

REL. IV 2

$$L = g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu = \left(\frac{dx}{d\tau} \right)^2$$

$$L = c^2 \left(1 - \frac{2\mu}{r}\right) \dot{t}^2 - \left(1 - \frac{2\mu}{r}\right)^{-1} \dot{r}^2 - r^2 \dot{\theta}^2 - r^2 \sin^2 \theta \dot{\phi}^2$$

EULER-LAGRANGE FOR r :

$$\frac{d}{dr} \left(\frac{\partial L}{\partial \dot{r}} \right) = \frac{d}{dr} \left[-2 \left(1 - \frac{2\mu}{r}\right)^{-1} \dot{r} \right]$$

$$= 2 \left(1 - \frac{2\mu}{r}\right)^{-2} \frac{2\mu}{r^2} \dot{r}^2 - 2 \left(1 - \frac{2\mu}{r}\right)^{-1} \ddot{r}$$

$$\frac{\partial L}{\partial r} = c^2 \frac{2\mu}{r^2} \dot{t}^2 + \left(1 - \frac{2\mu}{r}\right)^{-2} \frac{2\mu \dot{r}^2}{r^2} - 2r \dot{\theta}^2 - 2r \sin^2 \theta \dot{\phi}^2$$

$$\Rightarrow 2 \left(1 - \frac{2\mu}{r}\right)^{-2} \frac{2\mu}{r^2} \dot{r}^2 - 2 \left(1 - \frac{2\mu}{r}\right)^{-1} \ddot{r} = c^2 \frac{2\mu}{r^2} \dot{t}^2 + \left(1 - \frac{2\mu}{r}\right)^{-2} \frac{2\mu \dot{r}^2}{r^2} - 2r \dot{\theta}^2 - 2r \sin^2 \theta \dot{\phi}^2$$

$$\left(1 - \frac{2\mu}{r}\right)^{-2} \frac{2\mu}{r^2} \dot{r}^2 - \ddot{r} = c^2 \frac{\mu}{r^2} \left(1 - \frac{2\mu}{r}\right) \dot{t}^2 + \left(1 - \frac{2\mu}{r}\right)^{-1} \frac{\mu}{r^2} \dot{r}^2$$

$$- \left(1 - \frac{2\mu}{r}\right) r \dot{\theta}^2 - \left(1 - \frac{2\mu}{r}\right) r \sin^2 \theta \dot{\phi}^2$$

$$C.F. \quad \frac{d^2 r}{d\tau^2} + \Gamma_{rr}^r \frac{dr^m}{d\tau} \frac{dr^n}{d\tau} = 0$$

$$\Rightarrow \Gamma_{rr}^r = \left(1 - \frac{2\mu}{r}\right)^{-1} \frac{\mu}{r^2} - \left(1 - \frac{2\mu}{r}\right)^{-1} \frac{2\mu}{r^2} = - \left(1 - \frac{2\mu}{r}\right)^{-1} \frac{\mu}{r^2}$$

$$\Gamma_{tt}^r = c^2 \left(1 - \frac{2\mu}{r}\right) \frac{\mu}{r^2}$$

$$\Gamma_{\theta\theta}^r = - \left(1 - \frac{2\mu}{r}\right) r$$

$$\Gamma_{\phi\phi}^r = - \left(1 - \frac{2\mu}{r}\right) r \sin^2 \theta$$

ALL OTHER Γ_{xy}^r ARE ZERO.

EULER-LAGRANGE FOR Θ :

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) = \frac{\partial L}{\partial \theta}$$

LHS:

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) = \frac{d}{dt} \left(-r^2 \ddot{\theta} \right) = -4r \dot{r} \dot{\theta} - 2r^2 \ddot{\theta}$$

RHS:

$$\frac{\partial L}{\partial \theta} = -r^2 2 \sin \theta \cos \theta \dot{\phi}^2$$

$$\Rightarrow -4r \dot{r} \dot{\theta} - 2r^2 \ddot{\theta} = -r^2 2 \sin \theta \cos \theta \dot{\phi}^2$$

