

REL. IV 2

$$L = g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu = \left( \frac{ds}{dt} \right)^2$$

$$L = c^2 \left( 1 - \frac{2M}{r} \right) \dot{t}^2 - \left( 1 - \frac{2M}{r} \right)^{-1} \dot{r}^2 - r^2 \dot{\Theta}^2 - r^2 \sin^2 \Theta \dot{\Phi}^2$$

EULER-LAGRANGE FOR  $r$ :

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{r}} \right) = \frac{d}{dt} \left[ -2 \left( 1 - \frac{2M}{r} \right)^{-1} \dot{r} \right]$$

$$= 2 \left( 1 - \frac{2M}{r} \right)^{-2} \frac{2M}{r^2} \dot{r}^2 - 2 \left( 1 - \frac{2M}{r} \right)^{-1} \ddot{r}$$

$$\frac{\partial L}{\partial r} = c^2 \frac{2M}{r^2} \dot{t}^2 + \left( 1 - \frac{2M}{r} \right)^{-2} \frac{2M}{r^2} \dot{r}^2 - 2r \dot{\Theta}^2 - 2r \sin^2 \Theta \dot{\Phi}^2$$

$$\Rightarrow 2 \left( 1 - \frac{2M}{r} \right)^{-2} \frac{2M}{r^2} \dot{r}^2 - 2 \left( 1 - \frac{2M}{r} \right)^{-1} \ddot{r} = c^2 \frac{2M}{r^2} \dot{t}^2 + \left( 1 - \frac{2M}{r} \right)^{-2} \frac{2M}{r^2} \dot{r}^2 - 2r \dot{\Theta}^2 - 2r \sin^2 \Theta \dot{\Phi}^2$$

$$\left( 1 - \frac{2M}{r} \right)^{-2} \frac{2M}{r^2} \dot{r}^2 - \ddot{r} = c^2 \frac{2M}{r^2} \left( 1 - \frac{2M}{r} \right)^{-1} \dot{t}^2 + \left( 1 - \frac{2M}{r} \right)^{-1} \frac{2M}{r^2} \dot{r}^2$$

$$- \left( 1 - \frac{2M}{r} \right) r \dot{\Theta}^2 - \left( 1 - \frac{2M}{r} \right) r \sin^2 \Theta \dot{\Phi}^2$$

C.F.  $\frac{d^2 r}{dt^2} + \Gamma_{\mu\nu}^r \frac{dx^\mu}{dt} \frac{dx^\nu}{dt} = 0$

$$\Rightarrow \Gamma_{rr}^r = \left( 1 - \frac{2M}{r} \right)^{-1} \frac{M}{r^2} - \left( 1 - \frac{2M}{r} \right)^{-1} \frac{2M}{r^2} = - \left( 1 - \frac{2M}{r} \right)^{-1} \frac{M}{r^2}$$

$$\Gamma_{tt}^r = c^2 \left( 1 - \frac{2M}{r} \right) \frac{M}{r^2}$$

$$\Gamma_{\Theta\Theta}^r = - \left( 1 - \frac{2M}{r} \right) r$$

$$\Gamma_{\Phi\Phi}^r = - \left( 1 - \frac{2M}{r} \right) r \sin^2 \Theta$$

ALL OTHER  $\Gamma_{xy}^r$  ARE ZERO.

EULER-LAGRANGE FOR  $\Theta$ :

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\Theta}} \right) = \frac{\partial L}{\partial \Theta}$$

LHS:

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\Theta}} \right) = \frac{d}{dt} (-r^2 \dot{\Theta}) = -4r\dot{r}\dot{\Theta} - 2r^2\ddot{\Theta}$$

RHS:

$$\frac{\partial L}{\partial \Theta} = -r^2 2 \sin \Theta \cos \Theta \dot{\Phi}^2$$

$$\Rightarrow -4r\dot{r}\dot{\Theta} - 2r^2\ddot{\Theta} = -r^2 2 \sin \Theta \cos \Theta \dot{\Phi}^2$$

~~REARRANGE~~  
/ REARRANGE  
/:  $2r^2$

$$\ddot{\Theta} + \frac{2}{r} \dot{r}\dot{\Theta} - \sin \Theta \cos \Theta \dot{\Phi}^2 = 0$$

COMPARE WITH GEODESIC EQUATION, CONCLUDE:

$$\Gamma_{r\Theta}^{\Theta} = \Gamma_{\Theta r}^{\Theta} = \frac{1}{r}$$

$$\Gamma_{\Phi\Phi}^{\Theta} = -\sin \Theta \cos \Theta$$

ALL OTHER  
 $\Gamma_{XY}^{\Theta}$  IS ZERO.

