

Uncovering the Dynamic CO₂ Gas Uptake Behavior of CALF-20 (Zn) under Varying Conditions via Positronium Lifetime Analysis

Ahmed G. Attallah,* Volodymyr Bon, Eric Hirschmann, Maik Butterling, Andreas Wagner, Radosław Zaleski,* and Stefan Kaskel*

Carbon dioxide (CO₂) is a major greenhouse gas contributing to global warming. Adsorption in porous sorbents offers a promising method for CO₂ capture and storage. The zinc-triazole-oxalate-based Calgary framework 20 (CALF-20) demonstrates high CO2 capacity, low H2O affinity, and low adsorption heat, enabling energy-efficient and stable performance over multiple cycles. This study examines CO2 adsorption mechanism in CALF-20 using positron annihilation lifetime spectroscopy (PALS), in situ powder X-ray diffraction (PXRD), and gas adsorption experiments under varying temperatures and humidity levels. Variable-temperature PALS experiments demonstrate that CO₂ molecules are spatially localized within the CALF-20 cages, leaving temperature- and pressure-dependent gaps. CO2 begins at cage centers, forming 1D chains, and ultimately adheres to pore walls. Interestingly, positronium intensity correlates with the Langmuir-Freundlich isotherm, reflecting gas uptake behavior. Moreover, under pure relative humidity (RH), water molecules form isolated clusters or small oligomers at low RH, transitioning to hydrogen-bonded networks above 35 %RH, significantly altering free volumes. In humid CO2 conditions, competitive interactions arise: CO2 initially disrupts water propagation, but higher RH leads to extensive water networks filling the framework. The synergy between in situ-PALS, in situ-PXRD, and gas adsorption techniques provides comprehensive insights into CALF-20's potential for efficient CO2 capture under varying conditions.

1. Introduction

During the world climate conference in 2021, 153 countries agreed on a Glasgow Climate Pact keeping the 1.5 °C target of global warming temperature within reach, with plans to revisit and tighten their emissions targets by 2030.[1] The persistent rise in human-driven CO2 emissions continues to intensify climate change, underscoring the critical need for effective carbon capture strategies.^[2] As temperatures rise and environmental challenges unfold, the demand for innovative approaches to reduce carbon emissions grows steadily. Separation and capture of greenhouse gases from dilute emissions are among the seven most important separation processes, which are responsible for 45–55% of energy consumption in industry.[3] Various approaches, such as amine scrubbing, membrane separation, cryogenic distillation, and chemical looping,[4] have been explored to address this challenge, and the adsorption of CO2 from air under different environmental conditions emerges as a promising pathway due to its costeffectiveness and operational simplicity.[5,6]

A. G. Attallah, E. Hirschmann, M. Butterling^[+], A. Wagner Institute of Radiation Physics Helmholtz-Zentrum Dresden – Rossendorf 01328 Dresden, Germany E-mail: a.elsherif@hzdr.de



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 $^{[+]}$ Present address: Delft University of Technology, Faculty of Applied Sciences, Department of Radiation Science & Technology, Mekelweg 15, JB, Delft NL-2629, Netherlands

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A. G. Attallah
Physics Department
Faculty of Science
Minia University
Minia 61519, Egypt
V. Bon, S. Kaskel
Chair of Inorganic Chemistry I
Technische Universität Dresden
D-01062 Dresden, Germany
E-mail: stefan.kaskel@tu-dresden.de
R. Zaleski
Institute of Physics
Maria Curie-Sklodowska University
Lublin 20-03 1, Poland
E-mail: radek@zaleski.umcs.pl

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It is desirable to choose adsorbents that have a high CO_2 capacity, a low N_2 and H_2O affinity, and a relatively low CO_2 adsorption heat to conserve energy while retaining good CO_2 purity, recovery, and productivity. Various types of adsorbents are being investigated for CO_2 capture, such as activated carbon, zeolites, silica gel, III metal oxides, and metal-organic frameworks (MOFs). The main challenges of well-established industrial adsorbents such as zeolites and silica gels are quite high adsorption enthalpy and co-adsorption of water, which is in most cases present in the gas mixture. On the other hand, activated carbons tend to exhibit non-polar surface and therefore low selectivity towards CO_2 .

MOFs are a class of hybrid nanoporous crystalline materials that are constructed of metal ions or clusters connected by organic linkers using modular building principles.^[13] MOFs have attracted much attention as potential adsorbents for CO₂ capture due to their high surface area, tunable pore size and shape, adjustable functionality, and structural diversity.[14-16] Different approaches are proposed for enhancing CO2 capture in MOFs including the introduction of coordinative unsaturated metal sites, amine functionalization, modulation of hydrophobicity of the pores, and the mixture thereof.[14,16-18] At the forefront of this strive is Calgary Framework 20 (CALF-20) (Zn), a zinc-based MOF, first reported by Shimizu et al.,[19] distinguished by its hydrophobicity, high scalability, and robust structural integrity. With a surface area of ca. 530 m² g^{-1[19,20]} and an average pore size of 6-7 Å,[20-22] CALF-20 (Zn) demonstrates a remarkable capacity and selective CO₂ adsorption^[23] at desired conditions.^[21] Its unique framework, composed of 1,2,4-triazole and oxalate ligands confers stability in harsh environments.[21] The necessity of capturing CO₂ under fluctuating conditions is emphasized by the varying concentrations of CO2 found in diverse industrial emissions and atmospheric contexts. CALF-20 (Zn) exhibits exceptional versatility in maintaining operational efficacy across a range of temperatures and pressures, positioning it as an indispensable tool in global efforts to mitigate CO₂ emissions.^[5] Its capabilities are boosted through the incorporation of functional groups that amplify its affinity to CO2, thereby optimizing the adsorption process.^[2] Within the PrISMa (process-informed design of tailor-made sorbent materials) platform, which includes the comprehensive screening of materials including CO₂ source, destination, capture technology, and geographical regions, CALF-20 was identified among the top-performing materials for further study in terms of more detailed process modelling and investigation for technological development.^[24] Besides that, recently the guest-induced flexibility in CALF-20 was discovered. [20,25] It was found that under humid conditions, the structure undergoes a phase transition from α -CALF-20 to β -CALF-20, accompanied by a minor contraction and adsorption of water molecules. [20] This transition was studied by a combination of in situ PXRD and total scattering. However, the authors pointed out the challenges associated with characterizing CO_2 adsorption properties of β -CALF-20 phase due to the reversibility and fast kinetics of α - β - α transitions. Further in-depth structural studies result in a number of phases, most of which are close to β -phase in terms of porosity.[25]

To this end, studying gas adsorption dynamics on the highly scalable CALF-20 MOFs is essential to gain insights into the complex processes at play. This is important for $\rm CO_2$ capture un-

der varying conditions, where the adsorption behavior may shift. This approach enables the elucidation of the dynamic interactions between gas molecules and the porous structure of CALF-20, potentially offering a clearer understanding of CALF-20's performance in different environmental conditions.^[26]