~~cancel~~

/ REARRANGE

$/ : 2r^2$

$$\ddot{\theta} + \frac{2}{r} \dot{r} \dot{\theta} - \sin \theta \cos \theta \dot{\phi}^2 = 0$$

COMPARE WITH GEODESIC EQUATION, CONCLUDE:

$$\Gamma_{r\theta}^\theta = \Gamma_{\theta r}^\theta = \frac{1}{r} \quad !$$

ALL OTHER

Γ_{xy}^θ IS ZERO.

$$\Gamma_{\phi\phi}^\theta = -\sin \theta \cos \theta$$

EULER-LAGRANGE FOR ϕ :

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\phi}} \right) = \frac{\partial L}{\partial \phi}$$

LHS:

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\phi}} \right) = \frac{d}{dt} \left(-2r^2 \sin^2 \theta \dot{\phi} \right) =$$

$$= -4r \dot{r} \sin^2 \theta \dot{\phi} - 4r^2 (\sin \theta)(\cos \theta) \dot{\theta} \dot{\phi} - 2r^2 (\sin^2 \theta) \ddot{\phi}$$

RHS:

$$\frac{\partial L}{\partial \dot{\varphi}} = 0$$

$$LHS = RHS \quad / : -2 \quad \cancel{\text{cancel}}$$

$$\Rightarrow 2r^2 \sin^2 \theta \dot{\phi} + 2r^2 (\sin \theta)(\cos \theta) \ddot{\theta} \dot{\phi} + r^2 \sin^2 \theta \ddot{\phi} = 0 \quad / \text{REARRANGE} \\ : r^2 \sin^2 \theta$$

$$\ddot{\phi} + \frac{2}{r} \dot{\theta} \dot{\phi} + 2 \cot \theta \ddot{\theta} \dot{\phi} = 0$$

COMPARE WITH GEODESIC EQUATION, CONCLUDE:

(THERE ARE TWO
 Γ 'S FOR $\ddot{\theta}$:
 $\Gamma_{\theta\phi}^\theta$ & $\Gamma_{\phi\theta}^\theta$)

$$\Gamma_{\theta r}^\theta = \Gamma_{r\theta}^\theta = \cancel{\frac{2}{r}} \downarrow ! \quad \Gamma_{\theta\phi}^\theta = \Gamma_{\phi\theta}^\theta = \cancel{2 \cot \theta}$$

ALL OTHER Γ_{XY}^θ IS ZERO.

RELATIVITY / 2 (IV)

$$g_{\mu\nu} u^\mu u^\nu = \left(\frac{dt}{d\tau}\right)^2 = c^2$$

$$g_{\mu\nu} u^\mu u^\nu d\tau^2 = c^2 dt^2$$

$$g_{\mu\nu} u^\mu u^\nu d\tau^2 =$$

~~$$= g_{00} u^0 u^0 d\tau^2 + g_{11} u^1 u^1 d\tau^2 + g_{22} u^2 u^2 d\tau^2 + g_{33} u^3 u^3 d\tau^2$$~~

$$= g_{00} u^0 u^0 d\tau^2 + g_{11} u^1 u^1 d\tau^2 + g_{22} u^2 u^2 d\tau^2 + g_{33} u^3 u^3 d\tau^2 = c^2 dt^2$$

SINCE METRIC IS DIAGONAL.

$$u^\mu = \frac{dt}{d\tau} \left(c + \frac{dr}{dt}, \frac{d\theta}{dt}, \frac{d\phi}{dt} \right) \text{ WHERE}$$

CONSIDER MOVEMENT
OF SATELLITE

$d\tau$: SATELLITE PROPER TIME
 dt : SATELLITE COORDINATE TIME

$$c^2 d\tau^2 = \left[1 - \frac{2\mu}{r} \right] c^2 dt^2 + r^2 \left(\frac{d\theta}{dt} \right)^2 dt^2$$

