## FORMULARI ESTAT SOLID

| CS (a)          | $\vec{a}_1 = a\hat{i}; \vec{a}_2 = a\hat{j}; \vec{a}_3 = a\hat{k}$  | $\vec{b}_1 = \frac{2\pi}{a}\hat{i}; \vec{b}_2 = \frac{2\pi}{a}\hat{j}; \vec{b}_3 = \frac{2\pi}{a}\hat{k}$  |  |  |
|-----------------|---|--|--|--|
| BCC (a)         | $\vec{a}_{1} = \frac{a}{2} (-\hat{i} + \hat{j} + \hat{k}); \vec{a}_{2} = \frac{a}{2} (\hat{i} - \hat{j} + \hat{k})$ $\vec{a}_{3} = \frac{a}{2} (\hat{i} + \hat{j} - \hat{k})$ | $\vec{b}_{1} = \frac{2\pi}{a} (\hat{j} + \hat{k}); \vec{b}_{2} = \frac{2\pi}{a} (\hat{i} + \hat{k}); \vec{b}_{3} = \frac{2\pi}{a} (\hat{i} + \hat{j})$                                 |  |  |
| FCC (a)         | $\vec{a}_1 = \frac{a}{2}(\hat{j} + \hat{k}); \vec{a}_2 = \frac{a}{2}(\hat{i} + \hat{k}); \vec{a}_3 = \frac{a}{2}(\hat{i} + \hat{j})$  | $\vec{b}_{1} = \frac{2\pi}{a} (-\hat{i} + \hat{j} + \hat{k}); \vec{b}_{2} = \frac{2\pi}{a} (\hat{i} - \hat{j} + \hat{k})$ $\vec{b}_{3} = \frac{2\pi}{a} (\hat{i} + \hat{j} - \hat{k})$ |  |  |
| Hexagonal (a,c) | $\vec{a}_1 = a \hat{i} ; \vec{a}_2 = \frac{a}{2} (\hat{i} + \sqrt{3} \hat{j}) ; \vec{a}_3 = c \hat{k}$  | $\vec{b}_1 = \frac{2\pi}{a} \left( \hat{i} - \frac{1}{\sqrt{3}}  \hat{j} \right);  \vec{b}_2 = \frac{4\pi}{\sqrt{3}  a}  \hat{j};  \vec{b}_3 = \frac{2\pi}{c}  \hat{k}$                |  |  |

$$\vec{R} = n_i \vec{a}_i; \vec{G} = h \vec{b}_1 + k \vec{b}_2 + l \vec{b}_3 \qquad \vec{b}_i \cdot \vec{a}_j = 2\pi \delta_{ij} \qquad e^{i\vec{G} \cdot \vec{R}} = 1$$

| Estructura cristal·lina | Xarxa<br>Bravais | Base atòmica  | Factor estructura   |
|-------------------------|------------------|---|---|
| Diamant                 | FCC              | $\vec{r}_1 = (0,0,0)_{c.s.}; \vec{r}_2 = (1/4,1/4,1/4)_{c.s}$   | $S_{\text{diamant}} = S_{fcc} \cdot \left[ 1 + e^{-i\frac{\pi}{2}(h+k+l)} \right]$      |
| CsCl                    | CS               | $\vec{r}_{Cs^+} = (0,0,0)_{c.s}$<br>$\vec{r}_{Cl^-} = (1/2,1/2,1/2)_{c.s}$  | $S = f_1 + f_2 e^{-i\pi(h+k+l)}$  |
| NaC1                    | FCC              | $r_{Na^{+}}^{\rightarrow} = (1/2, 1/2, 1/2)_{c.s}$<br>$r_{Cl}^{\rightarrow} = (0, 0, 0)_{c.s}$                                | $S = S_{fcc} \cdot \left[ f_1 + f_2 e^{-i\pi(h+k+l)} \right]$                           |
| Zinc Blenda             | FCC              | $\vec{r}_{Zn^{+}} = (0,0,0)_{c.s}$<br>$\vec{r}_{S^{-}} = (1/4,1/4,1/4)_{c.s}$   | $S = S_{fcc} \cdot \left[ f_{Zn} + f_S e^{-i\frac{\pi}{2}(h+k+l)} \right]$              |
