

Nonequilibrium quantum criticality of interacting Dirac fermions

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Background

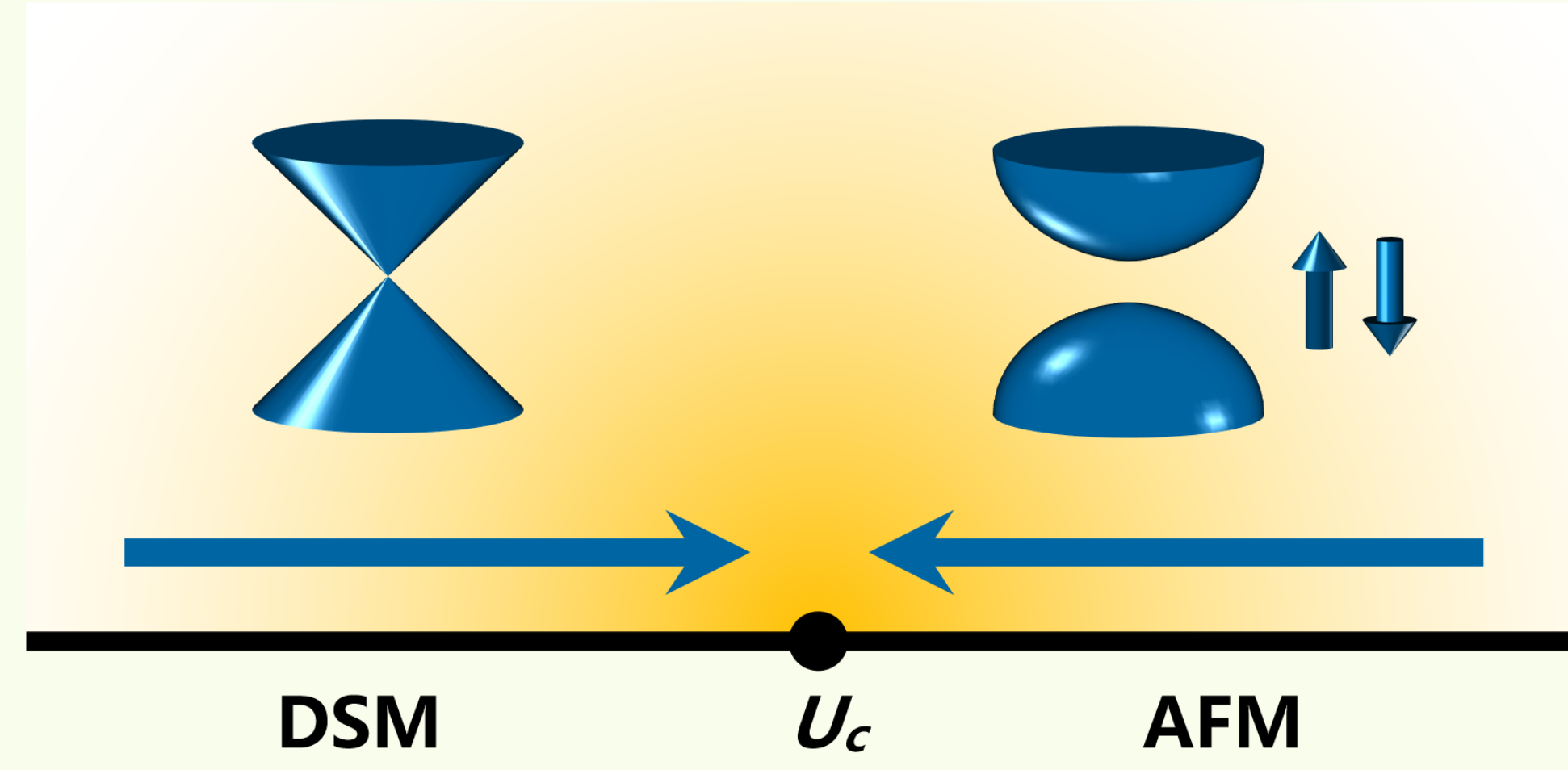
Current status in this field

- Nonequilibrium scaling has been developed in various classical or quantum systems, but rarely studied in strongly interacting fermionic systems[1–3].
- Dirac fermions with linear dispersion matter in phase transition widely in graphene, Weyl/Dirac semimetal, surface of topological insulator[4, 5].

This work

- For the first time, we demonstrate the nonequilibrium dynamics of chiral Heisenberg universality class for Dirac fermions.

