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# Random Self-similar Trees: Emergence of Scaling Laws

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1   **Abstract** The hierarchical organization and emergence of scaling laws in  
2   complex systems – geophysical, biological, technological, and socio-economic  
3   – have been the topic of extensive research at the turn of the 20th century.  
4   Although significant progress has been achieved, the mathematical origin  
5   of and relation among scaling laws for different system attributes remain  
6   unsettled. Paradigmatic examples are the Gutenberg-Richter law of seismology  
7   and Horton’s laws of geomorphology. We review the results that clarify the  
8   appearance, parameterization, and implications of scaling laws in hierarchical  
9   systems conceptualized by tree graphs. A recently formulated theory of random  
10   self-similar trees yields a suite of results on scaling laws for branch attributes,  
11   tree fractal dimension, power-law distributions of link attributes, and power-  
12   law relations between distinct attributes. Given the relevance of power laws  
13   to extreme events and hazards, our review informs related theoretical and  
14   modeling efforts and provides a framework for unified analysis in hierarchical  
15   complex systems.

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17 similarity · Horton's laws · Tokunaga model · Hierarchical Branching  
18 Process

19 **Article Highlights**

- 20 • Theory of random self-similar trees provides a unifying framework for  
21 studying scaling laws in complex systems  
22 • Hierarchical Branching Process explains power laws for system attributes,  
23 system fractal dimension, and other scalings  
24 • A one-parameter critical Tokunaga model closely fits the key data and  
25 scalings of geomorphology

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34 **Code availability.** The algorithms for the reviewed models are described in  
35 the text. The codes are available upon request.

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## 86 1 Motivation

87 The emergence of extremes in complex natural systems – tectonic, hydrological,  
 88 climatic, biological – has been a topic of extensive research, recognizing the  
 89 catastrophic impact that hazards produced by these systems exert on popula-  
 90 tion, economy and the environment. The complexity science, developed and  
 91 proliferated at the turn of the 20th century, recognizes several fundamental  
 92 traits shared by natural extremes and hazards of different origin: (i) They are  
 93 generic in systems that have *hierarchical organization*. Notably, such a hierarchy  
 94 “*need not be manifest in the object but may arise in the construction of a model*”  
 95 (Badii and Politi, 1999). (ii) The hazard-generating systems exhibit *scalings*,  
 96 often expressed as power-law distributions and/or power-law relations among  
 97 system’s attributes (Barenblatt, 1996). (iii) The scalings are often connected  
 98 to *self-similarity* – a property of the system to retain its statistical properties  
 99 after being zoomed in or out via a suitable transformation (Mandelbrot, 1982;  
 100 Barenblatt, 1996). The results reviewed in this work originated in analysis of  
 101 systems whose hierarchical organization is particularly evident, and is com-  
 102 monly represented by a *tree graph*, and whose scalings and self-similarity are  
 103 well established empirically.

104 An example of such a system is the Earth lithosphere that generates *earth-*  
 105 *quakes*. The two staples of statistical seismology are the power-law distribution  
 106 of earthquake moments, which is equivalent to the celebrated Gutenberg-Richter  
 107 law of earthquake magnitudes (Gutenberg and Richter, 1954), and power-law  
 108 temporal decay of event rate within aftershock series (Omori, 1894; Utsu et al.,  
 109 1995). Multiple other physical and statistical scalings of earthquake attributes  
 110 are summarized in (Ben-Zion, 2008, Table 2). Hierarchical representation of  
 111 seismicity by branching processes is also well explored, starting from the pio-  
 112 neering works of Kagan (1973), Kagan and Knopoff (1976, 1981), and Vere-Jones  
 113 (1976). A very popular Epidemic Type Aftershock Sequence (ETAS) model of  
 114 earthquake dynamics (Ogata, 1998), is a Galton-Watson branching process with  
 115 a power-law offspring distribution and space-time-magnitude marks (Saichev  
 116 et al., 2005; Baró, 2020). A trajectory of this process is a tree graph, whose  
 117 vertices represent individual earthquakes and edges – triggering processes. Alter-  
 118 native tree representations of seismicity are discussed in (Baiesi and Paczuski,  
 119 2004; Holliday et al., 2008; Zaliapin et al., 2008; Yoder et al., 2013; Zaliapin  
 120 and Ben-Zion, 2013). More conceptually, the lithosphere can be thought as “*a*  
 121 *hierarchy of blocks separated by boundary zones, with densely fractured nodes*  
 122 *at junctions and intersections*” (Keilis-Borok, 2002). This hierarchy spans a  
 123 wide range from the seven major tectonic plates of continental size to nearly  
 124  $10^{25}$  grains of rocks. The earthquakes are produced by complex dynamics and  
 125 interaction of these blocks (Burridge and Knopoff, 1967; Allegre et al., 1982;  
 126 Gabrielov et al., 1990; Zaliapin et al., 2003; Soloviev and Ismail-Zadeh, 2003).

127 Another prime hierarchical system prone to natural hazards is geomorpho-  
 128 logical landscape evolution, which is associated with and in part driven by  
 129 mass-movement processes like *sediment transport, rockfalls, debris flows*, and  
 130 *landslides*. One of the system’s fundamental outputs is the network of stream

131 channels that spans the continental earth in the form of permanent river and  
132 delta networks and ephemeral drainage pathways that extend to the grain  
133 scales. This network is naturally linked to such hazards as *floods* (Gupta et al.,  
134 1994, 2007; Tessler et al., 2015) and *coastal and hillslope erosion* (Roering  
135 et al., 1999). Geomorphological networks are conventionally represented by  
136 trees (for converging river channels) or directed acyclic graphs (for diverging  
137 deltaic systems or braided rivers) (Sapozhnikov and Foufoula-Georgiou, 1996;  
138 Lashermes et al., 2007; Passalacqua et al., 2010; Tejedor et al., 2017, 2015a,b),  
139 and their scaling laws have been recognized since the groundbreaking work  
140 of Horton (Horton, 1945); see (Rodriguez-Iturbe and Rinaldo, 2001, and refs.  
141 therein).

142 Geomorphic hazards like *landslides*, *avalanches*, and *forest-fires* are char-  
143 acterized by scaling laws similar to those in seismicity (Malamud et al., 1998;  
144 Turcotte et al., 2002; Malamud et al., 2004a,b), and have been successfully  
145 examined within a hierarchical framework (Turcotte, 1999; Turcotte et al.,  
146 1999, 2002). Biological hazards, such as spread of *human, animal, and plant*  
147 *epidemics*, are naturally modeled by time-oriented trees where vertices repre-  
148 sent infected subjects. The discussed phenomenology is relevant in other areas  
149 beyond hazard studies where hierarchical organization and related scalings  
150 have been reported. These areas include computer science (Flajolet et al., 1979;  
151 Drmota and Prodinger, 2006), statistical physics of fracture (Zaliapin et al.,  
152 2003; Davidsen et al., 2007; Herrmann and Roux, 2014), vascular analysis  
153 (Kassab, 2000), brain studies (Cassot et al., 2006), ecology (Grant et al., 2007),  
154 scaling of biomass in river streams (Barnes et al., 2007; Gangodagamage et al.,  
155 2007), fractal hydraulic conductivity (Neuman, 1990; Molz et al., 1997), and  
156 allometric scaling laws in biology (West et al., 1997; Turcotte et al., 1998).

157 The multitude of systems traditionally studied via a prism of tree represen-  
158 tation and associated hierarchical dynamics calls for a unifying framework to  
159 address the following questions:

160           *What is a self-similar tree?*

161           *How to model a self-similar tree?*

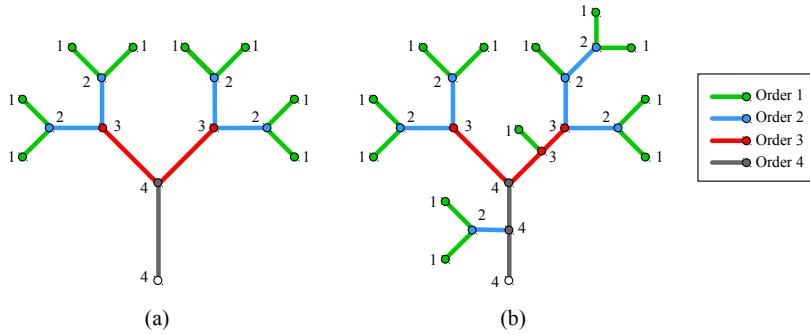
162           *How to test for self-similarity with limited data (in a single tree)?*

163           *What does self-similarity imply for the scalings of observed attributes?*

164 This survey summarizes the currently available answers to these questions.  
165 We show that the key manifestation of self-similarity is *Horton's laws* that  
166 describe scaling of various tree attributes. A geomorphologic origin of the  
167 Horton's laws and the fact that tree representation of river networks is direct  
168 and intuitive affected our choice of examples and terminology. Our main  
169 results however have universal applicability and are formulated in generic  
170 graph-theoretic terms.

## 171   **2 Introduction**

172 In a pioneering study “*of streams and their drainage basins*,” Robert E. Horton  
173 took the first steps toward exploring “*the problems of the development of land*

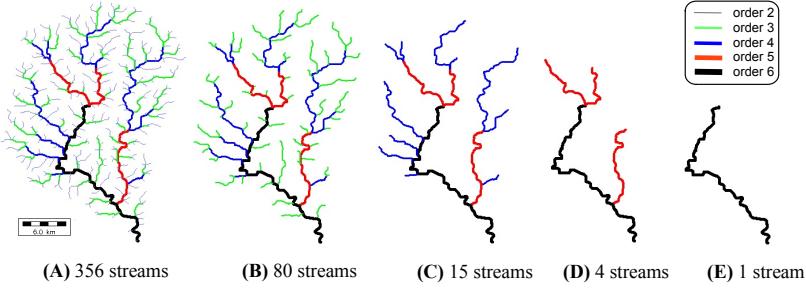


**Fig. 1** Horton-Strahler orders in a binary tree. Different colors correspond to different orders of vertices and edges, as indicated in legend. (a) Perfect binary tree – orders are inversely proportional to vertex/edge depth. (b) General binary tree – orders represent vertex importance in the hierarchy, from leaves (smallest order) to the root (largest order).

174 forms, particularly drainage basins and their stream nets, along quantitative  
 175 lines” (Horton, 1945). Starting from William Playfair’s shrewd observation  
 176 of “a nice adjustment of [stream] declivities” that produces “system of valleys,  
 177 communicating with one another”, Horton revealed deeper regularities in orga-  
 178 nization of river streams. He introduced the concept of stream order and  
 179 formulated two fundamental laws of the composition of stream-drainage nets  
 180 (Horton, 1945, p. 291). *The Law of Stream Numbers* postulates a geometric  
 181 decay of the numbers  $N_K$  of streams of increasing order  $K$ , with the exponent  
 182  $R_B$ ; see Sect. 3.3, Eq. (8). *The Law of Stream Lengths* postulates a geometric  
 183 growth of the average length  $L_K$  of streams of increasing order  $K$ , with the  
 184 exponent  $R_L$ . The initial Horton’s ordering scheme has been later adjusted  
 185 by Arthur Newell Strahler (1957) to its present form (which we call Horton-  
 186 Strahler orders, Sect. 3.2), preserving the laws of stream numbers and lengths.  
 187 Horton-Strahler orders are illustrated in Figs. 1,2.

188 During the 20th century, geometric dependence on the order has been  
 189 documented for multiple physical and combinatorial stream attributes, including  
 190 upstream area, magnitude (number of upstream sources), the total channel  
 191 length, the longest stream length, link slope, mean annual discharge, and  
 192 energy expenditure (Rodriguez-Iturbe and Rinaldo, 2001, and refs. therein).  
 193 A geometric scaling of an arbitrary river stream attribute with order is called  
 194 *Horton’s law* and the respective geometric index is called *Horton exponent*;  
 195 see Sect. 3.3, Eq. (9). The Horton’s laws play an elemental role in studies of  
 196 drainage networks. Being important in their own right, Horton’s laws imply  
 197 power law tails for the empirical frequencies of link attributes (Sect. 3.5)  
 198 and power law relations between different attributes (Sect. 3.4). A celebrated  
 199 example is the upstream contributing area  $A_{(i)}$  of a link  $i$  and the length  $\Lambda_{(i)}$  of  
 200 the longest channel from the link  $i$  to the basin divide. Each of these attributes  
 201 has a power-law empirical frequency,

$$\#\{i : A_{(i)} \geq a\} \propto a^{-\beta_A}, \quad \#\{i : \Lambda_{(i)} \geq l\} \propto l^{-\beta_\Lambda}, \quad (1)$$



**Fig. 2** Stream network of Beaver Creek, Floyd County, KY. (A) Streams (branches) of orders  $K = 2, \dots, 6$  are shown by different colors (see legend on the right). Streams of order  $K = 1$  (source streams) are not shown for visual convenience. Accordingly, this is the first Horton pruning of the network. (B)–(E) Consecutive Horton prunings of the river network; uses the same color code for branch orders as panel (A). The basin has order  $K = 6$  since it is completely eliminated in 6 Horton prunings. The channel extraction is done using RiverTools software (<http://rivix.com>).

202 where  $i$  spans a large collection of links. The expression  $a \propto b$  means that  $a$  is  
 203 proportional to  $b$ , that is  $a = \text{Const.} \times b$ . Furthermore, the two quantities are  
 204 related via a power-law  $A_{(i)} \propto A_{(i)}^{\mathbf{h}}$ , with  $\mathbf{h} \approx 0.6$ . This relation is known as the  
 205 Hack's law (Hack, 1957); it is often reported for the area and the length of the  
 206 longest stream in a basin (Rigon et al., 1996; Rodriguez-Iturbe and Rinaldo,  
 207 2001).

208 Intriguingly, the key parameters of these and other scaling laws can often be  
 209 expressed via the Horton exponents  $R_B$  and  $R_L$ . For example, De Vries et al.  
 210 (1994) and La Barbera and Rosso (1989) have shown, under some simplifying  
 211 assumptions, that

$$\beta_A = 1 - \frac{\log R_L}{\log R_B} \quad \text{and} \quad \beta_A = \frac{\log R_B}{\log R_L} - 1, \quad (2)$$

212 and the fractal dimension  $\mathbf{d}$  of a large tree is given by

$$\mathbf{d} = \frac{\log R_B}{\log R_L}. \quad (3)$$

213 This yields simple relations among the examined quantities:

$$\beta_A = 1 - \mathbf{h}, \quad \beta_A = \mathbf{d} - 1, \quad \text{and} \quad \mathbf{h} = \frac{1}{\mathbf{d}}. \quad (4)$$

214 Despite their recognized importance, the Horton's laws remain an empirical  
 215 finding and their origin and apparent ubiquity remain unsettled. A first attempt  
 216 at rigorous explanation of the Horton's laws and related scalings was made by  
 217 Ronald L. Shreve (Shreve, 1966), who claimed that "*the statistical nature and*  
*218 remarkable generality of Horton's law of stream numbers suggest the speculation*  
*219 that the law of stream numbers arises from the statistics of a large number*  
*220 of randomly merging channels in somewhat the same fashion that the law of*  
*221 perfect gases arises from the statistics of a large number of randomly colliding*

*gas molecules.*" To substantiate this claim, Shreve examined a "topologically random population of channel networks, defined as a population within which all topologically distinct networks with given number of first-order streams are equally likely". This model is equivalent to the critical binary Galton-Watson process with a given progeny (Burd et al., 2000; Pitman, 2006; Kovchegov and Zaliapin, 2020). Shreve's calculations imply that in this model the Horton's law of stream numbers hold with  $R_B = 4$ . Although not attempted by Shreve, it can be shown (Burd et al., 2000) that the law of stream lengths also holds here with  $R_L = 2$  under assumption of constant or equally distributed edge lengths. This corresponds to

$$\beta_A = 1/2, \quad \beta_A = 1, \quad \mathbf{d} = 2, \quad \text{and} \quad \mathbf{h} = 1/2. \quad (5)$$

Albeit insightful and mathematically tractable, the random topology model deviates from the observations. This has been explicitly noted by De Vries et al. (1994) who examined the observed area scaling and have shown that  $\beta_A \approx 0.45 \neq 0.5$ , and by Peckham (1995) who has shown in a detailed analysis of river networks that  $R_B \approx 4.5 \neq 4$ . This called for developing alternative modeling approaches.

Versatile modeling efforts of the past decades have proven challenging to develop an approach that would be mathematically tractable and flexible enough to fit a range of observations. One end of the modeling spectrum is occupied by conceptual models, such as the Peano fractal basin (Rodriguez-Iturbe and Rinaldo, 2001, Sect. 2.4) that has already appeared in Horton's work under a different name (Horton, 1945, Fig. 25), or Scheidegger's lattice model (Scheidegger, 1967; Takayasu et al., 1988). These models provide an invaluable insight into the origin of the observed scalings; they however lack realistic dendritic patterns and values of scaling exponents because of a stiff geometry. On the other end are simulation approaches that are successful in generating visually appealing networks and closely fitting selected exponents, but can be analytically opaque. The Optimal Channel Network (OCN) model (Rinaldo et al., 1992; Rigon et al., 1993; Balister et al., 2018) is a particularly recognized simulation technique. Following the energy expenditure minimization principle, the model constructs random drainage basins on a planar lattice (or more general graph) and fits a variety of observed scaling laws. We refer to (Rodriguez-Iturbe and Rinaldo, 2001) for a comprehensive discussion of these and other models.

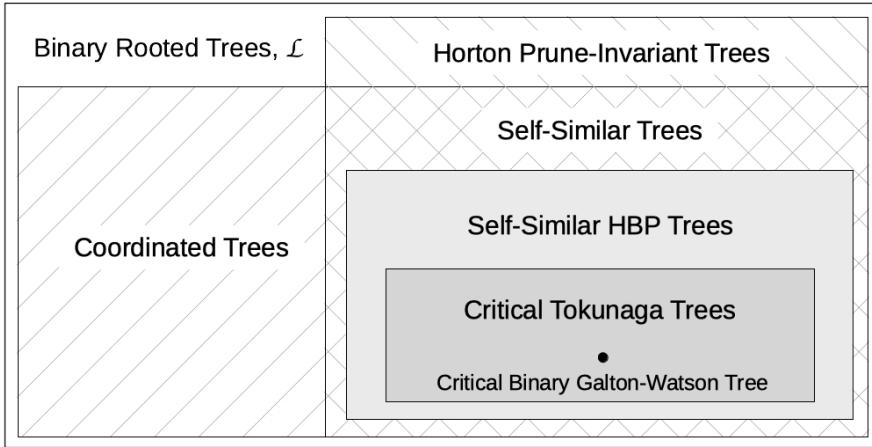
Despite the progress achieved by the modeling efforts of the 20th century, the following essential questions remain unanswered:

*What are sufficient conditions for the Horton's laws?*

*What are the values of the Horton exponents? and*

*How are the Horton exponents for different stream attributes related to each other and to other basin parameters?*

There is a consensus that Horton's laws are connected to the self-similar structure of a basin, which is generally understood as invariance of basin's statistical structure under changing the scale of analysis (zooming in or out).



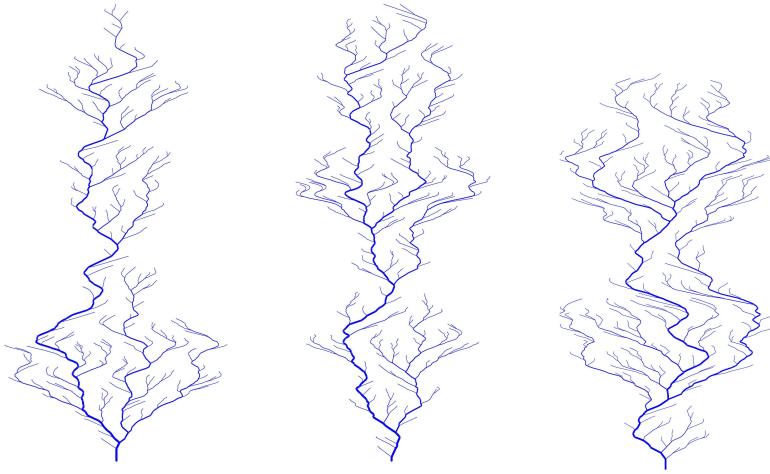
**Fig. 3** Tree spaces examined in this survey: A Venn diagram. We work with binary rooted trees with edge lengths from the space  $\mathcal{L}$  formally defined in Sect. 3.1. The **Self-Similar Trees** belong to the intersection of **Coordinated Trees** (Sect. 4.2) and **Horton Prune-Invariant Trees** (Sect. 4.3). The **Hierarchical Branching Process (HBP)** of Sect. 8 generates a particularly symmetric (infinite-dimensional) family of self-similar tree distributions, which we propose as a default model for applications. A one-parametric family of **Critical Tokunaga Trees** (Sect. 9) is a model proposed for river networks. It includes the celebrated **Critical Binary Galton-Watson Tree** with exponential edge lengths. The combinatorial part of this tree, being conditioned on the tree size, is equivalent to the Shreve's random topology model.

265 However, a consensus is still lacking about a suitable formal definition of tree self-  
 266 similarity. Three alternative definitions have been studied: the Toeplitz property  
 267 of the Tokunaga coefficients (Peckham, 1995; Newman et al., 1997); invariance  
 268 of a distribution with respect to the Horton pruning (cutting source streams)  
 269 in Galton-Watson trees (Burd et al., 2000); and statistical self-similarity of  
 270 basin attributes (Gupta and Waymire, 1989; Peckham and Gupta, 1999). This  
 271 triggers the questions:

272     *How are the alternative definitions of tree self-similarity related? and*  
 273     *Is self-similarity (any version) sufficient for selected Horton's laws?*

274 Answers to these questions require a rigorous toolbox that would go beyond  
 275 the conventional heuristic approaches, which, albeit able to suggest quick routes  
 276 to useful results, may lead to contradictions. For example, the classical works  
 277 of De Vries et al. (1994), La Barbera and Rosso (1989) and many later studies  
 278 adopted the assumption of *an ideal basin* that obeys the exact Horton's laws  
 279 of stream numbers and lengths with  $R_L < R_B$  (Rodriguez-Iturbe and Rinaldo,  
 280 2001, Sect. 2.5). The mean size (the number of links) of a basin of order  $K$  in  
 281 this model is asymptotically given by (Rodriguez-Iturbe and Rinaldo, 2001,  
 282 Eq. (2.91))

$$\frac{R_B^K - R_L^K}{R_B - R_L}. \quad (6)$$



**Fig. 4** Examples of HBP trees (Sect. 8). The trees are generated by the critical Tokunaga process with  $c = 2.3$  and order  $K = 5$  (Sect. 9). The line width is proportional to the contributing area approximated by  $\sum \ell_i^2$ , where the sum is taken over all upstream edges. The figure accurately represents the tree combinatorial structure; the edge lengths are scaled for a better planar embedding. We notice that the HBP generates trees with no planar embedding. The current figure uses an ad-hoc embedding; accordingly, the related purely geometric properties, such as junction angles or spacing between channels, are not a part of the model.

Observe that the mean basin size also equals twice the mean number  $N_1$  of leaves, which is  $2N_1 = 2R_B^{K-1}$ , where the equality holds because of the exact Horton's law for stream numbers. Equating the two expressions for the mean basin size and considering the limit of large  $K$  one obtains  $R_B = 2R_L$ . If in addition one expects the fractal dimension of a basin to be  $d = 2$ , then  $R_B = R_L^2$ , in accordance with (3). The two constraints lead to the unique solution  $R_B = 4$ ,  $R_L = 2$  that corresponds to the Shreve's random topology model (or, more precisely, to the critical binary Galton-Watson model, since the basin size is random). However, the initial assumption of the exact Horton's law does not hold in the Shreve's model. Moreover, the mean tree size in this model scales as  $\frac{4}{3}R_B^{K-1}$ ; see Sect. 9, Eqs. (72), (73). In general, the ideal basin assumption is unrealistic in analysis of river networks as follows from the results of the present paper (see Appendix G). This prompts one to carefully validate the results yielded by heuristic approaches and calls for developing formal techniques. (We must notice that despite the mentioned contradiction, the key heuristic results obtained in the classical works are valid and are reproduced by formal techniques.)

We answer the questions posed above and develop a rigorous toolbox of working with branching structures capitalizing on a self-consistent mathemat-

302 ical theory of random self-similar trees recently developed by the authors  
 303 (Kovchegov and Zaliapin, 2020; Kovchegov et al., 2021). The main goal of the  
 304 current work is to present the theory in relation to the empirical and modeling  
 305 constraints accumulated in the studies of river networks.

306 The theory builds on the self-similarity concepts developed by Horton  
 307 (1945), Strahler (1957), Hack (1957), Shreve (1966, 1969), Tokunaga (1978),  
 308 Mandelbrot (1982), Tarboton et al. (1988), La Barbera and Rosso (1989), Gupta  
 309 and Waymire (1989), Tarboton et al. (1989), Leopold et al. (1992), Rinaldo  
 310 et al. (1992), Rigon et al. (1993), Tarboton (1996), Maritan et al. (1996),  
 311 Turcotte (1997), Gupta and Waymire (1998), and many others. Technically, its  
 312 impetus is provided by the works of Peckham (1995), Newman et al. (1997),  
 313 Turcotte et al. (1998), Peckham and Gupta (1999), Burd et al. (2000), Veitzer  
 314 and Gupta (2000), McConnell and Gupta (2008).

### 315 2.1 How to Use This Survey

316 The survey has a three-fold goal: (i) to outline the key technical tools for  
 317 examining scaling laws in trees; (ii) to present a number of scaling results  
 318 for familiar tree attributes; and (iii) to propose the critical Tokunaga process  
 319 as a model for river networks. To help a reader to promptly find the desired  
 320 material, we briefly (and informally) summarize the main results of the work  
 321 below in Sect. 2.2. Section 2.3 discusses the organization of material throughout  
 322 the paper. Finally, Table 1 provides cross references to relevant equations and  
 323 figures for each of the attribute and scaling exponent examined in the work.

### 324 2.2 The Main Results: A Brief Overview

325 Here we take a short stroll through the main staples of the theory of random  
 326 self-similar trees and give an overview of the results presented in this work,  
 327 before these are expanded in detail in the sections that follow.

328 We work with systems represented by binary tree graphs. The action takes  
 329 place on the space  $\mathcal{L}$  of all such trees, with a root and positive edge lengths. In  
 330 essence, the survey examines a series of consecutively narrower subspaces of  
 331 trees with consecutively stronger symmetries related to scaling laws – Fig. 3  
 332 illustrates the examined hierarchy.

333 The key element of the theory is the operation of *Horton pruning*  $\mathcal{R}$   
 334 that removes (using the hydrological terminology) the source streams from  
 335 a basin. The number of Horton prunings necessary to remove a link from a  
 336 basin defines the link's Horton-Strahler order, hence the pruning name. An  
 337 alternative (equivalent) counting approach to assigning the Horton-Strahler  
 338 orders is described in Sect. 3.2. Figures 1, 2 illustrate Horton-Strahler orders  
 339 and Horton pruning in simple binary trees and in the stream network of Beaver  
 340 Creek, KY.