| "FCC"                   | CS               | $\vec{r}_1 = (0,0,0)_{c.s.}; \vec{r}_2 = (0,1/2,1/2)_{c.s}$<br>$\vec{r}_3 = (1/2,0,1/2)_{c.s}; \vec{r}_4 = (1/2,1/2,0)_{c.s}$ | $S_{fcc} = f \cdot \left[ 1 + e^{-i\pi(h+k)} + e^{-i\pi(h+l)} + e^{-i\pi(k+l)} \right]$ |
| "BCC"                   | CS               | $\vec{r}_1 = (0,0,0)_{c.s}; \vec{r}_2 = (1/2,1/2,1/2)_{c.s}$  | $S_{bcc} = f \left[ 1 + e^{-i\pi(h+k+l)} \right]$                                       |

$$\begin{split} & \overline{k} \cdot \left(\frac{1}{2} \, \overrightarrow{G}\right) = \left(\frac{1}{2} \, \overrightarrow{G}\right)^2 \Rightarrow 2 \operatorname{d} \sin \theta = n \, \lambda \\ & F = N \int_{\operatorname{cel·la}} dV \, n(\vec{r}) \, e^{-i \vec{G} \cdot \vec{r}} = N \cdot S_G \end{split} \qquad \begin{aligned} & d = \frac{2\pi}{|\vec{G}|} \quad I = \operatorname{degen.} |S^2| \\ & S_{\tilde{G}} = \sum_{j} f_{j} e^{-i \vec{G} \cdot \vec{r}_{j}} \right] \quad f_{j} = \int_{\operatorname{cel·la}} dV \, n_{j}(\vec{\rho}) \, e^{-i \vec{G} \cdot \vec{\rho}} \\ & S_{\operatorname{cel·la}} \end{aligned} \qquad \begin{aligned} & N_{1D} = k_{F} \cdot 2_{(\operatorname{spin})} \cdot D(k) \\ & N_{2D} = (\pi \, k_{F}^{2}) \, 2_{(\operatorname{spin})} \cdot D(k) \\ & N_{2D} = (\pi \, k_{F}^{2}) \, 2_{(\operatorname{spin})} \cdot D(k) \\ & N_{3D} = \frac{4\pi \, k_{F}^{3}}{3} \, 2_{(\operatorname{spin})} \cdot D(k) \end{aligned} \\ & D(E) = \frac{dN}{dk} \cdot \frac{dk}{dE} \quad D_{3D}(E) = \left(\frac{V}{2\pi^{2}}\right) \left(\frac{2m}{\hbar^{2}}\right)^{3/2} E^{1/2} \quad \text{(densitat d'estats monoparticulars)} \quad E_{F} = \frac{\hbar^{2}}{2m} \left(3\pi^{2} \, \rho\right)^{2/3} \\ & \rho = \frac{N}{V} \quad N = \int_{0}^{\infty} D(E) \, f(E) \, dE \quad U = \int_{0}^{\infty} D(E) \, f(E) \, E \, dE \quad \Delta U \approx \frac{\pi^{2}}{6} (k_{B} \, T)^{2} D(E_{F}) \quad C = \frac{\pi^{2}}{2} \left(k_{B} \, \frac{T}{E_{F}}\right) N \, k_{B} \end{aligned} \\ & f_{0} = \frac{1}{e^{\frac{E-\mu}{k_{B}T}} + 1} \quad \mu \approx E_{F} \left[1 - \frac{1}{3} \left(\frac{\pi \, k_{B} \, T}{2E_{F}}\right)^{2}\right] \end{aligned}$$
Sòlid: 
$$\Phi_{\vec{k}}(\vec{r}) = \sum_{\vec{G}} C_{\vec{k} - \vec{G}} e^{i(\vec{k} - \vec{G})\vec{r}} \quad \text{Eq. Central:} \left[\frac{\hbar^{2} |\vec{k} - \vec{G}|^{2}}{2m} - E_{\vec{k}} \right] C_{\vec{k} - \vec{G}} + \sum_{\vec{E} \in \mathcal{L}} V_{\vec{G}'' - \vec{G}} C_{\vec{k} - \vec{G}''} = 0 \end{aligned}$$

Sòlid:  $\phi_{\vec{k}}(\vec{r}) = \sum_{\vec{G}} C_{\vec{k}-\vec{G}} e^{i(\vec{k}-\vec{G})\vec{r}}$  Eq. Central:  $\left[ \left[ \frac{\hbar^2 |\vec{k}-\vec{G}|^2}{2m} - E_{\vec{k}} \right] C_{\vec{k}-\vec{G}} + \sum_{\vec{G}''} V_{\vec{G}''-\vec{G}} C_{\vec{k}-\vec{G}''} = 0 \right]$ 