For a comprehensive insight into the CO₂ adsorption dynamics of CALF-20 (Zn), positron annihilation lifetime spectroscopy (PALS) stands out as an invaluable tool.[27-30] PALS can differentiate the subtleties of pore filling and emptying processes,[28,29,31] providing insights into the adsorption mechanisms and molecular packing within the porous materials, e.g., MOFs. This non-destructive technique has proven to be sensitive to conformational, structural, and microenvironmental transformations, which are crucial in understanding the molecular organization and transport properties within CALF-20 (Zn). More details about porosimetry by PALS are available in Appx. S1 (Supporting Information). Briefly, PALS utilizes the annihilation characteristicsquantified by lifetime and intensity-of ortho-positronium (o-Ps), a quasi-stable bound state formed by an electron and a positron with parallel spins (Appx. S1, Supporting Information), to probe the microstructural properties of materials.^[28,32-36] In porous materials, o-Ps annihilates with a lifetime determined by probability of its bound positron annihilation with an electron of parallel spin from the pore walls (in empty pores) or from adsorbents and guests within the pores. This process, known as "pick-off" annihilation, occurs at a rate that depends on the size of the pores and the accommodated substances. Consequently, the lifetime of o-Ps is exceptionally sensitive to the local electronic environment (due to adsorption, for example) and can provide detailed information about changes in size and shape of unoccupied spaces. Meanwhile, o-Ps intensity indicates the formation and trapping probabilities of o-Ps, which are influenced by pore density, connectivity, surface area, and chemistry.[37-40] In this regard, PALS can be employed as a complementary tool alongside other chemical and spectroscopic methods to characterize CO₂ adsorption in CALF-20 under environmental conditions, particularly in scenarios where some of these methods face limitations related to temperature, pressure, or humidity. Additionally, the processes of adsorption and pore filling can result in the formation of closed or restricted pores that traditional gas adsorption methods may fail to detect, whereas o-Ps can effectively identify these structures. Therefore, we believe that PALS is an informative new addition for understanding the adsorption dynamics at different conditions of CALF-20, which can lead to improvements in the material's design and functionality, making it more effective in reducing CO₂ emissions.

To advance our understanding of CO_2 capture, this study investigates the adsorption dynamics on CALF-20 (Zn) across various conditions taking advantage from PALS capabilities. This entails conducting precise treatments including a control of the sample temperature between 253 K and 373 K under vacuum, followed by in situ CO_2 adsorption experiments and full-pressure scans within the same temperature range during PALS measurements. Furthermore, we executed comprehensive full-scale humid CO_2 adsorption experiments, spanning humidity levels from 11% to 98%. Additionally, our investigation is complemented by in situ-PXRD (powder X-ray diffraction) analysis and physisorption experiments, conducted under identical conditions allowing us to compare the results with in situ gas adsorption data obtained via

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Table 1. Variations of o-Ps parameters and the corresponding pore sizes in CALF-20 at different temperatures under vacuum.

T [K]	$ au_{\mathrm{cage}}$ [ns]	$ au_{ m grain}$ [ns]	d _{cage} [nm]	d _{grain} [nm]	I _{cage} [%]	I _{grain} [%]
253	4.34 ± 0.03	33.75 ± 0.47	0.60 ± 0.002	1.53 ± 0.01	13.11 ± 0.06	9.91 ± 0.05
273	4.47 ± 0.04	34.53 ± 0.49	0.61 ± 0.002	1.55 ± 0.01	12.81 ± 0.06	10.26 ± 0.05
295	4.54 ± 0.04	33.92 ± 0.47	0.61 ± 0.003	1.54 ± 0.01	12.49 ± 0.06	10.54 ± 0.05
320	4.75 ± 0.04	34.15 ± 0.47	0.63 ± 0.003	1.55 ± 0.01	12.11 ± 0.06	11.00 ± 0.05
350	4.72 ± 0.05	33.28 ± 0.44	0.62 ± 0.003	1.53 ± 0.01	11.83 ± 0.06	11.40 ± 0.05
373	4.95 ± 0.05	33.45 ± 0.46	0.64 ± 0.003	1.53 ± 0.01	11.40 ± 0.06	11.64 ± 0.06

PALS, thus offering a holistic understanding of CALF-20's adsorption behavior.

2. Results and Discussions

Prior to any experimental procedures, the as-received sample was annealed at 423 K for 10 h under a dynamic vacuum of ca. 1.5×10^{-6} mbar, as detailed in the method section.

2.1. Thermal Effect on CALF-20 Porosity

One primary objective of this study was to investigate CO_2 adsorption dynamics in CALF-20 under different temperatures. In order to distinguish temperature effects from adsorption induced deformation, we initially conducted PALS measurements^[41] at varying temperature over a range from 253 to 373 K in a vacuum $(1.5 \times 10^{-6} \text{ mbar})$, with 2 h intervals per step. The activated sample was first heated to 373 K, then gradually cooled to 253 K at a heating/cooling rate of 2.5 K min⁻¹.

PALS analysis using PALSFit^[42] revealed that there are two resolvable o-Ps components with distinct lifetimes (ca. 4.5-4.34 ns and 33.3-34.5 ns, Table 1) in CALF-20. The short-lived o-Ps lifetime (τ_{cage}) corresponds to square cross-section channels of about 0.60-0.64 nm, closely matching the crystallographically calculated cages of CALF-20 in the α -phase (Appx. S2 and Figure S1, Supporting Information). The long-lived o-Ps lifetime (τ_{grain}) corresponding to gaps of around 1.5 nm is commonly interpreted as annihilation between grains (shown in Figure S2, Supporting Information) as noted in Table 1. This component may also include a fraction of o-Ps annihilating from delocalized Bloch states as discussed below. The analysis indicates also that about onequarter (I_{cage} + I_{grain} in Table 1) of positrons in CALF-20 form o-Ps. The results in Table 1 suggest a weak dependence of o-Ps parameters, and thus cages in CALF-20 remain in the α -phase over the 253 to 373 K temperature range in a vacuum. The cage size, as probed by o-Ps in Table 1, negligibly increases with increasing the temperature between 253 and 373 K. Therefore, any greater changes observed during CO2 dosing in the next section are likely attributable to the adsorption process.

2.2. Isothermal CO₂ Adsorption at Varying Temperatures

To explore the CO₂ adsorption course in CALF-20, we investigated its temperature-dependent adsorption using PALS. In the

following, we have discussed the lifetimes and intensities separately. The discussion of hysteresis noticed in some curves and variation of $\tau_{\rm grain}$ are detailed in Appx. S3 and Figure S3 (Supporting Information), respectively.

2.2.1. o-Ps Lifetime and Cage Filling

As CO_2 is introduced at different temperatures, significant changes occur in the o-Ps lifetimes and intensities in the cages. For clarity, the entire course of changes can be best traced in the pressure range of 5–1000 mbar at room temperature (**Figure 1a** and Appx. S3, Supporting Information). The change in τ_{cage} consists of three stages:

- Lack of any change in lifetime below a certain pressure threshold value.
- Sigmoidal decrease in lifetime that can be well described by equation

$$\tau_{\text{cage}} = \frac{\tau_{\text{empty}} - \tau_{\text{filled}}}{1 + e^{\frac{\ln\left(\frac{p/p_{1/2}}{k}\right)}{k}}} \% + \tau_{\text{filled}}$$
(1)

where p is CO₂ pressure, $\tau_{\rm empty}$, $\tau_{\rm filled}$, $p_{/_2}$, and k are fitting parameters meaning o-Ps lifetimes before and after sigmoidal change, CO₂ pressure at the middle of the sigmoidal change, and width of sigmoidal change, respectively.

1) Slow logarithmic decrease in τ_{cage} with p.

This pattern is observed at all studied temperatures; however, the sigmoidal change occurs at a very low pressure of several mbar at 253 K, while it spreads above the maximum ${\rm CO_2}$ pressure of 1000 mbar at 350 K and 373 K. Nevertheless, Equation 1 was fitted for all temperatures (Figure 1a and Table S1, Supporting Information). Given only the partial curves at T=350 and 373 K, fitting was challenging. Therefore, we used fixed values from Figure S9 (Supporting Information) in all fits by extrapolating $\tau_{\rm filled}$, assuming a linear change with temperature, and weighted the values with errors to ensure consistency.

