



# Boundary Integral Equations

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## A Motivating Example

Consider the Boundary Value Problem given by:

$$\begin{aligned}-\frac{d^2u}{dx^2} &= f(x) \\ u(0) &= u(1) = 0\end{aligned}$$

How can we avoid the error produced by approximating the derivative operators and discretizing on a mesh grid?

What if we tried integrating it?

$$\begin{aligned}-\frac{du}{dx} &= \int_0^x f(y)dy + C_1 \\ -u(x) &= \int_0^x \int_0^y f(s)ds + C_1x + C_2\end{aligned}$$

With a little help from the following theorem:

$$\int_a^x \int_a^s f(y)dy = \int_a^x f(y)(x-y)dy$$

And applying the boundary conditions:

$$-u(x) = \int_0^x (x-s)f(s)ds + x \int_0^1 (s-1)f(s)ds$$

And changing the bounds of integration:

$$= \int_0^x s(1-x)f(s)ds + \int_x^1 x(1-s)f(s)ds$$



# The Green's Function!!

The previous solution is often rewritten as

$$u(x) = \int_0^1 G(x, s) f(s) ds$$
$$G(x, s) = \begin{cases} x(1-s) & x < s \\ s(1-x) & x > s \end{cases}$$

(Note: this is an exact formula)

How is this evaluated numerically?

$$= \int_0^x s(1-x) f(s) ds + \int_x^1 x(1-s) f(s) ds$$

Apply quadrature for each desired point

## Benefits:

- Integration is more stable than derivatives with numerical methods
- Error is only dependent on the quadrature method being used
- Error is localized and does not rely on previous steps

Unfortunately this method requires more overhead to set up, and requires a Green's function to exist for the given problem

How do you even find the Green's function?



# Linear Algebra Recap

Any linear differential equation can be written in the form

$$L[u] = f$$

How can we invert a differential operator?

Some form of integration!

$$u(x) = L^{-1}[f] = \int_a^b G(x, s) f(s) ds$$

But where does this come from?

Introducing: PHYSICS!!!

$$L[u] = \delta(x)$$

What if we moved it around and introduced the boundary conditions?

$$L[u] = \delta(x - x_0)$$

$$u(a) = u(b) = 0$$

Solving this problem would result in the Green's Function

$$u(x) = G(x, x_0)$$

(Note: Dependent on the operator  $L$  and the B.C.s)

Convolving this with the forcing function allows the point charge to be “moved around” and capture the behavior of the solution



# Sturm-Liouville Problems

Consider the well-known problem:

$$\frac{d}{dx} \left( p(x) \frac{du(x)}{dx} \right) + q(x)u(x) = f(x) \quad a < x < b$$

Using variation of parameters:

$$G(x, s) = \begin{cases} \frac{u_1(s)u_2(x)}{pW}, & a \leq x \leq s \\ \frac{u_1(x)u_2(s)}{pW}, & s \leq x \leq b \end{cases}$$

Where  $u_1$  and  $u_2$  are linearly independent solutions of the associated homogeneous problem with the respective homogeneous boundary conditions

Notes:

- The Wronskian is never zero on the interval because they are solutions to the ODE
- $pW$  will always be constant for ODEs of this form

But what about non-homogeneous boundary conditions?

$$u(x) = u_h(x) + u_p(x)$$

$$L[u_h] = 0 \quad u_h(a) = \alpha \quad u_h(b) = \beta$$

$$L[u_p] = f(x) \quad u_p(a) = 0 \quad u_p(b) = 0$$

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# Numerical Examples!!!!

(Pretty graphs & actual numbers)

Trapezoidal Quadratures use 100 nodes, Gaussian Quadratures use 10

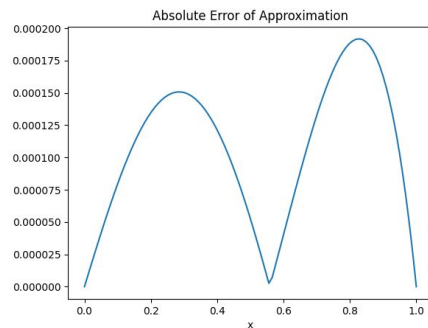
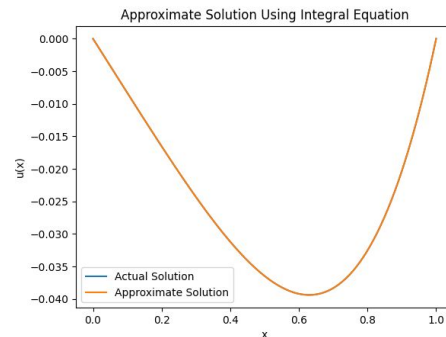
# Starting Easy

$$u''(x) = x^2$$
$$u(0) = u(1) = 0$$

Green's Function:

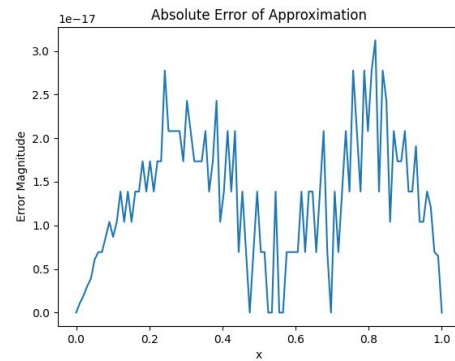
$$G(x, s) = \begin{cases} s(x-1) & 0 \leq x \leq s \\ x(s-1) & s \leq x \leq 1 \end{cases}$$

Solution:



Trapezoidal Error:

Gauss-Legendre Error:



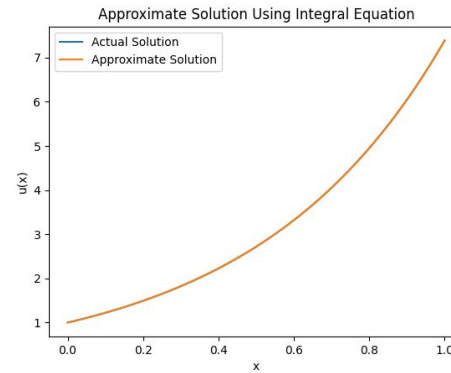
# Nonhomogeneous B.C.s

$$u''(x) = 4e^{2x}$$
$$u(0) = 1 \quad u(1) = e^2$$

Note: Same Green's function

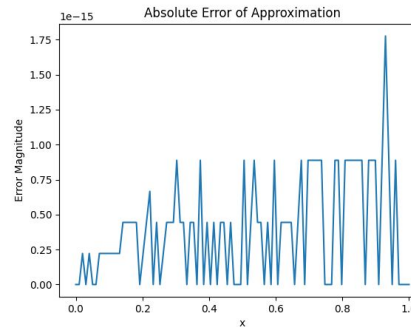
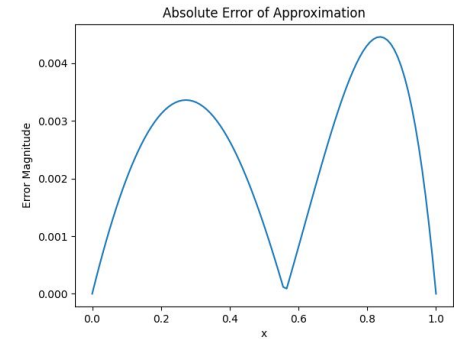
Particular Solution:

$$u_p(x) = x(e^2 + 1) + 1$$



Solution

Trapezoidal Error



Gauss-Legendre Error



# Meaner Problem

$$\frac{d}{dx} \left( x^2 \frac{du}{dx} \right) - 2u(x) = 4x^2$$

$$u(1) = 2$$

$$u(2) = -1$$

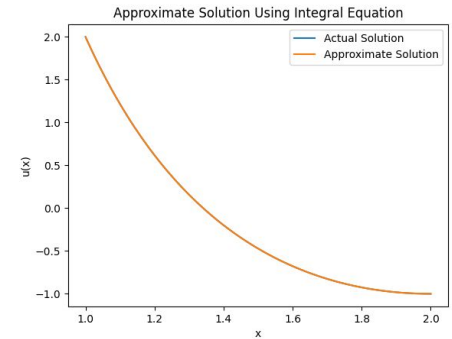
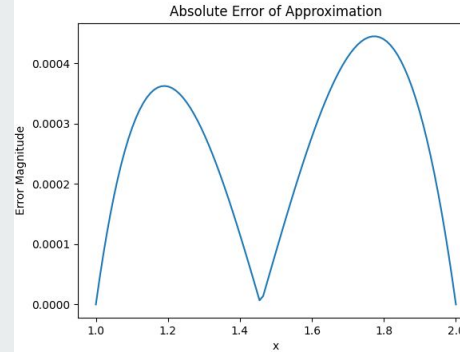
Green's Function:

$$G(x, s) = \begin{cases} \frac{1}{21}(s - s^{-2})(x - 8x^{-2}), & 1 \leq x \leq s \\ \frac{1}{21}(x - x^{-2})(s - 8s^{-2}), & s \leq x \leq b \end{cases}$$