341 *Self-similarity* is defined as distributional invariance with respect to the  
 342 Horton pruning in trees that satisfy the coordination property. *Coordination*

**Table 1** Tree attributes and scaling exponents examined in the study, with corresponding equations and figures.

Attribute	Description	Eqns.	Figs.
$A_{(i)}$	Contributing area of vertex $i$	(1),(14),(99),(100)	15,16a
$\bar{A}_k$	Empirical average of the contributing areas of order- $k$ branches	(17),(18)	
$A_k$	Mean contributing area of an order- $k$ branch	(37),(38),(59), (76),(77)	5c,10c
$L_{[i]}$	Length of branch $i$	(105)	16b
$\bar{L}_k$	Empirical average of the lengths of order- $k$ branches	(17),(18)	
$L_k$	Mean length of an order- $k$ branch	(51),(60)	5d,9
$L_k^{\text{tot}}$	Mean total channel length upstream of an order- $k$ branch		5c
$\Lambda_{(i)}$	Length of the longest stream to the divide from vertex $i$ (height)	(101),(102),(103)	15
$\bar{\Lambda}_k$	Empirical average of the heights of order- $k$ branches	(17),(18)	
$\Lambda_k$	Mean length of the longest stream to the divide from an order- $k$ branch (height)	(95),(96),(97) (98)	5d
$M_{(i)}$	Magnitude of vertex $i$		9
$\bar{M}_k$	Empirical average of the magnitudes of order- $k$ branches	(17),(18)	
$M_k$	Mean magnitude of an order- $k$ branch	(34),(36),(40),(42), (59),(72),(73),(122)	5b,10b
$N_k[T]$	Number of branches of order $k$ in a tree $T$	(8),(10), (78),(79),(80),(82)	5a,10a
$\mathcal{N}_k[K]$	Mean number of branches of order $k$ in a tree of order $K$	(35),(36),(43),(44), (59),(72),(73),(122), (123)	
$T_{i,j}$	Tokunaga coefficients	(20)	
$T_k$	Tokunaga sequence	(23),(24)	
$\bar{S}_k$	Empirical average of the number of edges (vertices) in order- $k$ branches	(17),(18)	
$S_k$	Mean number of edges (or vertices) in an order- $k$ branch	(33),(50),(60),(118)	5b,10d
Exponent	Description	Eqns.	Figs.
$d$	Fractal dimension of a tree	(3),(4),(5), (85),(87)	11a,12a,14a
$h$	Hack's exponent	(4),(5),(88), (96),(97),(98)	11b,12b,14b, 16
$R_B$	Horton exponent for mean branch numbers $\mathcal{N}_k[K]$	(7),(8),(10),(13), (27),(43),(44),(59), (70),(78),(79),(120), (100),(102)	5a,10a,13a
$R_M$	Horton exponent for mean branch magnitudes $M_k$	(7),(40),(42),(70)	5b,10b
$R_A$	Horton exponent for mean branch contributing areas $A_k$	(7),(12),(48),(70), (100),(102)	5c,10c
$R_L$	Horton exponent for mean branch lengths $L_k$	(7),(12),(51),(60), (70),(85)	5d
$R_S$	Horton exponent for mean combinatorial branch lengths $S_k$	(7),(13),(50),(70), (100),(102)	5b,10d
$R_{\Lambda}$	Horton exponent for mean lengths $\Lambda_k$ of the longest stream to the divide (heights)	(7),(102)	5d,15
$\beta_A$	Exponent of the power law exceedance frequency of branch contributing areas $A_{(i)}$	(1),(2),(4),(5),(14), (99),(100), (102),(103)	16a
$\beta_{\Lambda}$	Exponent of the power law exceedance frequency of lengths of the longest stream to the divide (heights) $\Lambda_{(i)}$	(1),(2),(4),(5),(15), (101),(102),(104)	

means that the (random) structure of a river basin is determined by its order. For example, a basin with outlet of order three and a sub-basin of order three within a basin of order nine have, statistically, the same structure. This assumption is in the heart of analyses based on the Horton-Strahler orders; it has been imposed, explicitly or implicitly, in the mainstream studies of river networks (Horton, 1945; Rodriguez-Iturbe and Rinaldo, 2001; Shreve, 1966; De Vries et al., 1994; Peckham, 1995; Peckham and Gupta, 1999; Tarboton, 1996). A distribution that satisfies the coordination property is called coordinated. The *Horton pruning*  $\mathcal{R}$  is a natural model for the change of resolution in a river network (Fig. 2). Indeed, better observations lead to detecting smaller streams, which increases the basin order. Pruning a basin by order is roughly equivalent to decreasing the resolution of stream detection. The Horton prune-invariance requires that the statistical structure of trees remains the same after zooming in or out (Sect. 4.3, Eq. (22)).

Self-similar distributions are abundant on spaces of rooted trees. Each self-similar distribution corresponds to a unique sequence of non-negative *Tokunaga coefficients*  $T_k$ ,  $k \geq 1$ , equal to the mean number of tributaries of order  $K - k$  within a stream of order  $K$ , for any  $K$  (Sects. 4.1, 4.4). At the same time, an arbitrary sequence of Tokunaga coefficients  $T_k$  corresponds to an infinite number of self-similar distributions (with the same mean numbers of side-tributaries). The well-established models such as Peano basin or Shreve's topologically random model are self-similar in the above sense.

A foundational result (Sect. 7, Thm. 1 and Cor. 1) states that self-similarity implies Horton's law for the mean branch numbers  $N_K$  with exponent  $R_B$  and for the mean branch magnitudes  $M_K$  with exponent  $R_M = R_B$ . Furthermore, a conventional hydrological assumption of equally distributed link lengths (Rodriguez-Iturbe and Rinaldo, 2001; Tarboton et al., 1989) yields Horton's laws for the mean branch contributing areas  $A_K$ , mean number  $S_K$  of links within a stream, and mean total stream length  $L_K$ . The corresponding Horton exponents are uniquely expressed via  $T_k$ . Section 7 and Appendix C examine the Horton's laws for mean branch attributes in the most general situation, with and without the equally distributed link length assumption. The Horton's laws for the mean attributes imply Horton's laws (with the same exponents) for the random attributes obtained by averaging over branches of a given order in a single tree (Sect. 3.7).

Horton's laws imply power-law frequencies of stream attributes and power-law relations between different attributes (Sects. 3.4, 3.5). This includes the power law frequencies of (1) for link contributing areas and the length of the longest channel from a link to the basin divide, as well as for the length of a random stream in a basin. This, in turn, leads to the conventional expressions (3) and (4) that involve the Hack's law and basin fractal dimension. Table 1 lists the attributes and exponents examined in this work, with references to the related equations and figures.

A self-similar Hierarchical Branching Process (HBP) introduced in Sect. 8 generates a particularly symmetric distribution of trees for a given Tokunaga sequence  $T_k$ . The HBP trees obey the strongest forms of Horton's laws for

**Table 2** Selected scaling exponents (1st column) in critical Tokunaga model expressed via the model parameter  $c \geq 1$  (2nd column), fractal dimension  $d$  (3rd column), and Hack's exponent  $h$  (4th column). Columns 5–7 show the values of the exponents for  $c = 2.0, 2.3, 2.5$ . Column 8 shows the values estimated in the OCN model. Columns 9 summarizes estimations in the observed river networks. The agreement of the exponents of the critical Tokunaga model with  $c = 2.3$  with those observed from real basins is noted.

Exponent	Expressed via			Critical Tokunaga model			OCN <sup>†</sup>	Real basins <sup>‡</sup>
	$c$	$d$	$h$	$c = 2.0^*$	$c = 2.3$	$c = 2.5$		
$R_B = R_M = R_A$	$2c$	$2d/(d-1)$	$2^{1/(1-h)}$	4	4.6	5.0	4	4.1 – 4.8
$R_S = R_L$	$c$	$2^{1/(d-1)}$	$2^{h/(1-h)}$	2	2.3	2.5	2	2.1 – 2.7
$d = \frac{\log R_B}{\log R_L}$	$\log_c(2c)$	$d$	$h^{-1}$	2	1.832	1.756	2	1.7 – 2.0
$h = \frac{\log R_L}{\log R_B}$	$\log_{2c} c$	$d^{-1}$	$h$	0.5	0.546	0.569	0.57	0.5 – 0.6
$\beta_A$	$\log_2 2$	$1 - d^{-1}$	$1 - h$	0.5	0.454	0.431	0.43	0.4 – 0.5
$\beta_A$	$\log_c 2$	$d - 1$	$h^{-1} - 1$	1	0.832	0.756	0.8	0.65 – 0.9

\* Equivalent to the critical binary Galton-Watson branching process with i.i.d. exponential edge lengths.

† Mean values estimated in simulated OCN basins. According to Rodriguez-Iturbe and Rinaldo (2001); Cieplak et al. (1998).

‡ According to Rodriguez-Iturbe and Rinaldo (2001); De Vries et al. (1994); Peckham (1995); Maritan et al. (1996); Rigon et al. (1996); Tarboton et al. (1988).

389 multiple stream attributes. A fast and simple recursive simulation algorithm  
 390 allows one to generate networks of realistic size within seconds (Fig. 4). Multiple  
 391 additional symmetries and a well-developed theoretical framework make the  
 392 process an efficient modeling tool.

393 A special subfamily of HBPs, a one-parameter *critical Tokunaga model*, is  
 394 specified by  $T_k = (c - 1)c^{k-1}$  for some  $c \geq 1$  (Sect. 9). This model yields a  
 395 simple relation among the Horton exponents:

$$2c = R_B = R_M = R_A > R_S = R_L = R_A = c, \quad (7)$$

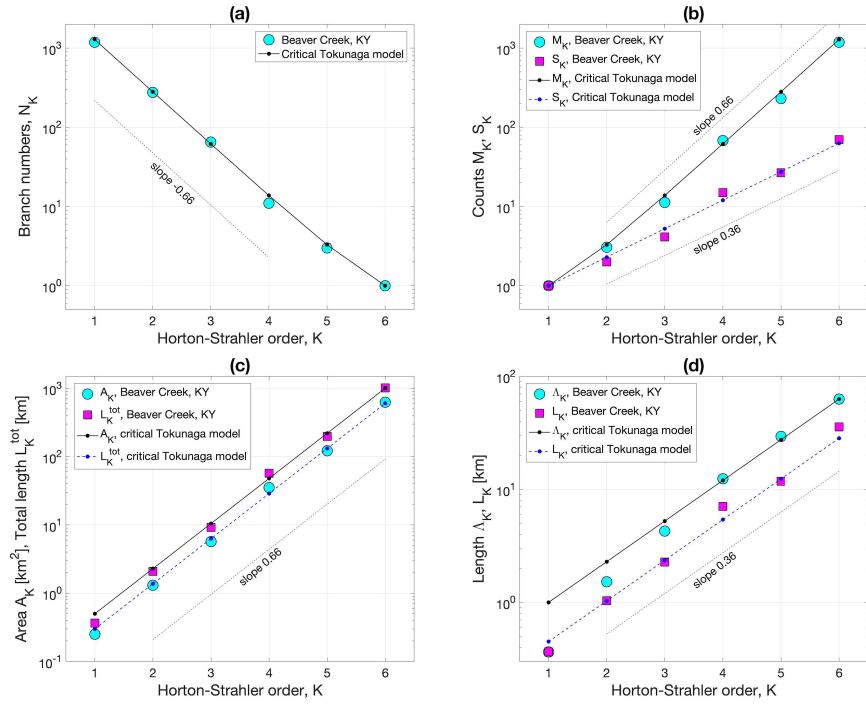
396 where we list, left to right, the Horton exponents for the stream counts ( $R_B$ ),  
 397 magnitude ( $R_M$ ), area ( $R_A$ ), number of links in a stream ( $R_S$ ), stream length  
 398 ( $R_L$ ), and lengths of the longest channel to the divide ( $R_A$ ). The critical  
 399 Tokunaga model provides a close fit to the data and scalings reported in river  
 400 studies over the past decades. We illustrate this in the Beaver Creek basin  
 401 of Fig. 2. Figure 5 shows seven Horton's laws and their respective fits by the  
 402 critical Tokunaga model with  $c = 2.3$ . Specifically, we consider the following  
 403 stream attributes parameterized by stream order  $K = 1, \dots, 6$ : the stream  
 404 numbers  $N_K$  (panel a), the mean magnitude  $M_K$  and the mean number  $S_K$  of  
 405 links in a stream (panel b), the mean contributing area  $A_K$  and the mean total  
 406 channel length  $L_K^{\text{tot}}$  upstream (panel c), the mean length  $\Lambda_K$  of the longest  
 407 channel to the divide and the mean stream length  $L_K$  (panel d). The fitting  
 408 lines correspond to the theoretical model predictions (see figure caption) that  
 409 we derive in Sects. 9, 11. We notice that the model predictions in panels (a) and  
 410 (b) span the entire range of orders, while those in panels (c) and (d) only give  
 411 asymptotic behavior at large orders. This explains some fairly large deviations  
 412 between the data and fitting lines that one can notice at small orders in panel  
 413 (d).

414 Table 2 summarizes the expressions for the Horton exponents and main  
 415 scaling constants in the critical Tokunaga model and compare them with the  
 416 respective quantities estimated in data and the OCN model.

417 It may seem remarkable that a model with a single parameter provides such  
 418 a close fit to the variety of Horton's laws (and other attributes, as can be seen  
 419 from the further discussion). This hints at deep symmetries in the structure  
 420 of trees that describe river networks. The theory of random self-similar trees  
 421 explains the mathematical origin of these symmetries and provides one with  
 422 tools for future exploration.

### 423 2.3 Survey Organization

424 The rest of the paper is organized as follows. Section 3 presents main concepts  
 425 and definitions. Tree representation of a stream network and graph-theoretic  
 426 terminology (Fig. 6) used throughout the paper are introduced in Sect. 3.1.  
 427 Sections 3.2, 3.3 define Horton-Strahler stream orders and related Horton's laws  
 428 (in their simplest form) for stream attributes. The remainder of this section



**Fig. 5** Horton's laws in the stream network of Beaver Creek, Floyd County, KY of Fig. 2. Symbols correspond to the observed attributes. Lines and dots show theoretical fit by a critical Tokunaga model (Sect. 9) with  $c = 2.3$ . (a) Stream numbers  $N_K$ . The model fit is given by (73); it has asymptotic slope  $-\log_{10}(2c) \approx -0.66$ , which is achieved here at small orders. (b) Stream magnitudes  $M_K$  (cyan circles) and number of links  $S_K$  in a stream (magenta squares). The fit for  $M_K$  is given by (73); it has asymptotic slope  $\log_{10}(2c) \approx 0.66$ , which is achieved here at large orders. For this model,  $S_K = c^{K-1}$ , which corresponds to the slope  $\log_{10} c \approx 0.36$ . (c) Contributing areas  $A_K$  (cyan circles) and total upstream channel length  $L_K^{\text{tot}}$  (magenta squares). The fitting lines have theoretical slope  $\log_{10}(2c) \approx 0.66$ ; see (76). The intercept is selected so that the fitting line coincides with an observed quantity at  $K = 6$ . (d) Lengths  $\Lambda_K$  of the longest stream to the divide (cyan circles) and lengths  $L_K$  of streams (magenta squares). The fitting lines have theoretical slope  $\log_{10}(c) \approx 0.36$ ; see (60) and (95).

429 discusses essential heuristic implications of the Horton's laws for power-law  
430 frequencies of and relations among stream attributes.

431 Section 4 introduces the key technical tools of self-similarity analysis –  
432 Tokunaga coefficients, coordination of tree measures, Horton pruning, and  
433 Horton prune-invariance.

434 Section 5 recalls basic facts from the theory of generating functions and  
435 complex analysis that are used to establish our main results. Here we formalize  
436 the notion of Horton's laws for mean stream attributes by considering three  
437 consecutively stronger versions – (R), (Q), and (G) – of geometric variation.  
438 Later we refer to those as the *root* Horton's law, *quotient* Horton's law, and  
439 *geometric* Horton's law, respectively.

440 The key stream attributes and their relations are presented in Sect. 6. This  
 441 includes the mean number  $S_K$  of edges in a branch of order  $K$ , the mean  
 442 magnitude  $M_K$  of a branch of order  $K$ , the mean number  $\mathcal{N}_k[K]$  of branches  
 443 of order  $k$  in a tree of order  $K$ , the mean length  $L_K$  of a branch of order  $K$ ,  
 444 and the mean contributing area  $A_K$  of a branch of order  $K$ .

445 Section 7 presents the main results of this work – Horton’s laws for mean  
 446 branch attributes in a self-similar tree. We start in Sect. 7.1 with the geometric  
 447 Horton’s laws for the mean branch magnitudes  $M_k$  (Theorem 1) and mean  
 448 branch numbers  $\mathcal{N}_k[K]$  (Corollary 1). These laws hold in any self-similar tree  
 449 under a mild constraint  $\limsup_{k \rightarrow \infty} T_k^{1/k} < \infty$ ; they form a foundation for  
 450 further development. Section 7.2 establishes quotient Horton’s laws for  $S_k$ ,  $L_k$ ,  
 451 and  $A_k$  under additional Assumption 1, which seems practically appealing to  
 452 most applications. Section 7.3 discusses a special case of Horton’s laws under  
 453 the hydrologically relevant constraint of unit Horton exponent for edge lengths,  
 454 which is a substantial generalization of the condition of equally distributed edge  
 455 lengths that is well-documented in hydrologic observations. Section 8 introduces  
 456 a self-similar Hierarchical Branching Process (HBP) – a computationally simple  
 457 and analytically tractable model that generates trees with arbitrary Tokunaga  
 458 sequences and obeys the (strongest) geometric Horton’s laws for the mean  
 459 branch counts, magnitudes, and lengths. An important one-parametric family  
 460 of self-similar HBP – critical Tokunaga process with  $T_k = (c - 1)c^{k-1}$  for  $c \geq 1$   
 461 – is examined in Sect. 9. We propose this as a conceptual model for river stream  
 462 networks, as it provides a very close fit to the attributes and scalings reported  
 463 in observations (see Table 2 and Fig. 5). A classical model with slightly relaxed  
 464 constraints on the Tokunaga sequence,  $T_k = ac^{k-1}$ , is discussed in Appendix E.

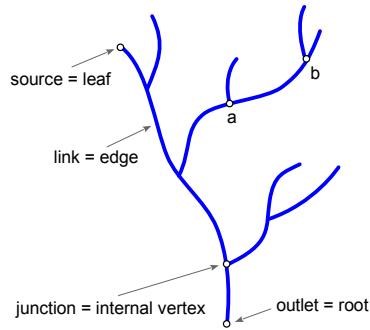
465 Fractal dimension and Hack’s law in self-similar HBP trees are examined in  
 466 Sects. 10, 11, respectively. We show in particular that the fractal dimension  $\mathbf{d}$   
 467 and Hack’s exponent  $\mathbf{h}$  are reciprocal to each other, which has been heuristically  
 468 known for trees with exact Horton’s laws since the 1980s (La Barbera and  
 469 Rosso, 1989; De Vries et al., 1994; Peckham, 1995).

470 Section 12 illustrates the origin of power-law exceedance frequencies of edge  
 471 and branch attributes in a tree that satisfies selected Horton’s laws. Specifically,  
 472 we consider the vertex contributing area  $A_{(i)}$ , the length  $L_{(i)}$  of the longest  
 473 stream from a vertex to the divide, and the length  $L_{[i]}$  of a randomly selected  
 474 branch. Section 13 provides concluding remarks. Proofs and most general  
 475 results, which may not be of prime interest in applied analyses, are given in  
 476 Appendices.

### 477 3 Horton’s Laws and Their Implications: A Heuristic Approach

#### 478 3.1 Tree Representation of River Networks

479 River studies commonly represent a stream network that drains a single basin  
 480 (watershed, catchment) as a rooted binary tree with planar embedding. The  
 481 basin *outlet* (point furthest downstream) corresponds to the tree root, *sources*



**Fig. 6** Hydrologic vs. graph theoretic terminology. The figure shows, clockwise from the bottom, the basin outlet (tree root), a stream junction (internal vertex), link (edge), and a source (leaf). Vertex  $b$  is an offspring of  $a$ ; and  $a$  is the parent for  $b$ .

482 (points furthest upstream) to leaves, *junctions* (points where two streams meet)  
 483 to internal vertices, and *links* (stream segments between two successive nodes)  
 484 to edges. A node  $j$  immediately upstream of a node  $i$  is called an *offspring* of  $i$ ,  
 485 and  $i$  is called the *parent* of  $j$ . Any node  $j$  upstream of  $i$  is called a *descendant* of  
 486  $i$ , and any node  $j$  downstream of  $i$  is called an *ancestor* of  $i$ . Figure 6 illustrates  
 487 this correspondence.

488 In this work we use the graph-theoretic nomenclature, which provides a  
 489 better link to the other systems examined using their tree representation (see  
 490 Sect. 1). We assume that all examined trees belong to the space  $\mathcal{T}$  of finite  
 491 binary rooted trees, or to the space  $\mathcal{L}$  of trees from  $\mathcal{T}$  with positive edge lengths  
 492 (Kovchegov and Zaliapin, 2020, Sect. 2.1). The models discussed in this work  
 493 do not deal with planar embedding of trees, which is an important separate  
 494 problem. We notice at the same time, that a suitable (non-physical) planar  
 495 embedding of a given tree can be readily developed. One such embedding  
 496 approach is used to illustrate synthetic HBP trees in Fig. 4.

### 497 3.2 Horton-Strahler Orders

498 The importance of vertices and their parental edges is measured by the Horton-  
 499 Strahler order  $K \geq 1$  (Horton, 1945; Strahler, 1957). We agree that each vertex  
 500 and its parental edge (the unique edge that connects this vertex to its parent,  
 501 or the immediate downstream edge) have the same order. The order assignment  
 502 is done in a hierarchical fashion, from the leaves towards the root (that is,  
 503 from the sources downstream). Specifically, each leaf (and its parental edge)  
 504 is assigned order  $K = 1$ . When two edges of the same order  $K$  merge at a  
 505 vertex, the vertex is assigned order  $K + 1$ . When two edges with different  
 506 orders  $K_1 > K_2$  merge at a vertex, the largest order prevails and the vertex  
 507 is assigned order  $K_1$ . The connected sequence of vertices and their parental  
 508 edges of the same order  $K$  is called a *branch* of order  $K$ . The Horton-Strahler

orders are illustrated in Figs. 1, 2. We denote by  $N_K = N_K[T]$  the number of branches of order  $K$  in a finite tree  $T$ .

The Horton-Strahler order of a tree is that of its root, or equivalently, the maximal order of its vertices (edges, branches). We show below that multiple fundamental regularities in the structure and dynamics of river networks are expressed in terms of the Horton-Strahler orders.

### 3.3 Horton's Laws

The observed stream counts  $N_K$  in a large basin are closely approximated by (Horton, 1945)

$$\frac{N_K}{N_{K+1}} = R_B \Leftrightarrow N_K \propto R_B^{-K} \quad (8)$$

for some *Horton exponent*  $R_B \geq 2$ . The lower bound on  $R_B$  follows immediately from the definition of Horton-Strahler orders, since it takes at least two branches of order  $K$  to create a single stream of order  $K + 1$ . It has been noticed by (Strahler, 1957, p. 914) that the value of the empirical ratio  $R_B$  in river streams is between 3 and 5, and is usually close to 4. This has been strongly corroborated in numerous observational studies, e.g. Kirchner (1993), Shreve (1966), Leopold et al. (1992), Peckham (1995), Tarboton (1996), Turcotte (1997), Gupta and Waymire (1998), Zanardo et al. (2013), Rodriguez-Iturbe and Rinaldo (2001), and Mesa (2018).

In hydrogeomorphology, a geometric scaling of any branch attribute with order, similar to that of Eq.(8), is called *Horton's law*. Horton's laws are documented for multiple physical and combinatorial quantities, including upstream area, magnitude (number of upstream sources), the total channel lengths, link slope, mean annual discharge, energy expenditure, etc. (Rodriguez-Iturbe and Rinaldo, 2001). These quantities often increase with order (unlike the branch counts  $N_K$  that decrease with order), which justifies a slightly different form of the respective Horton's laws. Specifically, consider the values  $Z_K$  obtained by averaging a selected attribute  $Z$  over branches of order  $K$ . Horton's law with exponent  $R_Z \geq 1$  states that  $Z_K$  scale as

$$\frac{Z_{K+1}}{Z_K} = R_Z \Leftrightarrow Z_K \propto R_Z^K. \quad (9)$$

In both cases (8) and (9) the law is formulated in such a way that the Horton exponent is greater than unity. Informally,  $Z_K$  may represent a particular way to measure the branch "size", and the law (9) states that the order  $K$  of a branch is proportional to its logarithmic size  $\ln(Z_K)$ .

Horton's laws play an elemental role in statistical modeling of river basins, which rests upon empirical regularities that describe the frequencies of and relations among the key geometric and physical characteristics of individual streams. Remarkably, many such regularities heuristically follow from Horton's laws and are parameterized by the respective Horton exponents. Below we discuss several key power laws that are commonly observed in river networks.