The decrease in o-Ps lifetime in MOFs often results from complex interactions of o-Ps with metal nodes. [36] However, in CALF-20 with Zn, where no unpaired electrons are present, there is no evidence of such effects—such as chemical quenching, spin conversion, Ps inhibition, or Ps complex formation. Instead, the predominating relevant Ps interaction is the pick-off process, which is directly related to the size of the free volumes. Hence,



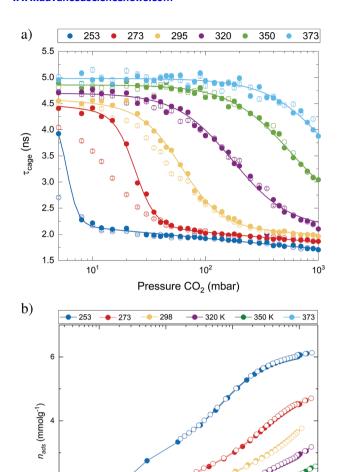


Figure 1. a) Pressure dependence of lifetimes of ortho-positronium trapped in cages of CALF-20 during $\rm CO_2$ sorption at different temperatures, where sigmoidal functions were fitted to adsorption data. b) Variable temperature $\rm CO_2$ physisorption isotherms in CALF-20. Adsorption data are solid circles and desorption data are open circles.

Pressure CO2 (mbar)

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the smaller lifetime implies varying degrees of cage filling with CO_2 at different temperatures. Because the molecular diameter of CO_2 (ca. 0.3×0.6 nm)^[43] is comparable with the size of the CALF-20 cage, even a single CO_2 molecule acts as a blockage of the pore and forms a potential barrier for o-Ps. So, the o-Ps probe detects an average distance between the blockages, which can consist of one or more molecules. In this context, the sigmoidal change observed in Figure 1a is the shortening of the distance between blockages from the infinity (open channel) to a certain minimum distance of Van der Waals radius due to filling cages with subsequent CO_2 molecules. This distance between molecules can be calculated assuming a cuboid shape of free volume (Figure 2a and Table S3, Supporting Information). At pressure just above the pore filling step in the τ_{cage} isotherm (when lifetime decreased by 99% of its total change) this gap is 0.45-

0.50 nm long, i.e., of a size that can accommodate one or two ${\rm CO_2}$ molecules. Below 298 K, where the step in $au_{\rm cage}$ is observed at the pressure below 60 mbar further increase in pressure causes a reduction in the average gap length ($L_{\rm gap}$), which rate seems to be independent of temperature. Therefore, at 1000 mbar L_{gap} increases with temperature at a rate of 1.5 pm K⁻¹ starting from 0.38 nm at 253 K (Figure 2a). These two effects, L_{pap} decreases with pressure and increases with temperature, reflect changes of the CO₂ arrangement in the cages of CALF-20. There is another possibility that o-Ps exerts pressure on CO₂ molecules, repels them, and thus L_{gap} increases. This effect, known as "bubble" formation, is observed in bulk liquids.[44] However, as explained in the "main cages and Ps bubble" (Supporting Information), Ps bubble is not expected in the microporous CALF-20 cages. Therefore, it is not justified to ascribe the changes in the gap size to the interaction of CO₂ with o-Ps.

The sigmoidal change related to cage filling with CO₂ can be described by two parameters; i) $p_{1/2}$, the CO₂ pressure at which $\tau_{\rm cage}$ is halfway from the maximum value ($\tau_{\rm empty}$) to the minimum value ($\tau_{\rm filled}$) when the $\tau_{\rm cage}$ change is finished and ii) k, the dimensionless parameter describing the rate of the $\tau_{\rm cage}$ change, where

- I. $p_{\frac{1}{2}}$ changes exponentially with temperature from tens of mbar at 253 K to over 1000 mbar at 373 K (Figure 2b).
- II. k increases linearly (i.e., width increases exponentially like $p_{1/2}$ because pressure in Equation 1 is introduced in the logarithmic scale, so for greater $p_{1/2}$ the same k value means greater width of the sigmoidal change) with temperature in the range of 253–320 K and then it seems to stabilize for higher temperatures (Figure S8, Supporting Information). It should be remembered that the sigmoidal change above 320 K is incomplete within the measured pressure range and these results can be inaccurate.

2.2.2. Ps Intensities and Adsorption Mechanism

The intensities of both o-Ps components undergo the sigmoidal change with increasing CO₂ pressure. The intensity of the intergranular component (I_{grain}) changes in almost the same way as the o-Ps lifetime in cages (Figures S5 and S6, Supporting Information), following a function like that in Equation 1. The strong correlation between pressure dependences of fitted au_{cage} and I_{grain} (Figure 2b) confirms that of o-Ps annihilating between grains originate from inside the cages. Filling the cages with CO₂ results in blocking their outlets, preventing o-Ps migration to the intergranular spaces. Possibly, at higher temperatures the mobility of CO₂ increases and it can easily migrate to the center of the cage, leaving longer open ends that increase o-Ps intensity outside grains when cages are filled (I_{filled} in Table S2, Supporting Information). This effect can be also related to the greater kinetic energy of o-Ps at higher temperatures. Alternatively, I_{grain} in fact contains signals from o-Ps in delocalized Bloch states, which spread over the periodic structure of CALF-20 instead of localization in a single channel. In such a case, their formation and annihilation are affected by the presence of CO₂ within the cages. This is thoroughly explained below.

More intriguing is the change in o-Ps intensity in cages (I_{cage}) with CO_2 pressure (Figure S6, Supporting Information). One

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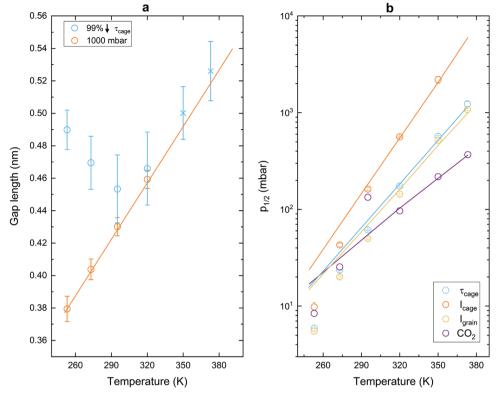


Figure 2. Temperature dependence of a) the length of the gap between CO_2 molecules blocking the cages of CALF-20 calculated from the o-Ps lifetime assuming cage side obtained without gas when it is reduced by 99% of total change during the sigmoidal change and at 1000 mbar (circles), values for 350 and 373 K are extrapolated (crosses) and not applicable for 1000 mbar. b) CO_2 pressure at which the fitted o-Ps lifetime in cages (τ_{cage}), fitted intensity of the cage (τ_{cage}) and intergranular (τ_{cage}) components, and τ_{cage} 0 uptake during the physisorption isotherm are halfway between the maximum values and the minimum values.

would expect it mirrors $I_{\rm grain}$, since o-Ps, unable to escape the cages, will annihilate with the lifetime of $au_{\rm cage}$ instead of $au_{\rm grain}$. Although the expected increase in I_{cage} occurs and is approximately equal to the decrease in I_{grain} (Figure S5, Supporting Information), it is clearly shifted towards higher pressures (Figure 2b). This shift ranges from almost two times greater pressure at 253 K to over three times at 320 K, and possibly more at higher temperatures, but this exceeds 1000 mbar and cannot be determined with sufficient accuracy. An additional channel of Ps migration outside grains is expected to justify high $I_{\rm grain}$ from a relatively large CALF-20 crystals (Appx. S5 and Figure S12, Supporting Information). The origin of additional "intergranular" o-Ps may be o-Ps from delocalized (hot) Bloch states [39,45] spreading outside the grains, where disappearance of the periodic structure of CALF-20 causes their localization. If this is the main source of I_{grain} , its decrease is related to $\tau_{\rm cage}$, reflecting a disruption of the periodic structure. In contrast, "intragranular" o-Ps can hardly get outside grains and changes in $I_{\rm cage}$ are only slightly related to the migration of o-Ps and reflect other processes in the cages that are hardly related to o-Ps migration. The localized and delocalized o-Ps states are discussed in the "o-Ps parameters during in situ CO₂ adsorption" (Supporting Information).

