Superposed Rigid Body Motion (SRBM)

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1 Invariance under SRBM

The balance laws and constitutive functions holds for all motions, including a SRBM. The balance laws and constitutive response functions are form-invariant under a SRBM.

$$T = f(E) \Longrightarrow T^+ = f(E)^+$$

where $\tilde{\mathcal{I}}$ is Cauchy stress tensor and $\tilde{\mathcal{E}}$ is Green–Lagrange strain tensor. Note that the superscript + indicates rigid body motion.

2 Rigid-Body Motion

From the definition of rigid-body motion,

$$\underline{x}^{+} = Q(t) \left(\underline{x} - \underline{x}_{0} \right) + \underline{d}(t) = Q(t)\underline{x} + \underline{c}(t),$$

where

$$QQ^T = Q^TQ = 1,$$
 $\det Q = 1.$ (1)

From the derivative of L.H.S. of (1),

$$\dot{Q}\dot{Q}^T + \dot{Q}\dot{Q}^T = 0 \Longrightarrow \dot{Q}\dot{Q}^T = -\left(\dot{Q}\dot{Q}^T\right)^T.$$

Let $\Omega = \dot{Q}Q^T$, clearly $\Omega = -\Omega^T$, Ω is skew-symmetric tensor. From the definition of dual vector,

$$\Omega \underline{x} = \underline{\omega} \times \underline{x}.$$

$$\therefore \dot{\underline{Q}} = \Omega \underline{Q}$$

$$\underline{v}^+ = \dot{\underline{x}}^+ = \dot{\underline{Q}}\underline{x} + \dot{\underline{Q}}\dot{\underline{x}} + \dot{\underline{c}}(t) = \Omega \underline{Q}\underline{x} + \underline{Q}\underline{v} + \dot{\underline{c}}(t) = \underline{\omega} \times (\underline{x}^+ - \underline{c}) + \underline{Q}\underline{v} + \dot{\underline{c}}(t)$$

$$\Longrightarrow \underline{v}^{+} = \underline{\Omega} \left(\underline{x}^{+} - \underline{c} \right) + \underline{Q}\underline{v} + \dot{\underline{c}}(t)$$

$$\underline{x}^{+} = \underline{Q}(t) + \underline{c}(t) \Longrightarrow \underline{F}^{+} = \frac{\partial \underline{x}^{+}}{\partial \underline{X}} = \frac{\partial \underline{x}^{+}}{\partial \underline{x}} \frac{\partial \underline{x}}{\partial \underline{X}} = \underline{Q}\underline{F}$$

$$J^{+} = \det \underline{F}^{+} = \det \underline{Q} \det \underline{F} = \det \underline{F} = J \Longrightarrow dV^{+} = dV \text{ (isochoric)}.$$

Recall

$$da = da\hat{n} = JF^{-T}dA = JF^{-T}dA\hat{N},$$

then

$$d\underline{a}^+ = \hat{n}^+ da^+ = d\underline{x}_1^+ \times d\underline{x}_2^+ = J^+ \left(\underline{F}^+\right)^{-T} \hat{N} dA = J \left(\underline{Q}\underline{F}\right)^{-T} \hat{N} dA = J \underline{Q}\underline{F}^{-T} \underline{N} dA = \underline{Q} \hat{n} da$$
 and

$$(da^{+})^{2} = d\underline{a}^{+} \cdot d\underline{a}^{+} = \hat{n}^{+} da^{+} \cdot \hat{n}^{+} da^{+} = \underline{Q} \hat{n} da \cdot \underline{Q} \hat{n} da = \hat{n} \cdot \underline{Q}^{T} \underline{Q} \hat{n} (da)^{2} = (da)^{2}$$
$$\therefore da^{+} = da, \qquad \qquad \hat{n}^{+} = \underline{Q} \hat{n}.$$

2.1 Cauchy-Green deformation tensor

$$C = E^T E \Longrightarrow C^+ = (E^+)^T E^+ = (QE)^T (QE) = E^+ \underbrace{Q^T Q}_{=1} E = E^T E = C.$$

