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APDE: Charpit: F(p,q,u,x,y) = 0 with u_x = p, u_y = q, \dot{x} = F_p, \dot{y} = F_q. Then via F_x, F_y, \& p_y = q_x \rightarrow p_\tau = -F_x - pF_u, q_\tau = -F_y - qF_u, u_\tau = pF_p + qF_q. Also, \frac{du_0}{ds} = p_0 \frac{dx_0}{ds} + q_0 \frac{dy_0}{ds}; F_0 = 0 - last 2 needed to show u defined on Γ. Proof of Charpit: Write \frac{dF}{d\tau} = F_p \frac{dp}{d\tau} + F_q \frac{dq}{d\tau} + F_u \frac{du}{d\tau} + F_x \frac{dx}{d\tau} + F_y \frac{dx}{d\tau}
           F_y \frac{dy}{d\tau} = 0. Then write \phi := u_s - px_s - qy_s. Take \phi_\tau to get \phi_\tau = F_s - \phi F_u + \partial_s (u_\tau - pF_p - qF_q) \to 0
          \phi_{\tau}^{\sigma} = -\phi F_u. Then we have u_{\tau} = x_{\tau}u_x + y_{\tau}y_x, and u_s = x_s p + y_s q which means p := u_x, q := u_x
          u_{y}. Max Principle: For -\Delta u = f \leq 0 \rightarrow \max u \in \partial D. First show contradiction assuming
          LU = f < 0, then try some auxiliary function \psi = U + \alpha(T_{\text{max}})g(x_i, y_i) s.t. L\psi < 0 so \max \psi = 0
          \max_{\epsilon \partial D} \psi. Gets \max_{i,j}; change to -\alpha for \min_{i,j}. Laplacian: In 2D: r^{-1}(rf_r)_r + r^{-2}f_{\theta\theta}. In 3D: r^{-2}(r^2f_r)_r + r^{-2}\sin^{-2}(\theta)f_{\phi\phi} + r^{-2}\sin^{-1}(\theta)(\sin(\theta)f_{\theta})_{\theta} Green's f'n Circle: For G = 0|_{\partial D} we
          have G = \frac{-1}{4\pi} \left( \frac{1}{|x-\xi|} - \frac{1}{|\xi||x-\xi'|} \right) Riemann: For u_{xy} + au_x + bu_y + cu = f we have \int_D RLu - uL^*R
          = \int_D \partial_x \left( Ru_y + auR \right) + \partial_y \left( -uR_x + buR \right) = \int_{\partial D} dy \left( Ru_y + Rau \right) + dx \left( uR_x - buR \right). Expand over triangle going B-P-A (B at bottom right) \rightarrow need R_x = bR@y = \eta, R_y = aR@x = \xi, R(P) = 1, L^*R = 0.
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         Also ensure IVP on \int_B^P dy R u_y \to R u|_B^P - \int_B^P dy \, u R_y. Riemann Invariants: If we have \frac{d}{dx} [u-v] = -f on y = x + c_1, and \frac{d}{dx} [u+v] = f on y = -x + c_2, then we have: u - v + \int_{-c_1}^x ds \, f(s, s + c_1) = k_1,
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          and u+v-\int_{c_2}^x ds \ f(s,-s+c_2) = k_2 for constants k_1,k_2. R-H: Derived via P_x\psi+Q_y\psi=R\psi
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          \int_{D} (P\psi)_{x} + (Q\psi)_{y} \left( = \int_{\Gamma} \psi P dy - \psi Q dx \right) = \int_{D} P\psi_{x} + Q\psi_{y} + R\psi = \int_{D_{1}+D_{2}} P\psi_{x} + Q\psi_{y} + R\psi, \text{ where } \int_{D_{i}} = \int_{D_{i}} (P\psi)_{x} + (Q\psi)_{y} + \psi (R - P_{x} - Q_{y}). \text{ So } \int_{\Gamma} \psi P dy - \psi Q dx = \int_{\Gamma + C_{1} - C_{2}} \psi P dy - \psi Q dx \text{ and so } \int_{\Gamma} \psi P dy + Q\psi_{y} 
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          \int_{C_1+C_2} \psi P dy - \psi Q dx = 0 \to dy/dx = [Q]_-^+/[P]_-^+ Canonical: For au_{xx} + 2bu_{xy} + cu_{yy} = f, we need
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          Cauchy-Kowalevski s.t. first derive defined: x' := \frac{dx}{ds} s.t. on \Gamma p'_0 = x'_0 u_{xx} + y'_0 u_{xy}, q'_0 = x'_0 u_{xy} + y'_0 u_{yy}. Use these 3, solve det A!=0 s.t. ay'_0^2 - 2bx'_0y'_0 + cx'_0^2 \neq 0. Solve quadratic s.t. b^2 > ac \rightarrow h, b^2 < ac \rightarrow e, b^2 = ac \rightarrow e
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          ac \to p. H: \lambda_1, \lambda_2 \to \xi, \eta. E: \lambda = \lambda_R \pm i\lambda_I; \lambda_R \to \xi, \lambda_I \to \eta. P: \lambda_1 \to \xi, choose \eta independent e.g. xy, x^2.
