

# Modeling Spontaneous Emission with a Density Matrix and a Quantum Trajectories Approach

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# Problem

- ▶ Quantum computer needs quantum bits (qubits)
- ▶ Ion trap model for qubit
- ▶ Qubit is a two level system, represented by density matrix:

$$\rho = \sum_i p(i) |\psi_i\rangle \langle \psi_i| = \begin{pmatrix} \rho_{ee} & \rho_{eg} \\ \rho_{ge} & \rho_{gg} \end{pmatrix}$$

- ▶  $|\psi_i\rangle$  is either  $|e\rangle$  (excited state) or  $|g\rangle$  (ground state)
- ▶ Following approaches can be extended to quantum systems with  $n$  levels
- ▶ Ion may spontaneously emit a photon and lose energy, transitioning to a lower energy level

# Approach

- ▶ Losing energy to the environment is a non-unitary time evolution
- ▶ Lindblad equation models time evolution of density matrix interacting with environment:

$$\frac{\partial \rho}{\partial t} = -\frac{i}{\hbar}[H, \rho] + \sum_k [L_k \rho, L_k^\dagger] + [L_k, \rho L_k^\dagger]$$

- ▶  $H$  is the system Hamiltonian, and  $L_k$  are Lindblad operators modeling interaction with environment (emission)

# Approach

- ▶ When does the emission happen?
- ▶ Rewritten version of Lindblad equation:

$$\frac{\partial |\psi\rangle}{\partial t} = -\frac{i}{\hbar} H |\psi\rangle - L^\dagger L |\psi\rangle$$

- ▶ Apply emission operator  $L$  at random times
- ▶ Take average of  $|\psi\rangle \langle \psi|$  over many possible trajectories (emission times) to obtain estimate for  $\rho$
- ▶  $|\psi\rangle$  is an  $n$  vector,  $\rho$  is an  $n \times n$  matrix: fewer differential equations to solve

# Linear Multi-step Methods

- ▶ Initial value problem:

$$y' = f(t, y) \text{ s.t. } y(t_0) = y_0$$

- ▶ Linear multi-step method:

$$y_{n+s} = \sum_{k=0}^{s-1} a_k \cdot y_{n+k} + h \cdot \left( \sum_{k=0}^s b_k \cdot f(t_{n+k}, y_{n+k}) \right)$$

- ▶ Discrete-variable method: provide approximations to exact solutions at discrete points  $y(t_i) \approx y_i$
- ▶ Coefficients  $a_0, \dots, a_{s-1}$  and  $b_0, \dots, b_s$  completely determine the method

# Adams-Moulton Method

- ▶ Implicit Adams method for non-stiff system
- ▶ Coefficients  $a_i$  are fixed:

$$a_{s-1} = -1, \quad a_{s-2}, \dots, a_0 = 0$$

- ▶ Coefficients  $b_i$  determined by interpolating polynomial passing through points  $p(t_{n+i}) = f(t_{n+i}, y_{n+i}) \forall i \in \{0, \dots, s-1\}$  using Lagrange polynomial basis

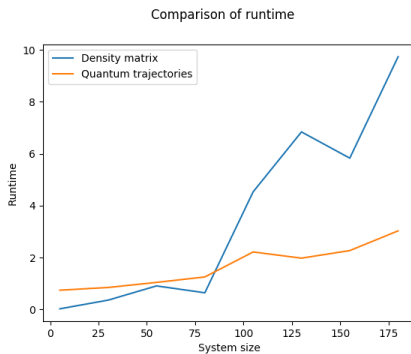
# Adams-Moulton Method

- ▶  $p$  is a locally good approximation for  $y' = f(t, y)$
- ▶ Update step:

$$y_{n+s} = y_{n+s-1} + \int_{t_{n+s-1}}^{t_{n+s}} p(t) dt$$

- ▶ Replacing  $f(t, y)$  by the interpolant  $p$  incurs an error of  $O(h^{s+1})$
- ▶ Implicit Adams method has order  $s + 1$

# Findings

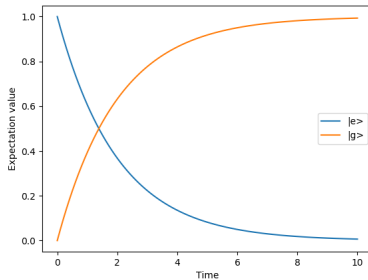


- ▶ Quadratic time complexity for density matrix approach
- ▶ Linear for quantum trajectories

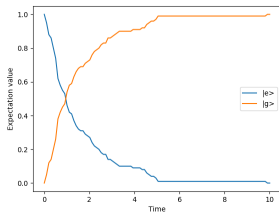


# Findings

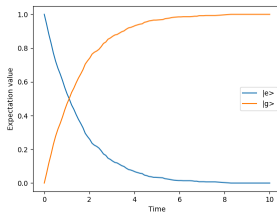
Density Matrix: Two Level System



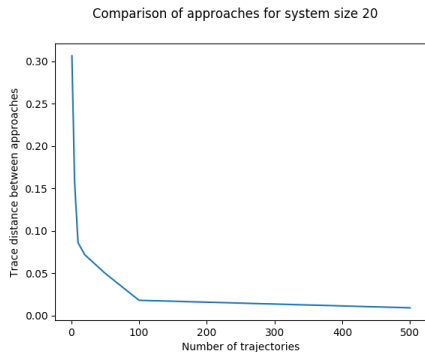
Quantum Trajectories: Two Level System, 100 Trajectories



Quantum Trajectories: Two Level System, 1000 Trajectories



# Findings



- Trace distance:

$$T(\rho, \sigma) = \frac{1}{2} \text{Tr}|\rho - \sigma| = \frac{1}{2} \sum_i |\lambda_i|$$

# Findings

- ▶ Above findings confirm the theoretical equivalence of the two approaches as number of trajectories is increased
- ▶ Different Hamiltonians model different physical systems (i.e. qubit, truncated harmonic oscillator, etc.)
- ▶ Varying Hamiltonian does not change running time or difference between approaches much (i.e.  $H = 0$ ,  $H = \text{constant}$ ,  $H = H(t)$ )
- ▶ Approaches can be used to solve systems which cannot be solved analytically (i.e. perturbed harmonic oscillator)

# References

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- 6 *Quantum Noise*, C. W. Gardiner and P. Zoller
- 7 *Quantum toolbox in Python (QuTiP)*
- 8 *Quantum trajectories for time-dependent adiabatic master equations*
- 9 *Solving Ordinary Differential Equations I*, Ernst Hairer, et al.