*Xarxa buida*:  $\left| E_{\vec{k}-\vec{G}} = \frac{\hbar^2}{2 \, m} |\vec{k} - \vec{G}|^2 \right|$  *Fortament lligats* (LCAO):  $E(\vec{k}) = E_0 - \alpha - \gamma \sum_{m} e^{i \vec{l} \cdot \vec{\rho}_m}$ Feblement lligats: No degen.  $E_{\vec{k}-\vec{G}} = \frac{\hbar^2}{2m} |\vec{k}-\vec{G}|^2 + V_0 + \frac{\sum_{\vec{G}''\neq 0} |V_{\vec{G}''}|^2}{\frac{\hbar^2}{2m} |\vec{k}-\vec{G}|^2 - \frac{\hbar^2}{2m} |\vec{k}-\vec{G}-\vec{G}''|^2}$ Quasi degen:  $E_{\pm} = \frac{E_{\vec{k}-\vec{G}}^0 + E_{\vec{k}-\vec{G}-\vec{G}'}^0}{2} \pm \sqrt{\left(\frac{E_{\vec{k}-\vec{G}}^0 - E_{\vec{k}-\vec{G}-\vec{G}'}^0}{2}\right)^2 + |V_{\vec{G}'}|^2}$  Degenerat:  $E_{\pm} = \frac{\hbar^2}{2m} |\vec{k} - \vec{G}|^2 \pm |V_{\vec{G}'}|$  (1 pla bisector,  $V_0 = 0$ ) Propietats de transport:  $\frac{-\hbar^2}{2m}\nabla^2 + V(\vec{r}) \left| \phi_{n\vec{k}}(\vec{r}) = E_n(\vec{k}) \phi_{n\vec{k}}(r) \right| \Rightarrow \phi_n(\vec{r},t) = \sum_{\vec{r}} g(\vec{k}) \phi_{n\vec{k}}(r) e^{-iE_n(\vec{k})\frac{t}{\hbar}}$  $\vec{v}_n(\vec{k}) = \frac{1}{\hbar} \vec{\nabla}_{\vec{k}} E_n(\vec{k}) \Big|_{\vec{k} = \vec{k}} \qquad \vec{F} = \hbar \frac{d\vec{k}}{dt} \qquad \vec{F} = -e \left[ \vec{\epsilon} + \vec{v} \times \vec{B} \right] \qquad \left( \frac{1}{m^*} \right) = \frac{1}{\hbar^2} \frac{\partial^2 E}{\partial k \cdot \partial k} \qquad a_i = \left( \frac{1}{m^*} \right) F_j$  $\frac{\partial f}{\partial t} + \vec{v} \cdot \vec{\nabla}_{\vec{r}} f + \frac{\vec{F}}{\hbar} \cdot \vec{\nabla}_{\vec{k}} f = \left(\frac{\partial f}{\partial t}\right) \quad \text{(estacionari, t relaxació)} \approx \Rightarrow f = f_0 - \tau \vec{v} \cdot \vec{\nabla}_{\vec{r}} f - \tau \frac{\vec{F}}{\hbar} \cdot \vec{\nabla}_{\vec{k}} f$ (homogeni, linealització)  $f(\vec{k}) \approx f_0(\vec{k}) - \tau(\vec{k}) \frac{\vec{F}}{\hbar} \cdot \vec{\nabla}_{\vec{k}} f_0(\vec{k})$  (no val per camps magnètics)  $\vec{J} = -\frac{e}{V} \sum_{\vec{k}} f(\vec{k}) \vec{v}(\vec{k}) \qquad \sigma_{ij} = \frac{-e^2}{V} \sum_{\vec{k}} \tau(\vec{k}) \frac{\partial f_0}{\partial E} v_i v_j \qquad \qquad J_x = e^2 \frac{\epsilon}{3V} \tau(k_F) v^2(k_F) D(E_F)$   $, \text{ metalls:} \qquad \sigma = \frac{e^2 \tau_F n}{*}$  $\vec{F} = -e[\vec{\epsilon} + \vec{v} \times \vec{B}] \Rightarrow f(\vec{k}) = f_0(\vec{k}) + \frac{e\tau}{1 + \omega^2 \tau^2} \frac{\partial f_0}{\partial E} \vec{v} \cdot \left[ \vec{\epsilon} - \frac{e\tau}{m^*} (\vec{\epsilon} \times \vec{B}) \right] \qquad \omega_c = \frac{eB}{m^*}$  $\vec{J} = \sigma_0 \vec{A} \qquad \vec{A} = \frac{e\tau}{1 + \omega^2 \tau^2} \left| \vec{\epsilon} - \frac{e\tau}{m^*} (\vec{\epsilon} \times \vec{B}) \right| \qquad \frac{1}{\tau} = \frac{1}{\tau_{varya}} + \frac{1}{\tau_{impureses}} \quad \text{Efecte Hall:} \quad R_H = \frac{-e\tau}{m^* \sigma_0} = -\frac{1}{ne}$ <u>Moviment dels ions:</u>  $V = \sum_{s} \frac{C}{2} (u_{s+1} - u_s)^2$   $m \frac{d^2 u_s}{dt^2} = C[u_{s+1} + u_{s-1} - 2u_s]$   $u_s(x, t) = A e^{i(kx - \omega t)}$  $\omega = \sqrt{\frac{4C}{m}} \left| \sin \left( \frac{ka}{2} \right) \right|$   $v_g = \frac{d\omega}{dk} = a\sqrt{\frac{C}{m}} \cos \left( \frac{ka}{2} \right)$   $\omega_{\text{max}} = \frac{2v_{so}}{a}$   $v_{so} = \frac{d\omega}{dk} \Big|_{z=z} = a\sqrt{\frac{C}{m}}$ Base diatòmica:  $\omega^2 = C \left( \frac{m+M}{mM} \right) \pm C \left[ \left( \frac{m+M}{mM} \right)^2 - \frac{2(1-\cos ka)}{mM} \right]^{\frac{1}{2}}$ En 3D:  $F_{s,\alpha} = -\sum_{r} \sum_{\beta} \phi_{\alpha\beta}(s,r) u_{r,\beta}$   $\phi_{\alpha\beta} = \frac{\partial^2 V}{\partial u_{\beta} \partial u_{\beta}}$   $m_s \ddot{u}_{s,\alpha} = -\sum_{r} \sum_{\beta} \phi_{\alpha\beta}(s,r) u_{r,\beta}$  $Det \left| \phi_{\alpha\beta}(s,r) - m_s \omega^2 \delta_{sr} \delta_{\alpha\beta} \right| = 0$ Fonons:  $H = \frac{p^2}{2m} + \frac{1}{2}k x^2 = \hbar \omega (N + \frac{1}{2})$   $\langle n_{\vec{k}s} \rangle = \frac{\sum_{n=0}^{\infty} n e^{-n\hbar \omega_s(\vec{k})/k_B T}}{\sum_{n=0}^{\infty} e^{-n\hbar \omega_s(\vec{k})/k_B T}} = \frac{1}{e^{\hbar \omega_s(\vec{k})/k_B T} - 1}$  $\langle E_{\vec{k}s} \rangle = \langle n_{\vec{k}s} \rangle \hbar \omega_s(\vec{k}) = \frac{\hbar \omega_s(\vec{k})}{2\pi^2} \qquad N = \frac{V}{6\pi^2} k^3 \qquad D_s(\omega) = \frac{dN}{d\omega} = \left(\frac{V k^2}{2\pi^2}\right) \frac{dk}{d\omega} - \frac{V}{2\pi^2} \left(\frac{V k^2}{2\pi^2}\right) \frac{dk}{d\omega} = \frac{V}{2\pi^2} \left(\frac$  $U = \sum_{a \text{ homogeneous}} \int d\omega D_p(\omega) \left( \frac{1}{2} + \frac{1}{e^{\hbar \omega / k_B T} - 1} \right) \hbar \omega \qquad C_{\text{xarxa}} = k_B \sum_s \int d\omega D_s(\omega) \frac{x^2 e^x}{\left(e^x - 1\right)^2} \qquad x = \frac{\hbar \omega}{k_B T}$ Debye:  $\omega = v_s k$   $D(\omega) = \frac{V \omega^2}{2\pi^2 v_s^3}$   $\omega_D = \left(\frac{6\pi^2 v_s^3 N}{V}\right)^{1/3}$   $\theta_D = \frac{\hbar \omega_D}{k_B}$   $C = 9Nk_B \left(\frac{T}{\theta_D}\right)^3 \int_0^{x_D} \frac{x^4 e^x}{(e^x - 1)^2} dx$ Einstein:  $\omega = \omega_E$   $D(\omega) = 3N\delta(\omega - \omega_E)$   $\theta_E = \frac{\hbar \omega_E}{k_B}$   $C = 3Nk_B \frac{e^{\theta_E/T}}{(1 + \theta_E/T - 1)^2} \left(\frac{\theta_E}{T}\right)$