2.2 Green–Lagrange strain tensor

$$E = \frac{1}{2} (C - 1) \Longrightarrow E^+ = \frac{1}{2} (C^+ - 1) = \frac{1}{2} (C - 1) = E.$$

2.3 Velocity gradient tensor

$$\dot{\mathcal{L}}^{+} = \frac{\partial \dot{\mathcal{V}}^{+}}{\partial \dot{x}^{+}} = \dot{\Omega} + \dot{\mathcal{Q}} \frac{\partial \dot{\mathcal{V}}}{\partial \dot{x}} \frac{\partial \dot{\mathcal{X}}}{\partial \dot{x}^{+}} \stackrel{*}{=} \dot{\Omega} + \dot{\mathcal{Q}} \dot{\mathcal{L}} \dot{\mathcal{Q}}^{T}$$

$$\dot{\mathcal{D}}^{+} = \frac{1}{2} \left(\dot{\mathcal{L}}^{+} + \dot{\mathcal{L}}^{+T} \right) = \frac{1}{2} \left(\dot{\mathcal{Q}} + \dot{\mathcal{Q}} \dot{\mathcal{L}} \dot{\mathcal{Q}}^{T} + \dot{\mathcal{Q}}^{T} + \dot{\mathcal{Q}} \dot{\mathcal{L}}^{T} \dot{\mathcal{Q}}^{T} \right) = \dot{\mathcal{Q}} \dot{\mathcal{D}} \dot{\mathcal{Q}}^{T}$$

$$\dot{W}^{+} = \frac{1}{2} \left(\dot{\mathcal{L}}^{+} - \dot{\mathcal{L}}^{+T} \right) = \frac{1}{2} \left(\dot{\Omega} + \dot{\mathcal{Q}} \dot{\mathcal{L}} \dot{\mathcal{Q}}^{T} - \dot{\Omega}^{T} - \dot{\mathcal{Q}} \dot{\mathcal{L}}^{T} \dot{\mathcal{Q}}^{T} \right) = \dot{\Omega} + \dot{\mathcal{Q}} \dot{\mathcal{W}} \dot{\mathcal{Q}}^{T}$$

$$\dot{\mathcal{Q}}^{-1} = \dot{\mathcal{Q}}^{T}, \, \dot{\chi} = \dot{\mathcal{Q}}^{T} \left(\dot{\chi}^{+} - \dot{\mathcal{C}}(t) \right)$$

2.4 Cauchy stress and tensor

We know that $\hat{n}^+ = Q\hat{n}$, and assume the component of stress vector \underline{t} in the direction of the outward normal \hat{n}^+ remains unchanged under SRBM so that

From the relation between stress vector and stress tensor, $t = T\hat{n}$ so that

2.5 Body force

$$\nabla \cdot T = \rho \left(\dot{v} - \dot{b} \right) \Longrightarrow \nabla^+ \cdot T^+ = \rho^+ \left(\dot{v}^+ - \dot{b}^+ \right)$$

From conservation of mass, $\rho_0 = \rho J = \rho^+ J^+ \Longrightarrow \rho = \rho^+$, and

$$\nabla^{+} \cdot \tilde{\mathcal{I}}^{+} = \frac{\partial T_{ij}^{+}}{\partial x_{j}^{+}} \hat{e}_{i} = \frac{\partial T_{ij}^{+}}{\partial x_{k}} \frac{\partial x_{k}}{\partial x_{j}^{+}} \hat{e}_{i} = \frac{\partial T_{ij}}{\partial x_{k}} Q_{jk} \hat{e}_{i} = Q_{im} \frac{\partial T_{mn}}{\partial x_{k}} \underbrace{Q_{jn} Q_{jk}}_{=\delta_{nk}} \hat{e}_{i} = Q_{im} \frac{\partial T_{mk}}{\partial x_{k}} \hat{e}_{i}$$