EULER-LAGRANGE FOR  $\Phi$ :

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\Phi}} \right) = \frac{\partial L}{\partial \Phi}$$

LHS:

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\Phi}} \right) = \frac{d}{dt} (-2r^2 \sin^2 \Theta \dot{\Phi}) =$$

$$= -4r\dot{r}\sin^2 \Theta \dot{\Phi} - 4r^2 (\sin \Theta)(\cos \Theta) \dot{\Theta} \dot{\Phi} - 2r^2 (\sin^2 \Theta) \ddot{\Phi}$$

RHS:

$$\frac{\partial L}{\partial \dot{\phi}} = 0$$

$$LHS = RHS \quad \therefore -2 \quad \text{~~EE~~$$

$$\Rightarrow 2 + 2 \sin^2 \theta \dot{\phi} + 2 r^2 (\sin \theta)(\cos \theta) \dot{\theta} \dot{\phi} + r^2 \sin^2 \theta \ddot{\phi} = 0 \quad \text{REARRANGE: } r^2 \sin^2 \theta$$

$$\ddot{\phi} + \frac{2}{r} \dot{\phi} + 2 \cot \theta \dot{\theta} \dot{\phi} = 0$$

COMPARE WITH GEODESIC EQUATION, CONCLUDE:

$$\Gamma_{\phi r}^{\phi} = \Gamma_{r \phi}^{\phi} = \frac{2}{r} \quad \text{!}$$

$$\Gamma_{\theta \phi}^{\phi} = \Gamma_{\phi \theta}^{\phi} = 2 \cot \theta$$

(THERE ARE TWO  
ΓS FOR  $\ddot{\phi}$ :  
 $\Gamma_{\theta \phi}^{\phi}$  &  $\Gamma_{\phi \theta}^{\phi}$ )

ALL OTHER  $\Gamma_{XY}^{\phi}$  IS ZERO.

$$g_{\mu\nu} u^\mu u^\nu = \left( \frac{ds}{d\tau} \right)^2 = c^2$$

$$g_{\mu\nu} u^\mu u^\nu d\tau^2 = c^2 d\tau^2$$

$$g_{\mu\nu} u^\mu u^\nu d\tau^2 =$$

$$g_{00} u^0 u^0 d\tau^2 + g_{11} u^1 u^1 d\tau^2 + g_{22} u^2 u^2 d\tau^2 + g_{33} u^3 u^3 d\tau^2$$

$$= g_{00} u^0 u^0 d\tau^2 + g_{11} u^1 u^1 d\tau^2 + g_{22} u^2 u^2 d\tau^2 + g_{33} u^3 u^3 d\tau^2 = c^2 d\tau^2$$

SINCE METRIC IS DIAGONAL.

$$u^\mu = \frac{d\tau}{d\tau} \left( c, \frac{dr}{d\tau}, \frac{d\theta}{d\tau}, \frac{d\phi}{d\tau} \right) \text{ WHERE}$$

$d\tau$ : SATELLITE PROPER TIME

$d\tau$ : SATELLITE COORDINATE TIME

CONSIDER MOVEMENT OF SATELLITE

$$c^2 d\tau^2 = \left( 1 - \frac{2M}{r} \right) c^2 dt^2 - r^2 \left( \frac{d\theta}{d\tau} \right)^2 dt^2$$

