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NS: Inverse 2 \times 2: For A := [a, b; c, d], A^{-1} := \frac{1}{ad - bc}[d, -b; -c, a] Radial: r\dot{r} = x\dot{x} + y\dot{y}, \ \dot{\theta} = x\dot{x} + y\dot{y}
     \left[\tan^{-1}(y/x)\right]' = \frac{x\dot{y} - \dot{x}y}{x^2 + y^2} Classifications: Node: \lambda_i \in \mathbb{R}, \Pi\lambda_i > 0 Centre: \lambda_i = \pm ib Focus: \lambda_i = a \pm ib
     Hyperbolic: \operatorname{Re}(\lambda) \neq 0 \to \operatorname{hyperbolic}. If all \lambda < 0 for \operatorname{Spec}(Df(x_0)) then A-Stable Invariant Set:
     |\phi_t(S)| \subseteq S Lim Pts: \omega pt. if \lim_{t\to\infty} \phi(x) = p, i.e. flows tend to p. \alpha pt. if \lim_{t\to\infty} \phi(x) = p.
     Attracting Set: A set A \subseteq S if \exists neighbourhood U s.t. \phi(U) \subseteq U \forall t \geq 0, and A = \bigcap_{t>0} \phi(U) Dense
     Orbits: If \forall \epsilon > 0, x \in A with A an attracting set, \exists \tilde{x} \in \Gamma s.t. | x - \tilde{x}| < \epsilon. I.e. a dense orbit goes as close
     as needed to any point within A Attractor: An attracting set with a dense orbit. Lyapunov Stable:
    If \forall \epsilon > 0, \exists \ \delta > 0 s.t. \forall \ x \in B_{\delta}, t \geq 0, \phi(t) \in B_{\delta} (i.e. points stay close within region). Asymptotically
    Stable: If L-Stable and \exists \ \delta > 0 s.t.\phi(x) \to x_0 \forall x \in B_\delta Lyapunov F'n: V(x_0) = 0, V(x) > 0 \forall \ x \neq x_0.
     Then if V < 0 \to A-Stable, or if V \le 0 L-Stable. Stable Manifold: If spectrum of Df(x_0) has k eigends
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     with positive real parts, and n-k with negative, then \exists an n-k dim manifold tangent to E^s s.t. for all
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     t>0 \phi(W^s_{loc})\subseteq W^s_{loc}, and \forall x\in W^s_{loc}\phi(x)\to x_0 as time increases. Repeat for k-dim unstable manifold
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     but for negative time. Then, define e.g. global stable manifold by W^s(x_0) := \bigcup_{t \leq 0} \phi_t(W^s_{loc}). Note that we
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     search backwards in time for stable, and forwards for unstable! Centre Manifold: If x_0 not hyperbolic
     (0 real part), then E^c is the centre subspace. Then \exists W^c parallel to E^c, of class C^r, and invariant under
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     flow. Want bifurcation at \mu = 0, so with change of variables first find eigers v_1, v_2. Then, construct
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     P := [v_1, v_2] s.t. \vec{x} = P\vec{\xi}. NOTE: first v_i in P is always associated with \text{Re}(\lambda) = 0. Solve for \vec{\xi} and then
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     expand with \eta = h(\xi, \tilde{\mu}) Alt. Centre Manifold: If vector v_1 \sim E^c = [a, b]^T then we have y = bx/a
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     (e.g. [1,1]^T \to y = x. If bifurcation at \mu = \alpha then have \mu = \tilde{\mu} + \alpha s.t. bifurcation when \tilde{\mu} = 0. Then have
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     \dot{x}(x,y,\tilde{\mu})=\ldots etc. Next, set up y=h(x,\tilde{\mu})=bx/a+b_1\tilde{\mu}+b_2\tilde{\mu}^2+a_2x^2+c_2\tilde{\mu}x and proceed as usual, s.t.
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    y is along E^c. Transcritical Bifurcation: Always two points, change type at origin. E.g. \dot{x} = \mu x - x^2
    Saddle-Node: E.g. \dot{x} = \mu - x^2 Bifurcation begins to exist at origin. Supercritical: E.g. \dot{x} = \mu x - x^3,
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     where stable \to 2\times stable and one unstable. Subcritical: E.g. \dot{x} = -\mu x + x^3, where unstable \to 2\times
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     unstable and one stable. General co-dim 1: If \dot{x} = f then \dot{x} = \mu f_u + 0.5x^2 f_{xx} + x\mu f_{x\mu} + 0.5\mu^2 f_{\mu\mu}
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    Generally this is a saddle-node but if f_u = 0 we have \dot{x} = x\mu f_{x\mu} + 0.5x^2 f_{xx}, which is a transcritical. However if x = -x then \dot{x} = x(\mu f_{x\mu} + \ldots) + x^3(f_{xxx}/6 + \ldots) \to \text{pitchfork}. Saddle-node stable under
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     perturbations! Homoclinic Orbits Sum of roots of cubic = - coeff. of x^2
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     FPDE: Types: 1^{st}: \exists scale s.t. solution found, not so for 2^{nd}. Heat: \hat{T} = u(\hat{T}_{\infty} - \hat{T}_{-\infty}) + \hat{T}_{-\infty} Oil
     Spread: Dims: x = x_f + \varepsilon \xi, t = \tau Ground Spread: (1-s)\phi h_t + Q_x = 0; Q \sim -hh_x, 0 < x_s < x_f. Have
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     h(x_f) = 0, h_t(x_s) = 0, and hh_x|_{x=0,x_f} = 0 (i.e. no flux at centre and front), and h, hh_x cont. at joint.