$$= Q \left(\nabla \cdot \tilde{\mathcal{I}} \right)$$

$$\implies \nabla^{+} \cdot \tilde{\mathcal{I}}^{+} = Q \left(\nabla \cdot \tilde{\mathcal{I}} \right) = \rho^{+} \left(\dot{v}^{+} - \dot{b}^{+} \right) \Longrightarrow Q \rho \left(\dot{v} - \dot{b} \right) = \rho^{+} \left(\dot{v}^{+} - \dot{b}^{+} \right)$$

$$\implies (\dot{v}^{+} - \dot{b}^{+}) = Q \left(\dot{v} - \dot{b} \right)$$

3 Derivation of Balance Laws from Balance of Energy

Recall balance of energy

$$\dot{K} + \dot{U} = P + Q$$

where

$$K = \int_{V} \frac{1}{2} \rho y \cdot y dV$$
$$U = \int_{V} \rho \varepsilon dV$$

$$\begin{split} P &= \int_{V} \rho \underline{b} \cdot \underline{v} dV + \int_{S} \underline{t} \cdot \underline{v} da \\ Q &= \int_{V} \rho r dV - \int_{S} h da, h(\hat{n}) = \underline{q} \cdot \hat{n} \end{split}$$

Assume that $\rho^+ = \rho$, $\varepsilon^+ = \varepsilon$, $r^+ = r$, $\tilde{q}^+ \cdot \hat{n}^+ = \tilde{q} \cdot \hat{n} \Longrightarrow \tilde{q}^+ = Q\tilde{q}$

$$\dot{K} + \dot{U} = P + Q \Longrightarrow \rho \dot{\varepsilon} = \tilde{\chi} \cdot \tilde{D} + \rho r - \nabla \cdot q$$

From

$$\frac{D}{Dt} \int_{V} \frac{1}{2} \rho \underline{v} \cdot \underline{v} dV + \frac{D}{Dt} \int_{V} \rho \varepsilon dV = \int_{V} \rho \underline{b} \cdot \underline{v} dV + \int_{S} \underline{t} \cdot \underline{v} da + \int_{V} \rho r dV - \int_{S} \underline{q} \cdot \hat{n} da$$

$$\frac{D}{Dt} \int_{V^{+}} \rho^{+} \left(\frac{1}{2} \underline{v}^{+} \cdot \underline{v}^{+} + \varepsilon^{+} \right) dV^{+} = \int_{V^{+}} \rho^{+} \underline{b}^{+} \cdot \underline{v}^{+} dV^{+} + \int_{S^{+}} \underline{t}^{+} \cdot \underline{v}^{+} da^{+}$$

$$+ \int_{V^{+}} \rho^{+} r^{+} dV^{+} - \int_{S^{+}} \underline{q}^{+} \cdot \hat{n}^{+} da^{+}$$
(*)

and transport theorem

$$\frac{D}{Dt} \int_{V} \phi dV = \int_{V} \left[\frac{D\phi}{Dt} + (\nabla \cdot v) \phi \right] dV \Longrightarrow \frac{D}{Dt} \int_{V^{+}} \phi dV^{+} = \int_{V^{+}} \left[\frac{D\phi^{+}}{Dt} + (\nabla^{+} \cdot v^{+}) \phi^{+} \right] dV^{+}$$

$$\frac{D}{Dt} \int_{V^{+}} \rho^{+} \phi^{+} dV^{+} = \int_{V^{+}} \left[\dot{\rho}^{+} \phi^{+} + \rho^{+} \dot{\phi}^{+} + (\nabla^{+} \cdot v^{+}) \rho^{+} \phi^{+} \right] dV^{+}$$

$$= \int_{V^{+}} \left[\left(\dot{\rho}^{+} + \rho^{+} \nabla^{+} \cdot v^{+} \right) \phi^{+} + \rho^{+} \dot{\phi}^{+} \right] dV^{+}$$

