shift}t}$$

$$\omega_{shift} = \gamma G_z z_{slice}$$

### MRI project

#### **Rotating Frame**

$$B_{x} = RF(t)\cos(\omega_{shift}t)$$

$$B_{y} = RF(t)\sin(\omega_{shift}t)$$

$$B_z = G_x x + G_y y + G_z z$$

## Questions?

Small Tip Angle Approximation

# Putting Slice Selection with the Signal Equation

## Putting Slice Selection with the Signal Equation

Define our slice profile function:

$$p(z) = gF\{B_1(t+t/2)\} \mid_{f=\frac{g}{2\rho}G_z z}$$

Magnetization:

$$m_{xy,rot}(z,3\tau/2) = i\gamma m_0 F \left\{ B_1(t+\tau/2) \right\} \Big|_{f=\frac{\gamma}{2\pi}G_z z}$$

$$m_{xy,rot}(z,3t/2) = im_0 p(z)$$

## Putting Slice Selection with the Signal Equation

3D distribution of magnetization by substituting  $im_0 = m(x, y, z)$  and putting it into the signal equation :

$$s(t) = i m(x,y,z) p(z) \exp(-i2p(xk_x(t) + yk_y(t))) dxdydz$$

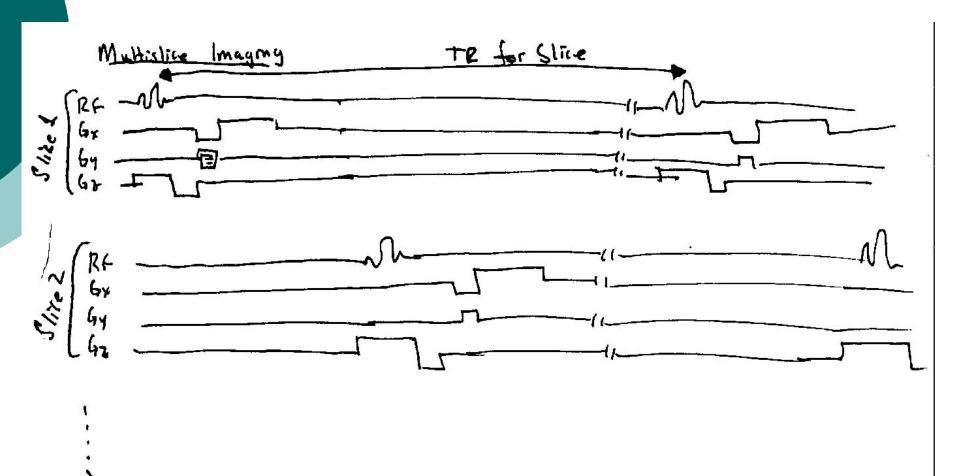
Here we are performing 2D imaging while integrating across the slice profile

## Larger Flip Angles

Analysis using small flip angle (e.g.  $\alpha < \pi/6$ ) approximation turns out perform well for large flip angles (e.g.  $\pi/2$ ) as well.

Here, a better approximation for the transverse component of the magnetization is:

$$m_{xy,rot}(z,3t/2) = im_0 \sin(p(z))$$

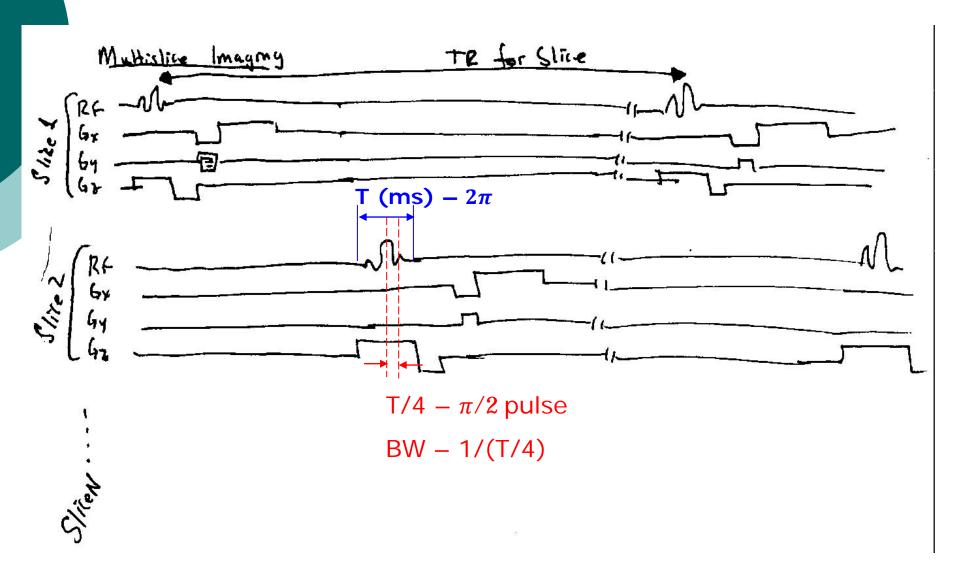


- Slice-selective, 2D spin-warp acquisition overall acquisition time is:
- $_{i}$  A)  $N_{y}^{*}N_{y}^{*}TR$
- $_{i}$  B)  $N_{y}^{*}TR$
- $C) N_z * TR$

