

PQMIV Q1 (I)

$$H(t) = H_0 + \Delta_s(t)$$

$$H_0 = \begin{pmatrix} E_1 & 0 \\ 0 & E_2 \end{pmatrix} \quad \Delta_s(t) = \begin{pmatrix} 0 & V_0 e^{i\omega t} \\ V_0 e^{-i\omega t} & 0 \end{pmatrix}$$

$$P(|1\rangle \rightarrow |2\rangle) = a(t)^* a(t)$$

$$a(t) = \frac{1}{i\hbar} \int_0^t e^{i\omega_{12} t'} \langle 1 | \Delta(t') | 2 \rangle dt'$$

$\omega_{12} = \frac{E_1 - E_2 + i\hbar\omega}{\hbar}$

$$\begin{aligned} \langle 1 | \Delta(t') | 2 \rangle &= (1)^+ V_0 \begin{pmatrix} 0 & e^{i\omega t'} \\ e^{-i\omega t'} & 0 \end{pmatrix} (0) \\ &= (1 \ 0) \begin{pmatrix} e^{i\omega t'} \\ 0 \end{pmatrix} = V_0 e^{i\omega t'} \end{aligned}$$

WE HAVE:

$$\begin{aligned} a(t) &= \frac{1}{i\hbar} \int_0^t e^{i\frac{E_1 - E_2}{\hbar} t'} V_0 e^{i\omega t'} dt' = \\ &= \frac{V_0}{i\hbar} \int_0^t e^{i(E_1 - E_2 + \hbar\omega)t'/\hbar} dt' = \\ &= \frac{V_0}{i\hbar} \frac{\hbar}{(E_1 - E_2 + i\hbar\omega)} \left(e^{i(E_1 - E_2 + \hbar\omega)t'/\hbar} - 1 \right) \\ &= -\frac{V_0}{E_1 - E_2 + i\hbar\omega} \left[e^{i(E_1 - E_2 + \hbar\omega)t'/\hbar} - 1 \right] \end{aligned}$$

QM IV Q1 (II)

$$L + X = E_1 - E_2 + \frac{e}{m} w$$

THEN:

$$d^* a = + \frac{\sqrt{\epsilon}^2}{X^2} \left(e^{-iXt/\frac{e}{m}} - 1 \right) \left(e^{iXt/\frac{e}{m}} - 1 \right)$$

$$= \frac{\sqrt{\epsilon}^2}{X^2} \left(1 - e^{-iXt/\frac{e}{m}} - e^{iXt/\frac{e}{m}} + 1 \right)$$

$$= \frac{\sqrt{\epsilon}^2}{X^2} \left[2 - \left(e^{-iXt/\frac{e}{m}} + e^{iXt/\frac{e}{m}} \right) \right]$$

$$= \frac{\sqrt{\epsilon}^2}{X^2} 2 \left[1 - \cos(Xt/\frac{e}{m}) \right]$$

$$= \frac{\sqrt{\epsilon}^2}{X^2} 2 \cdot 2 \sin^2\left(\frac{Xt}{2\frac{e}{m}}\right)$$

$$= \underline{\underline{\frac{4\sqrt{\epsilon}^2}{(E_1 - E_2 + \frac{e}{m} w)^2} \sin^2\left(\frac{(E_1 - E_2 + \frac{e}{m} w)t}{2\frac{e}{m}}\right)}}$$

PQM IV Q2 (I)

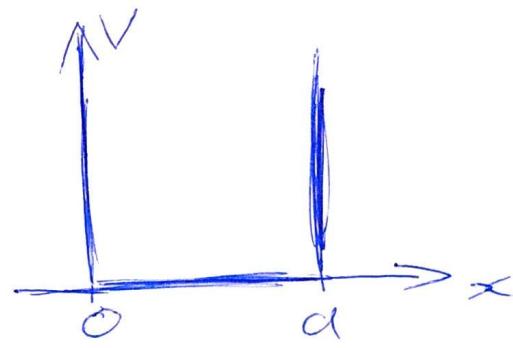
START FROM 1D CASE:

TDSE:

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V\psi = E\psi$$

"IN BOX"

$$\psi(0) = \psi(L) = 0$$



$$\Rightarrow \psi(x) = A \sin\left(\frac{n\pi x}{a}\right) \quad n = 1, 2, 3, \dots$$

$$E_n = -\frac{\hbar^2}{2m} \left(\frac{n^2\pi^2}{a^2} \right) = \frac{\pi^2 \hbar^2}{2ma^2} n^2$$

$$\int_0^a \psi^* \psi dx = 1 \Rightarrow A = \sqrt{\frac{2}{a}}$$

3D CASE:

(IE PARTICLE IN A BOX)

$$E_{klm} = \frac{\pi^2 \hbar^2}{2ma^2} (\ell^2 + m^2)$$

$$E_{\text{INITIAL}} = \frac{3\pi^2 \hbar^2}{2ma^2} \Rightarrow \ell = l = m = 1$$

$$E_{\text{LATE}} = \frac{6\pi^2 \hbar^2}{2ma^2} \Rightarrow \ell, l, m = 2, 1, 1$$

$$P(E_{\text{INITIAL}} \rightarrow E_{\text{LATE}}) = a^*(t) a(t)$$

$$a(t) = \frac{1}{it} \int_0^t e^{i\frac{E_{\text{LATE}} - E_{\text{INITIAL}}}{\hbar t} t'} \langle \text{LATE} | \Delta(t') | \text{INITIAL} \rangle dt'$$

$$a(t) = \frac{1}{it} \int_0^t e^{i\frac{E_{\text{LATE}} - E_{\text{INITIAL}}}{\hbar t} t'} \langle \text{LATE} | \Delta(t') | \text{INITIAL} \rangle dt'$$

$$\underline{H} = \underline{H}_0 + \underline{\Delta}$$

$$\underline{\Delta} = (eEX, 0, 0)$$

$$\langle 2, 1, 1 | \underline{\Delta} | 1, 1, 1 \rangle =$$

$$= \langle 2 | eEX | 1 \rangle + \langle 1 | 0 | 1 \rangle + \langle 1 | 0 | 1 \rangle$$

$$= \int_0^a \frac{2}{a} \sin\left(\frac{2\pi x}{a}\right) x \sin\left(\frac{\pi x}{a}\right) dx = \text{ALGEBRA STEPS/COMPUTER}$$

$$= -\frac{16a}{9\pi^2}$$

$$a(t) = \frac{1}{i\tau} \int_0^{t_i} e^{i \frac{E_{LATE} - E_{INIT}}{\tau} t'} \left(-\frac{16a}{9\pi^2} \right) dt'$$

$$= \frac{1}{i\tau} \left(-\frac{16a}{9\pi^2} \right) \frac{1}{i(E_{LATE} - E_{INIT})} \left(e^{i \frac{E_{LATE} - E_{INIT}}{\tau} t_i} - 1 \right)$$

~~$a(t) a(t) = \frac{16a}{9\pi^2}$~~

$$a^* a = \left(\frac{16a}{9\pi^2} \right)^2 \frac{1}{(E_{LATE} - E_{INIT})^2} \left(e^{i \frac{E_{LATE} - E_{INIT}}{\tau} t_i} - 1 \right) \left(e^{-i \frac{E_{LATE} - E_{INIT}}{\tau} t_i} - 1 \right)$$

$$= \left(\frac{16a}{9\pi^2} \right)^2 \frac{1}{E_{LATE} - E_{INIT}} + \sin^2 \left(\frac{E_{LATE} - E_{INIT}}{\tau} t_i \right)$$

PQM IV Q3 | (I)

$$P(|\text{Ground}\rangle \rightarrow |\text{LATE}\rangle) = a^* a$$

$$a(t) = \int_0^t e^{i \frac{E_{\text{LATE}} - E_{\text{Ground}}}{\hbar} t'} \langle \text{LATE} | \Delta(t') | \text{Ground} \rangle dt'$$

WE HAVE:

$$\Delta(t') = \alpha X = A + A^+$$

so:

$$\begin{aligned} \langle \text{LATE} | \Delta(t') | \text{Ground} \rangle &= \alpha \langle \text{LATE} | A + A^+ | \text{Ground} \rangle \\ &= \langle \text{LATE} | A | \text{Ground} \rangle + \langle \text{LATE} | A^+ | \text{Ground} \rangle \xrightarrow{\downarrow 0} \\ &= \langle n | A | 0 \rangle = \overline{J_{1n}} \end{aligned}$$

(A LOWERS n BY 1,
IF n IS NOT 1, IT IS NOT
LOWEDED TO 0, IE
RESULT WILL BE 0)

⇒ TRANSITION ONLY ALLOWED TO FIRST EXCITED STATE.