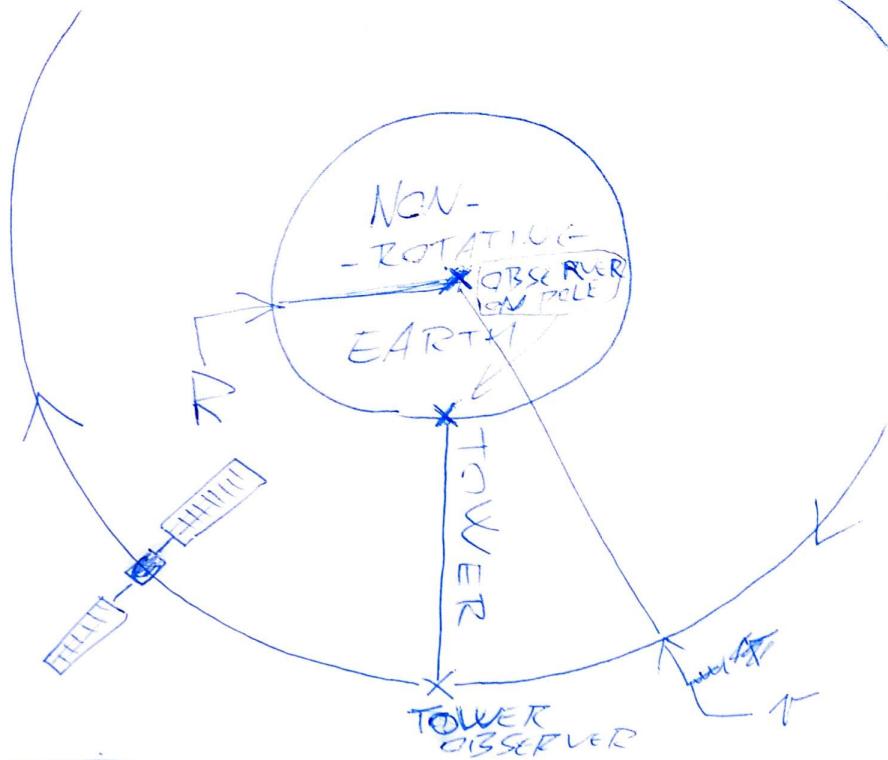
dr & $d\phi$
ARE ZERO
ON CIRCULAR
ORBIT.

$$\frac{d\theta}{dt} = \sqrt{\frac{GM}{r^3}} = \sqrt{\frac{\mu c^2}{r^3}}$$

$$c^2 d\tau^2 = \left(1 - \frac{2\mu}{r} \right) c^2 dt^2 - r^2 \frac{\mu c^2}{r^3} dt^2$$

$$= \left(1 - \frac{3\mu}{r} \right) c^2 dt^2$$

$$\Rightarrow \frac{d\tau}{dt} = \left(1 - \frac{3\mu}{r} \right)^{\frac{1}{2}}$$



REL
SHEET IV (2 ✓)

CONSIDER POLE & TOWER
OBSERVERS

$$c^2 dx^2 = c^2 dt^2 g_{00} = c^2 dt^2 \left(1 - \frac{2\mu}{R}\right)$$

$$\cancel{dt^2} \left(1 - \frac{2\mu}{R}\right)^{-\frac{1}{2}} = \frac{dt}{d\tau}$$

UNDERSTANDING IS
WEAK HERE. CAN I
THINK OF dt AS
PROPER TIME INCREMENT
OF SOMEONE IN CENTRE
OF EARTH, DISTANCE
FROM OBSERVER?

CONCLUDE:

$$\frac{\text{PERIOD IN SATELLITE PROPER TIME}}{\text{PERIOD IN POLE OBSERVER PROPER TIME}} = \frac{\Delta \tau_c}{\Delta \tau_o} = \frac{\Delta \tau_o}{\Delta t} \frac{\Delta t}{\Delta \tau_o}$$

COORDINATE TIME

$$= \left(1 - \frac{3\mu}{c^2}\right)^{\frac{1}{2}} \left(1 - \frac{2\mu}{R}\right)^{-\frac{1}{2}}$$

CONCEPTUAL WORK: & IN SCHWARTZCHILD LINE ELEMENT IS NOT THE SAME AS "NORMAL" RADIUS, IS IT?