(*) becomes

$$\begin{split} \int_{V^{+}} \left[\left(\dot{\rho}^{+} + \rho^{+} \nabla^{+} \cdot \underline{v}^{+} \right) \left(\frac{1}{2} \underline{v}^{+} \cdot \underline{v}^{+} + \varepsilon^{+} \right) + \rho^{+} \left(\underline{v}^{+} \cdot \dot{\underline{v}}^{+} + \dot{\varepsilon}^{+} \right) \right] dV^{+} \\ &= \int_{V^{+}} \rho^{+} \underline{b}^{+} \cdot \underline{v}^{+} dV^{+} + \int_{S^{+}} \underline{t}^{+} \cdot \underline{v}^{+} da^{+} + \int_{V^{+}} \rho^{+} r^{+} dV^{+} - \int_{S^{+}} \underline{q}^{+} \cdot \hat{n}^{+} da^{+} \end{split}$$

where

$$\begin{split} \int_{S^{+}} & \not z^{+} \cdot y^{+} da^{+} = \int_{S^{+}} t_{i}^{+} v_{i}^{+} da^{+} = \int_{S^{+}} T_{ij}^{+} n_{j}^{+} v_{i}^{+} da^{+} = \int_{V^{+}} \left(T_{ij}^{+} v_{i}^{+} \right)_{,j} dV^{+} \\ & = \int_{V^{+}} \left(\frac{\partial T_{ij}^{+}}{\partial x_{j}^{+}} v_{i}^{+} + T_{ij}^{+} \frac{\partial v_{i}^{+}}{\partial x_{j}^{+}} \right) dV^{+} = \int_{V^{+}} \left[y^{+} \cdot \left(\nabla^{+} \cdot \vec{x}^{+} \right) + \vec{x}^{+} \cdot \vec{L}^{+} \right] dV^{+} \end{split}$$

Recall

$$L_{ij} = \frac{\partial v_i}{\partial x_j} \Longrightarrow L_{ii} = \frac{\partial v_i}{\partial x_i} \Longrightarrow \nabla \cdot \underline{v} = \operatorname{tr} \underline{L}$$

and

$$\nabla^{+} \cdot \chi^{+} = \operatorname{tr} \mathcal{L}^{+} = \operatorname{tr} \left(\Omega + Q \mathcal{L} Q^{T} \right) = \operatorname{tr} \Omega + \operatorname{tr} \left(Q \mathcal{L} Q^{T} \right) = \operatorname{tr} \left(Q \mathcal{L} Q^{T} \right)$$

$$= Q_{ij} L_{jk} Q_{ik} = L_{jk} Q_{ik} Q_{ij} = \operatorname{tr} \left(\mathcal{L} Q^{T} Q \right) = \operatorname{tr} \mathcal{L} = \nabla \cdot \chi$$

$$T^{+} \cdot \mathcal{L}^{+} = \left(Q \mathcal{T} Q^{T} \right) \cdot \left(\Omega + Q \mathcal{L} Q^{T} \right) = \left(Q \mathcal{T} Q^{T} \right) \cdot \Omega + \left(Q \mathcal{T} Q^{T} \right) \cdot \left(Q \mathcal{L} Q^{T} \right)$$

$$\stackrel{\dagger}{=} \left(Q \mathcal{T} Q^{T} \right) \cdot \Omega + \operatorname{tr} \left(Q \mathcal{T} Q^{T} Q \mathcal{L} Q^{T} \right) = \left(Q \mathcal{T} Q^{T} \right) \cdot \Omega + \operatorname{tr} \left(Q \mathcal{T} \mathcal{L} Q^{T} \right)$$

$$= \left(Q \mathcal{T} Q^{T} \right) \cdot \Omega + \operatorname{tr} \left(\mathcal{T} \mathcal{L}^{T} Q^{T} Q \mathcal{L} Q^{T} \right) = \left(Q \mathcal{T} Q^{T} \right) \cdot \Omega + \operatorname{tr} \left(\mathcal{T} \mathcal{L}^{+} \right)$$