Particular Solution:

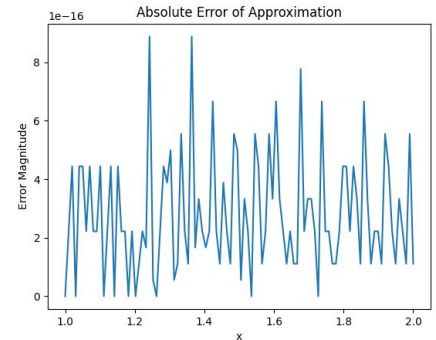
$$u_p(x) = -\frac{6}{7}x + \frac{20}{7}x^{-2}$$

Solution



Trapezoidal Error

Gauss-Legendre Error



# Independent Extension



Two dimensional PDEs (spooky scary)



## Green's Functions vs Fundamental Solutions

- In multiple dimensions, solutions still take the same form as before
- However boundary conditions become much harder to account for, especially on weird domain
- Instead we transfer the conditions onto an unknown density function

$$u(\mathbf{x}) = -\frac{1}{2\pi} \log |\mathbf{x} - \mathbf{x}_0|.$$

$$h(x) = \int_{\Gamma} \phi(x, y) \sigma(y) dl(y) \quad \text{where } u(x) = h(x) \quad x \in \Gamma$$

Let's look at Laplace's Equation

$$\begin{aligned} -\Delta u(x) &= 0 & x \in \Omega \\ u(x) &= g(x) & x \in \partial\Omega = \Gamma \end{aligned}$$



## Solving for the Density

- Will result in a dense linear system
- Becomes well conditioned when a double layer potential (dipole) is used

$$u(x) = \int_{\Gamma} \frac{n(y) \cdot (x - y)}{2\pi(x - y)^2} \sigma(y) dl(y)$$

When evaluated on the boundary, the jump discontinuity results in an extra term

$$g(x) = \frac{1}{2}\sigma(x) + \int_{\Gamma} \frac{n(y) \cdot (x - y)}{2\pi(x - y)^2} \sigma(y) dl(y)$$

(Fredholm Equation of the Second Kind)

Using any desired quadrature method (and parametrizing the boundary):

$$g(x_i) = \frac{1}{2}\sigma(x_i) + \sum_{j=1}^N \partial_{x_j} \phi(x_i, x_j) \sigma(x_j) w_j$$

$$\left(\frac{1}{2}I + D\right)\sigma = g$$

So we found the density, but what about the actual solution?

(Hint: It's already on this slide)

Use the same quadrature to evaluate the solution

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# More Numerical Examples!!!

(Even prettier graphs, but less numbers)

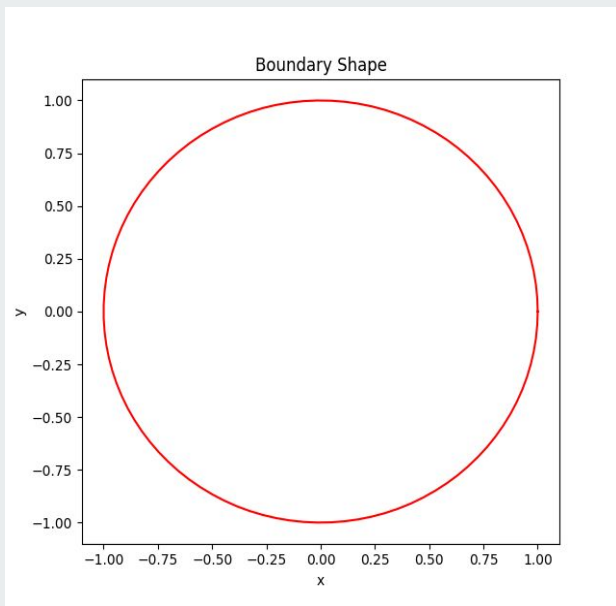
Note that all examples approximate Laplace's Equation with exact solution given by

$$\begin{aligned} -\Delta u(x) &= 0 \quad x \in \Omega \\ u(x) &= g(x) \quad x \in \partial\Omega = \Gamma \end{aligned}$$

