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APDE: Charpit: F(p,q,u,x,y) = 0 with u_x = p, u_y = q, \dot{x} = F_p, \dot{y} = F_q. Then via F_x, F_y, \& p_y = q_x \to q
             p_{\tau} = -F_x - pF_u, q_{\tau} = -F_y - qF_u, u_{\tau} = pF_p + qF_q. Also, u_0 = p_0 x_0 + q_0 y_0; F_0 = 0 - last 2 needed to
            show u defined on \Gamma. Max Principle: For -\Delta u = f \leq 0 \to \max u \in \partial D. First show contradiction
            assuming LU = f < 0, then try some auxiliary function \psi = U + \alpha(T_{\text{max}}) g(x_i, y_i) s.t. L\psi < 0 so
           \max \psi = \max_{\epsilon \partial D} \psi. Gets \max e_{i,j}; change to -\alpha for \min e_{i,j}. Laplacian: In 2D : r^{-1} (rf_r)_r + r^{-2} f_{\theta\theta}. In 3D : r^{-2} (r^2 f_r)_r + r^{-2} \sin^{-2}(\theta) f_{\phi\phi} + r^{-2} \sin^{-1}(\theta) (\sin(\theta) f_{\theta})_{\theta} Green's f'n Circle: For G = 0|_{\partial D} we
            have G = \frac{-1}{4\pi} \left( \frac{1}{|x-\xi|} - \frac{1}{|\xi||x-\xi'|} \right) Riemann: For u_{xy} + au_x + bu_y + cu = f we have \int_D RLu - uL^*R = \int_D RLu 
           \int_{D} \partial_{x} \left( Ru_{y} + auR \right) + \partial_{y} \left( -uR_{x} + buR \right) = \int_{\partial D} dy \left( Ru_{y} + Rau \right) + dx \left( uR_{x} - buR \right). Expand over triangle going B-P-A (B at bottom right) \rightarrow need R_{x} = bR@y = \eta, R_{y} = aR@x = \xi, R(P) = 1, L^{*}R = 0. Also
           ensure IVP on \int_B^P dy Ru_y \to Ru|_B^P - \int_B^P dy \ uR_y. Riemann Invariants: If we have \frac{d}{dx}[u-v] = -f
           on y = x + c_1, and \frac{d}{dx}[u + v] = f on y = -x + c_2, then we have: u - v + \int_{-c_1}^x ds \ f(s, s + c_1) = k_1.
           and u+v-\int_{c_2}^x ds\ f(s,-s+c_2)=k_2 for constants k_1,k_2. R-H: Derived via P_x\psi+Q_y\psi=R\psi
            \int_{D} (P\psi)_x + (Q\bar{\psi})_y (= \int_{\Gamma} \psi P dy - \psi Q dx)
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            = \int_{D} P\psi_{x} + Q\psi_{y} + R\psi = \int_{D_{1}+D_{2}} P\psi_{x} + Q\psi_{y} + R\psi, \text{ where } \int_{D_{i}} = \int_{D_{i}} (P\psi)_{x} + (Q\psi)_{y} + \psi (R - P_{x} - Q_{y})
           So \int_{\Gamma} \psi P dy - \psi Q dx = \int_{\Gamma + C_1 - C_2} \psi P dy - \psi Q dx and so \int_{C_1 + C_2} \psi P dy - \psi Q dx = 0 \to dy/dx = [Q]_-^+ / [P]_-^+
           Canonical: For au_{xx}+2bu_{xy}+cu_{yy}=f, we need Cauchy-Kowalevski s.t. first derive defined: x':=\frac{dx}{ds} s.t
           on \Gamma p'_0 = x'_0 u_{xx} + y'_0 u_{xy}, q'_0 = x'_0 u_{xy} + y'_0 u_{yy}. Use these 3, solve det A!=0 s.t. ay'_0^2 - 2bx'_0y'_0 + cx'_0^2 \neq 0. Solve quadratic s.t. b^2 > ac \rightarrow h, b^2 < ac \rightarrow e, b^2 = ac \rightarrow p. H: \lambda_1, \lambda_2 \rightarrow \xi, \eta. E: \lambda = \lambda_R \pm i\lambda_I; \lambda_R \rightarrow \xi, \lambda_I \rightarrow \eta. P: \lambda_1 \rightarrow \xi, choose \eta independent e.g. xy, x^2. Canonical Differentials: u_x = u_\xi \xi_x + u_\eta \eta_x, u_{xx} = u_\xi \xi_x + u_\eta \eta_x, u_{xx} = u_\xi \xi_x + u_\eta \eta_x.