$$= \left(Q \mathcal{T} Q^{T} \right) \cdot \Omega + \operatorname{tr} \left(\mathcal{T} \mathcal{L}^{T} Q^{T} Q \mathcal{L} Q^{T} \right) = \left(Q \mathcal{T} Q^{T} \right) \cdot \Omega + \operatorname{tr} \left(\mathcal{T} \mathcal{L}^{+} \right)$$

$$= \left(Q \mathcal{T} Q^{T} \right) \cdot \Omega + \mathcal{T} \cdot \mathcal{L}$$

$$\nabla^{+} \cdot \mathcal{Q}^{+} = \frac{\partial q_{i}^{+}}{\partial x_{i}^{+}} = \frac{\partial Q_{ik} q_{k}}{\partial x_{i}} \frac{\partial x_{j}}{\partial x_{i}^{+}} = Q_{ik} \frac{\partial q_{k}}{\partial x_{i}} Q_{ji}^{T} = \delta_{kj} \frac{\partial q_{k}}{\partial x_{i}} = \frac{\partial q_{j}}{\partial x_{i}} = \nabla \cdot \mathcal{Q}$$

then

$$\implies \left(\dot{\rho}^{+} + \rho^{+} \nabla \cdot \underline{v}\right) \left(\frac{1}{2} \underline{v}^{+} \cdot \underline{v}^{+} + \varepsilon^{+}\right) + \underline{v}^{+} \cdot \left(\rho^{+} \underline{\dot{v}}^{+} - \rho^{+} \underline{b}^{+} - Q \nabla \cdot \underline{T}\right) + \left(\rho^{+} \dot{\varepsilon}^{+} - \rho^{+} r^{+} - \underline{T} \cdot \underline{L} - Q \underline{T} \underline{Q}^{T} \cdot \underline{\Omega} + \nabla \cdot \underline{q}\right) = 0$$

Assume $\rho^+ = \rho$, $\varepsilon^+ = \varepsilon$, $r^+ = r$, and recall $(\dot{v}^+ - \dot{v}^+) = Q(\dot{v} - \dot{v})$ so that

$$\Longrightarrow (\dot{\rho} + \rho \nabla \cdot \underline{v}) \left(\frac{1}{2} \underline{v}^{+} \cdot \underline{v}^{+} + \varepsilon \right) + \underline{v}^{+} \cdot \underline{Q} \left(\rho \dot{\underline{v}} - \rho \underline{b} - \nabla \cdot \underline{\overline{v}} \right)$$

$$+ \left(\rho \dot{\varepsilon} - \rho r - \underline{T} \cdot \underline{L} - \underline{Q} \underline{T} \underline{Q}^{T} \cdot \underline{\Omega} + \nabla \cdot \underline{q} \right) = 0$$

$$(\bigstar)$$

for all $\underline{x}^+ = Q(t)\underline{x} + \underline{c}(t)$

$$^{\dagger} \underline{A} \cdot \underline{B} = A_{ij} B_{ij} = \operatorname{tr} \left(\underline{A} \underline{B}^T \right)$$

• No SRBM: Q = 1, c = 0, $\Omega = 0$, $x^+ = x$, and $x^+ = y$ (\bigstar) becomes

$$\Longrightarrow (\dot{\rho} + \rho \nabla \cdot \underline{v}) \left(\frac{1}{2} \underline{v} \cdot \underline{v} + \varepsilon \right) + \underline{v} \cdot (\rho \dot{\underline{v}} - \rho \underline{b} - \nabla \cdot \underline{\tau})$$

$$+ \left(\rho \dot{\varepsilon} - \rho r - \underline{\tau} \cdot \underline{L} + \nabla \cdot \underline{q} \right) = 0$$

$$(2)$$

• Translation: Q = 1, $\dot{c}(t) = v_0 \hat{e}$, $\dot{Q} = 0$, $\Omega = 0$, $\dot{x}^+ = \dot{x}$, and $\dot{x}^+ = \dot{x} + v_0 \hat{e}$