547 3.4 Power-Law Relations Between Attributes

548 Suppose a stream attribute  $Z$  satisfies the Horton's law (9) with exponent  $R_Z$ ,  
 549 and the branch counts  $N_K$  satisfy the Horton's law (8) with exponent  $R_B$ .  
 550 Then, using each of the laws to express  $K$  and equating these expressions, we  
 551 find

$$Z_K \propto N_K^{-\alpha}, \quad \text{with} \quad \alpha = \frac{\ln R_Z}{\ln R_B}. \quad (10)$$

552 Similarly, suppose that the Horton's law (9) holds for selected river at-  
 553 tributes  $Z$  and  $Y$ , with exponents  $R_Z$  and  $R_Y$ , respectively. Then,  $Z_K$  and  $Y_K$   
 554 are connected via a power-law relation

$$Z_K \propto Y_K^\alpha, \quad \alpha = \frac{\log R_Z}{\log R_Y}. \quad (11)$$

555 Equations (10), (11) are a punctuated (by discrete orders) version of a  
 556 general power-law relation  $Z \propto Y^\alpha$  that is abound among hydrologic quantities.  
 557

558 It is common to relate an attribute of interest to the basin area  $A$ . A well  
 559 studied example is Hack's law that relates the length  $L$  of the longest stream  
 560 in a basin to the basin area  $A$  via  $L \propto A^h$  with  $h \approx 0.6$  (Hack, 1957; Rigon  
 561 et al., 1996; Rodriguez-Iturbe and Rinaldo, 2001). Assuming Horton's laws for  
 562 the area and length of the longest stream, the parameter  $h$  is expressed via  
 563 the respective Horton exponents as in Eq. (11):

$$h = \frac{\log R_L}{\log R_A}. \quad (12)$$

564 3.5 Power-Law Frequencies of Link Attributes

565 Consider empirical frequencies of an attribute  $Z$  calculated at every edge (link)  
 566 in a large tree (basin). We write  $Z_{(i)}$  for the value of  $Z$  calculated at the  $i$ -th  
 567 edge. Assume that Horton's law holds (i) for the examined attribute  $Z$ , with  
 568 exponent  $R_Z$  and (for simplicity) proportionality constant equal to one; (ii) for  
 569 the average number  $S_K$  of edges within a branch of order  $K$ , with exponent  
 570  $R_S$ ; and (iii) for branch counts  $N_K$  with exponent  $R_B$  as in (8). The number of  
 571 edges of order  $K$  in such a tree is given by  $N_K S_K$ . One can now heuristically  
 572 approximate the expected frequencies of  $Z_{(i)}$  by using the same value  $Z_K$  for  
 573 any edge of order  $K$  and considering a limit of an infinitely large tree:

$$\#\{i : Z_{(i)} \geq R_Z^K\} \approx \sum_{j=K}^{\infty} N_j S_j \propto \sum_{j=K}^{\infty} \left(\frac{R_S}{R_B}\right)^j \propto \left(\frac{R_S}{R_B}\right)^K.$$

574 As before, this is a punctuated (by discrete order) version of a general power  
 575 law relation

$$\#\{i : Z_{(i)} \geq z\} \propto z^{-\beta}, \quad \beta = \frac{\log R_B - \log R_S}{\log R_Z}. \quad (13)$$

576 Such power laws are reported for the upstream contributing area, stream  
 577 lengths to the divide, water discharge, or energy expenditure. For example,  
 578 analyses of Tarboton et al. (1989), Rodriguez-Iturbe et al. (1992), and Maritan  
 579 et al. (1996) on river basins extracted from digital elevation models (DEM's)  
 580 suggest

$$\#\{i : A_{(i)} \geq x\} \propto x^{-\beta_A} \quad \text{with } \beta_A \approx 0.45, \quad (14)$$

$$\#\{i : A_{(i)} \geq x\} \propto x^{-\beta_A} \quad \text{with } \beta_A \approx 0.8, \quad (15)$$

581 where  $A_{(i)}$  is the area upstream of link  $i$  and  $A_{(i)}$  is the distance from link  $i$   
 582 to the furthest source (or, equivalently, to the basin divide) measured along  
 583 the channel network. Our analysis below shows that these power laws hold in  
 584 self-similar trees and their exponents found via (13) fit the empirical exponents  
 585 found in observations.

### 586 3.6 Modeling Physical Characteristics of a Stream

587 Classical hydrologic and geomorphologic studies of the mid 20th century  
 588 revealed that the key physical characteristics of streams – such as stream  
 589 width, depth, slope, and flow velocity, can be modeled as power functions of  
 590 the stream magnitude (Leopold and Miller, 1956; Leopold et al., 1992; Dodov  
 591 and Foufoula-Georgiou, 2004a,b, 2005),(Rodriguez-Iturbe and Rinaldo, 2001,  
 592 Chapter 1). Specifically, data analysis suggests that physical characteristics  
 593 of streams scale with the stream discharge  $Q$  defined as the volume of water  
 594 flowing through a river stream. Such scalings are called *hydraulic-geometric*  
 595 *relations*. For example, the velocity  $v$  through the stream can be approximated  
 596 by  $v \propto Q^\alpha$ , etc. The discharge  $Q$ , in turn, is a power-law function of the basin  
 597 area:  $Q \propto A^\beta$  (see Sect. 3.4). The value of exponent  $\beta$  depends on a precise  
 598 definition of discharge (bankful, mean annual, etc.) Finally, the basin area  $A$  is  
 599 closely approximated by the basin magnitude  $M$ , since it is natural to think of  
 600 a stream network as a space-filling tree (see Sect. 7 for a formal treatment).  
 601 Combining these observations, we find that the stream velocity can be modeled  
 602 as a power-law function of the stream magnitude:

$$v \propto M^{\alpha\beta}.$$

603 Hydraulic-geometric relations exist for other physical characteristics of a stream,  
 604 including the average link slope  $s$  (e.g., Gupta and Waymire, 1989):

$$s \propto A^{-\theta} \propto M^{-\theta}, \quad \theta \approx 0.5.$$

In summary, rather unexpectedly, essential physical characteristics of a river network (e.g., stream velocity or link slope) can be estimated from purely combinatorial statistics of its tree representation (e.g., magnitude). Gupta (2017) asserts that “*Self-similarity in channel networks plays a foundational role in understanding the observed scaling, or power law relations, between peak flows and drainage areas*”. For example, the emergent scaling behavior opens up the opportunity to circumvent a large number of parameters governing production and transport of runoff along the stream channels and use basin’s combinatorial characteristics for developing flood frequency relations and flood forecasting in ungauged basins (Gupta et al., 1994, 1996, 2010, 2007; Gupta, 2017). Accordingly, the results presented in this survey can inform modeling efforts aimed at physical quantities of the streams and the related processes and hazards. Gupta and Mesa (2014) discussed an alternative approach for establishing Horton’s laws for river physical attributes (hydraulic-geometric variables) based on the Buckingham  $\pi$  theorem and asymptotic self-similarity of first and second kinds (Barenblatt, 1996).

### 621 3.7 Beyond Heuristics

We observe that the above discussion in Sects. 3.3, 3.4, and 3.5 is heuristic, only maintaining a physical (but not mathematical) level of rigor. The very definition of Horton’s law via Eqs.(8) and (9) is not instrumental for developing a useful theory. Indeed, since it is hard to expect that the exact equalities would hold in a range of practically interesting situations, one should accept an approximate nature of these statements and, hence, define what is meant by “approximate”.

The approach adopted in this survey (and in most of the studies reviewed herewith) asserts that, for any fixed  $k$ , the branch number ratios  $N_k/N_{k+1}$  converge to the Horton’s exponent  $R_B$  when the tree size increases. Similarly, the other branch attribute ratios  $Z_{k+1}/Z_k$  converge to the appropriate Horton exponents when both  $k$  and the tree size increase. These convergences involve random variables  $N_k$  and  $Z_k$  and hence should be understood in a proper probabilistic sense (Bhattacharya and Waymire, 2007). The difference between the treatment of the branch numbers  $N_k$  and other branch attributes  $Z_k$  is explained by the observation that  $N_k$  decreases with  $k$ , while all other branch attributes (examined here) increase with  $k$ .

As an intermediate step, we consider the *mean values* of the examined attributes with respect to the examined distribution of trees. Such a mean value should not be confused with the *average value* that is calculated over a collection of branches within a single random tree. For example, we consider below a random length  $L_{[i]}$  of the  $i$ th branch. One can average the observed random lengths over the  $N_k$  branches of a given order  $k$  in a single tree  $T$  to

645 obtain the empirical *average* of the branch lengths:

$$\bar{L}_k = \frac{1}{N_k} \sum_{i=1}^{N_k} L_{[i]}.$$

646 Importantly,  $\bar{L}_k$  is a random variable that takes on a new value for each  
647 realization of a random tree. Finally, we consider the *mean branch length*

$$L_k = \mathbb{E}[L_{[i]}] \quad \text{for any } i \text{ because of coordination,}$$

648 which is a constant that only depends on the examined tree distribution.

649 Most of our results are formulated for the *mean branch attributes*. Import-  
650 antly, the respective results for the random attributes readily follow from  
651 these mean results. To illustrate this implication, let  $Z_{[i]}$  denote a random  
652 value of the examined attribute calculated for branch  $i$  in a random tree  $T$ ,  
653  $\bar{Z}_k$  denote the empirical average of the attribute over the branches of order  $k$   
654 in  $T$ , and  $Z_k$  denote the mean value of the attribute for a random branch of  
655 order  $k$ . Horton's law for the deterministic mean attribute is defined as a limit  
656 statement:

$$\lim_{k \rightarrow \infty} \frac{Z_{k+1}}{Z_k} = R_Z. \quad (16)$$

657 Section 5 discusses this and two other (weaker and stronger) forms of Horton's  
658 law that can hold in deterministic sequences of branch attributes. We have

$$\mathbb{E}[\bar{Z}_k] = \mathbb{E}[\mathbb{E}[\bar{Z}_k | N_k]] = \mathbb{E}[Z_{[i]}] = Z_k.$$

659 In a tree  $T$  of order  $K$  the branch numbers are bounded from below by  
660  $N_k \geq 2^{K-k}$  for any  $k \leq K$ , so  $K \rightarrow \infty$  implies  $N_k \rightarrow \infty$  for any fixed  $k$  with  
661 probability one. Hence, the Weak Law of Large Numbers (Bhattacharya and  
662 Waymire, 2007) asserts that

$$\bar{Z}_k \xrightarrow{p} Z_k \quad \text{for any } k \text{ as } K \rightarrow \infty, \quad (17)$$

663 where  $\xrightarrow{p}$  denote convergence in probability. Accordingly,

$$\frac{\bar{Z}_{k+1}}{\bar{Z}_k} \xrightarrow{p} \frac{Z_{k+1}}{Z_k} \quad \text{for any } k \text{ as } K \rightarrow \infty, \quad (18)$$

664 where the (deterministic) fraction in the right-hand side converges to  $R_Z$  as in  
665 (16). Practically, statements (16) and (18) suggest that the empirical averages  
666  $\bar{Z}_k$  satisfy Horton's approximation of (9) in a sufficiently large tree (or a finite  
667 collection of such trees).

668 The above discussion applies to the branch magnitudes  $M_k$ , combinatorial  
669 and metric branch lengths  $S_k$  and  $L_k$ , branch contributing areas  $A_k$ , and the  
670 length of the longest stream to the divide (height)  $\Lambda_k$ . At the same time,  
671 the probabilistic limit results for the random branch counts  $N_k$  require more  
672 sophisticated techniques that are outside of the scope of this survey. Section 9.2  
673 reviews limit laws for the random branch numbers in the critical Tokunaga  
674 model of Sect. 9.

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**675 4 Self-Similarity of River Networks**
**676 4.1 Tokunaga Coefficients**

677 The Tokunaga coefficients complement the branch counts  $N_K[T]$  in describing  
 678 the structure of a tree  $T$ . The empirical Tokunaga coefficient  $t_{i,j}[T]$  with  $i < j$   
 679 is the average number of branches of order  $i$  that merge with a branch of order  
 680  $j$  in a finite tree  $T$ :

$$t_{i,j}[T] = \frac{N_{i,j}[T]}{N_j[T]}, \quad (19)$$

681 where  $N_{i,j}[T]$  is the number of instances when an order- $i$  branch merges with  
 682 an order- $j$  branch within  $T$ . The merging of branches of distinct orders is  
 683 referred to as *side-branching*, and a branch that merges into a branch of a  
 684 higher order is called a *side-branch*. Merging of two branches of the same order  
 685 is called *principal branching*.

686 Assume that we fix a distribution  $\mu$  on the space  $\mathcal{T}$  of finite rooted binary  
 687 trees. For example, one might consider a uniform distribution among trees  
 688 with a given number of leaves, leading to the critical binary Galton-Watson  
 689 random tree (Burd et al., 2000; Pitman, 2006). Then one can define the  
 690 Tokunaga coefficient  $T_{i,j}$  as the expected number of side-branches of order  $i$  per  
 691 a randomly selected branch of order  $j$  (Dodds and Rothman, 1999; Tokunaga,  
 692 1966, 1978; Burd et al., 2000). This definition serves well the purpose of our  
 693 study; we refer to Kovchegov and Zaliapin (2020) for a more general approach.

694 We can arrange the Tokunaga coefficients for trees of a given order  $K$  in  
 695 an upper triangular matrix

$$\mathbb{T}_K = \begin{bmatrix} 0 & T_{1,2} & T_{1,3} & \dots & T_{1,K} \\ 0 & 0 & T_{2,3} & \dots & T_{2,K} \\ 0 & 0 & \ddots & \ddots & \vdots \\ \vdots & \vdots & \ddots & 0 & T_{K-1,K} \\ 0 & 0 & \dots & 0 & 0 \end{bmatrix}. \quad (20)$$

696 For example, the Tokunaga coefficients calculated for the Beaver Creek basin  
 697 illustrated in Figs. 2, 5 are given by

$$\mathbb{T}_6 = \begin{bmatrix} 0 & 1.06 & 2.40 & 8.91 & 15.33 & 44.00 \\ 0 & 0 & 0.92 & 3.64 & 8.67 & 20.00 \\ 0 & 0 & 0 & 2.00 & 4.00 & 9.00 \\ 0 & 0 & 0 & 0 & 0.33 & 4.00 \\ 0 & 0 & 0 & 0 & 0 & 1.00 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}. \quad (21)$$

**698 4.2 Coordination**

699 We assume that the structure of a river basin is determined by its order. This  
 700 means, for example, that a basin with outlet of order three and a sub-basin

of order three within a larger basin of order nine have, statistically, the same structure. This assumption is in the heart of the Horton-Strahler orders, and is imposed, explicitly or implicitly, in the mainstream studies of river networks (Shreve, 1966, 1969; Peckham, 1995; Rodriguez-Iturbe and Rinaldo, 2001). We refer to this assumption as *coordination*. Under the assumption of coordination the Tokunaga matrix  $\mathbb{T}_K$  of Eq. (20) coincides with the upper-left  $K \times K$  submatrix of the Tokunaga matrix  $\mathbb{T}_M$  for any  $M \geq K$ , which explains the assumption name.

### 4.3 Tree Self-Similarity

Most generally, *self-similarity* is understood as statistical invariance of a river basin under rescaling Mandelbrot (1982); Turcotte (1997); Dodds and Rothman (2000); Rodriguez-Iturbe and Rinaldo (2001). A fundamental specific way to downscale a river basin of order  $K$  is to only consider its branches with highest orders between  $K - k + 1$  and  $K$  for a given  $k < K$ . This results in a coarser basin, whose order (being computed according to the rules of Sect. 3.2) is  $k$ .

Formally, we consider the operation of *Horton pruning*  $\mathcal{R} : \mathcal{L} \rightarrow \mathcal{L}$  that removes the leaves from a tree  $T$  together with their parental edges, followed by a *series reduction* that eliminates all degree two non-root vertices by merging the edges adjacent to them. It is readily seen that the Horton pruning reduces the tree order by 1. Moreover, the order of each branch is also reduced by 1 (with understanding that branches of order 1 are eliminated). We refer to Peckham (1995), Burd et al. (2000), and Kovchegov and Zaliapin (2016) for a comprehensive discussion.

A coordinated distribution  $\mu$  on the space  $\mathcal{T}$  of combinatorial trees is called *self-similar* if it is invariant with respect to the Horton pruning (Burd et al., 2000; Kovchegov and Zaliapin, 2016):

$$\mu(\mathcal{R}^{-1}(T)|T \neq \phi) = \mu(T) \quad \text{for any } T \in \mathcal{T}. \quad (22)$$

Informally, consider a forest of trees, where each tree  $T$  occurs multiple times according to its probability  $\mu(T)$ . The forest is self-similar if after pruning each tree by  $\mathcal{R}$  we obtain the same forest. This definition can be extended to trees with edge lengths from space  $\mathcal{L}$ . In this case, we allow the edge lengths to scale by a *scaling constant*  $\zeta > 0$  after pruning. We refer to (Kovchegov and Zaliapin, 2020, Sect. 3) for a formal treatment.

We use a conventional abuse of terminology by saying that a tree  $T$  is self-similar; this means that  $T$  is a random tree drawn from a self-similar distribution  $\mu$ .

736 4.4 Tokunaga Sequence

737 A coordinated self-similar measure necessarily satisfies the following Toeplitz  
 738 property (Kovchegov and Zaliapin, 2016, 2020): there exists a *Tokunaga se-*  
 739 *quence*  $\{T_k\}_{k=1,2,\dots}$  such that

$$T_{i,i+k} = T_k \quad \text{for all } i, k > 0. \quad (23)$$

740 In this case, the Tokunaga matrix (20) is a Toeplitz matrix:

$$\mathbb{T}_K = \begin{bmatrix} 0 & T_1 & T_2 & \dots & T_{K-1} \\ 0 & 0 & T_1 & \dots & T_{K-2} \\ 0 & 0 & \ddots & \ddots & \vdots \\ \vdots & \vdots & \ddots & 0 & T_1 \\ 0 & 0 & \dots & 0 & 0 \end{bmatrix}. \quad (24)$$

741 Sometimes we refer to  $T_k$  as *Tokunaga indices*, which creates no confusion with  
 742  $T_{i,j}$  of Sect. 4.1. A comprehensive discussion of this approach can be found  
 743 in the works of Kovchegov and Zaliapin (2016) and Burd et al. (2000). The  
 744 Tokunaga indices  $T_k$  for the Beaver Creek of Figs. 2, 5 can be approximated by  
 745 averaging the values along the diagonals of the initial Tokunaga matrix (21):

$$T_1 = 1.06, \quad T_2 = 3.51, \quad T_3 = 8.86, \quad T_4 = 17.665, \quad T_5 = 44.00. \quad (25)$$

746 We emphasize that self-similarity (Sect. 4.3) is a property of a distribution  
 747 of trees on  $\mathcal{T}$  or  $\mathcal{L}$  and hence, formally, cannot be applied to a single tree. In  
 748 applied analysis, however, one works with a single basin, or a finite sample of  
 749 basins. The Tokunaga sequence  $T_k$  provides a fundamental connection between  
 750 the properties of a distribution on an infinite collection of trees and easily  
 751 computed attributes of a single tree.

752 We show below that the Tokunaga indices  $T_k$  provide enough information  
 753 to find the mean values of all other branching attributes in random self-similar  
 754 trees.

755 4.5 Tokunaga Two-Parameter Model

756 The first model for river networks that explicitly describes the network structure  
 757 in terms of side-branch counts is due to Eiji Tokunaga (1978). It postulates

$$T_k = ac^{k-1}, \quad a, c > 0. \quad (26)$$

758 The observed river networks are closely approximated by the *Tokunaga*  
 759 *model* (Tokunaga, 1978; Peckham, 1995; Kovchegov et al., 2021). The estimated  
 760 parameters  $a \approx 1.1$  and  $c \approx 2.5$  of this model have shown to be independent  
 761 of (or only weakly dependent on) river's geographic location (Peckham, 1995;  
 762 Dodds and Rothman, 2000; Zanardo et al., 2013).

<sup>763</sup> McConnell and Gupta (2008) have shown that the Tokunaga model obeys  
<sup>764</sup> the quotient Horton's law for stream numbers (when tree size increases) with

$$R_B = \frac{(a + c + 2) + \sqrt{(a + c + 2)^2 - 8c}}{2}. \quad (27)$$

<sup>765</sup> This result revealed, for the first time, the emergence of Horton's law from  
<sup>766</sup> the tree side-branch structure. Our Theorem 1 below establishes the most  
<sup>767</sup> general statement of this type, showing that (almost) *any* Tokunaga sequence  
<sup>768</sup> implies the geometric Horton's laws for branch numbers and mean magnitudes.  
<sup>769</sup> A detailed treatment of the two-parameter Tokunaga model (26) is given in  
<sup>770</sup> Appendix E.

<sup>771</sup> Burd et al. (2000) demonstrated that the Shreve's random topology model  
<sup>772</sup> (Shreve, 1966, 1969), equivalent to the critical binary Galton-Watson tree with  
<sup>773</sup> a fixed progeny, is a special case of the Tokunaga model with  $(a, c) = (1, 2)$ :

$$T_k = 2^{k-1} \quad \text{for } k \geq 1.$$

<sup>774</sup> Accordingly, it satisfies the geometric Horton's law for mean branch numbers  
<sup>775</sup> with  $R_B = 4$ .

<sup>776</sup>

<sup>777</sup> For a long time, the critical binary Galton-Watson tree has remained the  
<sup>778</sup> only well-studied probability model for which self-similarity was rigorously  
<sup>779</sup> established, and whose Horton-Strahler ordering has received attention in  
<sup>780</sup> the literature (Shreve, 1966, 1969; Kemp, 1979; Tarboton et al., 1988; Wang  
<sup>781</sup> and Waymire, 1991; Barndorff-Nielsen, 1993; Yekutieli and Mandelbrot, 1994;  
<sup>782</sup> Peckham, 1995; Devroye and Kruszewski, 1994; Burd et al., 2000). Scott  
<sup>783</sup> Peckham has explicitly noticed, by performing a high-precision extraction of  
<sup>784</sup> river channels for Kentucky River, Kentucky and Powder River, Wyoming,  
<sup>785</sup> that the Horton exponents and Tokunaga parameters for the observed rivers  
<sup>786</sup> significantly deviate from those in the Galton-Watson model (Peckham, 1995).  
<sup>787</sup> He reported values  $R_B \approx 4.6$  and  $(a, c) \approx (1.2, 2.5)$  and emphasized the  
<sup>788</sup> importance of studying a broad range of Horton exponents and Tokunaga  
<sup>789</sup> parameters.

<sup>790</sup> The general interest to fractals and self-similar structures in natural sciences  
<sup>791</sup> during the 1990s resulted in a quest, mainly inspired and led by Donald Turcotte,  
<sup>792</sup> for Tokunaga self-similar trees of diverse origin (Gabrielov et al., 1999; Newman  
<sup>793</sup> et al., 1997; Ossadnik, 1992; Pelletier and Turcotte, 2000; Turcotte, 1997;  
<sup>794</sup> Turcotte et al., 1999, 1998; Yakovlev et al., 2005; Zanardo et al., 2013). As a  
<sup>795</sup> result, the Tokunaga model and respective Horton's laws, with a broad range  
<sup>796</sup> of parameters, have been empirically or rigorously found in numerous observed  
<sup>797</sup> and modeled systems, well beyond river networks.

798 **5 Generating Functions: A Tool for Establishing Horton's Laws for  
799 Mean Attributes**

800 This section summarizes the basic facts about generating functions that are  
801 used below to derive asymptotic behavior and Horton's laws for the mean  
802 branch attributes in a self-similar tree.

803 The *generating function*  $f(z)$  of a sequence  $a_k \geq 0$ ,  $k = 0, 1, \dots$ , of non-  
804 negative real numbers is defined as a *formal power series*

$$f(z) = \sum_{k=0}^{\infty} a_k z^k, \quad (28)$$

805 where  $z$  is a complex number,  $z \in \mathbb{C}$ . It is known (Wilf, 1992) that there exist  
806 such a real number  $r \geq 0$  that the series in the right hand side (rhs) of (28)  
807 converges (to the function  $f(z)$ ) for any  $|z| < r$  and diverges for any  $|z| > r$ .  
808 The number  $r$  (which can be infinite) is called the *radius of convergence* of the  
809 sequence  $a_k$ . The value of  $r$  puts notable constraints on the asymptotic behavior  
810 of  $a_k$ . In general, the smaller the radius of convergence, the faster the growth  
811 of the coefficients. Roughly speaking,  $0 < r < 1$  implies that the coefficients  
812  $a_k$  increase geometrically,  $r > 1$  that the coefficients decrease geometrically,  
813 and  $r = 1$  that the coefficient vary at a rate slower than geometric (e.g.,  
814 polynomially). The values  $r = 0$  and  $r = \infty$  imply a faster than geometric  
815 growth or decay, respectively.

816 The Cauchy-Hadamard theorem expresses the radius of convergence in terms  
817 of the series coefficients (Wilf, 1992):

$$\frac{1}{r} = \limsup_{k \rightarrow \infty} a_k^{1/k}. \quad (29)$$

818 We consider the following, consecutively stronger, forms of geometric growth  
819 ( $r > 1$ ) or decay ( $r < 1$ ) of the sequence  $a_k$ :

$$\frac{1}{r} = \lim_{k \rightarrow \infty} a_k^{1/k}, \quad (\text{R})$$

$$\lim_{k \rightarrow \infty} \frac{a_k}{a_{k+1}} = r, \quad (\text{Q})$$

$$\lim_{k \rightarrow \infty} a_k r^k = \alpha \quad (\text{G})$$

822 The three limits are related as follows:

$$(\text{R}) \Rightarrow (\text{Q}) \Rightarrow (\text{G}), \quad (30)$$

823 which means that the existence of (G) implies the existence of (Q), etc. In  
824 hydrogeomorphology, the familiar quotient limit (Q) is referred to as Horton's  
825 law for coefficients  $a_k$  with *Horton exponent*  $r$ . We also consider a weaker root  
826 limit (R) and a stronger geometric limit (G), and refer to them as the root and  
827 geometric Horton's laws, respectively. The limits (R), (Q) and (G) may or may

not exist, but if either of them does, then it has the same Horton exponent  $r$  as in (29) and ensures the existence of the weaker limit(s), according to (30), with the same Horton exponent. As has been mentioned above, we use a convention that Horton exponents are greater than or equal to unity. Hence, if  $r < 1$ , we consider the reciprocal quotient in (Q),  $\lim a_{k+1}/a_k = r^{-1}$ , and change the other limits accordingly to make the Horton exponent equal to  $r^{-1}$ .

Often, the radius of convergence for  $a_k$  can be easily found from an explicit form of  $f(z)$ . Indeed, if  $r > 0$ , then the function  $f(z)$  is analytic within the disk  $|z| < r$  and has at least one *singularity* on the circle  $|z| = r$ , that is it has to diverge for at least one point on that circle (Wilf, 1992, Thm. 2.4.2). Thus, the radius of convergence equals to the modulus of a singularity closest to the origin. Furthermore, recalling that  $a_k \geq 0$  we have

$$|f(z)| = \left| \sum_{k=0}^{\infty} a_k z^k \right| \leq \sum_{k=1}^{\infty} a_k |z|^k = f(|z|), \quad (31)$$

where the equality is only achieved for  $z = |z|$ . This means that the singularity closest to the origin lies on the real axis (although there might be other singularities with the same modulus.) This makes the search for such a singularity much easier: one can only consider the restriction of the function  $f(z)$  to the real axis. In other words, despite the use of complex analysis in establishing some of our results, the applied examination of suitable generating functions can be done in the real domain.

One can examine the function  $f(z)$  in (28) to obtain more precise information about the coefficients  $a_k$ . If  $f(z)$  has no singularities inside the circle  $|z| = \rho$ , then (Ahlfors, 1953)

$$a_k = \frac{1}{2\pi i} \oint_{|z|=\rho} \frac{f(z)dz}{z^{k+1}}. \quad (32)$$

We have mentioned that in general neither of the limits (R), (Q) and (G) must exist. However, if the singularity of  $f(z)$  nearest to the origin is simple enough, then these properties are satisfied.

### Proposition 1 (Geometric Horton's Law for a Simple Pole Sequence)

Suppose  $f(z) = \sum_{k=1}^{\infty} a_k z^k$  with  $a_k \geq 0$  is analytic in the disk  $|z| < \rho$  except for a single singularity that occurs at a positive real value  $r < \rho$ , which is a pole of multiplicity one (simple pole). Then the geometric Horton's law (G), and hence the quotient Horton's law (Q) and the root Horton's law (R), are satisfied for the coefficients  $a_k$  with Horton exponent  $r$ . Furthermore, if we define  $g(z) = f(z)(z - r)$ , then the coefficient in the geometric Horton's law is  $\alpha = -g(r)/r$ .