Model and protocol



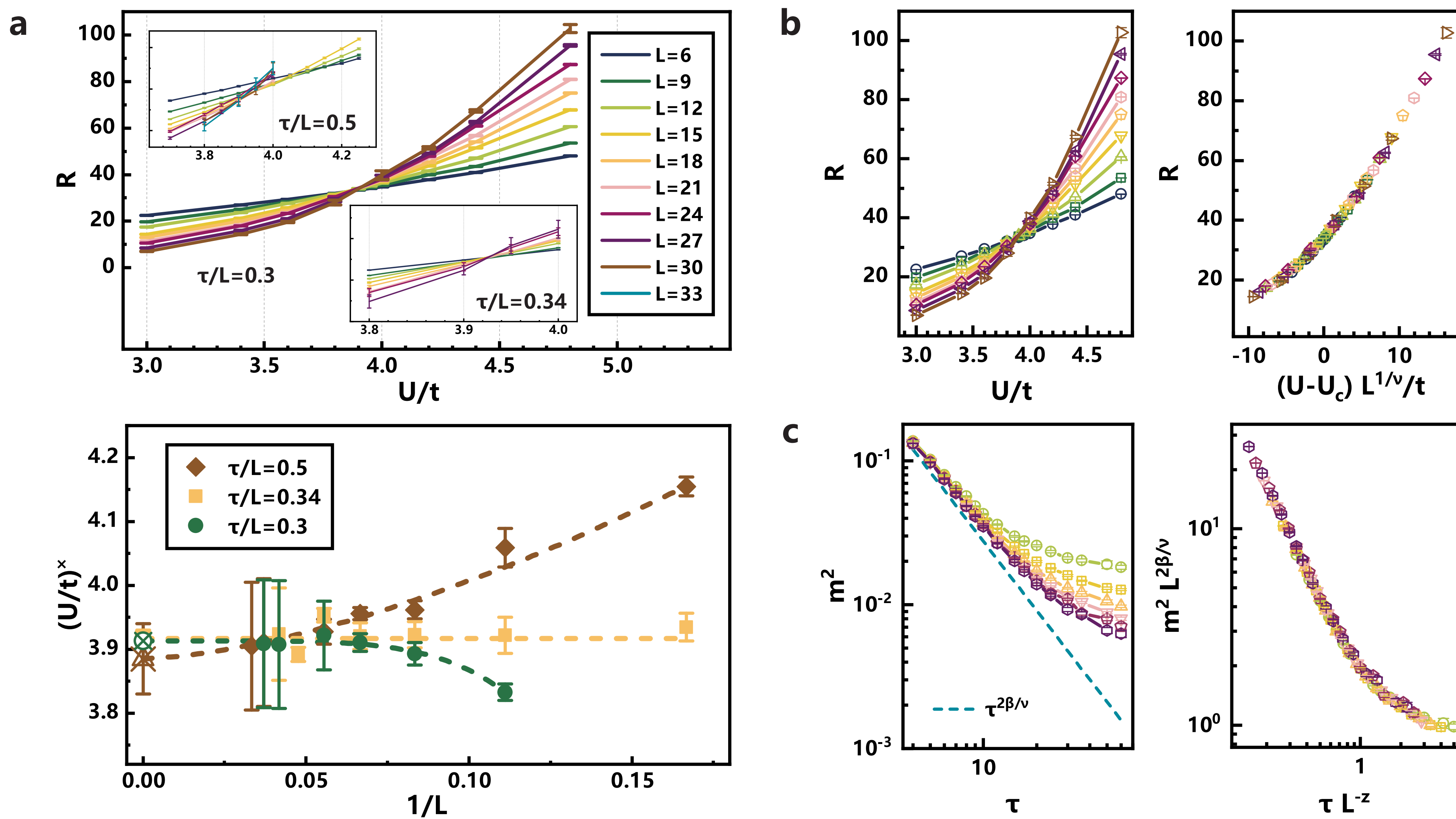
Half-filled honeycomb Hubbard model ($t = 1$)

$$H = -t \sum_{\langle ij \rangle, \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i \left(n_{i\uparrow} - \frac{1}{2} \right) \left(n_{i\downarrow} - \frac{1}{2} \right).$$

Steps

- Prepare initial states of DSM and AFM.
- Quench them to the quantum critical point (QCP) U_c .
- Explore the imaginary-time relaxation by determinant quantum Monte Carlo[6].

Imaginary-time relaxation dynamics with the AFM ordered initial state



We are concerned about these observable quantities:
 AFM structure factor $S(\mathbf{q}) = \frac{1}{L^2} \sum_{i,j} e^{i\mathbf{q} \cdot (\mathbf{r}_i - \mathbf{r}_j)} \langle m_i^{(z)} m_j^{(z)} \rangle$,
 AFM order parameter $m^2 = S(\mathbf{0})$,
 Correlation ratio $R = S(\mathbf{0}) / S(\Delta\mathbf{q})$.
 Imaginary-time correlation $A = \frac{1}{L^2} \sum_i \langle m_i^{(z)}(0) m_i^{(z)}(\tau) \rangle$.