$$\implies (\dot{\rho} + \rho \nabla \cdot \underline{v}) \left(\frac{1}{2} \underline{v} \cdot \underline{v} + \underline{v} \cdot (v_0 \hat{e}) + \frac{v_0^2}{2} + \varepsilon \right)$$

$$+ (\underline{v} + v_0 \hat{e}) \cdot (\rho \dot{\underline{v}} - \rho \underline{b} - \nabla \cdot \underline{T})$$

$$+ (\rho \dot{\varepsilon} - \rho r - \underline{T} \cdot \underline{L} + \nabla \cdot \underline{q}) = 0$$

$$(3)$$

$$(3) - (2) \Longrightarrow (\dot{\rho} + \rho \nabla \cdot \underline{v}) \left(\underline{v} \cdot (v_0 \hat{e}) + \frac{v_0^2}{2} \right) + v_0 \hat{e} \cdot \left(\rho \dot{\underline{v}} - \rho \underline{b} - \nabla \cdot \underline{T} \right) = 0 \forall v_0, \hat{e}$$

$$\Longrightarrow \begin{cases} \dot{\rho} + \rho \nabla \cdot \underline{v} = 0 \\ \rho \dot{\underline{v}} - \rho \underline{b} - \nabla \cdot \underline{T} = 0 \end{cases}$$

then (\bigstar) becomes

$$\rho \dot{\varepsilon} - \rho r - \tilde{\chi} \cdot \tilde{L} - \tilde{Q} \tilde{\chi} \tilde{Q}^T \cdot \tilde{\Omega} + \nabla \cdot \tilde{q} = 0$$

from $\rho \dot{\varepsilon} - \rho r - \tilde{\chi} \cdot \tilde{L} + \nabla \cdot \tilde{q} = 0$,

$$QTQ^T \cdot \Omega = 0 \Longrightarrow T = T^T$$

Summary.

$$\chi^{+} = Q(t)\chi + \zeta(t)$$

$$E = QE, \quad J^{+} = J, \quad \hat{n}^{+} = Q\hat{n}, \quad da^{+} = da, \quad dV^{+} = dV, \quad \rho^{+} = \rho, \quad E^{+} = E$$

$$\chi^{+} = Q(t)\chi + \zeta(t)$$

$$Q = QE, \quad Q^{+} = C, \quad E^{+} = E$$

$$\chi^{+} = Q(\chi^{+} - \zeta) + Q\chi + \dot{\zeta}(t), \qquad \qquad \Omega = \dot{Q}Q^{T}$$

$$\underline{L}^+ = \underline{\Omega} + \underline{Q}\underline{L}\underline{Q}^T, \quad \underline{D}^+ = \underline{Q}\underline{D}\underline{Q}^T, \quad \underline{W}^+ = \underline{\Omega} + \underline{Q}\underline{W}\underline{Q}^T, \quad \underline{t}^+ = \underline{Q}\underline{t}, \quad \underline{T}^+ = \underline{Q}\underline{T}\underline{Q}^T$$

4 Constitutive Equation

4.1 Ideal fluid

An ideal fluid has the following properties:

- It cannot sustain shearing motion.
- The stress vector acting on any surface in the fluid is always directed along the normal to the surface, i.e. \underline{t} is parallel to \underline{n} .

Assume that $\tilde{T} = \hat{T}(\rho, v)$. The constitutive equation holds for all motions, including SRBM.

$$T = \hat{T}(\rho, v) \Longrightarrow T^+ = \hat{T}(\rho^+, v^+)$$

A translation: Q = 1, $z(t) = v_0 t \Longrightarrow z^+ = z + v_0$

$$\Longrightarrow \mathcal{T} = \mathcal{T}^+ \Longrightarrow \hat{\mathcal{T}}(\rho, \underline{v}) = \hat{\mathcal{T}}(\rho^+, \underline{v}^+) \Longrightarrow \hat{\mathcal{T}}(\rho, \underline{v}) = \hat{\mathcal{T}}(\rho, \underline{v} + \underline{v}_0) \forall \underline{v}_0$$

$$\Longrightarrow \mathcal{T} = \hat{\mathcal{T}}(\rho) \Longrightarrow \mathcal{T}^+ = \hat{\mathcal{T}}(\rho)$$

$$\mathcal{T}^+ = Q \mathcal{T} Q^T \Longrightarrow \hat{\mathcal{T}}(\rho) = Q \hat{\mathcal{T}}(\rho) Q^T \forall Q$$