*Proof* See Appendix A.

Another useful result states that if we can write  $f(z) = g(z)h(z)$ , and the radius of convergence of  $g(z)$  is smaller than that of  $h(z)$ , then the coefficients

864 of  $f(z)$  satisfy the same Horton's laws as those of  $g(z)$ . A formal statement is  
 865 given below.

866 **Proposition 2 (Horton's Laws for Product Sequence)** *Consider complex  
 867 valued functions  $f(z)$ ,  $g(z)$ , and  $h(z) \neq 0$  that are analytic around the origin  
 868 with the following series expansions*

$$f(z) = \sum_{k=0}^{\infty} a_k z^k, \quad g(z) = \sum_{k=0}^{\infty} b_k z^k, \quad \text{and} \quad h(z) = \sum_{k=0}^{\infty} c_k z^k,$$

869 where  $a_k, b_k \in [0, \infty)$  and  $c_k \in \mathbb{R}$  for all  $k$ . Denote by  $r_a$ ,  $r_b$  and  $r_c$  the radii  
 870 of convergence of the sequences  $a_k$ ,  $b_k$  and  $c_k$ , respectively.

871 (i) Suppose that

$$f(z) = g(z)h(z) \quad \text{and} \quad r_c > r_b > 0.$$

872 Then, any of the Horton's laws (R), (Q), and (G) for the sequence  $b_k$   
 873 implies the same law for the sequence  $a_k$ , with Horton exponent  $R = 1/r_b =$   
 874  $1/r_a > 0$ .

875 (ii) Suppose, in addition, that  $h(z) \neq 0$  for  $|z| \leq r_b$ . Then any of the Horton's  
 876 laws (R), (Q), and (G) for the sequence  $a_k$  implies the same law for the  
 877 sequence  $b_k$ , with Horton exponent  $R = 1/r_a = 1/r_b > 0$ .

878 *Proof* See Appendix B.

879 Below, we find the generating functions for the sequences of mean branching  
 880 attributes relevant to our study. The respective radii of convergence provide  
 881 information on the asymptotic behavior of the examined sequences. In particu-  
 882 lar, the generating function  $M(z)$  for the mean branch magnitudes  $M_k$  has a  
 883 simple pole closest to the origin, and hence it satisfies the geometric Horton's  
 884 law by Proposition 1.

## 885 6 Self-Similar Trees: Main Attributes and Their Relations

886 We assume that a river basin is represented by a self-similar tree  $T$  with a  
 887 Tokunaga sequence  $T_k$ . This means, in particular, that each branch of Horton-  
 888 Strahler order  $j$  produces an mean of  $T_{j-i}$  side-branches of order  $i$  for each  
 889  $i$  such that  $1 \leq i < j$ . In this section we state the recursive relations for the  
 890 essential branch attributes: the mean number  $\mathcal{N}_k$  of branches of order  $k$ , the  
 891 mean number  $S_k$  of edges in a branch of order  $k$ , the mean magnitude (number  
 892 of descendant leaves)  $M_k$  of a branch of order  $k$ , the mean length  $L_k$  of a  
 893 branch of order  $k$ , and the mean contributing area  $A_k$  of a branch of order  $k$ .  
 894 These relations are mostly due to straightforward applications of the Wald's  
 895 formula (Bhattacharya and Waymire, 2007); it asserts that for a random sum  
 896 of  $N$  independent identically distributed (i.i.d.) random variables  $X_i$  we have

$$\mathbb{E}[X_1 + \dots + X_N] = \mathbb{E}[X_1]\mathbb{E}[N].$$

897 The expectations are taken with respect to a self-similar measure  $\mu$  on  $\mathcal{L}$ .

898 Let  $S_k$  denote the mean number of edges (or vertices) within a branch of order  
 899  $k$ . This attribute is also known as the mean number of links in a Strahler stream  
 900 of order  $k$  (e.g., Peckham, 1995). It equals the mean number of side-branches  
 901 that join this branch plus the branch starting vertex:

$$S_k = 1 + T_1 + \cdots + T_{k-1}, \quad k \geq 1. \quad (33)$$

902 The *mean magnitude*  $M_k$  is the mean number of leaves descendent to an  
 903 order  $k$  branch. It can be represented as the sum of magnitudes of two order  
 904  $k-1$  branches that created this branch (called *principal branches*), plus the  
 905 magnitudes of all the side-branches. Hence  $M_1 = 1$ , and

$$M_k = 2 M_{k-1} + \sum_{i=1}^{k-1} M_i T_{k-i}, \quad \text{for } k > 1. \quad (34)$$

906 The mean number  $\mathcal{N}_k[K]$  of branches of order  $k$  in a tree of order  $K$ , also  
 907 known as the mean total number of Strahler streams (Peckham, 1995), equals  
 908 twice the number of branches of order  $k+1$  plus the number of instances when  
 909 a branch of rank  $k$  joins a branch of a higher rank. Thus, for a tree of order  $K$   
 910 we have  $\mathcal{N}_K[K] = 1$  and

$$\mathcal{N}_k[K] = 2\mathcal{N}_{k+1}[K] + \sum_{i=k+1}^K \mathcal{N}_i[K] T_{i-k}, \quad \text{for any } k < K. \quad (35)$$

911 Comparing equations (34) and (35) we find

$$\mathcal{N}_{K-k+1}[K] = M_k \quad \text{for all orders } 1 \leq k \leq K. \quad (36)$$

912 Notice that we explicitly indicate the tree order  $K$  when working with the  
 913 mean number of branches  $\mathcal{N}_k[K]$ , and do not do that for  $S_k$  and  $M_k$ . This is  
 914 because the initial terms of the increasing sequences  $M_k$  and  $S_k$  coincide for  
 915 different values of  $K$ , which is not the case for  $\mathcal{N}_k[K]$  (e.g.,  $M_1[4] = M_1[5]$  but  
 916  $\mathcal{N}_1[4] \neq \mathcal{N}_1[5]$ ).

917 Let  $\ell_k$  denote the mean length of an edge of order  $k$ . Then the mean length  $L_k$   
 918 of an order  $k$  branch equals  $L_k = S_k \ell_k$ . By the *differential* contributing area  
 919 (as is opposed to the *total* contributing area) of an edge we understand the area  
 920 that drains directly to the edge (not via its upstream vertex). Assume that the  
 921 mean differential contributing area of an edge of order  $k$  equals  $\alpha_k$ . Then, for  
 922 the total mean contributing area  $A_K$  of a tree of order  $K \geq 1$ , we have

$$A_K = \sum_{k=1}^K \alpha_k S_k \mathcal{N}_k[K], \quad (37)$$

923 where  $S_k \mathcal{N}_k[K]$  is the number of edges of order  $k$  in a tree of order  $K$ .

Another important product that appears in (37) is  $\delta_k = \alpha_k S_k$ , which is the mean differential contributing area of a branch of order  $k$ . The total contributing area  $A_k$  can be expressed recursively by noticing that  $A_1 = \alpha_1$  and

$$A_k = 2A_{k-1} + \alpha_k S_k + \sum_{i=1}^{k-1} A_i T_{k-i} \quad \text{for } k \geq 2. \quad (38)$$

## 7 Horton's Laws in a Self-Similar Tree

This section establishes the main theoretical result of our work: geometric Horton's laws for the mean branch numbers  $\mathcal{N}_k[K]$  and mean magnitudes  $M_k$  in a self-similar tree (Sect. 7.1). Next, we show (Sect. 7.2) how Horton's laws for other mean attributes follow from these ones under additional assumptions.

### 7.1 Geometric Horton's Law for Mean Branch Numbers and Magnitudes

Consider a Tokunaga sequence  $T_k$  and its generating function  $T(z) = \sum_{k=1}^{\infty} T_k z^k$ . If we let  $t_1 = T_1 + 2$ , and  $t_k = T_k$  for  $k \geq 2$ , then  $t_k$  takes into account the *side-branching* and *principal branching*. Let

$$\hat{t}(z) = -1 + \sum_{k=1}^{\infty} t_k z^k = -1 + 2z + T(z).$$

Observe that  $\hat{t}(0) = -1$ , and since  $T_k \geq 0$  we have  $\hat{t}(1/2) = T(1/2) \geq 0$ . Furthermore, since

$$\frac{d}{dz} \hat{t}(z) = 2 + \sum_{k=1}^{\infty} k T_k z^{k-1} > 0$$

for all positive real values of  $z$ , the equation  $\hat{t}(z) = 0$  has a unique real root  $w_0$  of multiplicity one in the interval  $(0, 1/2]$ . Let  $r_T$  be the radius of convergence for  $T(z)$  and define  $R_T = r_T^{-1}$ . We notice that  $r_T > w_0$ . The following result of Kovchegov and Zaliapin (2016) ensures that  $w_0$  is the root of  $\hat{t}(z)$  closest to the origin; this fact will be used below.

**Lemma 1** Suppose  $\limsup_{k \rightarrow \infty} T_k^{1/k} < \infty$  and let  $w_0$  be the only real root of  $\hat{t}(z)$  in the interval  $(0, 1/2]$ . Then, for any other root  $w \in \mathbb{C}$  of  $\hat{t}(z)$ , we have  $|w| > w_0$ .

The generating function for the magnitudes  $M_k$  is obtained by multiplying both sides in (34) by  $z^k$  and summing over all  $k \geq 1$ :

$$M(z) = \sum_{k=1}^{\infty} M_k z^k = z + 2zM(z) + M(z)T(z).$$

948 Thus,

$$M(z) = \frac{z}{1 - 2z - T(z)} = -\frac{z}{\hat{t}(z)}. \quad (39)$$

949 The function  $M(z)$  is analytic with the exception of zeroes and singularities  
 950 of  $\hat{t}(z)$ . Lemma 1 asserts that  $w_0 \in (0, 1/2]$  is the closest to the origin root  
 951 of  $\hat{t}(z)$ ; recall that it has multiplicity one. Hence,  $w_0$  is a simple pole of  $M(z)$   
 952 and the only singularity of  $M(z)$  within a disk  $|z| < w_0 + \epsilon$  for a small enough  
 953  $\epsilon > 0$ . Consequently, the radius of convergence for  $M(z)$  is  $r_M = w_0$ . We define  
 954  $R_M = r_M^{-1}$ . Proposition 1 implies that the geometric Horton's law holds for  
 955  $M_k$ . We formulate this result in the following theorem.

956 **Theorem 1 (Geometric Horton's Law for Mean Branch Magnitudes)**

957 Suppose that  $r_T > 0$ , that is  $\limsup_{k \rightarrow \infty} T_k^{1/k} < \infty$ . Then, the geometric Horton's  
 958 law for mean branch magnitudes  $M_k$  holds with Horton exponent  $R_M = 1/w_0$ ,  
 959 where  $w_0$  is the only real root of the function  $\hat{t}(z) = -1 + 2z + \sum_{j=1}^{\infty} z^j T_j$  in the  
 960 interval  $(0, \frac{1}{2}]$ . Specifically, the geometric Horton's law states that

$$\lim_{k \rightarrow \infty} (M_k R_M^{-k}) = M < \infty, \quad (40)$$

961 where  $M$  is a positive real constant given by

$$M = -\frac{1}{w_0} \lim_{z \rightarrow w_0} \frac{z(z - w_0)}{\hat{t}(z)}. \quad (41)$$

962 The geometric Horton's law implies the quotient Horton's law

$$\lim_{k \rightarrow \infty} \frac{M_{k+1}}{M_k} = R_M. \quad (42)$$

963 Recalling (36) we notice that  $\mathcal{N}_1[K] = M_K$  and hence obtain the asymptotic  
 964 behavior for  $\mathcal{N}_k[K]$ .

965 **Corollary 1 (Geometric Horton's Law for Mean Branch Numbers)**

966 Under the assumption of Theorem 1 the geometric Horton's law holds for  
 967 the mean branch numbers  $\mathcal{N}_k[K]$  with Horton exponent  $R_B = R_M = 1/w_0$ .  
 968 Specifically, the geometric Horton's law states that

$$\lim_{K \rightarrow \infty} (\mathcal{N}_1[K] R_B^{-K}) = M < \infty, \quad (43)$$

969 where  $M$  is the same as in Thm. 1. This implies the quotient Horton's law: for  
 970 each positive integer  $j$  we have

$$\lim_{K \rightarrow \infty} \frac{\mathcal{N}_j[K]}{\mathcal{N}_{j+1}[K]} = R_B \quad \text{or} \quad \lim_{K \rightarrow \infty} \frac{\mathcal{N}_j[K]}{\mathcal{N}_1[K]} = R_B^{1-j}. \quad (44)$$

971 Informally, Theorem 1 and Corollary 1 ensure that the geometric Horton's  
 972 laws for mean branch magnitudes and mean branch counts hold with the same  
 973 Horton exponent in "any" self-similar tree, that is in any coordinated tree with  
 974 a well-defined Tokunaga sequence  $T_k$ . The assumption of non-zero radius of  
 975 convergence in  $T(z)$  eliminates obscure cases of super-exponential growth of  
 976  $T_k$ , such as  $T_k = k!$  or  $T_k = k^k$ .

977 7.2 Horton's Laws for Other Mean Branch Attributes

978 Horton's laws for other mean branch attributes are obtained by examining the  
 979 generating functions for the respective sequences and using the properties of  
 980 the series  $T_k$ ,  $\ell_k$ , and  $\alpha_k$ . The most general results that examine each type  
 981 of the Horton's law (root, quotient, and geometric) under the assumption  
 982  $\limsup_{k \rightarrow \infty} T_k^{1/k} < \infty$  are formulated in Appendix C. This section illustrates a  
 983 particular case of the quotient Horton's law (Q) for selected branch attributes  
 984 under the following more stringent yet practically appealing assumption. Infor-  
 985 mally, it suggests that sequences  $T_k$ ,  $\ell_k$  and  $\alpha_k$  behave "nicely".

**Assumption 1 (Quotient Horton's law for  $T_k$ ,  $\ell_k$ ,  $\alpha_k$ )** Assume that the  
 quotient Horton's law holds for  $T_k$ ,  $\ell_k$  and  $\alpha_k$ :

$$\begin{aligned} \lim_{k \rightarrow \infty} \frac{T_k}{T_{k+1}} &= c^{-1} > 0, & \lim_{k \rightarrow \infty} \frac{\ell_k}{\ell_{k+1}} &= \lambda^{-1} > 0, \\ \text{and} \quad \lim_{k \rightarrow \infty} \frac{\alpha_k}{\alpha_{k+1}} &= \alpha^{-1} > 0. \end{aligned} \quad (45)$$

986 Assumption 1 is satisfied for a multitude of natural choices for  $T_k$ ,  $\ell_k$ , and  $\alpha_k$ ,  
 987 including inverse polynomial, constant, polynomial, and geometric series. The  
 988 inequality  $r_T > w_0$  implies that  $c^{-1} > w_0$ .

The asymptotic behavior for the mean total contributing areas  $A_k$  follows  
 from that for  $M_k$ . First, we write the generating function for  $A_k$  via (38):

$$A(z) = \sum_{k=1}^{\infty} A_k z^k = 2zA(z) + \sum_{k=1}^{\infty} \alpha_k S_k z^k + A(z)T(z),$$

989 which yields, by (39),

$$A(z) = \frac{\sum_{k=1}^{\infty} \alpha_k S_k z^k}{1 - 2z - T(z)} = M(z) \left( \sum_{k=1}^{\infty} \alpha_k S_k z^{k-1} \right) = -\frac{D(z)}{\hat{t}(z)}. \quad (46)$$

990 Here  $D(z)$  is the generating function for the differential contributing areas  
 991  $\delta_k = \alpha_k S_k$  of branches of order  $k$ . The radius of convergence of  $D(z)$  and the  
 992 asymptotic behavior of  $\delta_k$  can be examined using Propositions 5,8 where  $\ell_k$   
 993 need to be replaced with  $\alpha_k$ .  
 994 Equation (46) implies the following convolution expression for  $A_k$  (Wilf, 1992):

$$A_k = \sum_{i=1}^k \alpha_{k+1-i} S_{k+1-i} M_i. \quad (47)$$

996 Observe that comparing equations (47) and (37) we arrive at  $N_{K-k+1}[K] = M_k$   
 997 that was first established in Eq. (36). We denote by  $r_A$  and  $r_D$  the radii of  
 998 convergence of  $A(z)$  and  $D(z)$ , respectively, and let  $R_A = r_A^{-1}$ ,  $R_D = r_D^{-1}$ .

999 **Proposition 3 [Quotient Horton's law for  $A_k$ ]**

1000 Suppose that Assumption 1 holds. Then

$$r_A = \min \{r_M, r_D\} = \min \{r_M, \alpha^{-1} \min\{1, r_T = c^{-1}\}\}.$$

1001 The quotient Horton's law holds with the Horton exponent

$$R_A = \max \{R_M, R_D\} = \max \{w_0^{-1}, \alpha \max\{1, c\}\},$$

1002 that is

$$\lim_{k \rightarrow \infty} \frac{A_{k+1}}{A_k} = R_A. \quad (48)$$

1003 We next examine the mean number  $S_k$  of edges within a branch of order  $k$ . The  
1004 most straightforward practical way to obtain the asymptotic of  $S_k$  is via direct  
1005 application of (33). The generating function approach clarifies the origin of the  
1006 respective Horton's laws. Multiplying both sides in (33) by  $z^k$  and summing  
1007 over  $k = 1, 2, \dots$ , we obtain the generating function  $S(z)$  of  $S_k$ :

$$S(z) = \sum_{k=1}^{\infty} S_k z^k = \sum_{k=1}^{\infty} z^k + \sum_{k=1}^{\infty} \left( \sum_{i=1}^{k-1} T_i \right) z^k = \frac{z(T(z) + 1)}{1 - z}. \quad (49)$$

1008 The function  $S(z)$  may become singular because of a singularity of  $T(z)$  or the  
1009 vanishing denominator  $(1 - z)$ . The singularity of  $S(z)$  closest to the origin is  
1010 the smallest of  $z = 1$  and the (only) singularity of  $T(z)$ . Let  $r_T, r_S$  denote the  
1011 radiiuses of convergence for the series  $T(z)$  and  $S(z)$ , respectively, and define  
1012  $R_T = r_T^{-1}$ ,  $R_S = r_S^{-1}$ .

1013 **Proposition 4 [Quotient Horton's law for  $S_k$ ]**

1014 Suppose that Assumption 1 holds. Then  $r_S = \min\{1, r_T = c^{-1}\}$  and the quotient  
1015 Horton's law holds with the Horton exponent  $R_S = \max\{1, c\} \geq 1$ :

$$\lim_{k \rightarrow \infty} \frac{S_{k+1}}{S_k} = R_S. \quad (50)$$

1016 Next, consider the generating function of  $L_k$  denoted by  $L(z)$ :

$$L(z) = \sum_{k=1}^{\infty} L_k z^k = \sum_{k=1}^{\infty} S_k \ell_k z^k.$$

1017 Let  $r_L$  denotes the radius of convergence for the series  $L(z)$ , and define  $R_L =$   
1018  $r_L^{-1}$ .

1019 **Proposition 5 [Quotient Horton's law for  $L_k$ ]**

1020 Suppose that Assumption 1 holds. Then  $r_L = \lambda^{-1} \min\{1, r_T = c^{-1}\}$  and the  
1021 quotient Horton's law holds with the Horton exponent  $R_L = \lambda \max\{1, c\}$ :

$$\lim_{k \rightarrow \infty} \frac{L_{k+1}}{L_k} = R_L. \quad (51)$$

1022 7.3 Quotient Horton's Laws in a River Basin

1023 Here we use the results of Sect. 7.2 to formulate Horton's laws for branch  
1024 attributes taking into account empirical constraints established for the observed  
1025 river basins.

1026 Observational studies suggest that the link lengths distribution in real rivers  
1027 is independent of the position of the link within a basin (Tarboton et al., 1989;  
1028 Rodriguez-Iturbe and Rinaldo, 2001). This corresponds to the assumption  
1029 that the edge lengths are i.i.d. random variables with the same mean. We  
1030 substantially relax this constraint in the following assumption.

1031 **Assumption 2 (Unit Quotient Horton's law for  $\ell_k$ )** *Assume that the quo-*  
1032 *tient Horton's law with exponent of unity holds for the mean edge lengths  $\ell_k$ :*

$$\lim_{k \rightarrow \infty} \frac{\ell_k}{\ell_{k+1}} = 1. \quad (52)$$

1034 Assumption 2 is trivially satisfied in the case of i.i.d. edge lengths, where  
1035  $\ell_k = \ell_{k+1}$  for all  $k$ . It also allows much wider variability of the edge lengths,  
1036 including, for example, different length distribution for different orders and  
1037 polynomial variation of the means  $\ell_k \sim k^{\text{Const.}}$ .

1038 Assumption 2 strongly supports the existence of the unit quotient Horton's  
1039 law for the mean differential contributing areas  $\alpha_k$  of edges. Specifically, let  
1040  $\xi_k$  and  $\nu_k$  be random variables that represent, respectively, the length and  
1041 differential contributing area of a randomly selected edge of order  $k$ . We have  
1042  $E[\xi_k] = \ell_k$  and  $E[\nu_k] = \alpha_k$ . Suppose there exist scalars  $b > 0$  and  $\sigma > 0$  such  
1043 that

$$\nu_k = b \xi_k^\sigma \quad (53)$$

1044 with  $\sigma$  and  $b$  being the same for all orders  $k \geq 1$ . Suppose, furthermore, that  
1045 the random variables  $\xi_k$  scale with order, that is

$$\frac{\xi_k}{\ell_k} \stackrel{d}{=} \frac{\xi_1}{\ell_1}, \quad (54)$$

1046 where  $\stackrel{d}{=}$  denotes equality of distributions. This gives, in particular,

$$\xi_{k+1} \stackrel{d}{=} \frac{\ell_{k+1}}{\ell_k} \xi_k. \quad (55)$$

Then

$$\alpha_{k+1} = E[\nu_{k+1}] = E[b \xi_{k+1}^\sigma] = E\left[b \left(\frac{\ell_{k+1}}{\ell_k}\right)^\sigma \xi_k^\sigma\right] = \left(\frac{\ell_{k+1}}{\ell_k}\right)^\sigma \alpha_k, \quad (56)$$

1047 which implies

$$\lim_{k \rightarrow \infty} \frac{\alpha_k}{\alpha_{k+1}} = \lim_{k \rightarrow \infty} \left(\frac{\ell_k}{\ell_{k+1}}\right)^\sigma = 1. \quad (57)$$

1048 The asymptotic of  $\alpha_k$  without the scaling assumption (55) is examined in  
1049 Appendix C.

1050 **Example 1 (Exponential edge lengths)** Suppose that the random edge length  
1051  $\xi_k$  is an exponential random variable with parameter  $1/\ell_k$  so that  $E[\xi_k] = \ell_k$ .  
1052 The scaling assumption (55) is satisfied and the expected differential contributing  
1053 area  $\nu_k$  of an order  $k$  edge is given by

$$\alpha_k = b E[\xi_k^\sigma] = b \int_0^\infty \frac{1}{\ell_k} x^\sigma e^{-x/\ell_k} dx = b \Gamma(\sigma + 1) \ell_k^\sigma.$$

1054 Accordingly, the relation (57) holds.

1055 The relation (53) between edge length and differential contributing area  
1056 and the scaling assumption (55) are sufficient but not necessary to obtain  
1057 the quotient Horton's law for  $\alpha_k$  under Assumption 2. To make our results  
1058 applicable to a range of specific situations, where these assumptions may or  
1059 may not hold, we make the following general assumption.

1060 **Assumption 3 (Unit Quotient Horton's law for  $\ell_k, \alpha_k$ )** Assume that the  
1061 quotient Horton's law holds for  $T_k$  and the quotient Horton's law with exponent  
1062 of 1 holds for the mean edge lengths  $\ell_k$  and mean differential edge contributing  
1063 areas  $\alpha_k$ :

$$\lim_{k \rightarrow \infty} \frac{T_k}{T_{k+1}} = c^{-1} > 0, \quad \lim_{k \rightarrow \infty} \frac{\ell_k}{\ell_{k+1}} = \lim_{k \rightarrow \infty} \frac{\alpha_k}{\alpha_{k+1}} = 1. \quad (58)$$

1064 **Proposition 6 [Quotient Horton's laws in a River Basin]**

1065 Suppose that Assumption 3 holds. Then the radii of convergence for the branch  
1066 attributes are related as

$$w_0 = r_M = r_A < r_L = r_S = c^{-1},$$

1067 and, accordingly,

$$c = R_L = R_S < R_B = R_M = R_A = w_0^{-1}.$$

In particular, the following quotient Horton's laws hold

$$\lim_{K \rightarrow \infty} \frac{\mathcal{N}_j[K]}{\mathcal{N}_{j+1}[K]} = \lim_{k \rightarrow \infty} \frac{M_{k+1}}{M_k} = \lim_{k \rightarrow \infty} \frac{A_{k+1}}{A_k} = R_B = w_0^{-1} > 1, \quad (59)$$

$$\lim_{k \rightarrow \infty} \frac{L_{k+1}}{L_k} = \lim_{k \rightarrow \infty} \frac{S_{k+1}}{S_k} = R_L = c. \quad (60)$$

*Proof* Observe that Assumption 1 is satisfied with  $\lambda = \alpha = 1$ , recall that  $w_0 < c^{-1}$ , and apply Theorem 1, Corollary 1, and Propositions 3,4,5. In particular,

$$\begin{aligned} r_A &= \min \{r_M, r_D\} = \min \{w_0, \min \{1, r_T = c^{-1}\}\} = w_0, \\ r_L &= \min \{1, c^{-1}\} > w_0, \quad r_S = \min \{1, c^{-1}\} > w_0. \end{aligned}$$

1068 Finally, we observe that the total length  $L_K^{\text{tot}}$  of a tree of order  $K$  is treated  
1069 similarly to the contributing area  $A_K$ , with  $\alpha_k$  replaced by  $\ell_k$  in definition (37).  
1070 This means that under the hydrology-motivated Assumption 3 of this section,  
1071 the sequence  $L_K^{\text{tot}}$  has the same Horton's law as  $A_K$ .

---

## 1072 8 Self-Similar Hierarchical Branching Processes (HBP)

1073 A flexible model that generates trees with arbitrary Tokunaga sequences  $T_k$  has  
 1074 been introduced and discussed by Kovchegov and Zaliapin (2018, 2020) and  
 1075 Kovchegov et al. (2021); it is called *Hierarchical Branching Processes* (HBP).  
 1076 Here we describe a self-similar version of the HBP.