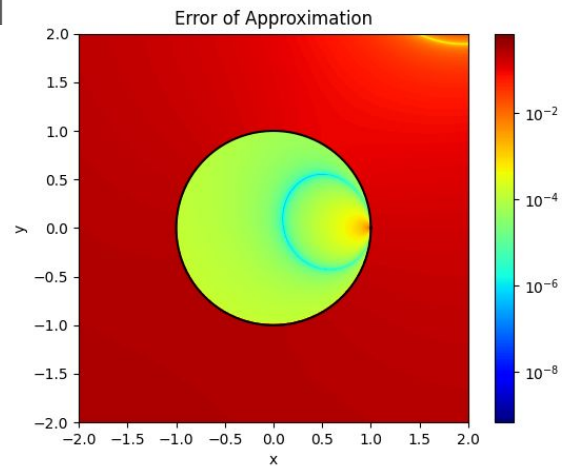
$$u(x) = -\frac{1}{2\pi} \log |x - (2, 3)|$$

Trapezoidal Quadratures use 1000 nodes, Gaussian Quadratures use 160 (10 panels of 16)

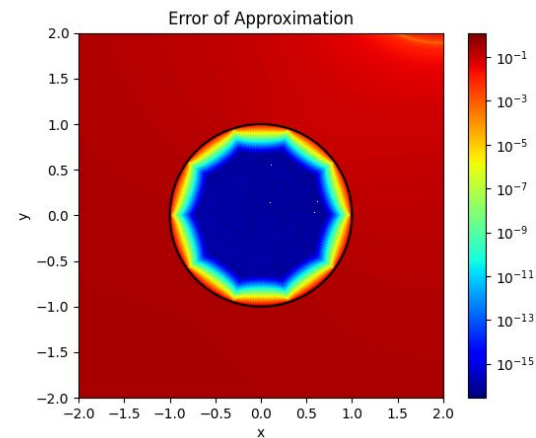
# Circle Boundary



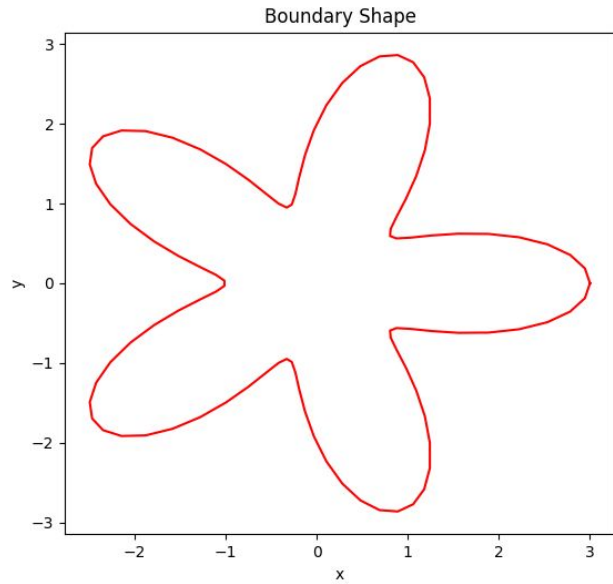
## Trapezoidal



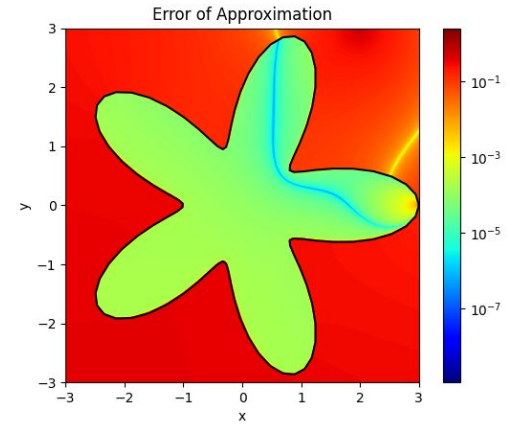
## Gauss-Legendre



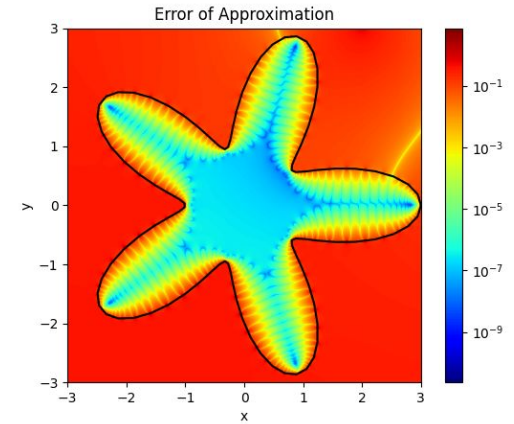
# Clover Boundary



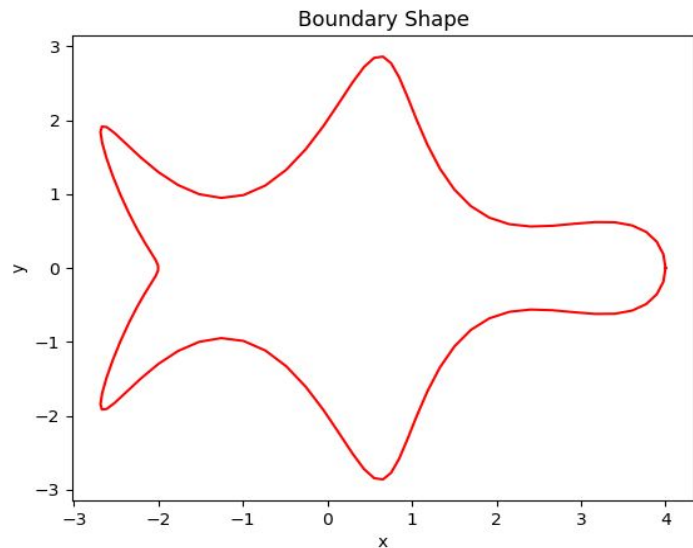
## Trapezoidal



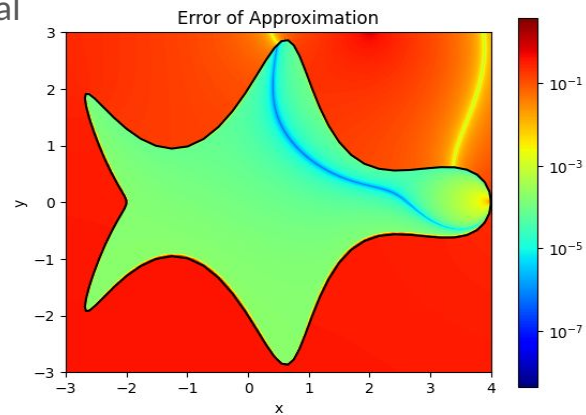
## Gauss-Legendre



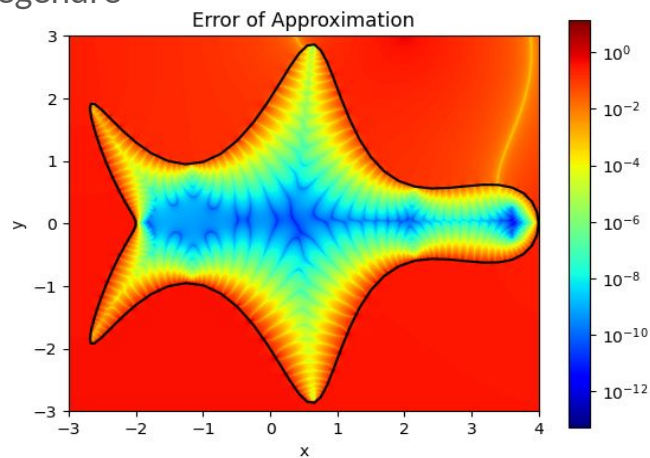
# Little Guy Boundary



## Trapezoidal

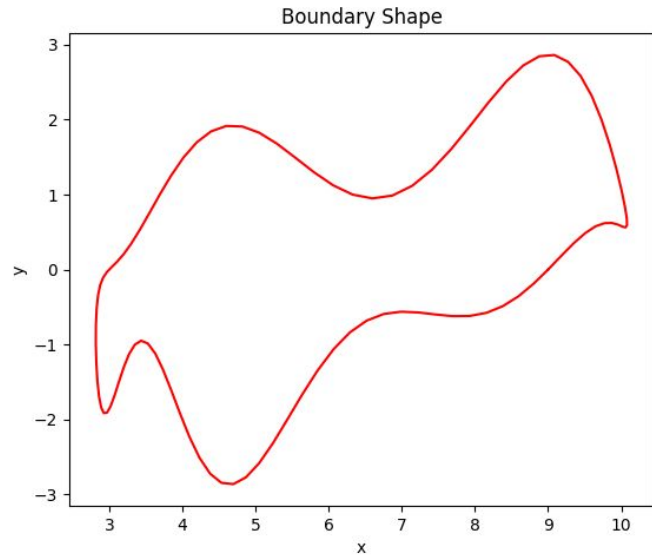


## Gauss-Legendre

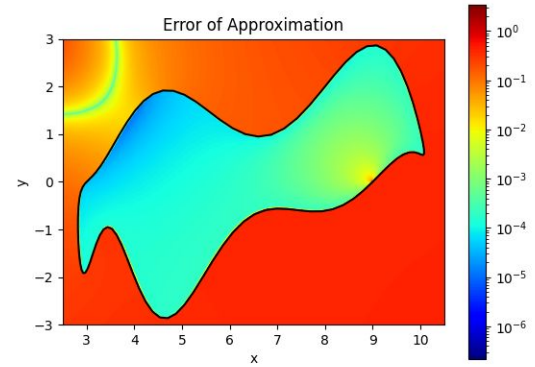




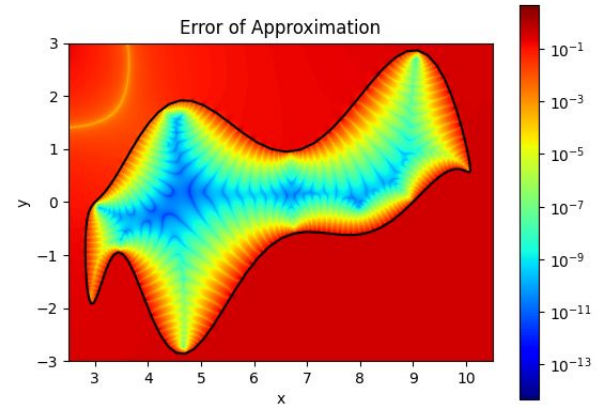