WORKING TO FIRST ORDER:

$$\frac{\Delta \tau_o}{\Delta \tau_c} \approx \left(1 - \frac{3\mu}{c^2}\right) \left(1 + \frac{\mu}{R}\right) \approx 1 - \frac{3\mu}{c^2} + \frac{\mu}{R}$$

$$= 1 - \frac{3GM}{c^2 r^2} + \frac{GM}{Rc^2}$$

$$3/ ds^2 = c^2 \left(1 - \frac{2\mu}{r}\right) dt^2 - \left(1 - \frac{2\mu}{r}\right)^{-1} dr^2 - r^2 d\theta^2 - r^2 \sin^2 \theta d\phi^2$$

$$u_{\text{ALICE}}^\mu = A \delta_0^\mu$$

ALICE ISN'T MASSLESS SO:

$$\begin{aligned} \frac{ds^2}{dt^2} &= \frac{c^2 dx^2}{dt^2} = g_{\mu\nu} \frac{dx^\mu}{dt} \frac{dx^\nu}{dt} \\ &= g_{\mu\nu} u^\mu u^\nu \\ &= g_{\mu\nu} u^\mu u^\nu = c^2 \left(1 - \frac{2\mu}{r}\right) A^2 \\ \Rightarrow A &= \left(1 - \frac{2\mu}{r}\right)^{-\frac{1}{2}} \end{aligned}$$

FOR MASSLESS PARTICLES, IE FOR THE PHOTON ALICE MEASURES:

$$p^\mu = c \frac{dx^\mu}{d\tau}$$

ALICE MEASURES PHOTON E TO BE:

$$\begin{aligned} E &= g_{\mu\nu} u^\mu p^\nu = g_{\mu\nu} A^\mu c \frac{dx^\nu}{d\tau} \\ &= c^2 \left(1 - \frac{2\mu}{r}\right) \left(1 - \frac{2\mu}{r}\right)^{-\frac{1}{2}} c \dot{x}^\nu \\ &= c^2 \left(1 - \frac{2\mu}{r}\right)^{\frac{1}{2}} c \dot{t} \end{aligned}$$

CONSIDER:

$$\left(\frac{ds}{d\tau}\right)^2 = L = c^2 \left(1 - \frac{2\mu}{r}\right) \dot{t}^2 - \left(1 - \frac{2\mu}{r}\right)^{-1} \dot{r}^2 - r^2 \dot{\theta}^2 - r^2 \sin^2 \theta \dot{\phi}^2$$

EULER-LAGRANGE EQ:

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{t}} = \frac{\partial L}{\partial t} \Rightarrow \frac{\partial L}{\partial \dot{t}} = \text{const} \Rightarrow \left(1 - \frac{2\mu}{r}\right) \dot{t} = \text{const}$$

REWRITE ALICE'S ENERGY MEASUREMENT:

$$E = c^2 \left(1 - \frac{2\mu}{\tau}\right)^{\frac{1}{2}} c \epsilon$$
$$= c^2 \left(1 - \frac{2\mu}{\tau}\right)^{-\frac{1}{2}} c \epsilon = h \nu$$

(WE MIGHT BE HAVING TOO MANY C-S)

~~REVIEW~~
LET'S HAVE AN OBSERVER WHO IS
WITH BOB WHEN HE EMITS A
PHOTON, BUT IS ALWAYS STATIONARY.
LETS CALL HIM CHARLIE.