1077 **Definition 1 (Self-similar Hierarchical Branching Process)** *We say that*  
 1078  *$S(t)$  is a self-similar hierarchical branching process with a Tokunaga sequence*  
 1079  *$\{T_k\}$ , and parameters  $p \in (0, 1)$ ,  $\gamma > 0$  and  $\zeta > 0$  if  $S(t)$  is a multi-type branch-*  
 1080 *ing process that develops in continuous time  $t > 0$  according to the following*  
 1081 *rules:*

- 1082 (i) *The process  $S(t)$  starts at  $t = 0$  with a single progenitor (root branch) whose*  
*Horton-Strahler order (type) is  $K \geq 1$  with probability  $\pi_K = p(1 - p)^{K-1}$ .*
- 1083 (ii) *Every branch of order  $j \leq K$  produces offspring (side branches) of every*  
*order  $i < j$  with rate  $\gamma \zeta^{1-j} T_{j-i}$ .*
- 1084 (iii) *A branch of order  $j$  terminates with rate  $\gamma \zeta^{1-j}$ .*
- 1085 (iv) *At its termination time, a branch of order  $j \geq 2$  splits into two independent*  
*branches of order  $j - 1$ .*
- 1086 (v) *A branch of order  $j = 1$  terminates without leaving offspring.*
- 1087 (vi) *Generation of side branches and termination of distinct branches are inde-*  
*pendent.*

### 1092 8.1 Properties of Self-similar HBP Trees

1093 The trajectories of the HBP are random trees from the space  $\mathcal{L}$  of binary trees  
 1094 with edge lengths. A random tree generated by an HBP is called an HBP tree.  
 1095 Each process parameter completely specifies a particular attribute of a random  
 1096 HBP tree: the Tokunaga sequence  $T_k$  specifies the combinatorial structure  
 1097 of a tree of a given order; the probability  $p$  specifies the frequencies of trees  
 1098 of different orders; the constant  $\zeta$  specifies the ratio of the mean lengths of  
 1099 branches of consecutive orders; and the rate  $\gamma$  specifies a unit of measurement  
 1100 for the tree edges. It has been shown by Kovchegov and Zaliapin (2018) that a  
 1101 random HBP tree  $T$  has the following properties:

1102 **Self-similarity.**  $T$  is a self-similar tree with the Tokunaga sequence  $T_k$  and  
 1103 scaling constant  $\zeta$ . This means that the distribution of  $T$  is invariant with  
 1104 respect to the Horton pruning as in (22), and the edge lengths of the random  
 1105 tree scale by  $\zeta^{-1}$  after each Horton pruning.

1106 **Side-branching.** The number  $N[b] \geq 0$  of side-branches within a branch  $b$  of  
 1107 order  $K$  has geometric distribution

$$\mathbb{P}(N[b] = k) = q(1 - q)^k, \quad \text{with } q = S_K^{-1}, \quad (61)$$

1108 where  $S_K$  is defined in (33). This implies, in particular,  $\mathbb{E}[N[b]] = S_K - 1$ .

1109 **Side-branch orders.** Let  $N_i[b] \geq 0$  be the number of side branches of order  $i$   
1110 within branch  $b$ . Conditioned on the total number  $N[b]$  of side branches, the  
1111 distribution of vector  $(N_1[b], \dots, N_{K-1}[b])$  is multinomial with  $N[b]$  trials  
1112 and success probabilities

$$\mathbb{P}(\text{side branch has order } i) = \frac{T_{K-i}}{S_K - 1}. \quad (62)$$

1113 **Branch and edge lengths.** The length of an order  $K$  branch has exponential  
1114 distribution with rate  $\gamma\zeta^{1-K}$ . The corresponding edge lengths  $\xi_K$  are i.i.d.  
1115 exponential random variables with rate

$$\gamma\zeta^{1-K} S_K. \quad (63)$$

1116 Accordingly,

$$\frac{\mathbb{E}[\xi_{k+1}]}{\mathbb{E}[\xi_k]} = \frac{\ell_{k+1}}{\ell_k} = \zeta \frac{S_k}{S_{k+1}}. \quad (64)$$

1117 This means that the quotient Horton's law holds for  $S_k$  if and only if the  
1118 quotient Horton's law holds for  $\ell_k$ , and in this case  $R_S = \zeta\lambda^{-1}$ . The same  
1119 equivalence holds for the root and geometric Horton's laws.

1120 **Geometric Horton's law for branch lengths.** Definition 1(iii) implies the  
1121 geometric Horton's law for the branch lengths  $L_k$  with  $R_L = \zeta$ . In fact,  
1122 here we have a stronger statement that holds for any  $k \geq 1$  (not only in  
1123 the limit of large  $k$ ):

$$L_k R_L^{-k} = \frac{\zeta^{k-1}}{\gamma} \zeta^{-k} = \frac{1}{\zeta\gamma}. \quad (65)$$

1124 It follows, in particular, that

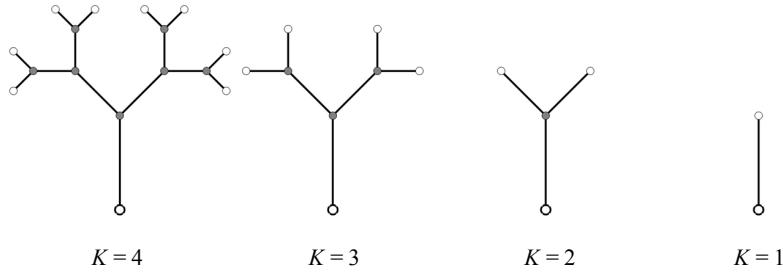
$$\frac{L_{k+1}}{L_k} = \zeta \quad \text{for any } k \geq 1.$$

1125 **Distribution of subtrees.** Consider a random HBP tree  $T$  and fix  $K$  such  
1126 that  $K > 1$  and  $K$  is less than or equal to the order of  $T$ . Select a uniform  
1127 random vertex  $v \in T$  of order  $K$ , and consider two planted trees  $T_a$  and  $T_b$   
1128 descendant to  $v$  in  $T$  that have  $v$  as their root. Informally, we consider the  
1129 pair of sibling trees at a random vertex of order  $K$ . It has been shown by  
1130 (Kovchegov and Zaliapin, 2020, Lem. 16) that the joint distribution of the  
1131 ordered statistics  $(K_1, K_2)$  of the orders  $(K_a, K_b)$  of these trees is given by

$$\mathbb{P}(K_1 = j, K_2 = m | K) = \begin{cases} S_K^{-1} & \text{if } j = m = K - 1, \\ T_{K-j} S_K^{-1} & \text{if } j < m = K. \end{cases} \quad (66)$$

1132 Moreover, the trees  $T_a$  and  $T_b$  are also HBP trees with the same parameters  
1133  $(\gamma, \zeta)$  as  $T$  and orders given by (66). This result is essential for a fast  
1134 recursive construction of  $T$  described in Sect. 8.2.

1135 **Independence of branches.** Distinct branches have independent structure.



**Fig. 7** Examples of perfect planted binary trees of orders  $K = 1, \dots, 4$ .

1136 8.2 Simulation of Self-similar HBP Trees

1137 This section describes three algorithms for constructing HBP trees that do  
1138 not involve time-dependent simulations. Each algorithm constructs a tree of a  
1139 given order  $K$ . To construct a random HBP tree, one first generates a random  
1140 order  $K \geq 1$  according to  $\pi_K$  of Definition 1(i) and then constructs a tree of  
1141 order  $K$  using either of the algorithms. Examples of HBP trees are shown in  
1142 Fig. 4.

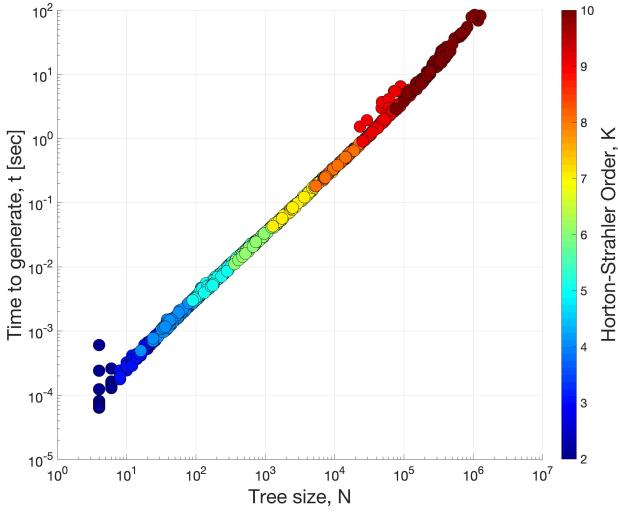
1143 8.2.1 Algorithm A: Recursion by Tree Depth

1144 This is the most straightforward algorithm that relies on the property (66)  
1145 of subtree distributions described above in Sect. 8. The tree is constructed  
1146 recursively, starting from the root and adding two principal subtrees at every  
1147 recursion step.

1148 Formally, a tree of order  $K = 1$  consists of two vertices (root and leaf)  
1149 connected by an edge of exponential length with rate  $\gamma$ . To construct a tree of  
1150 a given order  $K \geq 2$  we first use (66) to obtain the orders of its two principal  
1151 subtrees  $T_a$  and  $T_b$  rooted at the stem vertex farthest from the root. The stem  
1152 edge has exponential length with parameter  $\gamma\zeta^{1-K}S_K$ . To this stem we attach  
1153 two trees, each of which is generated using the same approach via recursion.

1154 Observe that the main branch of order  $K$  consists of a finite number of edges  
1155 with probability one, meaning that after a finite number of recursion steps  
1156 the two principal trees  $T_a$  and  $T_b$  will both have order  $K - 1$ . Extending this  
1157 argument to consequently smaller orders we find that this recursive procedure  
1158 does stop at a finite number of steps.

1159 The algorithm is linear with respect to the tree size  $N$  (number of tree  
1160 vertices) – it takes  $O(N)$  time units to generate a tree of size  $N$ , or, equivalently,  
1161  $O(R_B^K)$  time units to generate a tree of order  $K$ . Figure 8 illustrates the relations  
1162 among the tree size  $N$ , order  $K$ , and generation timing  $t$  in seconds for the  
1163 critical Tokunaga trees with  $c = 2$  and orders  $2 \leq K \leq 10$ . This corresponds to  
1164 sizes between  $N = 4$  and  $N = 1\,266\,454$ . We generated 100 trees of each order.  
1165 The time to create the largest trees using a 3.5 GHz desktop is about 100 sec.



**Fig. 8** Time  $t$  (in seconds) taken to generate HBP trees using the recursive Algorithm A of Sec. 8.2.1. The time  $t$  is shown as a function of tree size (number of vertices)  $N$ . Color represents the tree order  $K$  (see colorbar). This experiment corresponds to the critical Tokunaga process with  $\gamma = 1$  and  $c = 2$ , which is equivalent to critical binary Galton-Watson tree with exponential edge lengths. We generated 100 trees of each order  $2 \leq K \leq 10$ . The computations were performed in Matlab on an Apple Desktop 3.5 GHz 6-Core Intel Xeon E5 with 32GB memory.

For example, a tree of order  $K = 8$  has theoretical mean size  $2\mathcal{N}_1[8] = 21,846$  according to (73). In our simulations the mean size of order-8 trees is  $\bar{N} = 21\,445$ , and the mean running time is  $\bar{t} = 0.78$  sec.

It is clear from the algorithm description that the running time only depends on the Tokunaga sequence  $T_k$  and order  $K$  via the tree size, and is independent of the parameters  $\gamma$  and  $\zeta$ .

### 8.2.2 Algorithm B: Recursion by Tree Order

This algorithm is given in Kovchegov and Zaliapin (2020). It uses a recursion by tree order – we start with a perfect binary tree of order  $K$  and add its side-branches of smaller orders.

Formally, a tree of order  $K = 1$  consists of two vertices (root and leaf) connected by an edge of exponential length with rate  $\gamma$ . Assume now that we know how to construct a random tree of any order below  $K \geq 2$ . To construct a tree of order  $K$ , we start with a perfect (combinatorial) planted binary tree of depth  $K$ , which we call *skeleton*. The combinatorial shapes of such trees is illustrated in Fig. 7. All leaves in the skeleton have the same depth  $K$ , and all vertices at depth  $\kappa$  such that  $1 \leq \kappa \leq K$  have the same Horton-Strahler order  $K - \kappa + 1$ . The root (at depth 0) has order  $K$ . Next, we assign lengths to the branches of the skeleton. Observe that each branch in a perfect tree consists of

a single edge. To assign length to a branch  $b$  of order  $\kappa$ , with  $1 \leq \kappa \leq K$ , we generate a geometric number  $N[b]$  according to (61) with parameter  $q = S_\kappa^{-1}$  and then  $N[b] + 1$  i.i.d. exponential lengths  $\xi_{\kappa,i}$ ,  $i = 1, \dots, N[b] + 1$ , with the common rate  $\gamma\zeta^{1-\kappa}S_\kappa$  according to (63). The total length of the branch  $b$  is  $\xi_{\kappa,1} + \dots + \xi_{\kappa,N[b]+1}$ . The branch  $b$  has  $N[b]$  side branches that are attached along  $b$  with spacings  $\xi_{\kappa,i}$ , starting from the branch point closest to the root. The order assignment for the side branches is done according to (62). We generate side branches (each has order below  $K$ ) independently and attach them to the branch  $b$ . This completes the construction of a random tree of order  $K$ .

### 1195 8.2.3 Algorithm C: Random Attachment Model

1196 Here we construct a Markov tree process  $\{\Upsilon_K\}_{K=1,2,\dots}$  corresponding to the  
 1197 HBP  $S(t)$  following (Kovchegov and Zaliapin, 2020, Sect. 7.6). Each tree  $\Upsilon_K$  is  
 1198 distributed as a tree generated by the self-similar HBP with Tokunaga sequence  
 1199  $\{T_k\}$  and parameters  $(\gamma, \zeta)$ , conditioned on its Horton-Strahler order being  
 1200 equal to  $K$ , and with its edge lengths scaled by  $\zeta^{1-K}$ . This scaling is needed  
 1201 to ensure that  $\Upsilon_K \subset \Upsilon_{K+1}$ , when we consider each tree as a metric space of  
 1202 points connected by paths along the tree edges (Kovchegov and Zaliapin, 2020,  
 1203 Sect. 2.2). Accordingly, there exists the limit space, which informally can be  
 1204 considered an ‘‘infinite tree’’:

$$\Upsilon_\infty = \lim_{K \rightarrow \infty} \Upsilon_K = \bigcup_{K=1}^{\infty} \Upsilon_K.$$

1205 Section 10 uses this construction to find the fractal dimension of  $\Upsilon_\infty$ .

1206 Let  $\Upsilon_1$  be an I-shaped tree of Horton-Strahler order one, with the edge length  
 1207 distributed as an exponential random variable with parameter  $\gamma$ . Conditioned  
 1208 on  $\Upsilon_K$ , the tree  $\Upsilon_{K+1}$  is constructed according to the following transition rules.  
 1209 We attach new leaf edges to  $\Upsilon_K$  at the points sampled by an inhomogeneous  
 1210 Poisson point process with the intensity  $\rho_{j,K} = \gamma\zeta^{K-j}T_j$  along the edges  
 1211 of order  $j \leq K$  in  $\Upsilon_K$ . We also attach a pair of new leaf edges to each of  
 1212 the leaves in  $\Upsilon_K$ . The lengths of all the newly attached leaf edges are i.i.d.  
 1213 exponential random variables with parameter  $\gamma\zeta^K$  that are independent of the  
 1214 combinatorial shape and the edge lengths in  $\Upsilon_K$ . Finally, we let the tree  $\Upsilon_{K+1}$   
 1215 consist of  $\Upsilon_K$  and all the attached leaves and leaf edges.

By construction, a branch of order  $j$  in  $\Upsilon_K$  becomes a branch of order  $j+1$  in  $\Upsilon_{K+1}$  after the attachment of new leave edges. The length of order  $j$  branch in  $\Upsilon_K$  (and therefore, the length of order  $j+1$  branch in  $\Upsilon_{K+1}$ ) is exponential random variable with parameter  $\gamma\zeta^{K-j}$ . Therefore, in a tree  $\Upsilon_{K+1}$ , the number  $n_{1,j+1}(K+1)$  of side-branches of order one in a branch of order  $j+1$  has

geometric distribution:

$$\begin{aligned} \mathsf{P}(n_{1,j+1}(K+1) = r) &= \frac{\gamma\zeta^{K-j}}{\gamma\zeta^{K-j} + \rho_{j,K}} \left( \frac{\rho_{j,K}}{\gamma\zeta^{K-j} + \rho_{j,K}} \right)^r \\ &= \frac{1}{1+T_j} \left( \frac{T_j}{1+T_j} \right)^r, \quad r = 0, 1, 2, \dots \end{aligned} \quad (67)$$

1216 with the mean value

$$\mathsf{E}[n_{1,j+1}(K+1)] = \frac{\rho_{j,K}}{\gamma\zeta^{K-j}} = T_j.$$

Therefore, after  $i \geq 1$  rounds of attachments the mean number  $n_{i,j+i}(M)$  of side-branches of order  $i$  in a branch of order  $j+i$  in a tree  $\mathcal{T}_M$  (where  $M = K+i$  and  $K \geq j$ ) is

$$\mathsf{E}[n_{i,j+i}(M)] = T_j.$$

1217 **8.2.4 Comparison of the Algorithms**

1218 The Algorithms A and B are best suited for numerical simulations of HBP trees,  
1219 while the Random Attachment Model of Algorithm C has mainly a theoretical  
1220 value. The Algorithm A (Sect. 8.2.1) is slightly simpler than the Algorithm B  
1221 (Sect. 8.2.2), as it only involves generating a single edge length and merging  
1222 two trees. However, the recursion by tree depth used in Algorithm A could  
1223 make it computationally prohibitive. Heuristically, the expected value of the  
1224 tree depth  $\text{depth}(T)$  can be approximated by the sum of combinatorial lengths  
1225 (number of edges) of all orders:

$$\mathsf{E}[\text{depth}(T)] \approx \sum_{k=1}^K S_k \sim \text{Const.} \times c^K \quad \text{as } K \rightarrow \infty.$$

1226 The expression  $a_k \sim b_k$  as  $k \rightarrow \infty$  means that  $\lim_{k \rightarrow \infty} a_k/b_k = 1$ . This gives a  
1227 coarse estimate on the recursion depth that is required to successfully use  
1228 Algorithm A in generating large trees. For example, a critical binary Galton-  
1229 Watson tree corresponds to  $c = 2$ . Hence, the depth of a tree of order  $K = 10$   
1230 is about  $2^{10} = 1024$ . The mean size of such a tree is 349 525. All simulations  
1231 in this work have been done using Algorithm A.

1232 **9 Critical Tokunaga Tree: A Model for River Networks**

1233 Recall that the analysis of the observed river networks suggests that the  
1234 distribution of edge lengths is independent of their position within a tree (see  
1235 Sect. 7.3). Formally, this corresponds to the assumption that the edge lengths  
1236  $\xi_k$  are i.i.d. random variables. For the HBP model, this assumption is satisfied  
1237 only for a particular one-parametric class of trees, called *critical Tokunaga*  
1238 *process*, that we describe in this section. The critical Tokunaga trees enjoy

many additional symmetries as discussed by Kovchegov and Zaliapin (2018, 2019, 2020). The class is sufficiently broad and includes the critical binary Galton-Watson process with exponential edge lengths as a special case. Table 2 summarizes the main findings for the critical Tokunaga process and lists the values of its essential exponents fit to the observed river networks. Several examples of critical Tokunaga trees are shown in Fig. 4.

**Definition 2 (Critical Tokunaga process)** *For given  $\gamma > 0$  and  $c \geq 1$ , we say that a self-similar hierarchical branching process  $S(t)$  is a critical Tokunaga process with parameters  $(\gamma, c)$  if*

$$p = \frac{1}{2}, \quad \zeta = c, \quad \text{and} \quad T_k = (c - 1) c^{k-1}. \quad (68)$$

Similarly to the general HBP, we call a random tree generated by the critical Tokunaga process a critical Tokunaga tree. The parameter  $c$  completely determines the combinatorial structure of a random critical Tokunaga tree of a given order and scaling of the mean branch lengths. The rate  $\gamma$  specifies a unit of measurement.

It has been shown by (Kovchegov and Zaliapin, 2020, Thm. 14) that the critical Tokunaga process has the unit mean progeny, which explains the name critical, as in critical Galton-Watson process. These authors also have shown (Kovchegov and Zaliapin, 2020, Thm. 15) that the critical Tokunaga process with parameters  $(\gamma, c = 2)$  is equivalent to the critical binary Galton-Watson process with edge lengths distributed as independent exponential random variables with rate  $\gamma$ . One can also observe that  $c = 1$  corresponds to  $T_k = 0$ , which results in a perfect binary tree (no side branching).

## 9.1 Horton's Laws for Mean Attributes

According to Def. 2, the distribution of orders in the critical Tokunaga process is  $\pi_K = 2^{-K}$ . The mean number of edges within a branch of order  $K$  is  $S_k = 1 + T_1 + \dots + T_{k-1} = c^{k-1}$ . The edge lengths are i.i.d. exponential variables with common rate  $\gamma$ . Accordingly,  $\ell_k = \gamma^{-1}$  and  $L_k = \gamma^{-1} S_k = \gamma^{-1} c^{k-1}$ . While the HBP does not have formally defined areas, the discussion in Sect. 7.3 and Example 1 suggest that one can set  $\alpha_k = b\Gamma(\sigma + 1)\gamma^{-\sigma}$  for some  $b > 0$  and  $\sigma > 0$ .

The generating function  $\hat{t}(z)$  is given by

$$\hat{t}(z) = -1 + 2z + \sum_{k=1}^{\infty} (c - 1) c^{k-1} z^k = \frac{(1 - 2cz)(z - 1)}{1 - cz}. \quad (69)$$

The real root of  $\hat{t}(z)$  nearest to the origin is  $w_0 = (2c)^{-1}$ . Assumption 3 trivially holds, and Proposition 6 implies the existence of the quotient Horton's laws for  $\mathcal{N}_k[K]$ ,  $M_k$ ,  $A_k$ ,  $L_k$ , and  $S_k$ , with Horton exponents

$$c = R_L = R_S < R_M = R_N = R_A = 2c. \quad (70)$$

<sup>1273</sup> Stronger results are readily obtained by examining the generating functions.  
<sup>1274</sup> The generating function for the mean branch magnitudes  $M_k$  is

$$M(z) = -\frac{z}{\hat{t}(z)} = -\frac{z(1-cz)}{(1-2cz)(z-1)}, \quad (71)$$

<sup>1275</sup> and by Theorem 1

$$\mathcal{N}_{K-k+1}[K] = M_k = \frac{(2c)^k}{4c-2} + o((2c)^k). \quad (72)$$

<sup>1276</sup> An exact expression has been obtained using more powerful martingale tech-  
<sup>1277</sup> niques by (Kovchegov and Zaliapin, 2020, Eq. (133) of Cor. 4):

$$\mathcal{N}_{K-k+1}[K] = M_k = \frac{(2c)^k + 2c - 2}{4c - 2}. \quad (73)$$

<sup>1278</sup> The generating function  $D(z)$  for the mean differential contributing areas is

$$D(z) = \sum_{k=1}^{\infty} \alpha_k S_k z^k = d_0 \sum_{k=1}^{\infty} c^{k-1} z^k = \frac{d_0 z}{1-cz}, \quad (74)$$

<sup>1279</sup> with  $d_0 = b\Gamma(\sigma + 1)\gamma^{-\sigma}$ . This leads to

$$A(z) = -\frac{D(z)}{\hat{t}(z)} = -\frac{d_0 z}{(1-2cz)(z-1)}. \quad (75)$$

<sup>1280</sup> Proposition 1 gives the asymptotic form of  $A_k$ :

$$A_k = \frac{d_0(2c)^k}{2c-1} + o((2c)^k). \quad (76)$$

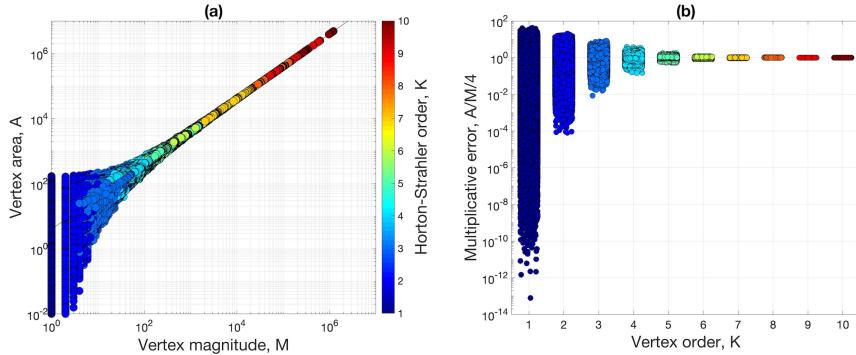
<sup>1281</sup> According to this discussion, the geometric Horton's law (not only the  
<sup>1282</sup> quotient Horton's law) holds for  $\mathcal{N}_k[K]$ ,  $M_k$  and  $A_k$ .

<sup>1283</sup> Comparing (72) and (76) we find

$$A_k \sim 2d_0 M_k \quad \text{as } k \rightarrow \infty. \quad (77)$$

<sup>1284</sup> Figure 9(a) shows the relation among the contributing areas and magnitudes of  
<sup>1285</sup> vertices in a critical Tokunaga tree with  $c = 2.5$  of order  $K = 10$ . This tree has  
<sup>1286</sup> 2 440 508 vertices. The asymptotic relation (77), which becomes in this case  
<sup>1287</sup>  $A = 4M$ , is closely followed for the high-order vertices,  $K \geq 6$ . At the same  
<sup>1288</sup> time, the low-order vertices may show pre-asymptotic behavior that results in  
<sup>1289</sup> substantial deviations from the asymptotic approximation  $A = 4M$ . Panel (b)  
<sup>1290</sup> shows the multiplicative error  $A/(4M)$  as a function of the vertex order. One  
<sup>1291</sup> notices extreme errors, up to several orders of magnitude, for the orders below  
<sup>1292</sup>  $K = 5$ . This experiment shows that the conventional approximation  $A \propto M$   
<sup>1293</sup> has a substantial error at low orders and may be misleading.

