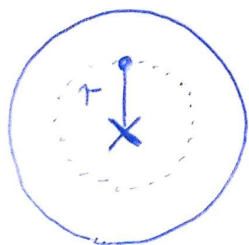


SDSG2

2.1.

POTENTIAL FROM PART OF SPHERE \rightarrow ↑
 NEWTON'S FIRST THM: NO FORCE DUE TO
 EXTERNAL SPHERICALLY
 SYMMETRIC MASS
 DISTRIBUTION



$\Rightarrow \Phi = \text{CONSTANT CONTRIB.}$
 FROM THIS BIT.

POTENTIAL FROM INSIDE SPHERE:

NEWTON SECOND THM: TREAT INSIDE AS
 POINT MASS.

$$f = -\nabla \Phi \Rightarrow \Phi = - \int_{\infty}^r f dr$$

$$= - \int_{\infty}^r \frac{m_{<r'} G}{r'^2} dr'$$

$$= - \int_{\infty}^r \frac{4}{3} \pi r'^3 \rho \frac{1}{r'^2} G dr' = \int_{\infty}^r \frac{4}{3} \pi \rho G r' dr' = \frac{4}{3} \pi \rho G \left[\frac{r'^2}{2} \right]_{\infty}^r$$

$$= \frac{2}{3} \pi \rho G r^2 = \frac{1}{2} \Omega^2 r^2 + C$$

$$\text{WHERE } \Omega^2 = \frac{4}{3} \pi \rho G$$

THE NOTATION

Ω SUGGESTS

THAT $\frac{4}{3} \pi \rho G$

SHOULD HAVE

DIMENSIONS SEC^{-2}

CHECK: $\sqrt{\frac{\text{kg}}{\text{m}^3} \cdot \text{N} \frac{\text{m}^2}{\text{kg}^2}} = \sqrt{\frac{\text{N}}{\text{m kg}}} = \sqrt{\frac{\text{kg} \frac{\text{m}}{\text{kg}^2 \text{sec}^2}}{\text{m kg}}} = \frac{1}{\text{SEC}} \text{ GOOD.}$

THE FACT WHICH WORKS

UNSTABLE IF: $\frac{d}{dr}(R^3 \ell) > 0$

ARISING QUESTIONS

RECALL: FOR CIRCULAR ORBIT:

$$-R\dot{\phi}^2 = \ell$$

• WHEN TAKING SECOND DERIVATIVE OF $V_{\text{EFFECTIVE}}$ DOESN'T WORK? IE

$$V_{\text{EFF}}(r) = U(r) + \frac{L^2}{2r^2}$$

$$\frac{dV_{\text{EFF}}}{dr} = \frac{dU(r)}{dr} + \frac{L}{r^2} L' - L^2 r^{-3}$$

CONSIDER:

$$\frac{dU(r)}{dr} = -f(r) = r\dot{\phi}^2 - \frac{L^2}{r^3} \rightarrow 0 \text{ FOR CIRCULAR}$$

$$\frac{dV_{\text{EFF}}}{dr} = r\dot{\phi}^2 + \frac{L L'}{r^2} - L^2 r^{-3} \rightarrow 0 \text{ BEC } L' = 0$$

ORBIT IS WHERE $\frac{dV_{\text{EFF}}}{dr} = 0$, IE $r\dot{\phi}^2 - L^2 r^{-3} = 0$

$$\frac{dV_{\text{EFF}}}{dr} = \cancel{r\dot{\phi}^2} - L^2 r^{-3} = 0$$

$$\dot{\phi}^2 = L^2 r^{-4}$$

$$L^2 = \dot{\phi}^2 r^4 \quad (\text{AS EXPECTED})$$

OH I SEE. SHOULDN'T THROW OUT THE $\frac{L L'}{r^2}$ TERM BEFORE TAKING SECOND DERIVATIVE.

• ~~IS $\frac{d}{dr}(L^2)$ EVER SATISFIED?~~ IS THE METHOD CORRECT THOUGH IN PRINCIPLE?
 $\frac{d}{dr}(L^2) = 2L L'$. L IS CONSTANT
 SO $L' = 0 \Rightarrow \frac{d}{dr}(L^2) = 0$ IF ANGULAR MOMENTUM IS CONSERVED IE ALWAYS
 SO HOW ARE THERE UNSTABLE ORBITS?

$$\frac{d}{dr}(R^3(-R\dot{\phi}^2)) > 0$$

$$\frac{d}{dr}(R^4 \dot{\phi}^2) < 0$$

RENAME $R \rightarrow r$
 USE: $L = r^2 \dot{\phi}$

$$\frac{d}{dr}(L^2) < 0$$

AS REQUIRED.