 $\hat{\mathcal{I}}(\rho)$ is 2^{nd} order isotropic tensor. Obviously, $\hat{\mathcal{I}}(\rho) = f(\rho)\mathbf{1}$ is a solution. Are there other solutions?

$$T = -p(\rho)$$
1.

Equation of motion

$$\rho\left(\underline{b} - \underline{\dot{v}}\right) = -\nabla \cdot \underline{T} = -T_{ij,j}\hat{e}_i = \left(p\delta_{ij}\right)_{,j}\hat{e}_i = p_{,i}\hat{e}_i = \nabla p$$

$$\Longrightarrow \boxed{\rho\underline{\dot{v}} = \rho\underline{b} - \nabla p}$$

which is famous Euler's Equation!

4.2 Viscous fluid

Assume that $\underline{T} = \overline{T}(\rho, \underline{v}, \underline{L}) = \overline{\underline{T}}(\rho, \underline{v}, \underline{D}, \underline{W})$, under SRBM $\underline{T}^+ = \overline{\underline{T}}(\rho^+, \underline{v}^+, \underline{D}^+, \underline{W}^+)$, also $\underline{T}^+ = \underline{Q}\underline{T}\underline{Q}^T$

$$\Longrightarrow Q\bar{T}(\rho, v, D, W)Q^T = \bar{T}(\rho^+, v^+, D^+, W^+)$$

and recall

$$\mathcal{D}^{+} = \mathcal{Q}\mathcal{D}\mathcal{Q}^{T}, \qquad \qquad \mathcal{W}^{+} = \mathcal{Q}^{+} + \mathcal{Q}\mathcal{W}\mathcal{Q}^{T}$$

$$\begin{split} \text{Translation: } & \ \underline{v}^+ = \underline{v} + \underline{v}_0, \ \underline{Q} = \underline{1}, \ \underline{\Omega} = \underline{0}, \ \underline{D}^+ = \underline{D}, \ \text{and} \ \underline{W}^+ = \underline{W} \\ \\ & \Longrightarrow \bar{\underline{T}}(\rho,\underline{v},\underline{D},\underline{W}) = \bar{\underline{T}}(\rho,\underline{v} + \underline{v}_0,\underline{D},\underline{W}) \forall \underline{v}_0 \end{split}$$

$$\Longrightarrow T = \bar{T}(\rho, D, W)$$

Rotation: $\underline{x}^+ = \underline{Q}(t) + \underline{c}(t)$; $\underline{c}(t) = 0$, $\underline{\Omega} = \underline{\Omega}_0 = \text{constant}$. At the instant: $\underline{Q} = \underline{1}$, $\underline{D}^+ = \underline{D}$, and $\underline{W}^+ = \underline{\Omega}_0 + \underline{W}$

$$\Longrightarrow \bar{T}(\rho, \underline{\mathcal{Q}}, \underline{\mathcal{W}}) = \bar{T}(\rho, \underline{\mathcal{Q}}, \underline{\mathcal{W}} + \underline{\Omega}_0) \forall \underline{\Omega}_0$$
$$\Longrightarrow T = \bar{T}(\rho, \underline{\mathcal{Q}})$$