Notably, the changes in $I_{\rm cage}$ and $I_{\rm grain}$ over the pressure course aids in predicting the adsorption mechanism within CALF-20. To explore the mechanism, we focus on their sigmoidal changes during ${\rm CO_2}$ adsorption. For clarity, we have selected the $I_{\rm cage}$ and

 $I_{\rm grain}$ curves at 295 K as a function of adsorption pressure and divided the sigmoidal part of $I_{\rm grain}$ (the fast-changing parameter) into three sections; A, B, and C, while point D is assigned to represent the change beyond the sigmoidal transition (**Figure 3**, left), to guide the discussion.

From vacuum until point A in Figure 3, left, both $I_{
m cage}$ and $I_{
m grain}$ are unchanged. Since o-Ps intensity correlates to pore surface, hence the unchanged intensities suggest no pronounced change in inner cage surface. This indicates that few CO₂ molecules are adsorbed at this pressure range and probably they locate in the center of the cage. As adsorption develops reaching point B (middle of the $I_{\rm grain}$ sigmoidal transition), $I_{\rm grain}$ declines but $I_{\rm cage}$ still unaffected, highlighting the shift mentioned above. Since the most probable binding sites of CO₂ in CALF-20 are in the middle of the cages, where the oxygen atoms of the oxalate groups in the framework serve as primary binding sites, [19] the introduction of CO₂ will disrupt the periodicity required to form the Bloch states and, therefore, reduces total Ps formation (Figure S7, Supporting Information), but only in the part that can localize outside grains (i.e., I_{grain}). This initiate changes in I_{grain} , while formation of Ps in the bulky part (I_{cage}) is not affected and the shift is noticed. At this point, gaps between CO₂ molecules are expected to get smaller inside the cage as au_{cage} decreases (Figure 1a). When advancing to point C, additional CO₂ molecules are introduced, which may form 1D chains along the channel. These chains further disrupt the periodicity of the cage. Consequently, the delocalized o-Ps



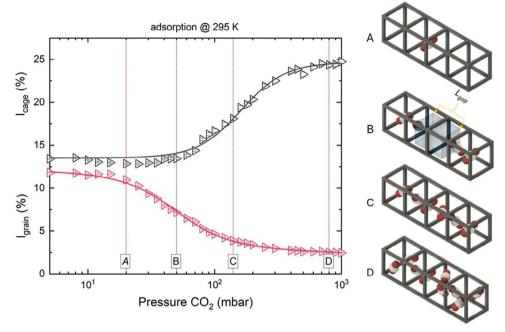


Figure 3. Left) changes of I_{cage} and I_{grain} as functions of adsorption pressure in CALF-20 at 295 K. Right) Sketch illustrating the anticipated CO₂ adsorption mechanism within CALF-20 based on o-Ps intensities.

does not actually form. $I_{\rm cage}$ starts to increase because, in addition to electrons from CALF-20, electrons from CO₂ contribute more and more to the formation of o-Ps. Finally, during stage D, $I_{\rm cage}$ reaches its maximum and $I_{\rm grain}$ approaches its minimum. In this region, $I_{\rm cage}$ increases by ca. 11.4%, while $I_{\rm grain}$ decreases by only 9.3% when compared to vacuum. The overcompensated increase in $I_{\rm cage}$ (only at 295 K and below) reveals an increase in the total o-Ps formation (also seen in Figure S7, Supporting Information). This supports the argument that the changes in $I_{\rm cage}$ and $I_{\rm grain}$ originate from independent sources. The obvious cause of the increase in o-Ps intensity is usually an increase in surface area. Hence, the increase of $I_{\rm cage}$ along stage D likely re-

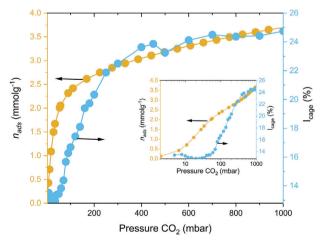


Figure 4. Experimental CO_2 physisorption isotherm and variation of o-Ps intensity in cages of CALF-20 from in situ PALS during CO_2 adsorption at 295 K. The same figure is plotted in the logarithmic scale in the inset.

sults from the filling of numerous free spaces with CO₂, which not only occupies the central regions of the cages but also fits into the free spaces between the organic "pillars" connecting the metallic nodes in CALF-20. Moreover, the increase in I_{cage} is not accompanied by an increase in τ_{cage} , which can be explained if o-Ps is trapped only in the central regions of the cages (between CO₂ molecules) rather than in the smaller spaces between the "pillars." This is justified because the smaller spaces between the pillars are connected to the larger cage centers, making it energetically favorable for o-Ps to migrate from the smaller to the larger free spaces. Additionally, a shift of the CO₂ molecules from the cage center to the cage sides at high pressures is also expected and it would lead to more empty spaces in the center that can host o-Ps, which may also contribute to the increase in I_{cage} (stage D in Figure 3). The discussion about the adsorption mechanism based on the intensities is summarized in the sketch in Figure 3, right.

Detailed discussions about variations of the total o-Ps intensity are provided in the "o-Ps intensity" section in Appx. S3 (Supporting Information).

2.2.3. Cage o-Ps Intensity as a CO2 Uptake Indicator

As discussed above, Figure 1a illustrates the pore filling mechanism in CALF-20, derived from in situ PALS data, and provides unique complementary information to the adsorption data (Figure 1b). In Figure 1b, the volumetric CO_2 adsorption experiments at the same temperatures and pressure range reveal a Langmuir-type isotherms with the uptake at 1000 mbar ranging from 1.4 mmol $\rm g^{-1}$ at 373 K to 6.0 mmol $\rm g^{-1}$ at 253 K. This shape of isotherm in Figure 1b and Figure S10 (Supporting Information), measured at the temperatures of PALS studies, shows

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Table 2. Fitting parameters obtained for the pressure dependence of I_{cage} at various temperatures T (Figure S6, Supporting Information): p is CO_2 pressure, $p_{1/2}$ is the pressure at the middle of the sigmoidal change, k is the width of sigmoidal change, I_{empty} is intensity in cages without CO_2 , I_{filled} (*fixed on maximum value for all T) is intensity in cages with CO_2 .

T [K]	p _½ [mbar]	$\Delta p_{1/2}$ [mbar]	k	Δk	I _{empty} [%]	ΔI_{empty} [%]	I _{filled} [%]	ΔI_{filled} [%]
253	9.8	1.1	0.341	0.111	14.04	0.57	24.67*	0.14
273	42.9	1.7	0.246	0.036	13.79	0.49	24.67*	0.14
295	162.0	4.9	0.398	0.026	13.51	0.42	24.67*	0.14
320	563.1	8.5	0.496	0.015	13.19	0.37	24.67*	0.14
350	2198.7	187.3	0.628	0.045	12.81	0.37	24.67*	0.14

an expected shape, which is in agreement with Shimizu and co-authors.^[19] The isosteric heats of adsorption in Figure S11 (Supporting Information), calculated using Clausius-Clapeyron equation^[46] overestimate the adsorption heat at low loadings (see Appx. S4, Supporting Information). At high loadings, the values are in range of 37–40 kJ mol⁻¹, which is in line with earlier reported data.