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     Expansions: Let \xi = z + \epsilon \eta for perturbations Scale: Try x = x_f + \epsilon \xi for groundwater Stefan: S_0 = \xi
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     C\left(T_{1}-T_{m}\right)/L, condition = \rho L\dot{s}=kT_{s}|_{s-}^{s+} 1ph Stefan: Bar = T_{h}|liq|_{s}sol|INS. Use T=T_{m}+(T_{1}-T_{m})u
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     s.t. S_0 u_t = u_{xx}, u = 1 @ x = 0, \{\dot{s} = -u_x, u = 0\} @ x = s, s(0) = 0. Sim. sol is s = \beta \sqrt{t}, f = f(x/\sqrt{t})
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    2ph Stefan: Use T = T_m + (T_1 - T_m)u s.t. S_0u_t = u_{xx} @ 0 < x < s, (S_0/\kappa)u_t = u_{xx} @ s < x < 1, u = 1 @ x = 0, u_x = 0 @ x = 1, {<math>\dot{s} = Ku_x|_{s_+} - u_x|_{s_-}, u = 0} @ x = s, \{s = 0, u = -\theta\} @ x = 0. Here
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     \theta := (T_m - T_0)/(T_1 - T_m), \kappa := c_1 k_1/(c_2 k_2), K := k_2/k_1 \text{ Sim. sol is } s = \beta \sqrt{t}, f = f(x/\sqrt{t}) \text{ 2-Dim:}
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    U_n = \hat{n} \cdot u = K(u_2)_n - (u_1)_n. If x = f(y,t) then \hat{n} := \nabla (x-f) = [1,-f_y]^T/\sqrt{1+f_y^2} Welding:
     Have 0 < s_2 < s_1. Have cold x = a, no flux x = 0. \theta = 1 in liquid. In mush \rho L\theta_t = J^2/\sigma, CoE
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     \to \theta \rho L \dot{s} + k T_x |_{s_-}^{s_+} = 0. Have \theta cont. (= 0) at s_1. I.e. we have S_0 u_t = u_{xx} + q, u_x = 0 @ x = 0, u = 0
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     -1 @ x = 1, \theta = 0 @ x = s_1. Also \theta_t = q in mush.
     FMM: Integral Constraint If J[y] = \int F dx with \int G dx = C then \tilde{J}[y] = \int F - \lambda G dx Hamilto-
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    nian: H:=y'F_{y'}-F\to H'=-F_x. If F=F(y,\dot{y}) then H=C Hamilton's Eqs: p:=F_{y'},q=y
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    and so p' = -H_q, q' = H_p Free Boundary: J[y, b] = \int_a^b F(x, y, y') dx where b free. Expand with y + \epsilon \eta, b + \epsilon \beta \to J = J_0 + \epsilon \left\{ \int_a^b \eta F_y + \eta' F_{y'} dx + \beta F(b, y(b), y'(b)) \right\} If y(b) = d \to d = y(b + \epsilon \beta) + \epsilon \eta(b + \epsilon \beta) = 0
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     y(b) + \epsilon(\beta y'(b) + \eta(b)) so \eta(b) = -\beta y'(b). IVP on integral so \beta [F - y'F_{y'}]_{x=b} + \int (\ldots) = 0 so F = y'F_y
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    at free boundary. Control: Have \int \xi h_x + \eta h_u dt = 0, \dot{\xi} = \xi f_x + \eta f_u. Sub for \eta, IVP s.t. \frac{d}{dt} \frac{h_u}{f_u} = h_x - f_x \frac{h_u}{f_u}
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    and \dot{x} = f Hamiltonian (Control): H := f \frac{h_u}{f_u} - h s.t. \dot{H} = \frac{h_u}{f_u} f_t - h_t \to \text{autonomous if } h_t = f_t = 0
     Fredholm Alt Integ Eqs. For y = f + \int K(x,t)y(t)dt we have ONE (N) has a unique sol y = 0 if f = 0,
     and adjoint has unique sol, or TWO (H) as sols y_1 \dots y_r iff \forall solutions to H^*, z_i, we have \langle f, z_i \rangle = 0. Ex
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    Solve y = f + \lambda \int \sin(x+t)y(t)dt. Unique sol iff (H) has trivial sol \to X_1 = \int y \cos(t) = \int \cos(t)y_H(t) \to \int \sin(x+t)y(t)dt.
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    solve [1, -\lambda \pi; -\lambda \pi, 1][X_1, X_2]^T = [0, 0]^T \to \text{unique sol if } \lambda \neq \pm 1/\pi. In this case X_1 = \int \cos(x) y_N(x) = \int \sin(x) dx
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     \lambda \pi X_2 + \int f(x) \cos(x), and similar for X_2. Invert matrix and solve. If non-unique sol, then find sols to
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    (H) first. If \lambda = 1/\pi then X_1 = X_2 = X so Ly = y - \pi^{-1}(\sin(x) + \cos(x)) \int \cos(x)y(x)dx with sols
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    y = c_1 (\sin(x) + \cos(x)) by inspection. Problem self adjoint so Ly = 0 = L^*w so need \int f(x)w(x) = 0
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    i.e. \int f(x)(\sin(x) + \cos(x)) = 0, repeat for \lambda = -1/\pi. Then y = y_p(x) + \sum_i y_{h,i}(x) Fred Diff Eq. For
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    nonunique sol to exist, need \langle Ly, w \rangle = \langle f, w \rangle \forall ws.t.L^*w = 0
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