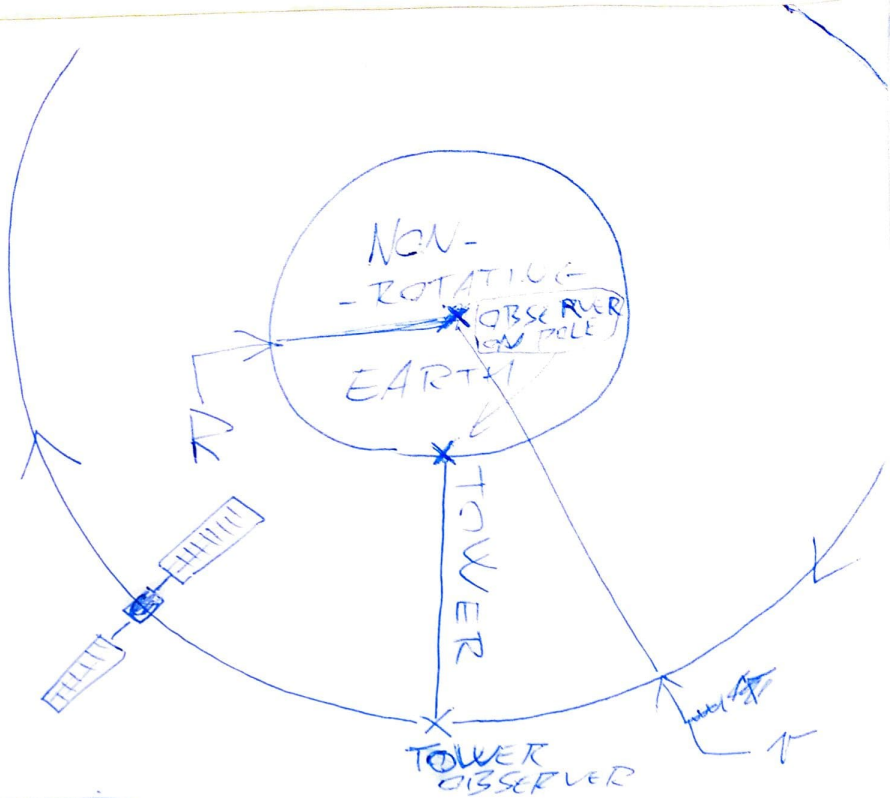
$dr$  &  $d\phi$  ARE ZERO ON CIRCULAR ORBIT.

$$\frac{d\theta}{d\tau} = \sqrt{\frac{GM}{r^3}} = \sqrt{\frac{Mc^2}{r^3}}$$

$$c^2 d\tau^2 = \left( 1 - \frac{2M}{r} \right) c^2 dt^2 - r^2 \frac{Mc^2}{r^3} dt^2$$

$$= \left( 1 - \frac{3M}{r} \right) c^2 dt^2$$

$$\Rightarrow \frac{d\tau}{dt} = \left( 1 - \frac{3M}{r} \right)^{\frac{1}{2}}$$





REL SURF IV (Z) V

CONSIDER POLE & TOWER OBSERVERS

$$c^2 d\tau^2 = c^2 dt^2 g_{00} = c^2 dt^2 \left(1 - \frac{2M}{R}\right)$$

$$\cancel{dt} \left(1 - \frac{2M}{R}\right)^{-\frac{1}{2}} = \frac{dt}{d\tau}$$

UNDERSTANDING IS WEAK HERE. CAN I THINK OF  $dt$  AS PROPER TIME INCREMENT OF SOMEONE IN CENTRE OF EARTH, DISTANCE  $R$  FROM OBSERVER?

CONCLUDE:

$$\frac{\text{PERIOD IN SATELLITE PROPER TIME}}{\text{PERIOD IN POLE OBSERVER PROPER TIME}} = \frac{\Delta \tau_c}{\Delta \tau_o} = \frac{\Delta \tau_o}{\Delta t} \frac{\delta t}{\delta \tau_o}$$

COORDINATE TIME

$$= \left(1 - \frac{3M}{r}\right)^{\frac{1}{2}} \left(1 - \frac{2M}{R}\right)^{-\frac{1}{2}}$$

CONCEPTUAL WORRY:  $r$  IN SCHWARTZSCHILD LINE ELEMENT IS NOT THE SAME AS "NORMAL" RADIUS, IS IT?

WORKING TO FIRST ORDER:

$$\frac{\Delta \tau_c}{\Delta \tau_o} \approx \left(1 - \frac{3M}{2r}\right) \left(1 + \frac{M}{R}\right) \approx 1 - \frac{3M}{2r} + \frac{M}{R}$$

$$= 1 - \frac{3GM}{2rc^2} + \frac{GM}{Rc^2}$$

3/

$$ds^2 = c^2 \left(1 - \frac{2M}{r}\right) dt^2 - \left(1 - \frac{2M}{r}\right)^{-1} dr^2 - r^2 d\theta^2 - r^2 \sin^2 \theta d\phi^2$$

$$u_{\text{ALICE}}^\mu = A \delta_0^\mu$$

ALICE ISN'T MASSLESS SO:

$$\frac{ds^2}{d\tau^2} = \frac{c^2 d\tau^2}{d\tau^2} = g_{\mu\nu} \frac{dx^\mu}{d\tau} \frac{dx^\nu}{d\tau}$$