Notably, the evolution of $\tau_{\rm cage}$ and $I_{\rm cage}$ provide the complementary information on the pore filling mechanism indicating nearly complete pore filling before the saturation is reached. This is supported by the behavior of $p_{1/2}$ in Figure 2b derived from fitting of $\tau_{\rm cage}$ and $I_{\rm cage}$ and from ${\rm CO_2}$ physisorption, where it increases with temperature but with different rates. The $I_{\rm cage}$ (Figure S6, Supporting Information) directly correlates with the ${\rm CO_2}$ adsorption isotherm (Figure 1b) as shown in Figure 4 for isotherms at 295 K as an example. This encouraged us to explore the $I_{\rm cage}$ in more detail. The shape of $I_{\rm cage}$ resembles that of a type I isotherm, $^{[7,19]}$ indicating it reflects the amount of ${\rm CO_2}$ uptake. It would be beneficial to analyze if the change in $I_{\rm cage}$ can be traced by a known adsorption isotherm model. In Figure S6 (Supporting Information), $I_{\rm cage}$ exhibits a sigmoidal shape, which can be best described by the following equation:

$$I_{\text{cage}} = \frac{I_{\text{empty}} - I_{\text{filled}}}{1 + \left(\frac{p}{p_{1/2}}\right)^k} + I_{\text{filled}}$$
(2)

here, the parameters have the same meaning as in Equation 1. Equation 2 can be rearranged in the following form (derivation is presented in Equations S2–S5, Supporting Information):

$$I_{\text{cage}} = \frac{I_{\text{filled}} \times R \times p^k}{1 + R \times p^k} + \frac{I_{\text{empty}}}{1 + R \times p^k}$$
(3)

with $R = (\frac{1}{p_{1/2}})^{k}$

Equation 3 can be regarded as a modified version of Langmuir–Freundlich (Sips) isotherm, [47] which takes the form:

$$q_e = \frac{q_{\rm m} \times K_{\rm s} \times C^n}{1 + K_{\rm s} \times C^n} \tag{4}$$

where q_e represents the amount of gas adsorbed at equilibrium, q_m is the maximum adsorption capacity, K_s is the proportional constant or Sips constant, C is the equilibrium gas concentration, and n is the dimensionless Sips exponent. By comparing

Equations 3 and 4 we observe that $I_{\rm cage}$ follows the Langmuir–Freundlich isotherm with an additional term, $\frac{I_{\rm empty}}{1+R\times p^k}$. This term, involving $I_{\rm empty}$, likely arises from the formation of o-Ps on the adsorption-free sites on the walls of CALF-20 at a given pressure. The terms in Equation 3 can be interpreted as representing ${\rm CO}_2$ adsorption in the centers of the cages, along with the gradual filling of the open framework walls by ${\rm CO}_2$ adsorbed. Additionally, the fitting exponent k in $I_{\rm cage}$ serves as a heterogeneity factor, similar to n in Sips isotherm. The reliability of the fitting parameters is shown in Table 2. However, at 373 K, the lack of sigmoidal behavior in $I_{\rm cage}$ across all pressure values (Figure S6, Supporting Information) leads to unreliable fitting, and thus, the corresponding values have been omitted from Table 2 and other related figures.

Although we illustrated that CO₂ adsorption isotherm (as seen by I_{cage}) adheres the Langmuir–Freundlich isotherm, Figure \$11, supporting information and other studies have indicated that the CO₂ adsorption is better described by a dual-site Langmuir isotherm.^[7] Notably, the Langmuir–Freundlich isotherm accounts for heterogeneous adsorption sites, whereas the Langmuir isotherm assumes uniform adsorption sites. The fitted heterogeneity factor k for I_{cage} is depicted in Figure S8 (Supporting Information). It follows very closely the k values for au_{cage} and I_{grain} , which confirms that despite the pressure shift, change in these parameters result from the same effect. Additionally, it is important to highlight the difference between I_{cage} and n_{ads} in Figure 4 below 250 mbar (Figure 4, inset). In particular, below 50 mbar I_{cage} does not change with increasing n_{ads} . This is possible if CO₂ molecules are distant from the CALF-20 skeleton, where thermalized positrons form o-Ps at this stage, i.e., CO2 molecules are presumably located in the cage center. In the range between 50 and 250 mbar, this behavior gradually changes to the state where CO₂ molecules start taking part in the o-Ps formation, i.e., they locate in the neighborhood of the CALF-20 framework more and more until both I_{cage} and n_{ads} follow the same dependence.

2.3. Humidity and Humid CO2 Adsorption

This section is focused on structural changes in CALF-20 during CO_2 adsorption in a humid environment, simulating CO_2 capture under ambient conditions. The experiment was conducted at RT with 970 mbar of CO_2 across relative humidity levels from 11% to 98%, and results were compared to CALF-20 exposed to humidity alone. The 0% RH was realized by using molecular



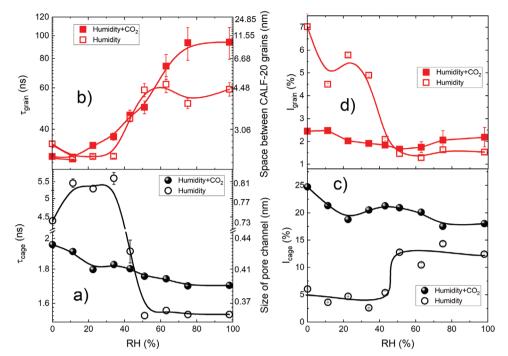


Figure 5. Impact of room temperature pure humidity exposure and 970 mbar CO_2 adsorption in CALF-20 as a function of relative humidity on: a) cage o-Ps lifetime (τ_{cage})—right axis shows the derived size of the square cross-section of cages, b) grain o-Ps lifetime (τ_{grain}), together with the corresponding size of gaps between calf-20 grains (right axis), c) o-Ps intensity in cages (I_{cage}), and d) o-Ps intensity between CALF-20 grains (I_{grain}). Lines are shown to guide the eye.

sieves to adsorb water from the sample chamber producing either dry air (reference for pure humidity experiment) or dry CO_2 (reference for humid CO_2 experiment). During the humid CO_2 exposure, both humidity and CO_2 were simultaneously introduced to the sample to mimic competition between $\mathrm{H}_2\mathrm{O}$ and CO_2 . The sample was reactivated in vacuum at 423 K for 4 h after each exposure, as detailed in the experimental section.

2.3.1. Water Vapor Adsorption

Adsorption: Adsorption of water vapor on CALF-20 sample at 298 K was conducted (Appx. S6 and Figure S13, Supporting Information (bottom)) in order to show the correlation between the PALS and adsorption data. The isotherm is similar to that reported by Shimizu and co-workers. [19] The adsorption branch indicates the steep water uptake in the pressure range $p/p_0=0$ –0.2. At higher pressure, less water is adsorbed in the pores. The desorption branch follows the adsorption showing a narrow hysteresis in the whole pressure range, originating from the slow adsorption kinetics of water vapor.