# The Blob



## Trapezoidal



## Gauss-Legendre



# Even More???

External Domain Problems!!! (WOOOO)

Just Kidding, but you can use this method to solve those as well

Other areas of extension:

- Non-smooth kernels (Helmholtz)
  - Neumann Boundary Conditions (Single Layer)
  - Evaluation close to the boundary
  - Different quadrature methods
  - Iterative solvers for the density
  - Interpolation of the density nodes
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# Sources

- [1] “Chapter 10: Integral equation formulations”. In: *Fast Direct Solvers for Elliptic PDEs*, pp. 105–114. DOI: 10.1137/1.9781611976045.ch10. eprint: <https://epubs.siam.org/doi/pdf/10.1137/1.9781611976045.ch10>. URL: <https://epubs.siam.org/doi/abs/10.1137/1.9781611976045.ch10>.
- [2] Jerry Farlow et al. *Differential equations & Linear Algebra*. 2nd ed. Pearson, 2018.
- [3] Dominique Habault. “Chapter 6 - Boundary Integral Equation Methods - Numerical Techniques”. In: *Acoustics*. Ed. by Paul Filippi et al. London: Academic Press, 1999, pp. 189–202. ISBN: 978-0-12-256190-0. DOI: <https://doi.org/10.1016/B978-012256190-0/50007-5>. URL: <https://www.sciencedirect.com/science/article/pii/B9780122561900500075>.
- [4] Richard Haberman. *Applied partial differential equations with Fourier series and boundary value problems*. 5th ed. Pearson, 2019.
- [5] John David Logan. *Applied mathematics*. 4th ed. Wiley-Interscience, 2006.
- [6] Gunnar Martinsson. “CBMS Conference on Fast Direct Solvers”. In: *Fast direct solvers for elliptic pdes*. Society for Industrial and Applied Mathematics, 2020.

# Questions?

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