

# Grad Project Equations in MCEN 5115: Rotating Inverted Pendulum

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April 11, 2016

## Normal Inverted Pendulum

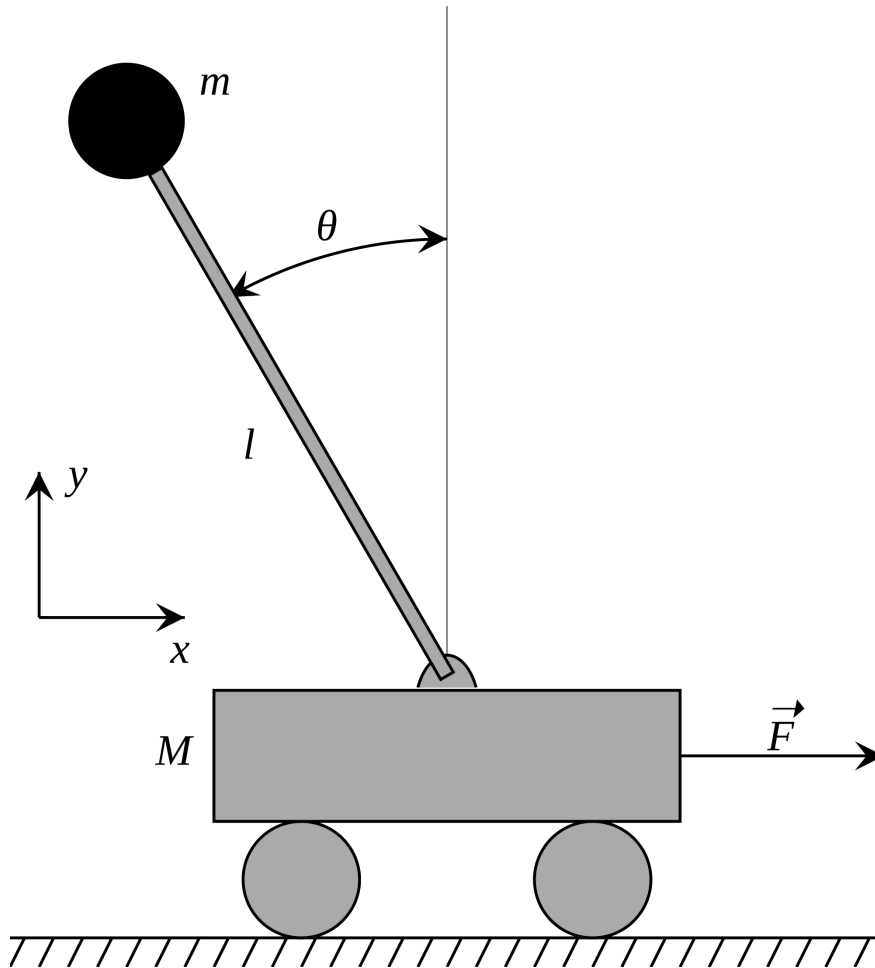


Figure 1: sup sup

Here we have the normal inverted pendulum equations for the image above:

$$(M + m)\ddot{x} - ml\ddot{\theta} + ml(\dot{\theta})^2 \sin(\theta) = F$$
$$l\ddot{\theta} - g \sin(\theta) = \ddot{x} \cos(\theta)$$

The  $x$  position for us is actually  $\omega$  because we are spinning. Also, the input force isn't  $F$ . What we have is a torque  $\tau = ||r|| ||F|| \sin(\theta)$ . Given that the  $\theta$  here is 90 degrees or 270 degrees, force can be equated to  $F = \frac{\tau}{r}$ .

Torque is some function of the motor current and the motor properties that we can put in later. We also note that  $x = \omega * r$  and thus  $\ddot{x} = \ddot{\omega}r$ .

Thus our equations become:

$$\begin{aligned}\frac{\tau}{r} &= ml(\dot{\theta})^2 \sin(\theta) - ml\ddot{\theta} + (M+m)\ddot{\omega}r \\ \ddot{\omega}r \cos(\theta) &= l\ddot{\theta} - g \sin(\theta)\end{aligned}$$

rewriting with  $\ddot{\omega}r = \frac{l\ddot{\theta} - g \sin(\theta)}{\cos(\theta)}$

$$\begin{aligned}\frac{\tau}{r} &= ml(\dot{\theta})^2 \sin(\theta) - ml\ddot{\theta} + \frac{(M+m)(l\ddot{\theta} - g \sin(\theta))}{\cos(\theta)} \\ &= ml(\dot{\theta})^2 \sin(\theta) + \ddot{\theta} \frac{((M+m)l - ml \cos(\theta))}{\cos(\theta)} - \frac{(M+m)g \sin(\theta)}{\cos(\theta)} \\ \ddot{\theta} &= \frac{\cos(\theta)}{(M+m)l - ml \cos(\theta)} \left( \frac{\tau}{r} - ml(\dot{\theta})^2 \sin(\theta) - \frac{(M+m)g \sin(\theta)}{\cos(\theta)} \right)\end{aligned}$$

hmmm, those equations suck to take partial derivatives.