RATIO OF FREQUENCY MEASUREMENTS BY ALICE & CHARLIE:

$$\frac{\nu_{\text{ALICE}}}{\nu_{\text{CHARLIE}}} = \frac{E_{\text{ALICE}}}{E_{\text{CHARLIE}}} = \sqrt{\frac{1 - \frac{2\mu}{\tau} |_{\text{CHARLIE}}}{1 - \frac{2\mu}{\tau} |_{\text{ALICE}}}}$$

$$1 + z = \frac{\nu_{\text{CHARLIE}}}{\nu_{\text{ALICE}}}$$

$$z = \frac{\nu_{\text{CHARLIE}}}{\nu_{\text{ALICE}}} - 1 = \sqrt{\frac{1 - \frac{2\mu}{\tau}}{1 - \frac{2\mu}{\tau_0}} - 1}$$

$$\text{IF } \tau_0 \rightarrow \frac{2GM}{c^2} \text{ IE } \tau_0 \rightarrow 2\mu \Rightarrow \frac{2\mu}{\tau_0} \rightarrow 1 \Rightarrow 1 - \frac{2\mu}{\tau_0} \rightarrow 0$$
$$\Rightarrow \sqrt{\frac{1 - \frac{2\mu}{\tau}}{1 - \frac{2\mu}{\tau_0}}} \rightarrow \infty \Rightarrow z \rightarrow \infty$$

AM I RIGHT IN THINKING THAT $1+z$ IS RATIO OF
PHOTONS MEASURED BY STATIONARY OBSERVERS AT
THE PLACE OF EMISSION & RECEPTION? IE WHAT I
HAVE CALCULATED $1+z$ IS NOT THE FREQUENCY RATIO
MEASURED BY ALICE & BOB.

KEL
SUCET IV

$$b = \frac{h}{c\epsilon} = \frac{r^2 \dot{\phi}}{c(1 - \frac{2M}{r})\dot{e}}$$

WE WANT

$$b = r \left(1 - \frac{2M}{r}\right)^{-\frac{1}{2}}$$

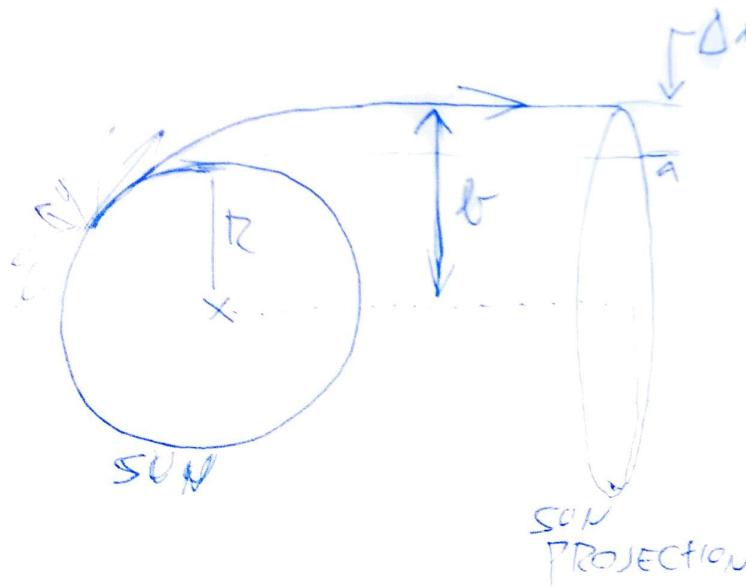
WE THEREFORE HOPE

$$\frac{r \dot{\phi}}{c \dot{e}} = \left(1 - \frac{2M}{r}\right)^{\frac{1}{2}}$$

REARRANGE:

$$\left(1 - \frac{2M}{r}\right) c^2 \dot{e}^2 - r^2 \dot{\phi}^2 = 0$$

THIS IS TRUE AT CLOSEST APPROACH, IE WHEN $\dot{r} = 0$.
(TOPIC 9, EQ48)



$$\Delta r = b - R = R \left(1 - \frac{2GM}{c^2 R}\right)^{\frac{1}{2}} - R$$

$$\approx 1500 \text{ m}$$

\Rightarrow DIAMETER IS BIGGER BY $\approx 3 \text{ km}$

REL SHEET IV | 5. THIS IS WRONG.