          Canonical Differentials: u_x = u_\xi \xi_x + u_\eta \eta_x, u_{xx} = u_{\xi\xi} \xi_x^2 + u_{\eta\eta} \eta_x^2 + 2u_{\xi\eta} \xi_x \eta_x + u_\xi \xi_{xx} + u_\eta \eta_{xx}. Repeat for
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          \partial_y, \partial_{yy} Green's Fn: DON'T USE GREENS THM USE NORMALS For u_{xx} + u_{yy} + au_x + bu_y + cu = f we
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          have \int_{D} GLu - uL^{*}G = \int_{D} (u_{x}G)_{x} + (u_{y}G)_{y} - (uG_{x})_{x} - (uG_{y})_{y} + (auG)_{x} + (buG)_{y} = \int_{D} \nabla \cdot (u_{n}G - uG_{n}) + (u_{y}G)_{y} + (u_{y}G
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          \nabla \cdot ((a\ b)^T \hat{n} G u) = \int_{\partial D} u_n G - u G_n + (a\ b)^T \hat{n} G. NB \hat{n} = (dy, -dx). Also note for quarter plane if we
          have G_x(0,y) = 0, G(x,0) = 0 then we have same sign at \xi_1 = (-x,y), opposite sign at \xi_2 = (x,-y).
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          and for the third we reflect \xi_2 across y axis so we have an opposite sign to \xi at \xi_3 = (-x, -y). Types:
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          Quasi: Coeffs don't depend on highest order derivs Semi: Coeffs depend on x, y. Causality: For
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          a n-dim prob, we have n characteristics. Shock intersects 2n. \exists k outgoing, 2n-k ingoing. Also
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          have n R-H relations, so 3n-k pieces of info. Unknowns are n components of \vec{u} on both sides of
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           shock & slope \Rightarrow 2n+1 unknowns. We demand 3n-k=2n+1 so k=n-1 outgoing character-
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          isitcs. d'Alembert: Consider triangle A-P-B with AB hypoteneuse. Via \xi = x + t, \eta = x - t we
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          get with R_{\eta}=0 on \xi=p, and R_{\xi}=0 on \eta=q, then via riemann f'n \phi(P)=-\int_{D}\frac{f}{4}. |J|=2
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          so \phi(r,s) = -\int_D \frac{f}{2} dx dt. Then have triangle ABP with AP: \eta = q := r - s \to x - t = r - s,
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          PB: \xi = p := r + s \to x + t = r + s, and AB: y = 0 so finally \phi(r,s) = -\frac{1}{2} \int_0^s dt \int_{r-s+t}^{r+s-t} dx f(x,t)
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          Integral Derivs \frac{d}{dt} \int_{b(t)}^{a(t)} dx f(x,t) = a'(t)f(a,t) - b'(t)f(b,t) + \int_{b(t)}^{a(t)} dt f_t(x,t)
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           SAM: Dists: Need linearity and continuity: \exists N, C \text{ s.t. } |(u,\phi)| \leq C \sum_{m \leq N} \max_{\in [-X,X]} |\phi^{(m)}|. OR
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          \lim_{n\to\infty}(u,\phi_n)=(u,\lim_{n\to\infty}\phi_n) for a sequence \phi_n\to 0 as n\to\infty. Orthog: \int_0^\pi\sin(kx)\sin(jx)=\frac{\pi}{2}\delta_{kj}
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          same for cos. S-L Operator For T := \alpha y'' + \beta y' + \gamma, multiply by \exp(\int dx \, \beta) to get T_{SL}
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