a. Estimation of quantum critical point U_c

$$R(g, \tau, L) = f_R(g L^{1/\nu}, \tau^{-1} L^z).$$

Fixed $\tau^{-1} L^z$, curves intersect at $g = U - U_c = 0$.
 \Rightarrow QCP: $U/t = 3.91 \pm 0.03$.

b. Fitting for ν

Fixing $\tau^{-1} L^z = 0.3$, $R(g) = f_R(g L^{1/\nu})$,
 $\Rightarrow \nu = 1.17 \pm 0.07$.

c. Fitting for β/ν

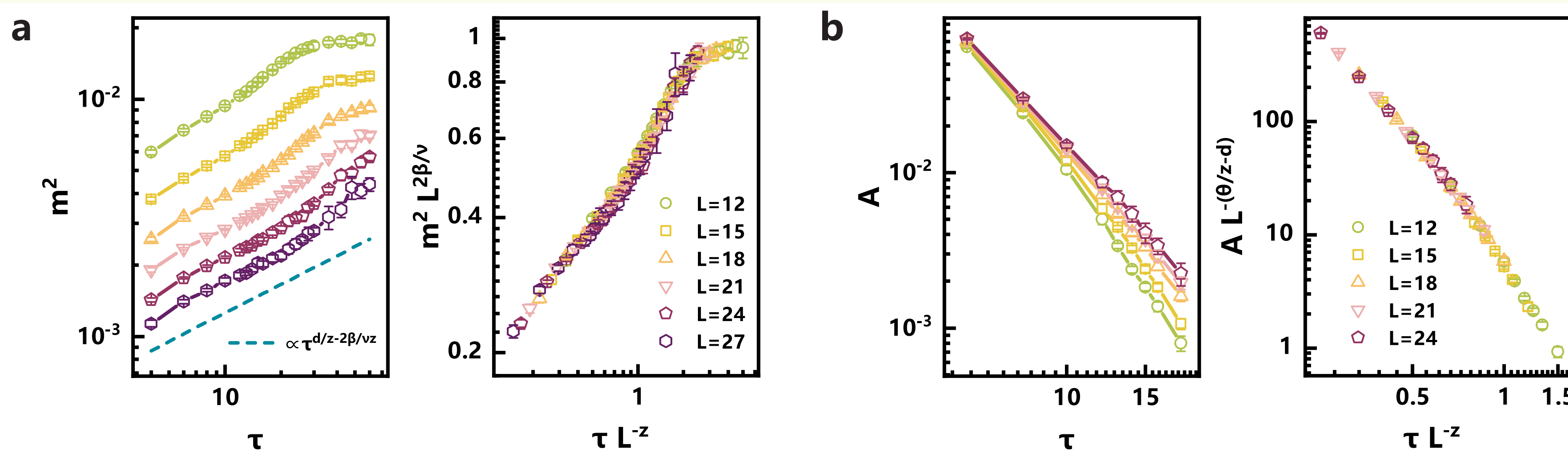
Saturated AFM initialized, $g = 0$ set,

$$m^2 = \tau^{-2\beta/\nu z} f_m(\tau^{-1} L^z),$$

$\Rightarrow \beta/\nu = 0.80 \pm 0.03$.

Limit $\tau L^{-z} \rightarrow 0$, $m^2 \sim \tau^{-2\beta/\nu z} + \tau^{-2\beta/\nu z} \mathcal{O}(\tau L^{-z})$.

Imaginary-time relaxation dynamics with the disordered initial state



a. Examination for β/ν

Initial DSM at $g = 0$. Here we take $\beta/\nu = 0.80$,

$$m^2 = L^{-d} \tau^{d/z - 2\beta/\nu z} f_m(\tau^{-1} L^z).$$

Compare with zero-order approximation $m^2 \sim \tau^{d/z - 2\beta/\nu z}$.

b. Fitting for θ

Random initial state at $g = 0$,

$$A = L^{\theta z - d} f_A(\tau^{-1} L^z).$$

$\Rightarrow \theta = -0.84 \pm 0.04$.

Discussion and conclusion

Comparison with previous works

model	method	U_c/t	β/ν	θ
honeycomb	QMC (present)	3.91(3)	0.80(3)	-0.84(4)
honeycomb	QMC[5]	3.85(2)	0.75(2)	-
honeycomb	QMC[4]	3.77(4)	0.8(1)	-
Gross-Neveu	$4 - \epsilon$ (1st order)[7]	-	0.945	-
Gross-Neveu	$4 - \epsilon$ (2nd order)[7]	-	0.985	-
Gross-Neveu	FRG[8]	-	1.008	-

Conclusion

- A negative critical initial slip exponent $\theta = -0.84(4)$ for chiral Heisenberg universality class.
- An efficient dynamic QMC method to determine the critical point and critical exponents.
- A new avenue for numerical studies on nonequilibrium quantum criticality of fermionic systems.



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