$$= g_{\mu\nu} u^\mu u^\nu$$

$$= g_{00} u^0 u^0 = c^2 \left(1 - \frac{2M}{r}\right) A^2$$

$$\Rightarrow A = \left(1 - \frac{2M}{r}\right)^{-\frac{1}{2}}$$

FOR MASSLESS PARTICLES, IE FOR THE PHOTON ALICE MEASURES:

$$p^\mu = c \frac{dx^\mu}{d\lambda}$$

ALICE MEASURES PHOTON E TO BE:

$$E = g_{\mu\nu} u^\mu p^\nu = g_{00} A c \frac{dx^0}{d\lambda}$$

$$= c^2 \left(1 - \frac{2M}{r}\right) \left(1 - \frac{2M}{r}\right)^{-\frac{1}{2}} c \dot{x}^0$$

$$= c^2 \left(1 - \frac{2M}{r}\right)^{\frac{1}{2}} c \dot{t}$$

CONSIDER:

$$\left(\frac{ds}{d\lambda}\right)^2 = L = c^2 \left(1 - \frac{2M}{r}\right) \dot{t}^2 - \left(1 - \frac{2M}{r}\right)^{-1} \dot{r}^2 - r^2 \dot{\theta}^2 - r^2 \sin^2 \theta \dot{\phi}^2$$

EULER-LAGRANGE EQ:

$$\frac{d}{d\lambda} \frac{\partial L}{\partial \dot{t}} = \frac{\partial L}{\partial t} \Rightarrow \frac{\partial L}{\partial \dot{t}} = \text{const} \Rightarrow \left(1 - \frac{2M}{r}\right) \dot{t} = E$$

REWRITE ALICE'S ENERGY MEASUREMENT:

$$E = c^2 \left( 1 - \frac{2M}{r} \right)^{\frac{1}{2}} c \dot{t}$$

$$= c^2 \left( 1 - \frac{2M}{r} \right)^{-\frac{1}{2}} c \dot{r} = h \nu$$

(WE MIGHT BE HAVING TOO MANY C-S)

LET'S HAVE AN OBSERVER WHO IS WITH BOB WHEN HE EMITS A PHOTON, BUT IS ALWAYS STATIONARY. LETS CALL HIM CHARLIE.

RATIO OF FREQUENCY MEASUREMENTS BY ALICE & CHARLIE:

$$\frac{\nu_{\text{ALICE}}}{\nu_{\text{CHARLIE}}} = \frac{E_{\text{ALICE}}}{E_{\text{CHARLIE}}} = \frac{\left( 1 - \frac{2M}{r} \right)_{\text{CHARLIE}}}{\left( 1 - \frac{2M}{r} \right)_{\text{ALICE}}}$$

$$1+z = \frac{\nu_{\text{CHARLIE}}}{\nu_{\text{ALICE}}}$$

$$z = \frac{\nu_{\text{CHARLIE}}}{\nu_{\text{ALICE}}} - 1 = \sqrt{\frac{1 - 2M/r}{1 - 2M/r_e}} - 1$$

$$\text{IF } r_e \rightarrow \frac{2GM}{c^2} \text{ IE } r_e \rightarrow 2M \Rightarrow \frac{2M}{r_e} \rightarrow 1 \Rightarrow 1 - \frac{2M}{r_e} \rightarrow 0$$

$$\Rightarrow \sqrt{\frac{1 - 2M/r}{1 - 2M/r_e}} \rightarrow \infty \Rightarrow z \rightarrow \infty$$

AM I RIGHT IN THINKING THAT  $1+z$  IS RATIO OF PHOTONS MEASURED BY STATIONARY OBSERVERS AT THE PLACE OF EMISSION & RECEPTION? IE WHAT I HAVE CALCULATED  $1+z$  IS NOT THE FREQUENCY RATIO MEASURED BY ALICE & BOB.