<sup>1294</sup> Figure 10 illustrates Horton's laws for  $N_k[T]$ ,  $M_k$ ,  $A_k$ , and  $S_k$  in three  
<sup>1295</sup> critical Tokunaga trees of order  $K = 9$  with  $\gamma = 1$ . We consider the perfect

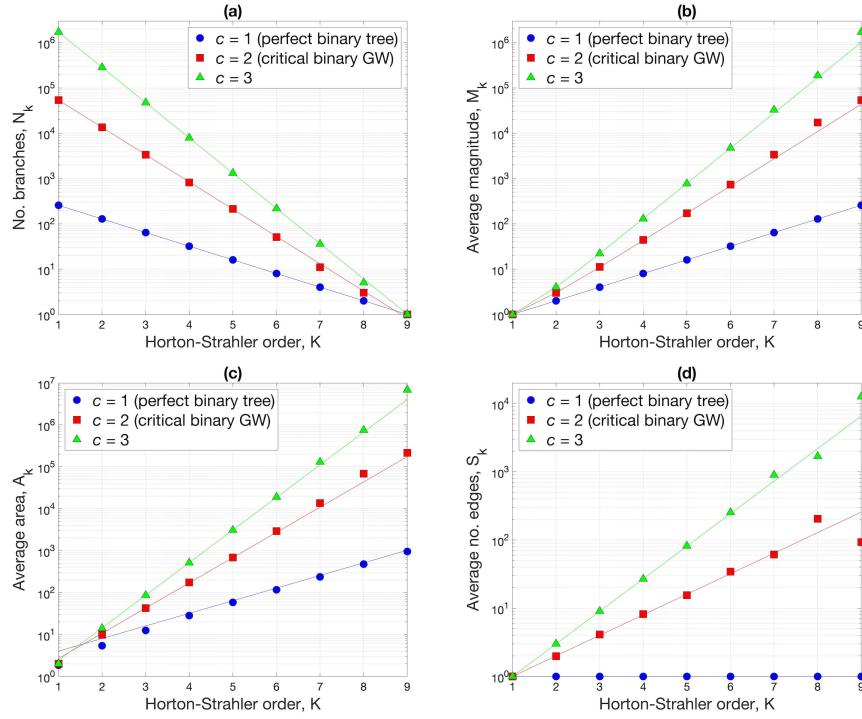


**Fig. 9** Relation between vertex contributing area  $A_{(i)}$  and vertex magnitude  $M_{(i)}$  in a critical Tokunaga tree with  $c = 2.5$  of order  $K = 10$ . (a)  $A_{(i)}$  as a function of  $M_{(i)}$ . The theoretical asymptotic relation (77),  $A = 4M$ , is shown by black line. Colorcode corresponds to vertex order (see colorbar). The vertical axis is trimmed at  $10^{-2}$ , although the minimal area is  $3 \times 10^{-13}$ . (b) Multiplicative error  $A_{(i)} / (4M_{(i)})$  of the asymptotic approximation: unity corresponds to a perfect fit. The horizontal coordinates are given with a uniform jitter. Colorcode is the same as in panel (a), it duplicates the horizontal coordinate.

binary tree ( $c = 1$ , blue circles), a critical binary Galton-Watson tree ( $c = 2$ , red squares), and a tree with  $c = 3$  (green triangles). Panel (a) illustrates Horton's law for the branch counts  $N_k[T]$ . Here the fitting lines start at  $N_1[T]$  and have a constant slope of  $-\log_{10}(2c)$  suggested by (78). We notice a very close fit for all examined orders.

Panel (b) refers to Horton's law for the average magnitudes  $M_k$ . The fitting lines are taken from (73) – the fit is ideal for the deterministic perfect binary tree (blue) and is very close for the two random trees (red and green). The only visible deviations from the theoretical quantities are observed for the high-order branches ( $K > 6$ ) that correspond to small-sample averaging. For instance, recall from panel (a) that we only have an average of 3 branches of order  $K = 8$  and 11 branches of order  $K = 7$  in a critical binary Galton-Watson tree (red squares).

Panel (c) illustrates Horton's law for the average contributing areas  $A_k$ . We assume here that the differential contributing areas  $\nu_k$  of edges are determined by the edge lengths  $\xi_k$  via  $\nu_k = \xi_k^\sigma$  with  $\sigma = 2$  and the contributing area of a branch is the sum of the differential contributing areas of all its descendant edges; see Sect. 7.3 for a discussion and examples of this approach. The fitting lines here correspond to the asymptotic expression (76) with  $d_0 = \gamma^{-\sigma} \Gamma(\sigma + 1) = \Gamma(3) = 2$ . We expect them to fit the observed values for the intermediate range of orders – when the asymptotic approximation already works yet the sample size (number of branches of a given order) is still large enough. In random trees (red and green), the fitting lines provide almost perfect fit to the data for orders  $2 \leq K \leq 6$  and show very small deviations at the higher orders. For the perfect binary tree (with random edge lengths and areas), where  $c = 1$ , the best fit is for the largest orders  $K > 5$ . The discrepancy at the low orders  $K < 4$  is related to the fact that the asymptotic expression (76)



**Fig. 10** Horton's laws in critical Tokunaga trees (Sec. 9). The figure illustrates three cases: a perfect binary tree,  $c = 1$  (blue circles); a critical binary Galton-Watson tree,  $c = 2$  (red squares); and  $c = 3$  (green triangles). Symbols correspond to the average attributes estimated in a single realization of a tree. Lines correspond to theoretical predictions (see below). (a) Branch counts  $N_k[T]$ . Lines start at  $N_1[T]$  and have the theoretical slopes  $-\log_{10}(2c)$  of (72). (b) Average branch magnitudes  $M_k$ . Lines show the theoretical means of (73). (c) Average contributing areas  $A_k$ . Lines show the theoretical means of (76). (d) Average number of edges  $S_k$ . Lines show the theoretical means of (33).

1323 suggests  $A_1 \approx 2d_0c/(2c - 1) = 4$ , while the actual mean here (for any  $c$ ) is  
1324  $A_1 = E[\xi^2] = 2$ , with  $\xi$  being an exponential random variable with rate  $\gamma = 1$ .  
1325 In general, observe that

$$A_1 \approx d_0 \frac{2c}{2c - 1} \rightarrow d_0, \quad \text{as } c \rightarrow \infty,$$

1326 so the asymptotic expression (76) does provide a good fit to the data at low  
1327 orders for large enough  $c$ , which we do observe for  $c = 2, 3$  (red and green).

1328 Panel(d) refers to Horton's law for the average number of edges  $S_k$ . The  
1329 fitting lines show the theoretical values  $S_k = 1 + T_1 + \dots + T_{k-1}$ . It follows from  
1330 the properties of the HBP model that this is merely an exercise in sampling  
1331 from the geometric distribution. The fit quality depends on the sample size  
1332 and is very good for orders  $K < 7$ . The deviations at higher orders are due  
1333 to small sample size. The perfect binary tree ( $c = 1$ , blue) has all branches  
1334 consisting of a single edge.

The condition  $T_{i,i+k} = T_k = a c^{k-1}$ , which is slightly more general than that of (68), was first introduced in hydrology by Eiji Tokunaga (1978) in a study of river networks, hence the process name. In the present work, the constraint  $a = c - 1$  is necessitated by the equality of the mean edge lengths, which requires the sequence  $\lambda_j$  to be geometric. The sequence of the Tokunaga coefficients then also has to be geometric, and satisfy  $a = c - 1$ . Interestingly, the constraint  $a = c - 1$  appears in the *random self-similar network* (RSN) model introduced by Veitzer and Gupta (2000), which uses a purely topological algorithm of recursive local replacement of the network generators to construct random self-similar trees. Results of Chunikhina (2018a,b) imply that the critical Tokunaga model with  $c = 2$  maximizes the entropy rate among the trees that satisfy the quotient Horton's law of stream numbers, and that the critical Tokunaga model with a fixed  $c$  maximizes the entropy rate among the trees that satisfy the quotient Horton's law for stream numbers with  $R_B = 2c$ .

## 9.2 Limit Laws for Random Branch Numbers

While this review mainly focuses on the mean branch numbers  $N_k[K]$ , much stronger, distributional, results are available for the random branch numbers  $N_k[T]$  in critical Tokunaga trees. This section assumes that we consider a critical Tokunaga tree with parameter  $c$ . The results refer to the combinatorial tree structure, and hence hold for an arbitrary measurement unit  $\gamma > 0$ .

Let  $\Delta_K$  be a critical Tokunaga tree of order  $K$ , then the following Weak Law of Large Numbers holds (Kovchegov and Zaliapin, 2020, Sect. 7.6.3, Cor. 5): for any  $k \geq 1$  we have

$$\frac{N_k[\Delta_K]}{N_{k+1}[\Delta_K]} \xrightarrow{p} R_B = 2c \quad \text{as } K \rightarrow \infty, \quad (78)$$

where  $\xrightarrow{p}$  denotes convergence in probability. Moreover, using the notations of Sect. 8.2.3 one can establish the following Strong Law of Large Numbers (Kovchegov and Zaliapin, 2020, Sect. 7.6.3, Thm. 16): for any  $k \geq 1$  we have

$$\frac{N_k[\Upsilon_K]}{N_{k+1}[\Upsilon_K]} \xrightarrow{\text{a.s.}} R_B = 2c \quad \text{as } K \rightarrow \infty, \quad (79)$$

where  $\xrightarrow{\text{a.s.}}$  denotes convergence with probability one (almost sure convergence). The difference between these two laws is that (78) considers a sequence of independent trees  $\Delta_K$ , while (79) refers to a sequence of trees  $\Upsilon_K$  related such that  $\mathcal{R}(\Upsilon_K) = \Upsilon_{K-1}$ . The following distributional geometric Horton's law also holds (Kovchegov and Zaliapin, 2020, Sect. 7.6.3, Cor. 16):

$$R_B^{1-K} N_1[\Upsilon_K] \xrightarrow{\text{a.s.}} V_\infty(\Upsilon_\infty) \quad \text{as } K \rightarrow \infty, \quad (80)$$

where  $V_\infty(\Upsilon_\infty)$  is a finite and positive value that depends on a particular realization of the Markov process  $\Upsilon_K$ . This random value allows different random trees to have very different sizes, while preserving the relative frequencies of

1369 branches of different orders. Specifically, the frequencies of branches of different  
1370 orders in a tree of order  $K$  are approximated by a geometric distribution:

$$\frac{\#\{\text{branches of order } k\}}{\#\{\text{branches}\}} = (R_B - 1)R_B^{-k}(1 + o(1)) \quad (81)$$

1371 for any  $k \geq 1$  as  $K \rightarrow \infty$ .

1372 Finally, one can prove the following Central Limit Theorem:

1373 **Theorem 2 (CLT for Branch Numbers in Critical Tokunaga Trees)**

1374 Let  $\Delta_K$  be a critical Tokunaga tree of order  $K > 0$  with parameter  $c \geq 1$ . Then,  
1375 for every integer  $k \geq 1$ ,

$$\sqrt{N_1[\Delta_K]} \left( \frac{N_{k+1}[\Delta_K]}{N_k[\Delta_K]} - \frac{1}{2c} \right) \xrightarrow{d} N(0, (c-1)(2c)^{k-3}) \quad \text{as } K \rightarrow \infty, \quad (82)$$

1376 where  $N(\mu, \sigma^2)$  denotes a Normal distribution with mean  $\mu$  and variance  $\sigma^2$ ,  
1377 and  $\xrightarrow{d}$  denotes convergence in distribution.

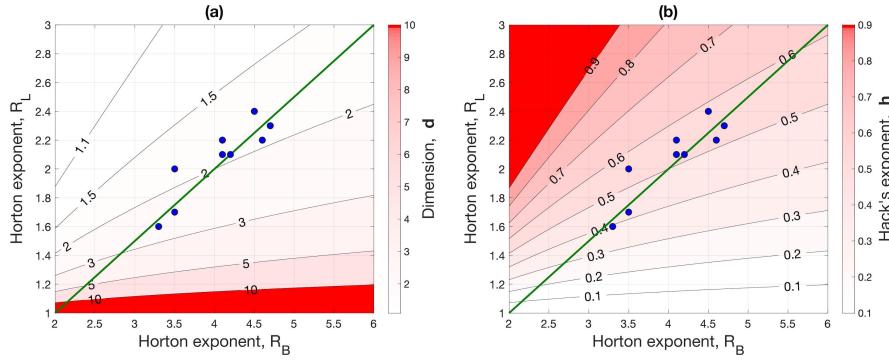
1378 In a special case of the critical Tokunaga tree with  $c = 2$ , which is equivalent  
1379 to the critical Galton-Watson process, similar limit results were established by  
1380 Wang and Waymire (1991), Yamamoto (2017), and Kovchegov and Zaliapin  
1381 (2020). These authors considered a random tree conditioned on the number  $N_1$   
1382 of leaves and established the conditional limits of  $N_{k+1}/N_k$  as  $N_1 \rightarrow \infty$ . We  
1383 notice that, trivially,  $K \rightarrow \infty$  implies  $N_1 \rightarrow \infty$  since a binary tree of order  $K$   
1384 has at least  $2^{K-1}$  leaves. Hence, the limit results for an increasing order  $K$   
1385 follow from those for an increasing tree magnitude  $N_1$ . Vice versa, it can be  
1386 shown that increasing  $N_1$  implies an increasing order  $K$ . In summary, the above  
1387 limit laws are equivalent under both forms of increasing tree size in critical  
1388 Tokunaga trees. This equivalence issue will be treated formally elsewhere.

1389 **10 Fractal Dimension of Self-Similar HBP Trees**

1390 Consider a self-similar HBP  $S(t)$  (Def. 1) with a Tokunaga sequence  $\{T_k\}$   
1391 satisfying  $\limsup_{k \rightarrow \infty} T_k^{1/k} < \infty$ , and parameters  $\gamma > 0$  and  $\zeta > 1$ . We use here  
1392 the Random Attachment Model representation of the HBP process discussed  
1393 in Sect. 8.2.3. The self-similarity of the HBP process (Sect. 8) suggests that the  
1394 limit space  $\Upsilon_\infty$  does not change its statistical properties after rescaling, which  
1395 corresponds here to the Horton pruning. Let  $\mathbf{d}$  denote its fractal dimension.  
1396 That the limit space includes at least the root branch  $\Upsilon_1$  suggests  $\mathbf{d} \geq 1$ .  
1397 Assume that  $\mathbf{d} > 1$ . Then, denoting the mean  $\mathbf{d}$ -dimensional volume of  $\Upsilon_\infty$  by  
1398  $\mathbf{vol}$ , we have

$$\mathbf{vol} = \sum_{k=1}^{\infty} t_k \frac{\mathbf{vol}}{\zeta^{\mathbf{d}k}}. \quad (83)$$

1399 This equation is obtained by splitting a tree  $\Upsilon_\infty$  into the subtrees attached  
1400 to its highest-order branch  $\Upsilon_1$ . There is an average of  $t_1 = T_1 + 2$  subtrees



**Fig. 11** Fractal dimension  $d = \max\{1, \frac{\log R_B}{\log R_L}\}$  (panel a) and Hack's exponent  $h = d^{-1}$  (panel b) of self-similar HBP tree in the limit of infinite size as a function of the Horton exponents  $R_B$  and  $R_L$ . Selected levels of  $d$  and  $h$  are shown by marked black lines. Green line corresponds to the critical Tokunaga process of Sect. 9, for which  $R_B = 2R_L$ . Blue circles depict the pairs  $(R_B, R_L)$  estimated in nine real river basins by Tarboton et al. (1988), see also (Rodriguez-Iturbe and Rinaldo, 2001, Table 2.1).

1401 distributed as  $\mathcal{T}_\infty$  scaled by  $\zeta^{-1}$ . In general, for each  $k$ , there will be an average  
 1402 of  $t_k$  subtrees distributed as  $\mathcal{T}_\infty$  scaled by  $\zeta^{-k}$ . Scaling the lengths by  $\zeta^{-k}$  in  
 1403 the  $d$ -dimensional space results in scaling the volume by  $\zeta^{-dk}$ . The **vol** term  
 1404 in (83) can be cancelled out, yielding

$$\hat{t}(\zeta^{-d}) = 0, \quad (84)$$

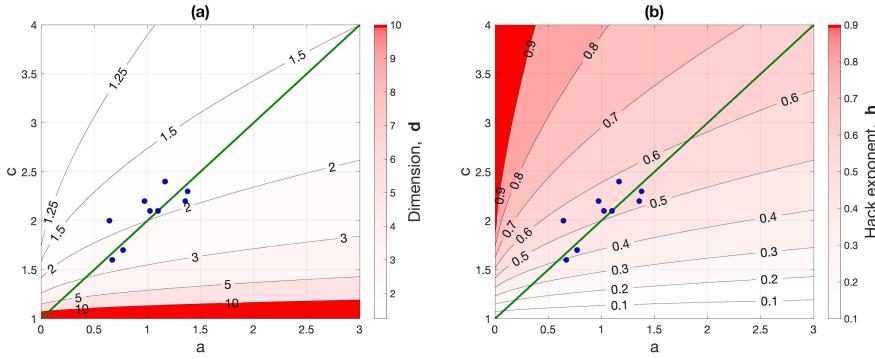
1405 and hence,  $\zeta^{-d} = w_0 = R_B^{-1}$ . Finally, we find:

$$d = \max\{1, d_0\}, \quad d_0 = -\frac{\log w_0}{\log \zeta} = \frac{\log R_B}{\log \zeta} = \frac{\log R_B}{\log R_L}. \quad (85)$$

1406 This expression has been first obtained for river networks by La Barbera  
 1407 and Rosso (1989). Figure 11(a) illustrates the fractal dimension of self-similar  
 1408 HBP trees for  $R_B \in [2, 6]$  and  $R_L \in [1, 3]$ , which are the ranges suitable for the  
 1409 studies of the observed river networks (Tarboton et al., 1988; Rodriguez-Iturbe  
 1410 and Rinaldo, 2001). The figure also shows the values  $(R_B, R_L)$  estimated for  
 1411 nine river basins by Tarboton et al. (1988). We notice the tendency of the  
 1412 estimated Horton exponents to cluster around the critical Tokunaga model for  
 1413 which  $R_B = 2R_L$  (green line). Figure 12(a) illustrates the fractal dimension for  
 1414 the two-parameter Tokunaga model (Appendix E) with parameters  $(a, c)$ . Here  
 1415  $d = \frac{\log R_B}{\log c}$  and  $R_B$  is given by (120). To add to this plot the values from the  
 1416 observed river basins, for which the pairs  $(R_B, R_L)$  are known, we use  $c = R_L$   
 1417 and find  $a$  by solving (120):

$$a = R_B - c - 2 + \frac{2c}{R_B}. \quad (86)$$

1418 Again, we see a tendency for the real basins to cluster around the critical  
 1419 Tokunaga model (green line) for which  $a = c - 1$ . Figure 13 shows the value



**Fig. 12** Fractal dimension  $d = \max\{1, \frac{\log R_B}{\log c}\}$  (panel a) and Hack's exponent  $h = d^{-1}$  (panel b) in the two-parameter Tokunaga model of Appendix E with parameters  $(a, c)$ . Here  $R_B = R_B(a, c)$  according to (120). Selected levels of  $d$  and  $h$  are shown by marked black lines. Green line corresponds to the critical Tokunaga process of Sect. 9, for which  $a = c - 1$  and  $R_B = 2c$ . Blue circles depict the pairs  $(a, c)$  that correspond to the values  $(R_B, R_L)$  estimated in nine real river basins by Tarboton et al. (1988), see also (Rodriguez-Iturbe and Rinaldo, 2001, Table 2.1).

of the Horton exponent  $R_B = R_B(a, c)$  and its reciprocal value  $w_0 = w_0(a, c)$  according to (120).

Recalling that  $R_B \geq 2$  and  $\zeta > 1$  we find that the dimension  $d$  can take any value  $d \geq 1$ . There exists an infinite collection of self-similar HBPs with a given value of  $d$ , since there are infinitely many ways to select a Tokunaga sequence  $T_k$  with a given  $w_0$ . Recalling (70) we find the dimension of a critical Tokunaga tree with parameter  $c > 1$  illustrated in Fig. 14(a):

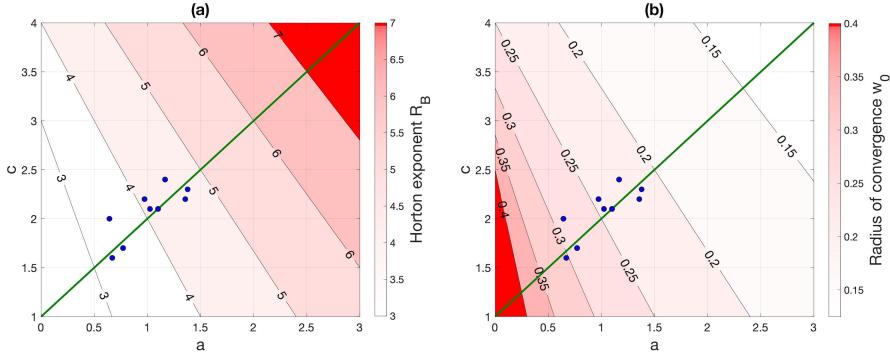
$$d = \frac{\log 2c}{\log c} = 1 + \log_c 2. \quad (87)$$

For the limit space  $\mathcal{T}_\infty$  to be embedded into a plane, one need to ensure that  $d \leq 2$ , which is equivalent to  $w_0 \geq \zeta^{-2}$  or  $R_B \leq \zeta^2$ , and in the family of critical Tokunaga trees to  $c \geq 2$ . We also observe that the condition  $w_0 = \zeta^{-2}$  corresponds to the space-filling tree with  $d = 2$  (see Newman et al. (1997)) and in the critical Tokunaga family this corresponds to  $c = 2$  (see Kovchegov and Zaliapin (2018)), which is the critical binary Galton-Watson tree.

## 11 Hack's Law in Self-Similar HBP Trees

One of the fundamental scaling laws of hydrology is the Hack's law (Hack, 1957; Mesa and Gupta, 1987; Rigon et al., 1996; Rodriguez-Iturbe and Rinaldo, 2001) that relates the lengths  $L$  of the longest stream in a river basin to the basin contributing area  $A$ :

$$L \propto A^h, \quad h \approx 0.6. \quad (88)$$



**Fig. 13** Horton exponent  $R_B$  and the radius of convergence  $w_0 = R_B^{-1}$  for  $M(z)$  in the two-parameter Tokunaga model of Appendix E with parameters  $(a, c)$ . Other notations are the same as in Fig. 12.

If  $T$  is the tree representing the stream network, then the length of the longest stream is the height of the tree  $T$ , denoted by  $\text{HEIGHT}(T)$  (Pitman, 2006; Kovchegov and Zaliapin, 2020). This section establishes the mean Hack's law in self-similar HBP trees.

Consider a tree  $T$  generated by a self-similar HBP with a Tokunaga sequence  $\{T_k\}$  satisfying  $\limsup_{k \rightarrow \infty} T_k^{1/k} < \infty$ , and parameters  $\gamma > 0$  and  $\zeta > 1$ . Let

$$\Lambda_k = \mathbb{E} [\text{HEIGHT}(T) \mid \text{ord}(T) = k] \quad (89)$$

that represents the mean length of *the longest river stream* in a basin with the Horton-Strahler order  $k$ . Notice that, since  $\text{HEIGHT}(\mathcal{R}(T)) \leq \text{HEIGHT}(T)$ ,

$$\zeta \Lambda_{k-1} = \mathbb{E} [\text{HEIGHT}(\mathcal{R}(T)) \mid \text{ord}(T) = k] \leq \mathbb{E} [\text{HEIGHT}(T) \mid \text{ord}(T) = k] = \Lambda_k. \quad (90)$$

Hence, since  $\Lambda_1 = \gamma^{-1}$ , we have  $\Lambda_k \geq \gamma^{-1} \zeta^{k-1}$ . Next, let

$$Y_1, Y_2, \dots, Y_{N_1[T]}$$

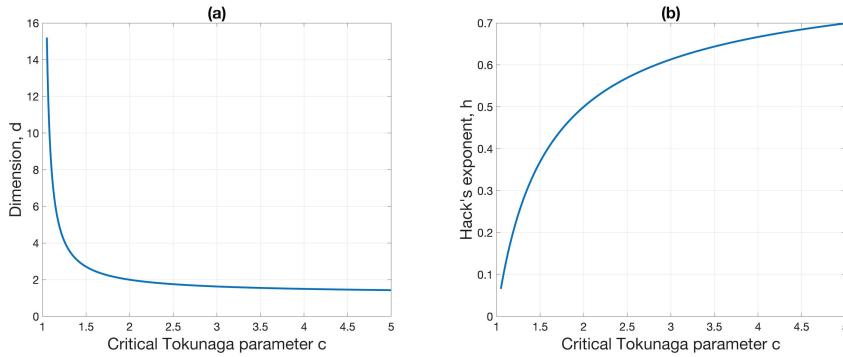
denote the leaf lengths in the tree  $T$ . Then, since

$$\text{HEIGHT}(T) \leq \text{HEIGHT}(\mathcal{R}(T)) + \max_{j=1, \dots, N_1[T]} Y_j,$$

we have,

$$\begin{aligned} \Lambda_k &\leq \mathbb{E} [\text{HEIGHT}(\mathcal{R}(T)) \mid \text{ord}(T) = k] + \mathbb{E} \left[ \max_{j=1, \dots, N_1[T]} Y_j \mid \text{ord}(T) = k \right] \\ &= \zeta \Lambda_{k-1} + \gamma^{-1} \mathbb{E} \left[ \sum_{j=1}^{N_1[T]} \frac{1}{j} \mid \text{ord}(T) = k \right] \end{aligned} \quad (91)$$

$$\begin{aligned} &\leq \zeta \Lambda_{k-1} + \gamma^{-1} \mathbb{E} [1 + \log(N_1[T])] \mid \text{ord}(T) = k \\ &\leq \zeta \Lambda_{k-1} + \gamma^{-1} + \gamma^{-1} \log(\mathbb{E}[N_1[T]] \mid \text{ord}(T) = k) \end{aligned} \quad (92)$$



**Fig. 14** Fractal dimension  $d = \frac{\log 2c}{\log c}$  (panel a) and Hack's exponent  $h = d^{-1}$  (panel b) in the critical Tokunaga model with parameter  $c > 1$ .