RECALL ORBIT EQ.:

$$\frac{d^2 u}{d\phi^2} + u = -\frac{f(r)}{h^2 u^2}$$

$$\text{with } u = \frac{1}{r}$$

$$h = r^2 \dot{\phi}$$

$$u = \frac{1}{r} = a^{-1} \exp(-b\phi)$$

$$\frac{1}{u^2 h^2} = \frac{1}{u^2 \bar{u}^4 \dot{\phi}^2} = u^2 \dot{\phi}^{-2}$$

$$\frac{b^2}{a} \exp(-b\phi) + \frac{1}{a} \exp(-b\phi) = \nabla \Phi u^2 \frac{1}{\dot{\phi}^2}$$

$$\frac{1}{a} (b^2 + 1) \exp(-b\phi) = \nabla \Phi \frac{u^2}{\dot{\phi}^2}$$

$$= \nabla \Phi \frac{1}{\dot{\phi}^2} a^{-2} \exp(-2b\phi)$$

$$\dot{\phi}^2 a (b^2 + 1) \exp(b\phi) = \nabla \Phi$$

$$\frac{h^2}{r^4} \overbrace{a (b^2 + 1) \exp(b\phi)}^r = \nabla \Phi$$

$$h^2 (b^2 + 1) r^{-3} = \nabla \Phi$$

$$\Rightarrow \Phi \propto \frac{1}{r^2}$$

CONSTANT OF
PROPORTIONALITY
CAN BE OBTAINED
IF WE KNOW h
OF THE ORBIT IS
WELL, CORRECT?

2.4/

$$u^2 = 2 \frac{E - \Phi}{h^2}$$

$$\Rightarrow \frac{d^2 \Phi}{du^2} = \frac{d^2}{du^2} \left(E - \frac{1}{2} u^2 h^2 \right) = -h^2$$

$$S(r) > 0 \quad \forall r \Rightarrow \Phi < 0 \quad \forall u$$

(USING: $\Phi(r) = - \int \frac{G S(r') dr'}{\sqrt{|r-r'|}}$)

WANT TO PROVE: THERE ARE 0 OR 2 ROOTS.

IE DISPROVE THAT THERE CAN BE 1 ROOT.

$$1 \text{ ROOT IF: } E - \Phi = 0$$

$$\frac{d^2 \Phi}{du^2} = -h^2 \Rightarrow \frac{d\Phi}{du} \text{ IS DECREASING AS } u \text{ INCREASES}$$

$$\Rightarrow \Phi \text{ IS CONCAVE. (AND THEN WHAT?)}$$

$$E = \underbrace{KE}_{\text{KINETIC ENERGY}} + \Phi \quad \Rightarrow E > \Phi \Rightarrow E - \Phi > 0$$

$$KE > 0$$

(IE PARTICLE IN
CENTRAL FORCE
FIELD DOES NOT
STAY STATIONARY)

$$\Downarrow$$

$$E - \Phi \neq 0$$

\Downarrow
CANNOT BE 1 ROOT

$$E - \Phi > 0 \Rightarrow u^2 = \frac{2}{h^2} (E - \Phi) \text{ MUST HAVE 2 ROOTS.}$$

THIS IS NOT WHAT WE WERE SUPPOSED TO BE PROVING.

SDSGZ

2.5 (I)

$$\nabla^2 \Phi = 4\pi G \rho$$

$$\text{TOTAL MASS} = \int_{\text{ALL SPACE}} \rho dV = \int_{\text{ALL SPACE}} \frac{1}{4\pi G} \nabla^2 \Phi dV$$

$$= \int_{r=0}^{\infty} \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} \left[\frac{1}{4\pi G} \nabla^2 \Phi \right] r^2 \sin \theta d\theta d\phi dr$$

$$= \frac{1}{G} \int_0^{\infty} \nabla^2 \Phi r^2 dr = \frac{1}{G} \int_0^{\infty} \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \Phi}{\partial r} \right) r^2 dr$$

$$= \int_0^{\infty} \frac{\partial}{\partial r} \left[r^2 \frac{\partial}{\partial r} \left(\frac{-M}{b + \sqrt{b^2 + r^2}} \right) \right] dr$$

$$= M \int_0^{\infty} - \frac{\partial}{\partial r} \left[r^2 \frac{\partial}{\partial r} \left(\frac{1}{b + \sqrt{b^2 + r^2}} \right) \right] dr$$

— TRUST OURSELVES THAT WE'RE GOOD IN
SETTING UP INTEGRALS

OR

— BRUTE FORCE ALGEBRA

OR

— PUT IT INTO COMPUTER
(THIS WAS MY CHOICE)

ARRIVE AT

$$= M$$

SDSG2
2.5(II)

$$S(0) = \frac{1}{4\pi G} \nabla^2 \Phi \Big|_{r=0}$$

$$= \frac{1}{4\pi G} \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \left(\frac{-M}{b + \sqrt{b^2 + r^2}} \right) \right) \Big|_{r=0}$$

①

$$= \frac{3M}{16\pi b^3}$$

[DONE BY COMPUTER,
RESULT CAN BE VERIFIED ON DESMOS:
[DESMOS.COM/CALCULATOR/ECNFBTM4MC](https://www.desmos.com/calculator/ECNFBTM4MC)]

$$S(r \gg b) = \frac{1}{4\pi G} \nabla^2 \Phi \Big|_{r \gg b} \quad \text{using: } \frac{r + \sqrt{b^2 + r^2}}{2} \approx r \quad \text{when } r \gg b$$

$$\approx \frac{1}{4\pi} \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \left(\frac{-M}{r} \right) \right) \Big|_{r \gg b}$$

$$= \frac{1}{4\pi} \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \left(\frac{-M}{r + \sqrt{b^2 + r^2}} \right) \right)$$

$$= \frac{bM}{2\pi r^4}$$

[VERIFIED BY DESMOS]
(LINK ABOVE ↑)