2.3.2. Ps Components during Pure Humidity Exposure

Cages: In the pure humidity run at RT, a step increase in $\tau_{\rm cage}$ of 1 ns between 0% and 11% RH is observed (**Figure 5a**). The decrease in o-Ps pick-off annihilation rate (increase in lifetime) shows that water somehow screens the interaction between o-Ps (already formed in the bulky parts) and the electrons on the

pore walls. This suggests that, unlike CO₂, adsorbed water is initially located not in the center but closer to the cage's skeleton. This observation is supported by the snapshots in ref.[48] and the shortest distances between the guest molecules (CO₂ or H₂O) inside the cage and the CALF-20's skeleton in refs.[19,49] For H₂O, the shortest distance was estimated to be 2.8 Å using radial distribution function (RDF), [49] whereas for CO₂, it was calculated to be 3.03 Å through atomistic grand canonical Monte Carlo simulations. [19] After that, CALF-20 exhibits two distinct regions of $\tau_{\rm cage}$ with 35–50% RH as a transition region (Figure 5a). Below 35% RH, $\tau_{\rm cage}$ fluctuates around \approx 5.5 ns. As noted by Oktavian et al.,[49] the RDF analysis suggests that at low RH, water molecules position themselves near the oxygen atoms of the oxalate linkers and further away the Zn atoms, resembling a hydrophobicity. Water molecules thus prefer direct cooperative contact, forming zero or at most one hydrogen bond, resulting in isolated water molecules or small oligomers being adsorbed. [22] This minimal hydrogen bond concentration at low RH slightly affects au_{cage} , with its variability screening o-Ps interactions with the few water molecules present in the cages. Above 35% RH, $\tau_{\rm cage}$ shows an abrupt reduction to ≈1.5 ns between 35% and 50% RH and stabilizes afterwards at this value regardless of the RH level. The steep transition in $\tau_{\rm cage}$ between 35 and 50% RH may indicate the onset of cage filling, aligning with water isotherm in Figure S13 (Supporting Information) (bottom), which starts to saturate at almost the same values of relative pressure. However, the $au_{
m cage}$ value of 1.5 ns at RH > 50% is shorter than the o-Ps lifetime of 1.87 ns in bulk water (Ps bubble) or 2.15 ns in water confined in SBA-3 with 2 nm pores. [30] The lifetime of 1.5 ns corresponds to the residual water-free spaces with size of ca. 0.3 nm that are

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present in the cages.^[27] These water-free spaces are expected because studies have demonstrated that at higher RH levels, water form complex hydrogen-bonding networks and quasi-1D molecular chaines within the MOF cage,^[26,50] and spaces may exist between water chains.

Between 0% and 35% RH, $I_{\rm cage}$ (Figure 5c) gently decreases reaching 3% at 35% RH. Like the discussion of $\tau_{\rm cage}$, this reduction in $I_{\rm cage}$ may originate from the low amount of adsorbed water. Above 35% RH, $I_{\rm cage}$ starts to rise, reaching 12–15% at RH \geq 50%. According to Llewellyn et al., [22] at high values of partial pressure p/p_0 (\propto RH), percolation occurs, causing neighboring water chains to interconnect. These interconnected water chains begin to form clusters around the inner pore surface, [22] potentially creating spaces in between. This increase in the surface area along with detachment of water molecules from walls, which gives o-Ps an opportunity to repel them, could explain the observed growth of $I_{\rm cage}$.

Inter-Grain Spaces: $\tau_{\rm grain}$ in Figure 5b in the absence of CO₂ also exhibits a three-region dependence on RH, i.e., 0-35%, 35-50%, and 50–98%. At RH below 35%, $\tau_{\rm grain}$ slightly decreases, likely due to the presence of some water molecules in the intergranular spaces, but mostly in the cages, where they affect the delocalized o-Ps reducing its intensity and, in consequence, the average lifetime of migrated and delocalized o-Ps. Next, it increases linearly up in the RH range between 35% and 50%, and then stabilizes around 60 ns. This increase in $\tau_{\rm grain}$ indicates an enlargement of the spacing between CALF-20 layers or grains. This might be a result of enhanced structural flexibility, leaving larger spaces between grains, but more likely water condenses (supported by the increasing branch of water adsorption isotherm in Figure \$13, Supporting Information) at high humidity level and fills the smaller spaces in the grain spaces leaving the larger spaces unfilled. This can lead to a longer lifetime for o-Ps annihilating in the residual larger spaces. Alternatively, the suppression of remigration of o-Ps to the cages is responsible for the increase in au_{grain} , which is described in detail in the "Humidity and Humid *CO*₂ in CALF-20" (Supporting Information).

 $I_{\rm grain}$ drops between the dry state (0% RH) and the first humidity point (11%). This drop is not compensated in $I_{\rm cage,}$ and it indicates lower o-Ps formation as previously discussed. Again, this would mean that the inner pore surface is altered and the number of available sites to form the delocalized o-Ps becomes less. Subsequently, $I_{\rm grain}$ varies between 4% and 6% at RH ranging from 11% to 35% (Figure 5d), likely due to some adsorbed water molecules between grains. At RH = 43%, $I_{\rm grain}$ falls sharply to about 2%, and continues to decrease to 1.2–1.5% at higher RH levels. Aligning with the discussion of $\tau_{\rm grain}$, this abrupt reduction in $I_{\rm grain}$ can be attributed to the presence of residual larger intergranular spaces, while the smaller intergranular spaces, which constitute the majority, have already been filled.

2.3.3. Ps Components during Humid CO₂

 $\it Cages$: In this part, $\it CO_2$ at 970 mbar is introduced together with humidity at RT. The 970 mbar pressure was specifically chosen to observe the most pronounced effects on the material's porosity and adsorption dynamics by PALS. While this pressure exceeds typical environmental conditions, it enhances the differential and the properties of the p

tiation of changes in o-Ps lifetimes and intensities, thereby facilitating a clearer interpretation of the underlying adsorption mechanisms. This approach provides insights that can later be extrapolated to more realistic conditions at lower CO₂ pressures. Once the RT humid CO_2 is applied, τ_{cage} decreases, but with distinct characteristics. Initially, it drops from 1.95 to 1.8 ns between 0 and 40% RH (Figure 5a). The corresponding cage size changes from 0.43 to 0.41 nm. Although it has been noted that humid CO_2 induces a phase change from α -CALF-20 to β -CALF-20 above 23% RH,^[20] from Ps lifetime we cannot unambiguously attribute the change in cage sizes solely to this phase change due to the progressive cage filling. In humid CO2 at low RH, the uptake of CO₂ is favored, [19] but water molecules coexist. The oxalate groups in the framework serve as primary binding sites for both H₂O and CO₂;^[26] however, their spatial distribution within the pores differs^[19,49] as stated earlier, i.e., H₂O molecules are located near the framework's skeleton, which possibly due to variations in molecular interaction of CO2 and H2O with the framework, interaction strengths, etc. Consequently, there will be a competitive interaction between H₂O and CO₂ in CALF-20, due to their shared preference for the same void regions and similar interaction energies with the framework. This competitive adsorption phenomenon at low RH is consistent with the recent findings of Krishna and van Baten, [51,52] who demonstrated through molecular simulations that CO2 molecules face reduced competition than expected from uniform adsorption models, such as the ideal adsorbed solution theory (IAST) due to the segregated adsorption of H₂O, highlighting the preference of CO₂ uptake at low RH. Worth noting, recent work by Moreton et al.^[53] of varied temperature (298 K, 308, and 318 K) CO2-H2O mixture adsorption on CALF-20 thin films have illustrated that the deviation from IAST theory is only pronounced at RT, while it is of less importance at higher temperatures. Aligning with the work of Krishna and van Baten, [51,52] the $au_{\rm cage}$ in Figure 5a at low water content (up to 40% RH) suggests that the interaction between water and CO₂ (or oppositely—the interaction screening) may promote partial pore filling in addition to a possible phase change. This implies that the blockages created by CO2, as described previously, are not entirely closed, allowing water molecules to intrude between them. Indeed, RDFs and computational snapshots in refs. [51,52] confirm that H₂O-H₂O pairs are typically located between CO₂ molecules at distances of approximately 3 Å, while CO₂-H₂O pairs are spaced farther apart at around 8 Å, and CO₂-CO₂ separations are approximately 6.5 Å. In this context, τ_{cage} serves as a probe for the average empty spaces within the pores. This suggests that within the pores, water molecules can rapidly exchange between the cages without adhering to each other or to the pore walls because the pores are very small. Due to the co-adsorption, H₂O occupies energetically favorable positions, which appear to be situated between the CO₂ molecules. Therefore, the distinguishability between the phase transition and pore filling cannot be made with a reasonable precision. An increased tendency for water molecules to form hydrogen bonds among themselves was observed at RH > 40% in CALF-20.[19,20,22,26] Consequently, more interconnected hydrogen-bond networks are formed involving a greater number of H₂O molecules. The strong hydrogen bonding for H₂O-H₂O pairs at high RH may lead to a substantial increase in the competition between CO₂ and H₂O, as also corroborated by Krishna and van Baten.[51,52] They reported that under

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these conditions, CO2 molecules are forced to share adsorption sites with chains of hydrogen-bonded H₂O molecules appearing at distances of about 3 Å, leading to a significant reduction in CO₂ uptake and further deviation from IAST predictions. Hence, we hypothesize that the formation of water chains, along with the increased preference for water uptake over CO₂ at high RH levels, is manifested in the reduction of au_{cage} to approximately 1.8 ns.