1448 by Wald's equation, the Coupon Collector Problem, and finally, the Jensen's  
 1449 inequality (Bhattacharya and Waymire, 2007). Recall the geometric Horton's  
 1450 law (43) for the leaf count in a self-similar process

$$\mathcal{N}_1[k] = M_k = M R_B^k + o(R_B^k).$$

1451 Hence, equations (90) and (91) imply

$$0 \leq \Lambda_k - \zeta \Lambda_{k-1} \leq \gamma^{-1} k \log R_B + \beta$$

1452 for some constant  $\beta$ , and

$$0 \leq \frac{\Lambda_k}{\Lambda_{k-1}} - \zeta \leq \gamma^{-1} \frac{k \log R_B + \beta}{\Lambda_{k-1}} \leq \frac{k \log R_B + \beta}{\zeta^{k-2}} \rightarrow 0 \text{ as } k \rightarrow \infty. \quad (93)$$

1453 Accordingly,

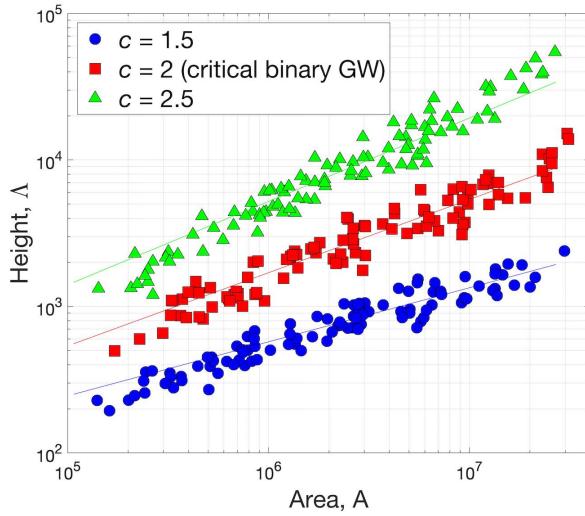
$$\log \Lambda_k = \sum_{j=2}^k \log \left( \frac{\Lambda_k}{\Lambda_{k-1}} \right) + \log \Lambda_1 = (k-1) \log \zeta + \sum_{j=2}^k \log(1+\mathcal{E}_j) - \log \gamma, \quad (94)$$

1454 where  $0 \leq \mathcal{E}_j \leq (k \log R_B + \beta) \zeta^{1-k}$ , and therefore,  $\sum_{j=2}^{\infty} \log(1+\mathcal{E}_j)$  converges to  
 1455 a constant. We therefore conclude that the geometric Horton's law holds for  
 1456  $\Lambda_k$  with Horton exponent  $R_A = R_L = \zeta$ :

$$\Lambda_k \sim \text{Const.} \times \zeta^k. \quad (95)$$

1457 This and the geometric Horton's law for the mean branch magnitudes  $M_k$   
 1458 implies the *Hack's law* for self-similar HBP:

$$\Lambda_k \sim \text{Const.} \times (M_k)^h, \quad \text{where } h = \frac{\log \zeta}{\log R_B}. \quad (96)$$



**Fig. 15** Hack's law in critical Tokunaga trees with parameters  $c = 1.5, 2, 2.5$  and  $\gamma = 1$ . Symbols correspond to individual simulated trees. Lines show theoretical slopes  $\mathbf{h} = \log(c)/\log(2c)$ . The differential contributing area  $\nu$  of an edge is calculated via the edge length  $\xi$  as  $\nu = \xi^2$ . For each  $(c, \gamma)$  we simulated 100 independent trees of different orders:  $11 \leq K \leq 14$  for  $c = 1.5$ ,  $9 \leq K \leq 11$  for  $c = 2$ , and  $8 \leq K \leq 10$  for  $c = 2.5$ . In each simulated tree, the area  $A_{(i)}$  and height  $\Lambda_{(i)}$  are reported for a random vertex from the stream of the highest order.

<sup>1459</sup> An asymptotic equivalence  $A_k \sim \text{Const.} \times M_k$  would imply the mean Hack's  
<sup>1460</sup> law in its classical form, relating the longest river channel to the basin's area:

$$\Lambda_k \sim \text{Const.} \times (A_k)^{\mathbf{h}}, \quad \text{where } \mathbf{h} = \frac{\log \zeta}{\log R_B}. \quad (97)$$

<sup>1461</sup> Our analysis in Sect. 7.2 asserts that such equivalence holds as soon as  $r_D >$   
<sup>1462</sup>  $r_M = w_0$ . This is so, for example, in the critical Tokunaga process of Sect. 9,  
<sup>1463</sup> or more generally under the hydrologic Assumption 3. The reciprocity of the  
<sup>1464</sup> Hack's law exponent  $\mathbf{h}$  and the fractal dimension  $\mathbf{d}$  has been heuristically  
<sup>1465</sup> established by Peckham (1995), La Barbera and Rosso (1989). Figures 11(b)  
<sup>1466</sup> and 12(b) show the exponent  $\mathbf{h}$  in the self-similar HBP as a function of  $(R_B, R_L)$   
<sup>1467</sup> and in the two-parameter Tokunaga model as a function of  $(a, c)$ , respectively.  
<sup>1468</sup> The nine observed river basins from Tarboton et al. (1988) (blue dots) have  
<sup>1469</sup> the Hack's exponent within the range  $0.39 < \mathbf{h} < 0.6$ .  
<sup>1470</sup> The Hack's law for the critical Tokunaga processes with parameter  $c > 1$  takes  
<sup>1471</sup> the form

$$\Lambda_k \sim \text{Const.} \times (A_k)^{\mathbf{h}}, \quad \text{where } \mathbf{h} = \frac{\log c}{\log(2c)} = \frac{1}{1 + \log_c 2}. \quad (98)$$

<sup>1472</sup> The value of  $\mathbf{h} = \mathbf{h}(c)$  is illustrated in Fig. 14(b) for  $1 < c < 5$ . Here, for  
<sup>1473</sup>  $c = 2.5$ ,  $\mathbf{h} \approx 0.57$ .... Figure 15 illustrates the Hack's law in simulated critical

<sup>1474</sup> Tokunaga trees with  $c = 1.5, 2, 2.5$  by showing the scatter between the tree  
<sup>1475</sup> contributing area  $A$  and height  $A$ .

## <sup>1476</sup> 12 Scaling Laws in Self-similar Trees

<sup>1477</sup> We have discussed in Introduction that Horton's laws imply a variety of power  
<sup>1478</sup> laws for the frequencies of edge attributes in a large tree and power-law relations  
<sup>1479</sup> between different attributes. This section illustrates these general observations  
<sup>1480</sup> with specific selected examples. We make here the hydrologic assumption of  
<sup>1481</sup> equality of the mean edge length  $\ell_k$  for all  $k$  and equality of the differential  
<sup>1482</sup> edge contributing areas  $\alpha_k$  for all  $k$ . We also assume a geometric Horton's  
<sup>1483</sup> law for  $T_k$  with  $r_T < 1$ . These assumptions are conventionally accepted in the  
<sup>1484</sup> hydrologic literature and are justified by field observations (Rodriguez-Iturbe  
<sup>1485</sup> and Rinaldo, 2001). The assumptions can be relaxed (with more technical  
<sup>1486</sup> work) if needed.

<sup>1487</sup> We observe that by Proposition 9 the areas  $A_k$  in this case satisfy geometric  
<sup>1488</sup> Horton's law with Horton exponent  $R_A$  such that  $R_A = R_B = R_M$ . Moreover,  
<sup>1489</sup> by Propositions 7,8 the combinatorial branch lengths  $S_k$  and metric branch  
<sup>1490</sup> lengths  $L_k$  also satisfy geometric Horton's laws with Horton exponents  $R_S$  and  
<sup>1491</sup>  $R_L$  such that  $R_S = R_L < R_B = R_M$ .

### <sup>1492</sup> 12.1 Power Laws for Edge Attributes

<sup>1493</sup> Recall (Sect. 3.5) that the geometric Horton's laws for branch counts  $N_k[T]$ ,  
<sup>1494</sup> edge counts  $S_k$ , and an arbitrary branch attribute  $Z_k$  imply a power-law  
<sup>1495</sup> frequency distribution for the edge attribute  $Z_{(i)}$  with power index given by  
<sup>1496</sup> (13). Section 7 establishes a variety of Horton's laws for a self-similar tree,  
<sup>1497</sup> including those for branch counts  $N_k[T]$  and edge counts  $S_k$ . This means that  
<sup>1498</sup> any attribute that satisfies the geometric Horton's law is expected to have  
<sup>1499</sup> a power-law frequency distribution when examined on individual edges. The  
<sup>1500</sup> well-studied hydrological examples are power law frequency statistics for link  
<sup>1501</sup> contributing areas (14) and distance to the divide (15); see Tarboton et al.  
<sup>1502</sup> (1989), Rodriguez-Iturbe et al. (1992), and Maritan et al. (1996).

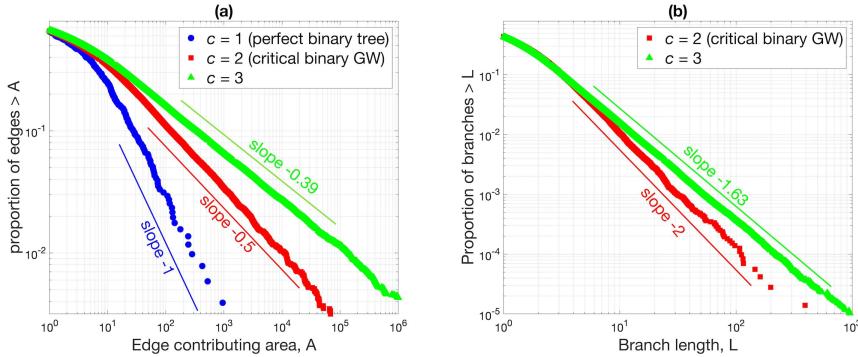
<sup>1503</sup> Consider a critical Tokunaga HBP with parameters  $\gamma = 1$  and  $c > 1$   
<sup>1504</sup> (Sect. 9). According to (70) we have  $R_B = R_A = 2c$  and  $R_S = c$ , hence  
<sup>1505</sup> establishing

$$\#\{i : A_{(i)} \geq x\} \propto x^{-\beta_A} \quad (99)$$

<sup>1506</sup> with

$$\beta_A = \frac{\log R_B - \log R_S}{\log R_A} = \frac{\log(2c) - \log c}{\log(2c)} = 1 - \frac{\log c}{\log(2c)} = 1 - h. \quad (100)$$

<sup>1507</sup> Here,  $c = 2.3$  corresponds to  $\beta_A \approx 0.45$  reported in analyses of river networks  
<sup>1508</sup> (Rodriguez-Iturbe et al., 1992; Rodriguez-Iturbe and Rinaldo, 2001). The last  
<sup>1509</sup> equality,  $\beta_A = 1 - h$ , holds in a general self-similar HBP tree under the



**Fig. 16** Power laws in critical Tokunaga trees with parameters  $c = 1, 2, 3$  and  $\gamma = 1$ . Symbols refer to the empirical counts. Lines show theoretical slopes. (a) Power law for the exceedance frequencies of edge contributing areas  $A_{(i)}$ . The lines show the theoretical slope  $\beta_A = -(1 - \mathbf{h}(c))$ , with convention  $\mathbf{h}(1) = 0$ . (b) Power law for the exceedance frequencies of branch lengths  $A_{[i]}$ . The lines show the theoretical slope  $-\mathbf{d}(c)$ .

assumptions listed in the beginning of Sect. 12. This relation is well-known in the analysis of river basins (Rodriguez-Iturbe and Rinaldo, 2001, Eq. (2.215)).

Figure 16(a) shows empirical exceedance frequencies for edge areas  $A_{(i)}$  calculated in three critical Tokunaga trees with  $c = 1, 2, 3$  and  $\gamma = 1$ . The distributions have power-law tails (seen as linear segments in the double logarithmic plot) with power indices given by  $\beta_A(c) = 1 - \mathbf{h}(c)$ , according to (100).

Similarly, the geometric Horton's law (95) for tree heights  $\Lambda_k$  implies a power-law distribution of the distances  $\Lambda_{(i)}$  from link  $i$  to the most distant source along the tree (i.e., heights of the edges using the graph-theoretic terminology):

$$\#\{i : \Lambda_{(i)} \geq x\} \propto x^{-\beta_A} \quad (101)$$

with

$$\beta_A = \frac{\log R_B - \log R_S}{\log R_A} = \mathbf{d} - 1 = \frac{\beta_A}{\mathbf{h}}. \quad (102)$$

This distribution can be alternatively derived by writing the Hack's law (97) on individual edges:

$$\#\{i : \Lambda_{(i)} \geq x\} = \#\{i : A_{(i)}^{\mathbf{h}} \geq \text{Const.} \times x\} \propto x^{-\beta_A/\mathbf{h}}. \quad (103)$$

The relation (102) is well-documented in the analysis of natural river basins as discussed by (Rodriguez-Iturbe and Rinaldo, 2001, Sect. 2.9.3, Eq. (2.185)). In a critical Tokunaga tree with parameter  $c$  we have

$$\beta_A = \frac{\log(2c) - \log c}{\log c} = \log_c 2. \quad (104)$$

Here for  $c = 2.3$  we have  $\beta_A \approx 0.83$ .

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 1528 12.2 Power Laws for Branch Attributes

1529 Tree self-similarity, and associated Horton's laws, also imply power laws for  
 1530 attributes calculated for random branches. The argument is very similar to  
 1531 that used to establish power laws for the edge attributes. We consider here  
 1532 the branch lengths  $L_{[i]}$ , where the lower bracketed index  $[i]$  indicates that we  
 1533 calculate the length of a uniformly randomly selected branch (and distinguishes  
 1534 this attribute from the mean length  $L_K$  of an order- $K$  branch). In a limit of a  
 1535 large tree,

$$\#\{i : L_{[i]} \geq R_L^K\} \propto \sum_{j=K}^{\infty} N_K = \sum_{j=K}^{\infty} R_B^{-j} \propto R_B^{-K}.$$

1536 This is a punctuated (by discrete order) version of a general power law relation  
 1537

$$\#\{i : L_{[i]} \geq z\} \propto z^{-\frac{\log R_B}{\log R_L}} = z^{-d}. \quad (105)$$

1538 Appendix H shows a rigorous derivation of the relative frequencies of branch  
 1539 lengths  $L_{[i]}$  in a critical Tokunaga tree, which leads to essentially the same result.  
 1540 Figure 16(b) shows the exceedance frequency for branch lengths  $L_{[i]}$  calculated  
 1541 in two critical Tokunaga trees with  $c = 2, 3$  and  $\gamma = 1$ . The distributions have  
 1542 power-law tails (expressed as linear segments in the double logarithmic plot)  
 1543 with power indices given by  $-d(c)$ .

 1544 **13 Discussion**

1545 A solid body of observational, modeling, and theoretical studies ascribe Hor-  
 1546 ton's laws, power-law distributions of tree attributes, and power-law relations  
 1547 between attributes to the self-similar structure of a tree that represents the  
 1548 examined system (Gupta and Waymire, 1989; Peckham, 1995; Gupta et al.,  
 1549 1996; Tarboton, 1996; Gupta and Waymire, 1998; Cieplak et al., 1998; Peckham  
 1550 and Gupta, 1999; Turcotte, 1997; Dodds and Rothman, 2000; Pelletier and  
 1551 Turcotte, 2000; Veitzer and Gupta, 2000; Gupta et al., 2007; Rodriguez-Iturbe  
 1552 and Rinaldo, 2001; Mesa, 2018). Here we review a recently formulated theory  
 1553 of random self-similar trees (Kovchegov and Zaliapin, 2020; Kovchegov et al.,  
 1554 2021) that suggests a rigorous treatment of the emergence of Horton's laws  
 1555 and related scalings in river networks and other dendritic systems.

1556 Self-similarity is defined here (Sect. 4) as invariance of a coordinated tree  
 1557 distribution with respect to the operation of Horton pruning (cutting the source  
 1558 streams); this definition is justified by the empirical and modeling evidence of  
 1559 the past decades. Horton's laws are rigorously defined as limit statements about  
 1560 random or mean values of the examined branch attributes (Sects. 3.7, 5, 9.2).  
 1561 We show that self-similarity guarantees the (strongest) geometric Horton's laws  
 1562 for mean branch numbers and mean magnitudes (Theorem 1 and Corollary 1).  
 1563 Horton's laws of different strengths for multiple other mean attributes follow  
 1564 under additional natural assumptions (Sect. 7). Each Horton's law for a mean

attribute (e.g., mean branch number  $\mathcal{N}_1[K]$ ) implies the respective Horton's law for its random counterpart (e.g., random branch number  $N_1[T]$ ). We have examined several commonly studied branch attributes (Table 1) whose scaling laws are well documented in the literature and have shown that the proposed self-similar model closely reproduces the scalings and exponents reported in observational studies (Table 2).

Our definition of tree self-similarity unifies several alternative definitions that have been introduced in studies of dendritic systems. Burd et al. (2000) define self-similarity in Galton-Watson trees as the Horton prune-invariance. This is a special case of our definition since the Galton-Watson trees are coordinated (Kovchegov and Zaliapin, 2020). Peckham (1995) and Newman et al. (1997) define self-similarity as Toeplitz property for the Tokunaga coefficients. This property, which only considers the Tokunaga coefficients and not the entire tree distribution, follows from and is weaker than our definition (Sect. 4.4). Moreover, Kovchegov and Zaliapin (2020) showed that the Toeplitz property alone, without coordination (Sect. 4.2), allows for a multitude of obscure measures that are hardly useful in practice. Gupta and Waymire (1989) and Peckham and Gupta (1999) suggested a concept of *statistical self-similarity* that requires a random stream attribute  $Z$  to have distribution that scales with order. It can be shown (Kovchegov and Zaliapin, 2020, Sect. 7) that (i) statistical self-similarity for some attributes (e.g., for any discrete attribute) may only hold asymptotically, and (ii) multiple attributes, including stream length, magnitude, and total basin length, are statistically self-similar in the limit of an infinitely large basin that is self-similar according to our definition.

The results reviewed herein contribute to a long-standing debate on the "inevitability" of Horton's laws in river networks (Shreve, 1966, 1969; Kirchner, 1993; McConnell and Gupta, 2008), and suggest that Horton prune-invariance is a useful paradigm for systems that exhibit such laws. The family of self-similar distributions is extremely rich and flexible. It includes the famous random topology model of Shreve (1966), which is equivalent to the critical binary Galton-Watson process with given progeny (Burd et al., 2000; Pitman, 2006), and closely fits the multitude of existing hydrologic observations summarized by Maritan et al. (1996), Turcotte (1997), Dodds and Rothman (1999), Dodds and Rothman (2000), Rodriguez-Iturbe and Rinaldo (2001), and Gupta et al. (2007); see Table 2. The self-similar family extends way beyond the hydrological constraints, allowing one to study self-similar trees with edge lengths that depend on the position within the hierarchy, having arbitrary fractal dimension  $d \in (1, \infty)$ , and Horton branch exponent  $R_B \in (2, \infty)$ . For instance, the HBP (Sect. 8) might be a suitable model for phylogenetic trees (Aldous, 2001; Blum and François, 2006) or dendritic structures generated by Diffusion Limited Aggregation (DLA) (Vicsek, 1984; Ossadnik, 1992; Newman et al., 1997). While this survey focuses on binary trees, the self-similarity definition and main results are readily extended to trees with multiple branching. The proposed approach emphasizes the importance of Tokunaga coefficients  $T_k$  that have been well-known in the literature (Tokunaga, 1978; Ossadnik, 1992; Peckham, 1995; Newman et al., 1997; Gabrielov et al., 1999; Turcotte, 1999;

Pelletier and Turcotte, 2000; Holliday et al., 2008; Yoder et al., 2013) and which also play a distinct role in the presented theory. Namely, each self-similar measure is characterized by an infinite sequence  $\{T_1, T_2, T_3, \dots\}$  of Tokunaga coefficients, and each such sequence corresponds to an infinite number of self-similar measures, which gives an idea of the richness of the self-similar family.

Notwithstanding this richness, the essential scalings established in the hydrological literature are closely fit by a one-parameter family of critical Tokunaga trees with  $T_k = (c - 1)c^{k-1}$  (Table 2). This empirical constraint has been known for long time (Tokunaga, 1978; Peckham, 1995; Veitzer and Gupta, 2000), with the special case  $c = 2$  corresponding to the Shreve's random topology model (Burd et al., 2000). However, only very recently a rigorous understanding has been gained of its theoretical importance. For example, this sequence is necessary to generate tree distributions that are time-invariant, critical, and having i.i.d. edge lengths (Kovchegov and Zaliapin, 2018, 2019, 2020; Kovchegov et al., 2021).

The critical Tokunaga model presents an ultimately symmetric class of trees characterized by coordination, Horton prune-invariance, criticality, time-invariance, and identically distributed link lengths (and hence local contributing areas). Despite these multiple constraints, this class is surprisingly rich, extending from perfect binary trees ( $c = 1$ ) to the famous Shreve's random topology model ( $c = 2$ ) to the structures reminiscent of the observed river networks ( $c \approx 2.3$ ) and beyond. While offering a convenient theoretical and modeling paradigm, the critical Tokunaga model is merely a subclass of a much broader family of self-similar trees that might better accommodate for various problem-specific data features. An applied study can use the self-similar theory to either focus on the symmetries of the critical Tokunaga family, or explore deviations from this stiff parameterization, both of which may have physical underpinnings.

This survey focuses on the results that concern the static structure of examined systems (e.g., river networks). Self-similarity of this structure might provide tangible constraints on the additional geometric attributes (e.g., channel slopes, junction angles, etc.) (Stark et al., 2009; Devauchelle et al., 2012), the dynamical processes that evolve along its static fabric (Mesa and Mifflin, 1986; Gupta et al., 1994, 1996; Menabde et al., 2001; Mantilla et al., 2006; Lashermes and Foufoula-Georgiou, 2007; Zaliapin et al., 2010; Gupta et al., 2010; Ramirez, 2012; Czuba and Foufoula-Georgiou, 2014) or control its formation and evolution (Seybold et al., 2007; Singh et al., 2015; Ranjbar et al., 2018).

The goal of this work was to review the recent results concerning tree self-similarity and present a simple model that explains a variety of scaling laws that are central for the studies of dendritic systems of diverse origin. Our prime illustration of the power of the proposed approach (Table 2, Fig. 5) uses empirical scaling laws of hydrogeomorphology that have been established and independently verified by multiple researchers since the 1940s. Accordingly, we intentionally avoided new data analyses and took all empirical constraints

from the existing literature. An original data analysis performed through the prism of the proposed modeling approach is a topic of future research.

The proposed approach to modeling dendritic systems based on random self-similar trees is subject to further testing and verification using data from diverse fields (see Sect. 1). There are several avenues for approaching such testing. One can test the foundational principle of the theory – tree self-similarity – that combines the tree coordination and Horton prune-invariance. The coordination property can be either rigorously verified (as in Galton-Watson trees) or heuristically accepted (as in river networks). The self-similarity is then tested either by checking the Horton prune-invariance property, which might be more suitable for theoretical models, or by verifying the Toeplitz property of the Tokunaga coefficients, which can be readily done for the empirical Tokunaga coefficients  $t_{i,j}$  of (19) via the ANOVA framework (Scheffe, 1999). Independently, one can check the theory's predictions. This includes (i) Horton's laws (e.g. Horton's law for branch numbers); (ii) Power law distribution of attributes calculated at individual vertices (e.g. vertex magnitudes); (iii) Power law relation between distinct attributes (e.g., Hack's law); and (iv) System's fractal dimension. Specifically, one would check if the above laws hold and if their scaling exponents are related in the way predicted by the theory (e.g., in critical Tokunaga trees the key exponents take only two distinct values). Finally, one can directly test if an observed tree (or forest of trees) can be approximated by a self-similar model. A likelihood approach to such direct statistical testing is being developed by the authors and will be presented elsewhere.

## A Proof of Proposition 1

We have, for any  $\Delta \in (0, r)$  (Ahlfors, 1953):

$$a_k = \frac{1}{2\pi i} \oint_{|z|=\Delta} \frac{f(z)dz}{z^{k+1}}. \quad (106)$$

By the Residue Theorem (Ahlfors, 1953), we obtain, for any  $\gamma \in (r, \rho)$

$$\frac{1}{2\pi i} \oint_{|z|=\gamma} \frac{f(z)dz}{z^{k+1}} = \text{Res}\left(\frac{f(z)}{z^{k+1}}; 0\right) + \text{Res}\left(\frac{f(z)}{z^{k+1}}; r\right) \quad (107)$$

$$= a_k + \text{Res}\left(\frac{f(z)}{z^{k+1}}; r\right). \quad (108)$$

Therefore,

$$a_k = \frac{1}{2\pi i} \oint_{|z|=\gamma} \frac{f(z)dz}{z^{k+1}} - \text{Res}\left(\frac{f(z)}{z^{k+1}}; r\right), \quad (109)$$

where

$$\left| \oint_{|z|=\gamma} \frac{f(z)dz}{z^{k+1}} \right| \leq \frac{\max_{|z|=\gamma} f(z)}{\gamma^k} = o(r^{-k}). \quad (110)$$

Consider  $g(z) = (z - r)f(z)$ . It is known that (Ahlfors, 1953)

$$\text{Res}(f(z); r) = g(r), \quad (111)$$

1685 and hence,

$$\text{Res} \left( \frac{f(z)}{z^{k+1}}; r \right) = \frac{g(r)}{r^{k+1}} = \frac{g(r)}{r} r^{-k}. \quad (112)$$

1686 Accordingly, we obtain

$$a_k = -\frac{g(r)}{r} r^{-k} + o(r^{-k}), \quad (113)$$

1687 which completes the proof.

## 1688 B Proof of Proposition 2

1689 First, we prove statement (i).

We begin with the root Horton's law (R). Suppose  $\lim_{k \rightarrow \infty} (b_k)^{1/k} = R > 0$ , then for a given  $\epsilon \in (0, 1/r_c) \subset (0, R)$ , there exist  $C_0, C_1 > 0$  such that

$$C_0(R - \epsilon)^k \leq b_k \leq C_1(R + \epsilon)^k \quad \text{for all } k = 0, 1, \dots$$

Then,

$$R - \epsilon \leq \liminf_{k \rightarrow \infty} \left( \sum_{j=0}^k b_{k-j} c_j \right)^{1/k} = \liminf_{k \rightarrow \infty} (a_k)^{1/k}$$

and

$$\limsup_{k \rightarrow \infty} (a_k)^{1/k} = \limsup_{k \rightarrow \infty} \left( \sum_{j=0}^k b_{k-j} c_j \right)^{1/k} \leq R + \epsilon$$

1690 since there exists  $D > 0$  such that  $c_k \leq D(1/r_c + \epsilon)^k$  for all  $k$ . Hence, since  $\epsilon$  can be taken  
1691 arbitrarily small,  $\lim_{k \rightarrow \infty} (a_k)^{1/k} = R$ .

Next, we consider the quotient Horton's law (Q). Suppose  $\lim_{k \rightarrow \infty} \frac{b_{k+1}}{b_k} = R$ . Then, by the Dominated Convergence Theorem,

$$\lim_{k \rightarrow \infty} \frac{a_k}{b_k} = \lim_{k \rightarrow \infty} \sum_{j=0}^k \frac{b_{k-j}}{b_k} c_j = \sum_{j=0}^{\infty} R^{-j} c_j < \infty$$

1692 as  $R^{-1} = r_b < r_c$ . Hence,  $\lim_{k \rightarrow \infty} \frac{a_{k+1}}{a_k} = R$ .

Finally, we consider the geometric Horton's law (G). Suppose  $\lim_{k \rightarrow \infty} b_k R^{-k} = \beta$  for some  $\beta > 0$ . Then, by the Dominated Convergence Theorem,

$$\lim_{k \rightarrow \infty} a_k R^{-k} = \lim_{k \rightarrow \infty} \sum_{j=0}^k R^{j-k} b_{k-j} R^{-j} c_j = \beta \sum_{j=0}^{\infty} R^{-j} c_j = \alpha.$$

1693 Statement (ii) follows from (i) if we write  $g(z) = f(z)\tilde{h}(z)$  with  $\tilde{h}(z) = \frac{1}{h(z)}$  analytic for  
1694  $|z| \leq r_b + \epsilon$  for some  $\epsilon > 0$ .