$$b = \frac{h}{c\dot{t}} = \frac{r^2 \dot{\phi}}{c(1 - \frac{2M}{r})\dot{t}}$$

WE WANT

$$b = r \left(1 - \frac{2M}{r}\right)^{-\frac{1}{2}}$$

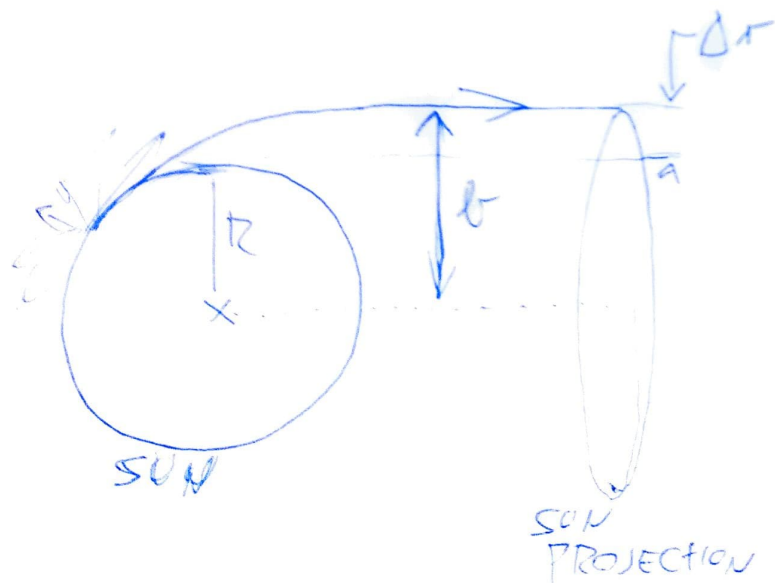
WE THEREFORE HAVE

$$\frac{r \dot{\phi}}{c \dot{t}} = \left(1 - \frac{2M}{r}\right)^{-\frac{1}{2}}$$

REARRANGE:

$$\left(1 - \frac{2M}{r}\right) c^2 \dot{t}^2 - r^2 \dot{\phi}^2 = 0$$

THIS IS TRUE AT CLOSEST  
APPROACH, IE WHEN  $\dot{r} = 0$ .  
(TOPIC 9. EQ 48)



$$\begin{aligned} \delta r &= b - R \\ &= R \left(1 - \frac{2GM}{c^2 R}\right)^{-\frac{1}{2}} - R \\ &\approx 1500 \text{ m} \end{aligned}$$

$\Rightarrow$  DIAMETER IS BIGGER  
BY  $\approx 3 \text{ km}$

[THIS IS PROPER NOT ALGEBRAIC  
[CORRECTING, SAYING IS BETTER  
THAN THIS]



REL IV  
SHEET 5

THIS IS WRONG.

TOPIC XI EQ IS:

$$C(r - r_0) = \frac{2}{3} \left[ \left( \frac{r^3}{2M} \right)^{\frac{1}{2}} - \left( \frac{r_0^3}{2M} \right)^{\frac{1}{2}} \right]$$

$$r_0 = r(r_0 = 0) = 2M$$

$$C r \Big|_{r=0} = \frac{2}{3} [2M - 0] = \frac{4}{3} M$$

(IE FALLING TO THE CENTRE)

$$\Rightarrow r = \frac{4}{3} \frac{M}{C}$$

NOT WHAT'S EXPECTED. I AM MISUSING & PROBABLY:  
IT IS NOT THE CONVENTIONAL RADIUS IN THE  
FORMULA I AM QUOTING.

REGION II REGION I

COLLAPSED  
STAR

$R_s$

MY METRIC IS  
ONLY VALID  
FOR ~~THE~~ VACUUM  
REGIONS.

WHY PARTICLES  
FALL INTO THE  
SINGULARITY,  
WHERE THE METRIC  
ISN'T EVEN VALID?  
(BEC. IT IS NOT  
VACUUM)

FRIEDMANN I EQ:  $\left(\frac{\dot{a}}{a}\right)^2 + \frac{Kc^2}{a^2} = \frac{8\pi G}{3} \rho + \frac{1}{3} \Lambda c^2$

"EMPTY UNIVERSE WITH VANISHING COSMOLOGICAL CONSTANT"

$\Rightarrow \rho = \Lambda = 0$

WE END UP WITH:

$$\left(\frac{\dot{a}}{a}\right)^2 = -\frac{Kc^2}{a^2}$$

$$\dot{a} = \sqrt{-K} c$$

$$a = \sqrt{-K} c t$$

RW LINE ELEMENT :  $ds^2 = c^2 dt^2 - a^2(t) \underbrace{\left[ \frac{dr^2}{1-Kr^2} + r^2 d\Omega^2 \right]}_{d\sigma^2}$