LETS DO THIS EXACTLY NOW.

$$\langle \text{LATE} | \Delta | \text{Ground} \rangle = \langle \text{LATE} | \frac{qV\varepsilon}{\sqrt{\pi}\lambda} \exp\left(-\frac{t^2}{\lambda^2}\right) \sqrt{\frac{\hbar}{2m\omega}} \hat{a} | \text{Ground} \rangle$$

$$= \frac{qV\varepsilon}{\sqrt{\pi}\lambda} \sqrt{\frac{\hbar}{2m\omega}} \exp\left(-\frac{t^2}{\lambda^2}\right) \overline{J_{1n}}$$

PQM IV Q3 (II)

$$a(t) = \int_{-\infty}^t e^{i \frac{E_{\text{LATE}} - E_{\text{GROUND}}}{\hbar} t'} \frac{q\varepsilon}{\sqrt{\pi}\varepsilon} \sqrt{\frac{\hbar}{2m\omega}} \exp\left(-\frac{t'^2}{\varepsilon^2}\right) dt'$$

$$E_{\text{LATE}} - E_{\text{GROUND}} = \hbar\omega\left(1 + \frac{1}{2}\right) - \hbar\omega\frac{1}{2} = \hbar\omega$$

$$\Rightarrow a(t) = \frac{q\varepsilon}{\sqrt{\pi}\varepsilon} \sqrt{\frac{\hbar}{2m\omega}} \int_0^t e^{i\omega t'} e^{-\frac{t'^2}{\varepsilon^2}} dt'$$

"FROM VERY EARLY TO VERY LATE TIMES" \Rightarrow

\Rightarrow LIMITS SHOULD BE: $-\infty, \infty$

~~limits~~

$$a(t) = \frac{q\varepsilon}{\sqrt{\pi}\varepsilon} \sqrt{\frac{\hbar}{2m\omega}} \int_{-\infty}^{\infty} e^{i\omega t'} e^{-\frac{t'^2}{\varepsilon^2}} dt' =$$

AFTER ~~LOTS~~ ^{SOME} OF ALGEBRA / COMPUTER

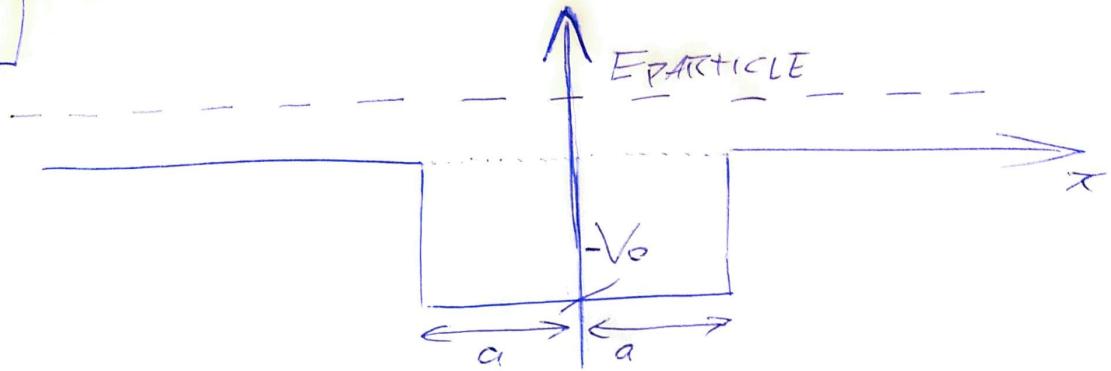
$$= \frac{q\varepsilon}{\sqrt{\pi}\varepsilon} \sqrt{\frac{\hbar}{2m\omega}} \cancel{\sqrt{\pi}} \times \exp\left(-\frac{1}{4} \frac{\varepsilon^2 \omega^2}{\hbar^2}\right)$$

$$= q\varepsilon \sqrt{\frac{\hbar}{2m\omega}} \exp\left(-\frac{1}{4} \frac{\varepsilon^2 \omega^2}{\hbar^2}\right)$$

$$a^* a = \underline{q^2 \varepsilon^2 \frac{\hbar}{2m\omega} \exp\left(-\frac{1}{2} \frac{\varepsilon^2 \omega^2}{\hbar^2}\right)}$$

I AM UNSURE IF I WAS RIGHT TO SEND THE LOWER
LIMIT TO $-\infty$ IN THE INTEGRAL ABOVE.