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           u_{\xi\xi}\xi_x^2 + u_{\eta\eta}\eta_x^2 + 2u_{\xi\eta}\xi_x\eta_x + u_{\xi}\xi_{xx} + u_{\eta}\eta_{xx}. Repeat for \partial_y, \partial_{yy} Green's Fn: DON'T USE GREENS THM USE NORMALS For u_{xx} + u_{yy} + au_x + bu_y + cu = f we have \int_D GLu - uL^*G = \int_D (u_xG)_x + (u_yG)_y - uL^*G
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            (uG_x)_x - (uG_y)_y + (auG)_x + (buG)_y = \int_D \nabla \cdot (u_nG - uG_n) + \nabla \cdot \left( (a\ b)^T \hat{n}Gu \right) = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}GG = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G = \int_{\partial D} u_nG - uG_n + (ab)^T \hat{n}G 
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            NB \hat{n} = (dy, -dx). Also note for quarter plane if we have G_x(0,y) = 0, G(x,0) = 0 then we have same
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            sign at \xi_1 = (-x, y), opposite sign at \xi_2 = (x, -y), and for the third we reflect \xi_2 across y axis so we
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           have an opposite sign to \xi at \xi_3 = (-x, -y). Types: Quasi: Coeffs don't depend on highest order derivs
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           Semi: Coeffs depend on x, y. Causality: For a n-dim prob, we have n characteristics. Shock intersects
           2n. \exists k \text{ outgoing, } 2n-k \text{ ingoing. Also have } n \text{ R-H relations, so } 3n-k \text{ pieces of info. Unknowns are } n
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            components of \vec{u} on both sides of shock & slope \Rightarrow 2n+1 unknowns. We demand 3n-k=2n+1 so
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           k = n - 1 outgoing characterisitcs. d'Alembert: Consider triangle A-P-B with AB hypoteneuse. Via
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           \xi=x+t, \eta=x-t we get with R_{\eta}=0 on \xi=p, and R_{\xi}=0 on \eta=q, then via riemann f'n \phi(P)=-\int_{D} rac{t}{4}
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           |J|=2 so \phi(r,s)=-\int_D \frac{f}{2}dxdt. Then have triangle ABP with AP:\eta=q:=r-s \to x-t=r-s, PB:\xi=p:=r+s \to x+t=r+s, and AB:y=0 so finally \phi(r,s)=-\frac{1}{2}\int_0^s dt \int_{r-s+t}^{r+s-t} dx f(x,t)
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           Integral Derivs \frac{d}{dt} \int_{b(t)}^{a(t)} dx f(x,t) = a'(t)f(a,t) - b'(t)f(b,t) + \int_{b(t)}^{a(t)} dt f_t(x,t)
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            SAM: Dists: Need linearity and continuity: \exists N, C \text{ s.t. } |(u,\phi)| \leq C \sum_{m < N} \max_{\in [-X,X]} |\phi^{(m)}|. OR
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           \lim_{n\to\infty}(u,\phi_n)=(u,\lim_{n\to\infty}\phi_n) for a sequence \phi_n\to 0 as n\to\infty. Orthog: \int_0^\pi\sin(kx)\sin(jx)=\frac{\pi}{2}\delta_{kj}
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            same for cos. S-L Operator For T := \alpha y'' + \beta y' + \gamma, multiply by \exp(\int dx \, \beta) to get T_{SL}
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