LET'S CHOOSE  $K=1 \Rightarrow a = \sqrt{-K} c t$

$$d\sigma^2|_{K=1} = \frac{dr^2}{1-r^2} + r^2 d\Omega^2$$

PARAMETRIZE BY:  $r = \sinh \chi$   $\chi: 0 \rightarrow \infty$

$$dr = \cosh \chi d\chi$$

$$d\sigma^2|_{K=1} = \frac{\cosh^2 \chi (d\chi)^2}{1 - \sinh^2 \chi} + \sinh^2 \chi d\Omega^2 = \frac{d\chi^2 + \sinh^2 \chi (d\theta^2 + \sin^2 \theta d\phi^2)}{1 - \sinh^2 \chi}$$

$ds^2$  BECOMES:

$$ds^2 = c^2 dt^2 + c^2 t^2 \left[ d\chi^2 + \sinh^2 \chi (d\theta^2 + \sin^2 \theta d\phi^2) \right]$$

AS WALTER

## GEOMETRY OF SPATIAL HYPERSURFACES?

$k=1 \Rightarrow$  HYPERSURFACE IS A 2-SPHERE EMBEDDED  
IN 4D (IE:  $(t, x, y, z) \in \text{SPACETIME}$ ).

- \* SHOW THAT METRIC\* IS MINKOWSKIAN.
- \* RECONCILE

} I DON'T KNOW  
THESE.

SOLUTION TO EINSTEIN'S EQUATIONS  
IF ISOTROPY & HOMOGENEITY ARE  
REQUIRED: FRIEDMANN EQUATIONS.

~~THESE, WITH  $\Lambda = 0$ :~~

$$\textcircled{I} \quad \left(\frac{\dot{a}}{a}\right)^2 + \frac{kc^2}{a^2} = \frac{8\pi G}{3} \rho + \frac{1}{3}\Lambda c^2 \quad (\text{I})$$

$$\left(\frac{\dot{a}}{a}\right) = -\frac{4\pi G}{3} \left(\rho + \frac{3p}{c^2}\right) + \frac{1}{3}\Lambda c^2 \quad (\text{II})$$

$$\rho > 0, p \geq 0$$

$$(\text{II}) \Big|_{\Lambda=0} \Rightarrow \frac{\ddot{a}}{a} = -A^2, \text{ where } A^2 = \frac{4\pi G}{3} \left(\rho + \frac{3p}{c^2}\right) > 0$$

$$\Rightarrow \ddot{a} + A^2 a = 0$$

SHM EQUATION FOR  $a$  WITH FREQUENCY  $A$ .

$\Rightarrow$  NOT STATIC.

STATIC, PRESSURELESS SOL. IF  $\Lambda > 0$ , BUT THIS IS UNSTABLE.



WE WANT  $\dot{a} = 0$  &  $\ddot{a} < 0$   
STATIC UNSTABLE.

$$(\text{I}) \Rightarrow \left(\frac{\dot{a}}{a}\right) = \pm \left( \frac{8\pi G}{3} \rho - \frac{kc^2}{a^2} + \frac{1}{3}\Lambda c^2 \right)^{\frac{1}{2}}$$

$$\dot{a} = \pm \sqrt{\dots}$$



$$\dot{a} = \pm a \sqrt{\frac{8\pi G}{3} S - \frac{kc^2}{a^2} + \frac{1}{3} \Lambda c^2}$$

FROM (II), IF  $\rho=0$ :

$$\frac{8\pi G}{3} S = -2 \frac{\ddot{a}}{a} + \frac{2}{3} \Lambda c^2$$

SUB THIS IN:

$$\dot{a} = \pm \sqrt{-2 \frac{\ddot{a}}{a} + \frac{2}{3} \Lambda c^2 - \frac{kc^2}{a^2} + \frac{1}{3} \Lambda c^2}$$

$$\pm \sqrt{-2 \frac{\ddot{a}}{a} + \Lambda c^2 - \frac{kc^2}{a^2}}$$

~~IF  $k=0$ , THEN  $\ddot{a}$  HAS TO BE NEGATIVE FOR  $\dot{a}$  TO BE PHYSICAL (IE NOT COMPLEX).~~

~~$\dot{a}$  CAN BE NEG~~

$$\text{WE WANT: } \dot{a}=0 \Rightarrow -2 \frac{\ddot{a}}{a} + \Lambda c^2 - \frac{kc^2}{a^2} = 0$$

WITH SUITABLE CHOICE OF  $k$  THIS SEEMS POSSIBLE.

I CANNOT SHOW FROM THIS THAT  ~~$\ddot{a}$~~   $\ddot{a} < 0$ ,  
SO I AM PROBABLY WRONG.