The presence of both CO₂ and humidity leads to smaller, but distinct, changes in $I_{\rm cage}$ than in pure water vapor indicating that the micropores are permanently filled with the guests. I_{case} is lower than in the dry CO2 at 970 mbar (0% RH)—due to less Ps formation because of changes in the chemistry of inner cage surface when water exist, but it hardly changes with humidity up to ca. 70%. Above this pressure (at 10³ mbar at 295 K in Figure S6, Supporting Information) it still remains clearly above the humidified sample. This confirms our previous interpretation that water molecules are initially well dispersed in the cages filled with CO₂ and only start to cluster at high humidity.

Inter-Grain Spaces: In the presence of humidity together with CO_2 , τ_{grain} behaves slightly different than in pure water adsorption experiment. When humidity and CO_2 coexist, au_{grain} starts changing already at 23% RH and increases linearly with RH until stabilizing at ca. 90 ns between 75% and 98% RH (Figure 5b). The transition differs from pure humidity not only due to the wider RH range, but also the different slope in au_{grain} changes and much greater value at high RH. It seems that the intergranular spaces are much larger at RH > 75%. The increased spaces may imply structural flexibility; however, the current results alone do not provide strong evidence to support this. Alternatively, different interactions between the humid gas and the delocalized o-Ps migrating to intergranular spaces may alter the ratio of o-Ps reentering the cages. This results in fewer o-Ps annihilating with the cage lifetime, while more o-Ps remain confined in the grain spaces, where they annihilate with longer lifetimes due to the larger size of these spaces.

In humid CO_2 adsorption, I_{grain} in Figure 5d starts at a low level to decrease with increasing RH until 50%. This decline may result from both hindered migration of o-Ps out of the cages and alterations in the surface chemistry of the outer particle surfaces facing the grains. It remains greater than in the humid sample without CO₂ at higher RH. This again could mean that more o-Ps are being confined in intergranular spaces due to the humid gas, where CO₂ accumulating between the grains, as water preferentially fills the cages at high RH. This is supported by $I_{\rm grain}$ at RH > 70% approaching its values observed in dry CO_2 (at 0% RH). However, positronium formation may be affected by other than geometrical factors, e.g., a high probability of Ps formation on the surface of water clusters or even near to their surface (if Ps can migrate outside to the inter-grain spaces).

2.4. Variable Temperature PXRD in the Controlled Atmospheres

To confirm the phase composition of the pristine sample at room temperature and under various temperatures and gas atmospheres explored by PALS and CO₂ physisorption, PXRD patterns were collected for the CALF-20 sample (Appx. \$7 and Figure S14, Supporting Information), which was exposed to dry and humid nitrogen and carbon dioxide gases. All PXRD data were analyzed using Pawley method, and the results are shown in Figure 6. The first measurement, conducted on the CALF-20 sample at 298 K in a dry nitrogen flow and prepared under ambient conditions, indicates the presence of the pure CALF-20- β phase (Figure 6i). Increasing the temperature to 320 K induces a phase transition from CALF-20- β to CALF-20- α , primarily reflected by an expansion along the b-axis and simultaneous contractions along the a- and c-axes. This finding is significant as it suggests that desorbing water and purging the CALF-20 cages require minimal energy, in accordance to observations of Chapman et al.[20] in their helium flow experiments at room temperature. At higher temperatures, only the CALF-20- α phase was maintained. These results also support the PALS findings in Table 1, where the CALF-20- α phase was observed at T > RT. In the case of wet nitrogen flow, only the CALF-20- α phase was observed across the entire temperature range (Figure 6iii). Notably, the observed variations in unit cell parameters are consistent with in silico data previously reported by Maurin et al. [54] When exposed to dry CO₂, no phase transition was initiated, and the CALF-20- α phase was observed across the entire temperature range (Figure 6ii). Interestingly, wet CO₂ also did not induce a phase transition, even at 298 K, [54] indicating strong host-guest interactions between CO₂ and the CALF-20- α phase (Figure 6iv).

The data obtained provide essential insights into the behavior of CALF-20 under various conditions, specifically:

- Under ambient conditions (298 K and 40–60% RH), CALF-20 exists in its β -phase (smaller pores, partially filled with water);
- At elevated temperatures (320 K and above), CALF-20-β transitions to CALF-20- α , regardless of the gas type or humidity;
- At ambient temperature in the presence of pure CO₂, interactions between CO₂ and CALF-20-α are stronger compared to water, resulting in no observed phase transition.

3. Conclusions

This study provides comprehensive insights into the thermal effects on the porosity of CALF-20(Zn) under dry and humid CO2 adsorption conditions, emphasizing the complex interplay between temperature, humidity, and guest interactions. In situ-PALS data reveal that o-Ps lifetimes in CALF-20 vary from 253 K to 373 K, with only a modest 4% increase in cage size under vacuum, indicating that structural integrity is preserved for effective CO2 adsorption.

The analysis shows a sigmoidal change in o-Ps lifetimes and intensities during CO₂ adsorption, highlighting the material's responsiveness to temperature and pressure variations. Notably, PALS results indicate the CO₂ adsorption results in some empty spaces or gaps inside the CALF-20 cages, that are temperatureand pressure-dependent. The correlation between PALS measurements and volumetric CO₂ adsorption confirms that CALF-20 exhibits Langmuir-type isotherms in the studied temperature range.

In humid environments, the competitive interactions between water and CO₂ significantly affect the course of adsorption on CALF-20. Depending on the humidity, three clearly distinct states of the system can be observed: dispersed water molecules, transition state, and water clustering. Above the transition, the formation of interconnected water chains around the inner pore



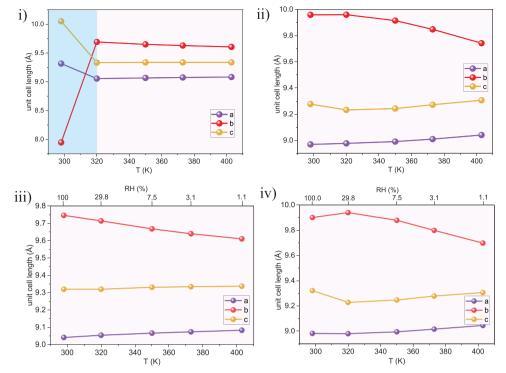


Figure 6. Unit cell axes evolution in CALF-20, measured at different temperatures in i) dry nitrogen; ii) dry carbon dioxide; iii) wet nitrogen; iv) wet carbon dioxide (conditions for α -CALF-20 are labelled yellow, β -phase light blue).

surfaces is observed, along with possibly enhanced structural flexibility and larger inter-grain spaces filled with condensed water. When CO_2 is introduced alongside humidity, the size of adsorbate-free volumes in the cages initially decreases but stabilizes as hydrogen-bond networks form among water molecules at high RH. This indicates competitive interactions for binding sites within CALF-20's framework. Variable temperature PXRD confirms that CALF-20 transitions from β -phase to α -phase at elevated temperatures (320 K and above), independent of gas or humidity conditions.