PQM IV Q4



~~RELEVANT~~
MOVE TO PARTICLE'S FRAME.

SWITCH ON PERTURBATION AT $t=0$,
SWITCH OFF AT $\frac{2a}{\omega}$

HERE I FIND TROUBLESOME TO CONCEPTUALLY
CONNECT ABSORPTION-EMISSION (TYPICAL
FERMI'S RULE APPLICATION AREA) TO
REFLECTION-TRANSMISSION.

PQM \Rightarrow Q5/I

$$H = \omega t_1 \left(A^+ A + \frac{1}{2} \right) + t_1 \underbrace{\left(f^*(\epsilon) A + f(\epsilon) A^+ \right)}_{\delta H(\epsilon)}$$

$$= H_0 + \delta H(\epsilon)$$

LET'S DEFINE:

$$U_0 = \exp \left(- \frac{i H_0 t}{\hbar} \right)$$

LET'S HAVE A STATE $|\psi(\epsilon)\rangle$ WHICH EVOLVES ACCORDING TO $H(\epsilon)$.

IF WE DEFINE $|\tilde{\psi}(\epsilon)\rangle$ THE FOLLOWING WAY,
IT'LL EVOLVE ACCORDING TO δH ONLY:

$$|\tilde{\psi}(\epsilon)\rangle = U^\dagger(t) |\psi(\epsilon)\rangle$$

SCHRÖDINGER EQ. FOR $|\tilde{\psi}(\epsilon)\rangle$:

$$i \frac{\hbar}{\hbar} \frac{d}{dt} |\tilde{\psi}(\epsilon)\rangle = i \frac{\hbar}{\hbar} \frac{d}{dt} \left[\exp \left(+ \frac{i H_0 t}{\hbar} \right) |\psi(\epsilon)\rangle \right]$$

$$= i \frac{\hbar}{\hbar} \underbrace{i \frac{\hbar}{\hbar} \exp \left(\frac{i H_0 t}{\hbar} \right) |\psi(\epsilon)\rangle}_{|\tilde{\psi}(\epsilon)\rangle} + i \frac{\hbar}{\hbar} \exp \left(\frac{i H_0 t}{\hbar} \right) \frac{d}{dt} |\psi(\epsilon)\rangle$$

SCHRÖDINGER FOR $|\psi(\epsilon)\rangle$: $i \frac{\hbar}{\hbar} \frac{d}{dt} |\psi(\epsilon)\rangle = (H_0 + \delta H) |\psi(\epsilon)\rangle$

USING THIS, REWRITE LIKE ABOVE.

$$= -H_0 |\tilde{\psi}(\epsilon)\rangle + \exp \left(\frac{i H_0 t}{\hbar} \right) (H_0 + \delta H) \exp \left(- \frac{i H_0 t}{\hbar} \right) |\tilde{\psi}(\epsilon)\rangle$$

$\cancel{\text{II}}$

← THESE IS COMMUTE &
CANCELED WITH THIS

$$= \exp \left(\frac{i H_0 t}{\hbar} \right) \delta H(t) \exp \left(- \frac{i H_0 t}{\hbar} \right) |\tilde{\psi}(\epsilon)\rangle$$

PQM \rightleftharpoons Q5(II)

$$= \exp(i\omega t) \mathcal{S} H(t) \exp(-i\omega t) |\tilde{\psi}(t)\rangle$$

$$= \underbrace{\exp(-i\omega t) \mathcal{S} H(t)}_{\text{THIS IS THE HAMILTONIAN GOVERNING THE TIME-DEPENDENT INTERACTION.}} |\psi(t)\rangle$$

$$\Rightarrow V_I(\epsilon) = \exp(i\omega t) \mathcal{S} H(\epsilon)$$

$$= \exp(i\omega t) \text{tr} (\hat{f}^* A + \hat{f} A^+)$$

$$= \text{tr} (\cancel{\hat{f}} e^{i\omega t} f(\epsilon) A + e^{i\omega t} \cancel{f} A^+)$$

$$= \text{tr} (\tilde{f}^* A + \tilde{f} A^+)$$

AS REQUIRED

(LARGELY INSPIRED BY MIT OCW NOTES)