$$\begin{aligned}\ddot{\theta} &= \dot{\theta}^2 \sin(\theta) - \frac{\tau}{rml} + \frac{(M+m)\ddot{\omega}r}{ml} \\ \ddot{\omega} &= \frac{l\ddot{\theta} - g \sin(\theta)}{r \cos(\theta)}\end{aligned}$$

Now let's pick some variables.  $x_1 = \dot{\theta}$ ,  $x_2 = \theta$ ,  $x_3 = \dot{\omega}$  and  $x_4 = \omega$ . We need to linearize around  $\omega = 0$  and  $\dot{\omega} = 0$ .

State matrix will look like:

$$\dot{x} = Ax = \begin{bmatrix} a_0 & a_1 & a_2 & a_3 \\ 1 & 0 & 0 & 0 \\ b_1 & b_2 & b_3 & b_4 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{bmatrix}$$

Here we have  $a_0 = 0$ ,  $a_1 = \frac{(M+m)g}{(M+m)l - ml}$ ,  $a_2 = 0$ , and  $a_3 = 0$ . Now we write out the equations in the other form.

$$l\ddot{\theta} = \frac{-\tau}{rm} + l(\dot{\theta})^2 \sin(\theta) + \frac{(M+m)\ddot{\omega}r}{m}$$

this gives:

$$\ddot{\omega}r \cos(\theta) = -\frac{\tau}{rm} + l(\dot{\theta})^2 \sin(\theta) + \frac{(M+m)\ddot{\omega}r}{m} - g \sin(\theta)$$

Simplifying we get:

$$\ddot{\omega} \left( r \cos(\theta) - \frac{(M+m)r}{m} \right) = \frac{-\tau}{rm} + l(\dot{\theta})^2 \sin(\theta) - g \sin(\theta)$$

etc:

$$\ddot{\omega} = \frac{-\tau}{r^2(m \cos(\theta) - (M+m))} + \frac{ml(\dot{\theta})^2 \sin(\theta)}{r(\cos(\theta) - (M+m))} - \frac{mg \sin(\theta)}{r(\cos(\theta) - (M+m))}$$

## small angle

This is not wrong, but we will be doing the small angle approximation then the jacobian. I don't need to do jacobian I need to do small angle approximation. For small  $\theta$   $\cos(\theta) \approx 1$  and  $\sin(\theta) \approx \theta$  Thus we get:

$$\ddot{\theta} = \frac{1}{Ml} \left( \frac{\tau}{r} - ml(\dot{\theta})^2\theta - (M+m)g\theta \right)$$

and:

$$\ddot{\omega} = \frac{-\tau}{-r^2M} + \frac{ml(\dot{\theta})^2\theta}{r(1-(M+m))} - \frac{mg\theta}{r(1-(M+m))}$$

Our state equations are thus:

$$\begin{bmatrix} \ddot{\theta} \\ \dot{\theta} \\ \ddot{\omega} \\ \dot{\omega} \end{bmatrix} = \begin{bmatrix} 0 & -(M+m)g & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & \frac{-mg}{r(1-(M+m))} & 0 & 0 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} \dot{\theta} \\ \theta \\ \dot{\omega} \\ \omega \end{bmatrix} + \begin{bmatrix} \frac{1}{Mlr} \\ 0 \\ \frac{1}{r^2M} \\ 0 \end{bmatrix} \tau$$

The output is then:

$$y = \begin{bmatrix} 0 & 1 & 0 & 0 \end{bmatrix} \begin{bmatrix} \dot{\theta} \\ \theta \\ \dot{\omega} \\ \omega \end{bmatrix}$$

Note here that  $a_0$  and  $b_1$  are zero. If this doesn't turn out to be good enough we can take the third part of the taylor series for the  $\dot{\theta}^2\theta$  term which will be non-zero unlike the first 2.