## 1695 C Horton's Laws for Mean Branch Attributes $S_k$ , $L_k$ , and $A_k$

1696 Recall that the mean branch length  $S_k$  is expressed via the Tokunaga coefficients as

$$S_k = 1 + \sum_{i=1}^{k-1} T_i, \quad k \geq 1, \quad (114)$$

1697 and the generating function for this sequence is given by (49). The following statement is  
1698 proven using Proposition 2.

**Proposition 7 [Asymptotic behavior of  $S_k$ ]**

- (a) If  $r_T > 1$  then  $r_S = 1$  and  $S_k \sim 1 + T(1) = \hat{t}(1) = 1 + \sum_{k=1}^{\infty} T_k$ . Accordingly, the geometric Horton's law (G) holds.
- (b) If  $r_T < 1$  then  $r_S = r_T$  and  $S_k$  have the same asymptotic as  $T_k$ . Namely, the same Horton's laws hold for  $T_k$  and  $S_k$ , with the same Horton exponent. In particular, if  $\lim_{k \rightarrow \infty} (T_{k+1}/T_k) = R_T > 1$ , then  $\lim_{k \rightarrow \infty} (S_{k+1}/S_k) = R_T$ .
- (c) If  $r_T = 1$  then  $r_S = 1$ . In this case the sequence  $S_k$  increases to infinity at a sub-exponential rate. The geometric Horton's law (G) does not hold. The quotient Horton's law (Q) and root Horton's law (R) may or may not hold depending on the form of  $T_k$ . See examples below.

**Example 2 (Relation between  $T_k$  and  $S_k$ )**

- (a) **Finite Tokunaga sequence.** Suppose  $T_1 > 0$  and  $T_k = 0$  for any  $k > 1$ . Then  $T(z) = T_1 z$  and  $r_T = \infty$  as in Prop. 7(a). According to (49),  $S(z) = z(T_1 z + 1)(1 - z)^{-1}$  and  $r_S = 1$ . In this case  $S_1 = 1$  and  $S_k = T_1 + 1$  for any  $k > 1$ .

- (b) **Harmonic Tokunaga sequence.** Suppose  $T_k = 1/k$  for any  $k$ . Then

$$T(z) = \sum_{k=1}^{\infty} k^{-1} z^k = -\ln(1 - z)$$

and  $r_T = 1$  as in Prop. 7(c). According to (49),  $S(z) = z(1 - \ln(1 - z))(1 - z)^{-1}$  and  $r_S = 1$ . In this case  $S_k = 1 + \sum_{i=1}^{k-1} k^{-1} \sim \ln(k)$  as  $k \rightarrow \infty$ . The quotient Horton's law (Q) holds with  $R_S = 1$ .

- (c) **Constant Tokunaga sequence.** Suppose  $T_k = 1$  for any  $k$ . Then  $T(z) = \sum_{k=1}^{\infty} z^k = z(1 - z)^{-1}$  and  $r_T = 1$  as in Prop. 7(c). According to (49),  $S(z) = z(1 - z)^{-2}$  and  $r_S = 1$ . In this case  $S_k = k$ . The quotient Horton's law (Q) holds with  $R_S = 1$ .

- (d) **Linear Tokunaga sequence.** Suppose  $T_k = k$  for any  $k$ . Then  $T(z) = \sum_{k=1}^{\infty} k z^k = z(1 - z)^{-2}$  and  $r_T = 1$  as in Prop. 7(c). According to (49),  $S(z) = z(z^2 - z + 1)(1 - z)^{-3}$  and  $r_S = 1$ . In this case  $S_k = \frac{k(k-1)}{2}$ . The quotient Horton's law (Q) holds with  $R_S = 1$ .

- (e) **Tokunaga sequence that does not satisfy Horton's law.** Let  $T_k = 2^{-k}$  if  $k$  is not a full square, and  $T_k = 2^{\sqrt{k}}$  otherwise. Here  $r_T = 1$  and even the root Horton's law (R) does not hold for  $T_k$ . We have  $S_{j^2} \sim 2^j$  and so  $T_{j^2}/S_{j^2} \rightarrow 1$ . At the same time,  $T_{j^2-1}/S_{j^2-1} \rightarrow 0$ . This means that the limit of  $T_k/S_k$  does not exist, and so the limit of  $S_{k+1}/S_k = 1 + T_k/S_k$ , which is equivalent to quotient Horton's law (Q), does not exist. The root Horton's law (R) also does not hold since  $(2^{\sqrt{k}})^{1/k} \rightarrow 1$ , while  $(2^{-k})^{1/k} = 2^{-1}$ .

- (f) **Geometric Tokunaga sequence.** Suppose  $T_k = (c - 1)c^{k-1}$  for any  $k$  with some  $c > 1$ . Then  $T(z) = (c - 1) \sum_{k=1}^{\infty} c^{k-1} z^k = (c - 1)z(1 - cz)^{-1}$  and  $r_T = c^{-1}$  as in Prop. 7(b). According to (49),  $S(z) = z(1 - cz)^{-1}$  and  $r_S = c^{-1}$ . In this case  $S_k = c^{k-1}$ .

The generating function of the mean branch lengths  $L_k$  is given by

$$L(z) = \sum_{k=1}^{\infty} L_k z^k = \sum_{k=1}^{\infty} S_k \ell_k z^k.$$

If we assume root Horton's law (R) for  $\ell_k$  with Horton exponent  $\lambda^{-1}$ , then as  $L_k = S_k \ell_k$ , we have  $r_L = \lambda^{-1} r_S$  whenever  $\lambda > 0$ , and  $r_L = \infty$  if  $\lambda = 0$ .

**Proposition 8 [Asymptotic behavior of  $L_k$ ]**

Suppose the root Horton's law holds for  $\ell_k$  with Horton exponent  $\lambda^{-1}$ ,  $\lambda \in (0, \infty)$ .

- (a) If  $r_T > 1$  then  $r_L = \lambda^{-1}$ . The same Horton's laws hold for  $\ell_k$  and  $L_k$ .
- (b) If  $r_T < 1$  then  $r_L = \lambda^{-1} r_T$ . If some Horton's laws hold for  $T_k$  and  $\ell_k$ , then the weakest of those holds for  $L_k$ .

- (c) If  $r_T = 1$  then  $r_L = \lambda^{-1}$ . The geometric Horton's law (G), the quotient Horton's law (Q) and the root Horton's law (R) may or may not hold for  $L_k$ .

**Example 3 (Relation between  $T_k$  and  $L_k$ )**

If  $\ell_k$  scales geometrically with  $k$ , i.e., there is a scalar  $\lambda > 0$  such that  $\ell_k = \ell_1 \lambda^{k-1}$ , then

$$L(z) = \ell_1 \lambda^{-1} \sum_{k=1}^{\infty} S_k \lambda^k z^k = \ell_1 \lambda^{-1} S(\lambda z).$$

Here, if  $r_T > 1$ , and therefore  $T(z)$  converges to a finite value  $T(1)$  at  $z = 1$ , we have

$$L_k = S_k \ell_k \sim \ell_1 \lambda^{k-1} (1 + T(1)) = \ell_1 \lambda^{k-1} \hat{t}(1).$$

In this case, the geometric Horton's law (G) holds with Horton exponent  $\lambda^{-1}$ . If  $T_k = (c-1)c^{k-1}$  for some  $c > 1$ , then  $r_T = c^{-1} < 1$  and  $S_k = c^{k-1}$ . Hence,  $L_k = \ell_1 (\lambda c)^{k-1}$  and the geometric Horton's law (G) holds for  $L_k$  with Horton exponent  $(\lambda c)^{-1}$ .

The generating function  $A(z)$  for the contributing areas is given by

$$A(z) = \frac{\sum_{k=1}^{\infty} \alpha_k S_k z^k}{1 - 2z - T(z)} = M(z) \left( \sum_{k=1}^{\infty} \alpha_k S_k z^{k-1} \right) = -\frac{D(z)}{\hat{t}(z)}. \quad (115)$$

**Proposition 9 [Asymptotic behavior of  $A_k$ ]**

Suppose the root Horton's law holds for  $\alpha_k$  with Horton exponent  $\alpha^{-1}$ ,  $\alpha \in (0, \infty)$ . Recall that  $\delta_k = \alpha_k S_k$ .

- (a) If  $r_T > 1$  then  $r_D = \alpha^{-1}$ . The same Horton's laws hold for  $\alpha_k$  and  $\delta_k$ . If  $\alpha^{-1} < w_0$  then  $r_A = \alpha^{-1}$  and  $A_k \sim \text{Const.} \times \alpha_k$ . If  $\alpha^{-1} > w_0$  then  $r_A = w_0$  and  $A_k \sim \text{Const.} \times M_k$ .
- (b) If  $r_T < 1$  then  $r_D = \alpha^{-1} r_T$ . If  $r_D < w_0$  then  $r_A = r_D$  and  $A_k \sim \text{Const.} \times \delta_k$ . If  $r_D < w_0$  and some of the Horton's laws hold for  $T_k$  and  $\alpha_k$ , the weakest of those holds for  $\delta_k$  and  $A_k$ . If  $r_D > w_0$  then  $r_A = w_0$  and  $A_k \sim \text{Const.} \times M_k$ .
- (c) If  $r_T = 1$  then  $r_D = \alpha^{-1}$ . The geometric Horton's law (G), the quotient Horton's law (Q) and the root Horton's law (R) may or may not hold for  $\delta_k$ . If  $\alpha^{-1} < w_0$  then  $r_A = \alpha^{-1}$  and  $A_k \sim \text{Const.} \times \delta_k$ . If  $\alpha^{-1} > w_0$  then  $r_A = w_0$  and  $A_k \sim \text{Const.} \times M_k$ .

## D Relation Between Edge Lengths and Differential Contributing Areas

Consider a random variable  $\xi_k$  representing the length of a randomly selected edge of order  $k \geq 1$ . Suppose there exists  $C \geq 0$  such that for all  $k \geq 1$ , the expectation  $E[\xi_k] = \ell_k$  and the standard deviation  $SD(\xi_k)$  satisfy

$$\frac{SD(\xi_k)}{\ell_k} \leq C.$$

Let  $\nu_k$  be a random variable representing the differential contributing area of the edge, and suppose there exist scalars  $b > 0$  and  $\sigma > 1$  such that  $\delta_k = b \xi_k^\sigma$  with  $\sigma$  and  $b$  being the same for all orders  $k \geq 1$ . Then, by Jensen's inequality and the size biasing method (Bhattacharya and Waymire, 2007), we have

$$\begin{aligned} \ell_k^\sigma &\leq E[\xi_k^\sigma] = \ell_k E[\zeta_k^{\sigma-1}] \leq \ell_k E[\zeta_k]^{\sigma-1} \\ &= \ell_k \left( \frac{E[\xi_k^2]}{\ell_k} \right)^{\sigma-1} = \ell_k^\sigma \left( 1 + \frac{\text{Var}(\xi_k)}{\ell_k^2} \right)^{\sigma-1} \leq (1 + C^2)^{\sigma-1} \ell_k^\sigma, \end{aligned} \quad (116)$$

where  $\zeta_k$  is a random variable distributed as  $P(\zeta_k \in A) = \frac{1}{\ell_k} E[\xi_k \mathbf{1}_A(\xi_k)]$ .

Hence, the expected contributing area  $\alpha_k = b E[\xi_k^\sigma]$  of an order  $k$  link satisfies

$$b \ell_k^\sigma \leq \alpha_k \leq b (1 + C^2)^{\sigma-1} \ell_k^\sigma.$$

1766 Accordingly, the root Horton's law holds for the lengths implies that for the areas:

$$\lim_{k \rightarrow \infty} \ell_k^{1/k} = \lambda^{-1} \Rightarrow \lim_{k \rightarrow \infty} \alpha_k^{1/k} = \lambda^{-\sigma}.$$

1767 Our analysis does not establish the quotient Horton's law, since

$$(1 + C^2)^{1-\sigma} \left( \frac{\ell_{k+1}}{\ell_k} \right)^\sigma \leq \frac{\alpha_{k+1}}{\alpha_k} \leq (1 + C^2)^{\sigma-1} \left( \frac{\ell_{k+1}}{\ell_k} \right)^\sigma.$$

1768 At the same time, the gap between the low and upper bounds above can be small, hence  
1769 implying the quotient Horton's law with a practical level of accuracy. For instance, if  
1770  $\sigma = 2$  and  $C = 1$ , then that gap is  $[1/2, 2]$ . One also can add to that the observation that  
1771 increasing/decreasing sequence of lengths corresponds to the increasing/decreasing sequence  
1772 of areas, which bounds the Horton exponent by 1 from below of above.

## 1773 E Tokunaga Two-parameter Model

1774 This section discusses a slight relaxation of the critical Tokunaga process constraint  $T_k =$   
1775  $(c - 1)c^{k-1}$ . Specifically, we consider the sequence of Tokunaga coefficients that has been  
1776 introduced in river studies by Tokunaga (1966, 1978, 1984):

$$\frac{T_{k+1}}{T_k} = c \quad \text{or} \quad T_k = a c^{k-1} \quad \text{for } a, c > 0. \quad (117)$$

1777 The trees satisfying (117) are usually called *Tokunaga trees*. The Tokunaga trees have been  
1778 shown to closely approximate multiple observed branching structures beyond river networks  
1779 (Peckham, 1995; McConnell and Gupta, 2008; Zanardo et al., 2013; Dodds and Rothman,  
1780 2000; Gabrielov et al., 1999; Kovchegov and Zaliapin, 2016; Newman et al., 1997; Ossadnik,  
1781 1992; Pelletier and Turcotte, 2000; Turcotte et al., 1998). A perfect binary tree is a Tokunaga  
1782 tree with  $a = 0$  and arbitrary  $c$ . The critical binary Galton-Watson tree corresponds to  
1783  $(a, c) = (1, 2)$  (Burd et al., 2000). The critical Tokunaga process of Sect. 4.1 corresponds  
1784 to a special case  $a = c - 1$ . The geometric behavior of the Tokunaga's indices allows one  
1785 to find an explicit form of the generating function  $T(z)$  and makes the branching analysis  
1786 particularly straightforward.

If  $c = 1$ , then  $S_k = 1 + \sum_{j=1}^{k-1} a = 1 + a(k - 1)$ . If  $c \neq 1$ , then

$$\begin{aligned} S_k &= 1 + \sum_{j=1}^{k-1} T_j = 1 + \sum_{j=1}^{k-1} a c^{j-1} = 1 + a \frac{c^{k-1} - 1}{c - 1} \\ &= \begin{cases} a(c-1)^{-1}c^{k-1} + O(1), & c > 1 \\ 1 + a(1-c)^{-1} + O(c^k), & c < 1 \end{cases} \end{aligned} \quad (118)$$

1787 Next, we have

$$\hat{t}(z) = -1 + 2z + \frac{az}{1 - cz} = \frac{-1 + (a + c + 2)z - 2cz^2}{1 - cz}, \quad (119)$$

whose two real roots are

$$w_{1,0} = \frac{(a + c + 2) \pm \sqrt{(a + c + 2)^2 - 8c}}{4c},$$

1788 with  $w_1 > w_0$ . The smallest root  $w_0$  has been reported in multiple works (e.g., Peckham,  
1789 1995; McConnell and Gupta, 2008). Here we have

$$R_B = w_0^{-1} = 2c w_1 = \frac{(a + c + 2) + \sqrt{(a + c + 2)^2 - 8c}}{2}. \quad (120)$$

<sup>1790</sup> Accordingly,

$$M(z) = -\frac{z}{\hat{t}(z)} = \frac{z(1-cz)}{2c(z-w_0)(z-w_1)}, \quad (121)$$

and, by Theorem 1,

$$\begin{aligned} \mathcal{N}_{K-k+1}[K] &= M_k = \frac{1}{w_0^k} \frac{(1-cw_0)}{2c(w_1-w_0)} + o\left(\frac{1}{w_0^k}\right) \\ &= \frac{1}{w_0^k} \frac{(2-c-a) + \sqrt{(2+c+a)^2 - 8c}}{4\sqrt{(2+c+a)^2 - 8c}} + o\left(\frac{1}{w_0^k}\right), \end{aligned} \quad (122)$$

<sup>1791</sup> The exact expression is derived in Appendix F:

$$\mathcal{N}_{K-k+1}[K] = M_k = w_0^{-k+1} \left[ 1 + (1-cw_1) \sum_{i=1}^{k-1} \left( \frac{w_0}{w_1} \right)^i \right]. \quad (123)$$

<sup>1792</sup> To examine the branch contributing areas  $A_k$  one need to make additional assumptions  
<sup>1793</sup> about the process. We adopt here the hydrological constraint of Assumption 2 of Sect. 7.3:

$$\lim_{k \rightarrow \infty} \frac{\ell_{k+1}}{\ell_k} = 1. \quad (124)$$

<sup>1794</sup> Combining this with (64) gives

$$\lim_{k \rightarrow \infty} \frac{\ell_{k+1}}{\ell_k} = \zeta \lim_{k \rightarrow \infty} \frac{S_k}{S_{k+1}} = 1. \quad (125)$$

<sup>1795</sup> The geometric form (117) of the Tokunaga coefficients implies that the quotient Horton's law  
<sup>1796</sup> holds for  $T_k$ . Proposition 4 then ensures that the quotient Horton's law also holds for  $S_k$  with  
<sup>1797</sup> Horton exponent  $R_S = \max\{1, c\} \geq 1$ . Hence,  $R_S = \zeta$  and  $\zeta \geq 1$ , which implies that the  
<sup>1798</sup> mean branch length is non-decreasing with order. The field observations (Rodriguez-Iturbe  
<sup>1799</sup> and Rinaldo, 2001) strongly suggest that the stream length increases geometrically with  
<sup>1800</sup> order. This implies  $\zeta > 1$  and hence  $\zeta = c > 1$ , which also means geometric growth of  $T_k$ .

<sup>1801</sup> The edge lengths in the HBP have exponential distribution. We use Example 1 to find

$$\lim_{k \rightarrow \infty} \frac{\alpha_{k+1}}{\alpha_k} = \lim_{k \rightarrow \infty} \left( \frac{\ell_{k+1}}{\ell_k} \right)^\sigma = 1. \quad (126)$$

<sup>1802</sup> Accordingly, Assumption 3 is satisfied here with  $c > 1$ , and Proposition 6 guarantees the  
<sup>1803</sup> existence of the quotient Horton's law for  $A_k$ .

<sup>1804</sup> One can obtain a stronger result by recalling

$$A(z) = -\frac{D(z)}{\hat{t}(z)}. \quad (127)$$

<sup>1805</sup> Proposition 8 applied to  $\alpha_k$  states that the radius of convergence for  $D(z)$  is  $r_D = r_T =$   
<sup>1806</sup>  $c^{-1} > w_0$ . This means that the asymptotic behavior of  $A_k$  is determined by the simple pole  
<sup>1807</sup> of  $\hat{t}(z)^{-1}$ . In other words,  $A_k \sim A M_k$ , where the proportionality constant is given by

$$A = \frac{D(w_0)(1-cw_0)}{2w_0c(w_1-w_0)}. \quad (128)$$

<sup>1808</sup> Accordingly, the geometric Horton's law holds for  $\mathcal{N}_k[K]$ ,  $M_k$ , and  $A_k$  with the Horton  
<sup>1809</sup> exponent  $R_B = R_M = R_A = w_0^{-1}$ .

<sup>1810</sup> **F Exact Form of Mean Magnitudes  $M_k$  in Tokunaga**  
<sup>1811</sup> **Two-parameter Model**

<sup>1812</sup> If  $T_k = a c^{k-1}$ , equation (34) implies

$$M_k = 2 M_{k-1} + a \sum_{i=1}^{k-1} c^{k-i-1} M_i \quad (k \geq 2) \quad (129)$$

<sup>1813</sup> and therefore,

$$M_{k+1} = (2 + a) M_k + a \sum_{i=1}^{k-1} c^{k-i} M_i \quad (k \geq 2). \quad (130)$$

Equations (129) and (130) yield the following recursion

$$M_{k+1} - c M_k = (2 + a) M_k - 2c M_{k-1}$$

<sup>1814</sup> which simplifies to

$$M_{k+1} - (2 + a + c) M_k + 2c M_{k-1} = 0. \quad (131)$$

<sup>1815</sup> The recurrence relation (131) is solved by finding the roots of its characteristic equation

$$x^2 - (2 + a + c)x + 2c = 0. \quad (132)$$

The roots of (132) equal  $R_B$  and  $2 + a + c - R_B$ . See (120). Therefore,

$$M_k = c_1 R_B^{k-1} + c_2 (2 + a + c - R_B)^{k-1},$$

where the initial conditions  $M_1 = 1$  and  $M_2 = 2 + T_1 = 2 + a$  yield

$$c_1 = \frac{R_B - c}{2R_B - 2 - a - c} \quad \text{and} \quad c_2 = \frac{R_B - 2 - a}{2R_B - 2 - a - c}.$$

<sup>1816</sup> Hence,

$$\mathcal{N}_1[K] = M_K = \frac{R_B - c}{2R_B - 2 - a - c} R_B^{K-1} + \frac{R_B - 2 - a}{2R_B - 2 - a - c} (2 + a + c - R_B)^{K-1}. \quad (133)$$

Notice that, by (120),

$$R_B > 2 + a + c - R_B > 0.$$

Recall that  $2 + a + c = 2c(w_0 + w_1)$  and  $R_B = w_0^{-1} = 2cw_1$ . Thus,  $2 + a + c - R_B = w_1^{-1}$ , and (133) can be rewritten as follows

$$M_K = \frac{2w_1 - 1}{2(w_1 - w_0)} w_0^{1-K} + \frac{1 - 2w_0}{2(w_1 - w_0)} w_1^{1-K}.$$

<sup>1817</sup> Finally, for  $1 \leq j \leq K$ , (133) yields

$$\mathcal{N}_j[K] = \mathcal{N}_1[K - j + 1] = \frac{R_B - c}{2R_B - 2 - a - c} R_B^{K-j} + \frac{R_B - 2 - a}{2R_B - 2 - a - c} (2 + a + c - R_B)^{K-j}. \quad (134)$$

1818 **G Exact Quotient Horton's Law for Mean Branch Counts,**  
1819 **Magnitudes**

1820 Assume that the quotient Horton's law for the mean branch counts  $\mathcal{N}_1[K]$ , and hence for  
1821 the mean magnitudes  $M_K$ , holds exactly, that is (using the fact that  $M_1 = 1$ ):

$$M_K = R_M^{K-1}. \quad (135)$$

1822 Then,

$$M(z) = \frac{z}{1 - R_M z}$$

1823 which leads to

$$t(z) = -\frac{z}{M(z)} = -1 + R_M z \quad \text{and} \quad T(z) = (R_M - 2)z.$$

1824 This implies that the only self-similar model with exact quotient Horton's law for the mean  
1825 branch counts and magnitudes corresponds to the Tokunaga sequence

$$T_1 = R_M - 2, \quad T_k = 0 \quad \text{for } k > 1.$$

1826 **H Power Law Frequency for Branch Lengths in Critical Tokunaga  
1827 Trees**

1828 Here we examine the frequencies of branch lengths  $L_{[i]}$  in critical Tokunaga trees. First, we  
1829 prove the following technical lemma that establishes a power law decay of a function that  
1830 later will be interpreted as the survival function of the branch lengths.

1831 **Lemma 2** Define

$$\phi(x) = \sum_{k=0}^{\infty} \alpha^{-k} \exp(-\beta^{-k} x), \quad x > 0$$

1832 for some constants  $\alpha > 1$  and  $\beta > 1$ . Then,

$$\phi(x) = b(x)x^{-\frac{\log \alpha}{\log \beta}}, \quad (136)$$

1833 where  $b(x)$  is a function bounded from above and from below by positive constants.

1834 **Proof** Consider

$$\phi(x) = \sum_{k=0}^{\infty} \alpha^{-k} e^{-\beta^{-k} x}, \quad x > 0.$$

1835 Let

$$\sigma(x) = \sum_{k=-\infty}^{\infty} \alpha^{-k} e^{-\beta^{-k} x}, \quad x > 0.$$

Since  $\lim_{x \rightarrow \infty} \frac{\phi(x)}{\sigma(x)} = 1$ , we have

$$\phi(x) = \sigma(x)(1 + o(1)) \quad \text{as } x \rightarrow \infty.$$

Next, observe that

$$\sigma(x) = \alpha \sigma(\beta x),$$

and therefore

$$\frac{(\sigma(\beta x))^{-\frac{\log \beta}{\log \alpha}}}{\beta x} = \frac{(\sigma(x))^{-\frac{\log \beta}{\log \alpha}}}{x}.$$

Hence,

$$\pi(y) = e^{-y} \left( \sigma(e^y) \right)^{-\frac{\log \beta}{\log \alpha}}$$

is a positive continuous periodic function with period  $\log \beta$  that is bounded away from  $+\infty$  and from 0. Observe that

$$\sigma(x) = \left( x \pi(\log x) \right)^{-\frac{\log \alpha}{\log \beta}}.$$

Accordingly,  $\phi(x)$  can be expressed as

$$\phi(x) = \sigma(x)(1 + o(1)) = b(x)x^{-\frac{\log \alpha}{\log \beta}},$$

where

$$b(x) = \left( \pi(\log x) \right)^{-\frac{\log \alpha}{\log \beta}} (1 + o(1))$$

is a positive function, bounded from above and from below, by positive quantities. This completes the proof.

Recall from Sect. 9.2, Eq. (81) that the empirical frequencies of branch orders in a critical Tokunaga tree of order  $K$  are approximately geometric:

$$\{\text{proportion of branches of order } k\} = (R_B - 1)R_B^{-k}(1 + o(1)) \quad (137)$$

for any  $k \geq 1$  as  $K \rightarrow \infty$ . The lengths of branches of order  $k$  are i.i.d. exponential random variables with rate  $\gamma\zeta^{1-k}$ . Accordingly, the relative proportion of the lengths of order- $k$  branches that exceed a given value  $x > 0$  is  $\exp(-\gamma\zeta^{1-k}x)$ . Taking into account the relative proportions of branches of different orders, we find the relative proportion of branches with length exceeding  $x$ :

$$\begin{aligned} \frac{\#\{i : L_{[i]} \geq x\}}{\text{total no. branches}} &= \sum_{k=1}^{\infty} (R_B - 1)R_B^{-k} \exp(-\gamma\zeta^{1-k}x) \\ &= (1 - R_B^{-1}) \sum_{k=0}^{\infty} R_B^{-k} \exp(-\gamma\zeta^{-k}x) \\ &= (1 - R_B^{-1})b(\gamma x)(\gamma x)^{-\frac{\log R_B}{\log \zeta}} \\ &\propto b(\gamma x)x^{-\frac{\log R_B}{\log \zeta}}, \end{aligned}$$

where  $b(x)$  is a positive function bounded from zero and infinity. Here the next to the last step uses Lemma 2 with  $\alpha = R_B$  and  $\beta = \zeta$ .

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### Conflict of Interest

The authors declare that they have no conflict of interest.

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