- Slice-selective, 2D spin-warp acquisition overall acquisition time:  $N_y *TR$
- For example, acquire a T1-weighted image
  - Ny 20 slices
  - □ TR 500 ms
  - 1 128 phase encoding lines in k-space
  - Total acquisition time  $N_y *TR = \sim 1$  minute

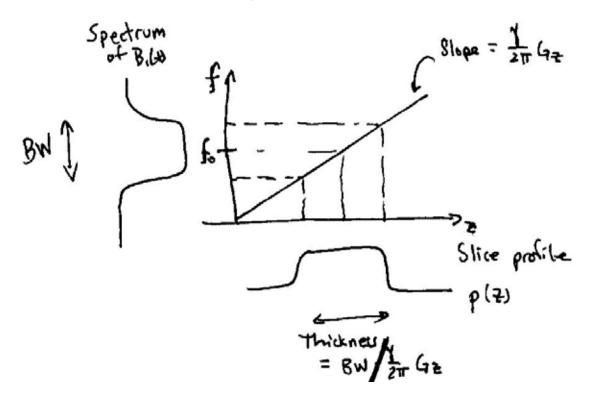
## Question?

Multi-slice imaging



- The most common way to image 3D volumes in MRI uses interleaved slice selective excitation.
  - Slice 1 is excited and part of the k-space data are acquired, then slice 2, 3...
  - After all have been acquired, come back to slice 1 to acquire additional parts of the k-space data
  - When one slice is excited, the others are not perturbed each slice has it's own TR.
  - I Slice selection allows the efficient use of longer TR's by simultaneously acquiring many slices.

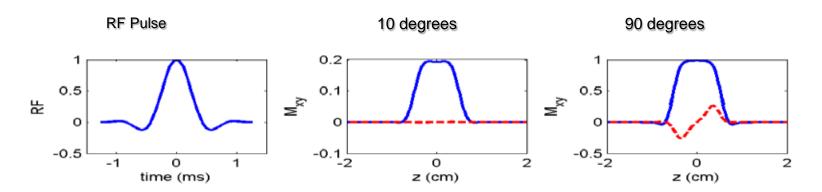
 $f = \frac{g}{2p}G_zz$  is the conversion between spectrum and the z



- K-space picture for rephasing
  - Assume RF pulse occurs instantaneously a the center of the pulse (at  $\tau/2$ )
  - Begin accumulation area in k-z after  $\tau/2$
  - By applying a negative gradient for the same duration as the last half of the pulse, the areas cancel and the k-space location in the z direction is returned to the origin.

## Larger Flip Angles

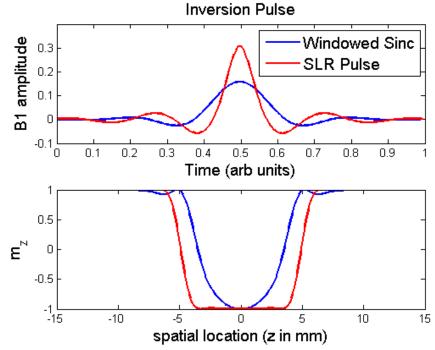
Due to the non-linearity of the Bloch equations, an exact solution would require numeric simulation of the RF pulse.



## Design of Large Flip Angle pulses

Suppose you wanted a slice selective 180 degree pulse. "sinc" pulses are not very good.

Here is an example of a 180 degree pulse designed using Fourier (small tip) methods and a large tip angle algorithm known as the Shinnar-LeRoux (SLR) method



## Example -sinc RF pulse

Consider an RF pulse:

Spectrum:

$$F\{B_1(s+t/2)\} = AT \operatorname{rect} \stackrel{\text{@}}{c} \frac{f}{BW} \stackrel{\ddot{o}}{\rightleftharpoons} \text{ where } BW = 1/T$$

Magnetization:

$$m_{xy,rot}(z,3t/2) = ign_0ATrect\overset{\text{@}}{c}\frac{zG_zg}{2pBW}\overset{\text{"o}}{\rightleftharpoons} = iam_0\text{rect}\overset{\text{@}}{c}\frac{z}{Dz}\overset{\text{"o}}{\rightleftharpoons}$$

Slice thickness: 
$$D_z = \frac{2pBW}{gG_z}$$

Flip angle (small angle):  $\alpha \cong \gamma AT$ 

#### The 3<sup>rd</sup> Dimension - Z

Spin-Warp imaging can be easily extended to 3D by using phase encoding in two dimensions (rather than 1) and frequency encoding in the remaining dimension.