• Newtonian fluid

 $\underline{T} = \underline{T}(\rho, \underline{D})$, linear function of \underline{D} with the coefficients depending on ρ so that

$$\bar{T}_{ij}(\rho, D_{ij}) = A_{ij}(\rho) + C_{ijkl}(\rho)D_{kl}$$

where $C_{ijkl} = C_{jikl} = C_{jilk}$ because of symmetry: $T_{ij} = T_{ji}$ and $D_{ij} = D_{ji}$

$$Q\bar{\mathcal{T}}(\rho, D)Q^T = \bar{\mathcal{T}}(\rho^+, D^+),$$
 $D^+ = QDQ^T$

$$\Longrightarrow Q\bar{T}(\rho, D)Q^T = \bar{T}(\rho, QDQ^T) \Longrightarrow Q_{ij}\bar{T}_{jk}(\rho, D_{pq})Q_{km}^T = \bar{T}_{im}(\rho, Q_{pr}D_{rs}Q_{sq}^T)$$

$$\Longrightarrow Q_{ij}A_{jk}(\rho)Q_{km}^T + Q_{ij}C_{jkrs}(\rho)D_{rs}Q_{km}^T = A_{im}(\rho) + C_{impq}Q_{pr}D_{rs}Q_{sq}^T$$

choose D = 0,

$$\Longrightarrow Q_{ij}A_{jk}(\rho)Q_{km}^T = A_{im}(\rho) \Longrightarrow A = -p(\rho)1$$

and

$$Q_{ij}C_{jkrs}D_{rs}Q_{km}^{T} = C_{impq}Q_{pr}D_{rs}Q_{sq}^{T} \Longrightarrow Q_{ij}Q_{km}^{T}C_{jkrs}D_{rs} = Q_{pr}Q_{sq}^{T}C_{impq}D_{rs}$$

$$\Longrightarrow Q_{ij}Q_{km}^{T}C_{jkrs} = Q_{pr}Q_{sq}^{T}C_{impq} \Longrightarrow Q_{ru}^{T}Q_{vs}Q_{ij}Q_{km}^{T}C_{jkrs} = Q_{ru}^{T}Q_{vs}Q_{pr}Q_{sq}^{T}C_{impq}$$

$$\Longrightarrow Q_{ru}^{T}Q_{vs}Q_{ij}Q_{km}^{T}C_{jkrs} = \delta_{pu}\delta_{vq}C_{impq} \Longrightarrow Q_{ij}Q_{mk}Q_{ur}Q_{vs}C_{jkrs} = C_{imuv}$$

or

$$C_{ijkl} = Q_{ip}Q_{jq}Q_{kr}Q_{ls}C_{pqrs}$$

which is 4th order isotropic tensor

$$C_{ijkl} = \lambda \delta_{ij} \delta_{kl} + \mu \delta_{ik} \delta_{jl} + \gamma \delta_{il} \delta_{jk}$$

owing to $C_{ijkl} = C_{jikl} = C_{jilk}$,

$$C_{ijkl} = \lambda \delta_{ij} \delta_{kl} + \mu \left(\delta_{ik} \delta_{jl} + \delta_{il} \delta_{jk} \right)$$

then

$$T_{ij} = -p\delta_{ij} + \lambda D_{kk}\delta_{ij} + 2\mu D_{ij}$$

Recall $D_{ij} = (v_{i,j} + v_{j,i})/2$ so that

$$\nabla \cdot \tilde{T} = T_{ij,j} \hat{e}_i = \left(-p \delta_{ij} + \lambda D_{kk} \delta_{ij} + 2\mu D_{ij} \right)_{,j} \hat{e}_i$$
$$= \left[-p_{,i} + (\lambda + \mu) v_{k,ki} + \mu v_{i,kk} \right] \hat{e}_i$$

Equation of motion

$$\rho\left(\dot{\underline{b}} - \dot{\underline{v}}\right) = -\nabla \cdot \underline{T} = -T_{ij,j}\hat{e}_{i}$$

$$\Longrightarrow \left[\rho\dot{\underline{v}} = \rho\underline{b} - \nabla p + (\lambda + \mu)\nabla\left(\nabla \cdot \underline{v}\right) + \mu\nabla^{2}\underline{v}\right]$$

which is famous Navier-Stokes Equation!