TOPIC EXPLANATION:

$$C(z - z_0) = \frac{2}{3} \left[\left(\frac{r^3}{2R} \right)^{\frac{1}{2}} - \left(\frac{r^3}{2R} \right)^{\frac{1}{2}} \right]$$

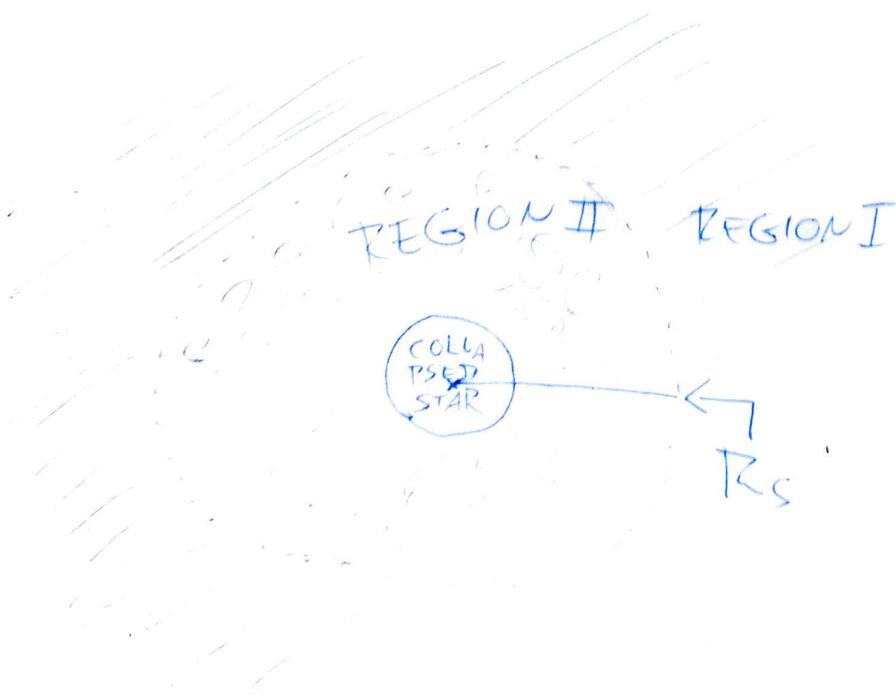
$$z_0 = r (z_0 = 0) = 2R$$

$$Cz = \frac{2}{3} [2R - 0] = \frac{4}{3} R$$

(IE FALLING
TO THE
CENTRE)

$$\Rightarrow z = \frac{4}{3} \underline{R}$$

NOT WHAT'S EXPECTED. I AM MISUSING & PROBABLY:
IT IS NOT THE CONVENTIONAL RADIUS IN THE
FORMULA I AM QUOTING.



MY METRIC IS
ONLY VALID
FOR ~~VACUUM~~ VACUUM
REGIONS.

WHY PARTICLES
FALL IN TO THE
SINGULARITY,
WHEN THE METRIC
ISN'T EVEN VALID?
(B.C. IT IS NOT
VACUUM)

REL. SHEET IV
Q 6 (I)

FRIEDMANN I EQ: $\left(\frac{\dot{a}}{a}\right)^2 + \frac{Kc^2}{a^2} = \frac{8\pi G}{3} S + \frac{1}{3} \Lambda c^2$

"EMPTY UNIVERSE WITH VANISHING COSMOLOGICAL CONSTANT"

$$\Rightarrow S = \Lambda = 0$$

WE END UP WITH:

$$\left(\frac{\dot{a}}{a}\right)^2 = -\frac{Kc^2}{a^2}$$

$$\dot{a} = \sqrt{-K} c$$

$$a = \sqrt{-K} c t$$

RW LINE ELEMENT: $ds^2 = c^2 dt^2 - a^2(t) \left[\frac{dr^2}{1-Kr^2} + r^2 d\Omega^2 \right]$

LET'S CHOOSE $K=1$, $\Rightarrow a = \frac{i c t}{d\Omega^2}$

$$d\Omega^2 \Big|_{K=1} = \frac{dr^2}{1-r^2} + r^2 d\Omega^2$$

PARAMETRIZE BY: $r = \frac{\sinh x}{\cosh x}$ $x: 0 \rightarrow \infty$

$$dr = \cosh x dx$$

$$d\Omega^2 \Big|_{K=1} = \frac{\cosh^2 x dx^2}{1-\sinh^2 x} + \sinh x d\Omega^2 = dx^2 + \sinh^2 x (d\theta^2 + \sin^2 \theta d\phi^2)$$

ds^2 BECOMES:

$$ds^2 = c^2 dt^2 + c^2 t^2 \left[dx^2 + \sinh^2 x (d\theta^2 + \sin^2 \theta d\phi^2) \right]$$

AS WANTED

REL SHEET IV

Q6 (II)

GEOMETRY OF SPATIAL HYPERSURFACES?