Overall, the integration of PALS, gas and water physisorption, and in situ PXRD analyses in this work elucidates the fundamental mechanisms governing CO_2 adsorption in CALF-20 while underscoring the importance of environmental factors in optimizing its application for carbon capture technologies. Future research should focus on long-term stability and performance under varied operational conditions to fully reveal CALF-20's potential.

4. Experimental Section

Materials: CALF-20 powder was purchased from Sigma Aldrich and used in adsorption experiments, in situ-PXRD and in situ PALS measurements after desolvation in dynamic vacuum for 10 h at 430 K.

Methods—PALS: Data Collection and Treatment: PALS measurements were conducted using a digital positron lifetime spectrometer equipped with an Acqiris DC 282 digitizer (10-bit vertical resolution) and three photomultiplier tubes (PMTs) from SCIONIX HOLLAND coupled to CeBr₃ scintillators. Two scintillators had dimensions of $\emptyset=51$ mm and h=10 mm, while the third scintillator had $\emptyset=51$ mm and h=25.4 mm. The scintillators were deliberately positioned to ensure comprehensive cover-

age around the sample holder and minimize the detection of backscattered photons, which can distort PALS spectra due to high efficiency of scintillators. The PMTs were arranged with precision: two were oriented horizontally at a 90° angle to each other, with their front faces separated by approximately 20 mm. The third PMT was vertically aligned, positioned above the sample holder. Customized data acquisition software facilitated simultaneous detection of start and stop signals from each of the three PMTs. This capability allowed for the collection of six spectra in a single measurement session, significantly reducing measurement time compared to traditional two-tube spectrometers. The sample holder was designed specifically to accommodate around 0.4 mL of powdered sample, surrounding a 20 µCi ²²Na positron source sealed within a 5 µm Kapton foil (DuPont). The coincidence box was set to a time range of 500 ns, while the PALS data acquisition software operated with a 600 ns time window and a channel width of 5 ps. The PALS spectra were resolved by fitting exponential decay curves to the histograms of recorded time differences, revealing lifetime components (τ_n) with their respective intensities (I_n) . Analysis showed that a reliable fit, characterized by minimal residual and a χ^2 value of 1.03–1.08, was achieved only when considering four lifetime components. The resolved lifetime components were categorized based on their lifetimes: τ_1 , corresponding to p-Ps annihilation ($\tau_1 \approx 0.12$ -0.18 ns); τ_2 , originating from unbound e^+ annihilation ($\tau_2\approx$ 0.35–0.50 ns); and the longest-lived components representing o-Ps annihilation within the pores ($\tau_3 > 1$ ns). It is believed that τ_4 reflects o-Ps annihilation between crystals in the intergranular spaces. These components are visualized in Figure S2 (Supporting Information).

The numerical spectra analysis was conducted using the PALSFit routine, [42] which separates the spectra from resolution functions, background, and source contribution. The conversion of the measured o-Ps lifetime to pore sizes is done by the EELViS code, [55] which is based on the extended Tao-Eldrup model. [40] Al and Sn reference samples with known lifetimes were employed to estimate the resolution functions and source contribution. Approximately 10.6% source contribution (distributed between Kapton foil (0.382 ns /10.4%) and glue (2.6 ns/0.2%)) and an

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average time resolution of 0.28 ns (full width at half maximum) were determined during the fitting process.

- i. Activation: The as-received sample underwent activation at 423 K in five slices (totaling 10 h, with 2 h per slice) under dynamic vacuum of 1.5×10^{-6} mbar to remove any impurities absorbed by the CALF-20 during sample transfer for PALS measurements. Subsequently, the evacuated sample was cooled to room temperature (RT) and PALS spectrum was measured again over 2 h. This RT measurement serves as a reference for the initial pore state.
- iii. Temperature-dependent measurements of empty pores: As mentioned in Section 2.1, PALS measurements were performed under vacuum conditions (1.5×10^{-6} mbar) over a temperature range of 253 to 373 K. The experimental procedure involved first activating the sample at 423 K for 10 h, followed by a temperature program starting at 373 K and cooling to 253 K, with heating and cooling rates of 2.5 K min $^{-1}$. Measurements were conducted at 2 h intervals at each temperature step during the temperature program.
- iii. Dry CO₂ adsorption: The dry CO₂ adsorption experiments have been carried out at different temperatures: 353 K, 323 K, 295 K, 273 K, and 253 K. At each T, the gas pressure was regulated by the software and varied from 5 mbar to 1000 mbar during the adsorption cycle and then reversed during the desorption cycle. The PALS measurements for each adsorption/desorption step were conducted continuously for 2 h at specific temperature and gas pressure conditions. After each adsorption scan at a specific temperature, the sample was annealed at 473 K for 4 h to purge the pore for the subsequent step.
- iv. Humidity and humid CO2: Again, to thoroughly remove any residual CO2 in the pores from the previous experiment, the sample underwent reactivation at 423 K for 4 h. Following this, the humidity experiments commenced at RT, consisting of two steps: first, exposing the sample to pure humidity, and then subjecting it to humid CO2 conditions at 970 mbar CO2. Humidity levels were adjusted using saturated salt solutions to achieve specific relative humidity (RH) levels of 11.3%, 22.5%, 34%, 43%, 51%, 63.5%, 75%, and 98% (water). At each RH value, the procedure involved closing the valve to CO2 to establish pure humidity, removing the salt solution, heating the sample under vacuum at 423 K for 4 h to restore clean pores, reintroducing the salt solution, and simultaneously opening the CO2 valve to initiate the humid CO2 experiment. Subsequently, the sample was left exposed overnight to ensure saturated humidity over the sample. The measurement then carried out by recording at least five slices, each lasting 2 h, and the average of these slices is presented as the result. Then, the activation procedure was repeated before the next step.

Methods—Volumetric Adsorption Measurements: Water vapor physisorption isotherms were measured using BELSORP-max apparatus (MICROTRAC MRB). A water-filled thermostat combined with water bath ($T=298\,$ K) was used to control the temperature. Helium gas (purity 99.999%) was used for dead volume measurement. The conditions of 0.5% of pressure change within 450 s was assumed as adsorption equilibrium. Carbon dioxide physisorption isotherms were measured using BELSORP-HP (MICROTRAC MRB), equipped with closed cycle helium cryostat and customized adsorption cell. The dead volume was measured using helium gas (purity 99.999%) and carbon dioxide (purity 99.999%) were used in adsorption experiments. The adsorption equilibration was defined as 1% of the pressure change within 450 s.

Methods—Variable Temperature PXRD in Controlled Gas Atmospheres: Powder diffraction patterns at variable temperatures were measured on PANALYTICAL Empyrean Powder X-ray diffractometer (α -1 system, ω -2 θ goniometer), equipped with Cu-K α 1 radiation (primary monochromator λ = 1.54059 Å), motorized antiscatter and receiving slits. A Pixcel 1D detector was used for the measurements of reflection intensities. HTK-1200N high temperature chamber (ANTON PAAR) was used to control the sample temperature. All measurements were conducted in a gas flow of 5 mL min⁻¹. PXRD patterns were collected in reflection geometry using ω -2 θ scans in the 2 θ range from 3 to 70° with 0.013° steps and 120 s per step. The sample was conditioned for 30 min after reaching the target temperature. The sample height correction program for the measurements in the inert atmosphere was applied to correct the thermal expansion of the sample holder. Nitrogen (purity 99.999%) and carbon dioxide (purity 99.999%) were used in experiments. The gases were humidified in a customized water-filled bubbler in order to achieve a saturated humidity level at 298 K.

Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

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Conflict of Interest

The authors declare no conflict of interest.

Data Availability Statement

The data that support the findings of this study are available from the corresponding author upon reasonable request.

Keywords

 $\mathsf{CALF-20},\,\mathsf{CO}_2$ capture and storage, environmental gas uptake, positron annihilation

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