$k = 1 \Rightarrow$ HYPERSURFACE IS A 2-SPHERE EMBEDDED
IN 4D (IE: (t, x, θ, ϕ)) SPACETIME.

- SHOW THAT METRIC* IS MINKOWSKIAN }
• RECONCILE } 
- I DON'T KNOW THESE.

Q8

SOLUTION TO EINSTEIN'S EQUATIONS
 IF ISOTROPY & HOMOGENEITY ARE
 REQUIRED: FRIEDMANN EQUATIONS.

THESE, WITH $\Lambda = 0$:

$$\textcircled{1} \quad \left(\frac{\dot{a}}{a}\right)^2 + \frac{kc^2}{a^2} = \frac{8\pi G}{3}S + \frac{1}{3}\Lambda c^2 \quad (\text{I})$$

$$\left(\frac{\ddot{a}}{a}\right) = -\frac{4\pi G}{3}\left(S + \frac{3\Lambda}{c^2}\right) + \frac{1}{3}\Lambda c^2 \quad (\text{II})$$

$$S > 0, p \geq 0$$

$$\text{(II)} \Rightarrow \frac{\ddot{a}}{a} = -A^2, \text{ WHERE } A^2 = \frac{4\pi G}{3}\left(S + \frac{3\Lambda}{c^2}\right) > 0$$

$$\Rightarrow \ddot{a} + A^2 a = 0$$

SHM EQUATION FOR a WITH FREQUENCY A .

\Rightarrow NOT STATIC.

STATIC, PRESSURELESS SOL. IF $\Lambda > 0$, BUT THIS IS UNSTABLE.

~~(1)~~ WE WANT $\dot{a} = 0$ & $\ddot{a} < 0$
STATIC UNSTABLE.

$$\text{(I)} \Rightarrow \left(\frac{\dot{a}}{a}\right) = \pm \sqrt{\frac{8\pi G}{3}S - \frac{kc^2}{a^2} + \frac{1}{3}\Lambda c^2}$$

$$\dot{a} = \pm \sqrt{\frac{8\pi G}{3}S - \frac{kc^2}{a^2} + \frac{1}{3}\Lambda c^2}$$

$$\ddot{a} = \pm a \sqrt{\frac{8\pi G}{3} S - \frac{kc^2}{a^2} + \frac{1}{3} \mathcal{L} c^2}$$

FRAM(II), IF $n=0$:

$$\frac{8\pi G}{3} S = -2 \frac{\ddot{a}}{a} + \frac{2}{3} \mathcal{L} c^2$$

SUB THIS IN:

$$\begin{aligned}\ddot{a} &= \pm \sqrt{-2 \frac{\ddot{a}}{a} + \frac{2}{3} \mathcal{L} c^2 - \frac{kc^2}{a^2} + \frac{1}{3} \mathcal{L} c^2} \\ &= \pm \sqrt{-2 \frac{\ddot{a}}{a} + \mathcal{L} c^2 - \frac{kc^2}{a^2}}\end{aligned}$$

~~If $k=0$, then \ddot{a} has to be negative for \dot{a} to be physical (i.e. not complex).~~

~~it can be zero~~

$$\text{WE WANT: } \dot{a}=0 \Rightarrow -2 \frac{\ddot{a}}{a} + \mathcal{L} c^2 - \frac{kc^2}{a^2} = 0$$

WITH SUITABLE CHOICE OF k THIS SEEMS POSSIBLE.

I CANNOT SHOW FROM THIS THAT ~~$\ddot{a} < 0$~~ ,
SO I AM PROBABLY WRONG.