

Solution to Exercise 16.2 from Spin Dynamics (2nd, Malcolm H. Levitt)

Zheng Zuo

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The Problem 16.2 in book Spin Dynamics (p.451-452)

Pulse sequence for 2QF-COSY (double-quantum-filtered COSY)

Calculation of density product operator(s)

(i, ii, iii) The product operators in cosine version pulse sequence (cycle $m=0$)

(iv) The product operators for the sine version pulse sequences

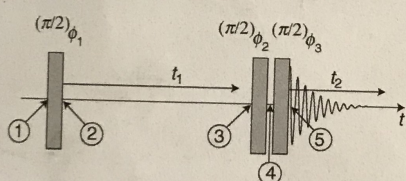
Construction of phase-sensitive 2D peaks according to States Method

(v) Using initial density operator \hat{I}_{2z}

The form of DQF-COSY peaks

The Problem 16.2 in book Spin Dynamics (p.451-452)

16.2 This exercise investigates the *double-quantum-filtered COSY (2QF-COSY) experiment*. The 'cosine' pulse sequence is as follows:



The phase table is as follows:

Cycle counter m	ϕ_1	ϕ_2	ϕ_3	ϕ_{rec}
0	$3\pi/2$	0	$\pi/2$	0
1	$3\pi/2$	0	π	$3\pi/2$
2	$3\pi/2$	0	$3\pi/2$	π
3	$3\pi/2$	0	0	$\pi/2$

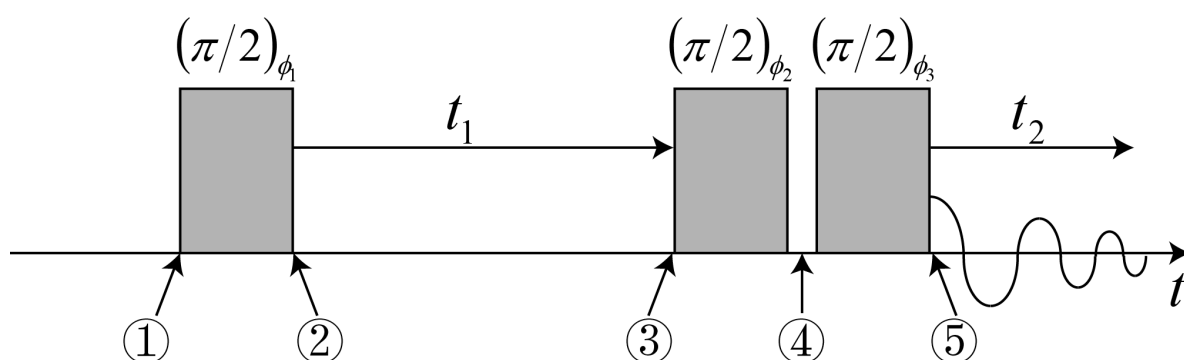
Assume a homonuclear AX spin system. The short delay between the last two $\pi/2$ pulses may be ignored.

- Take the first step in the phase cycle procedure ($m = 0$). For simplicity, start with an initial spin density operator $\hat{\rho}_1 \sim \hat{I}_{1z}$. What is the spin density operator at time point ④?
- Convert the x - and z -operators at time point ④ into shift operators. Which coherence orders are excited at this point?
- The phase cycle has the effect of suppressing all signal components that do not pass through (± 2) -quantum coherences at time point ④. Select the operators that have order ± 2 and apply the final $\pi/2$ pulse. What is the spin density operator at time point ⑤? To bring out the meaning

of the final expression, replace all products of trigonometric functions by single trigonometric functions.

- (iv) Repeat the calculation for the 'sine' component of the States procedure, in which the first pulse has phase π instead of phase $3\pi/2$. Use the same initial density operator.
- (v) Repeat the calculations using the full form of the initial spin density operator, $\hat{\rho}_{(1)} \sim \hat{I}_{1z} + \hat{I}_{2z}$. Sketch the form of the two-dimensional spectrum and remark on the shapes of the diagonal peaks and cross-peaks. What are the favourable properties of the double-quantum-filtered COSY spectrum? Are there any disadvantages of double-quantum-filtered COSY compared with ordinary COSY?

Pulse sequence for 2QF-COSY (double-quantum-filtered COSY)



Cycle counter m	ϕ_1	ϕ_2	ϕ_3	ϕ_{rec}
0	$3\pi/2$	0	$\pi/2$	0
1	$3\pi/2$	0	π	$3\pi/2$
2	$3\pi/2$	0	$3\pi/2$	π
3	$3\pi/2$	0	0	$\pi/2$

Calculation of density product operator(s)

(i, ii, iii) The product operators in cosine version pulse sequence (cycle $m=0$)

$$\begin{aligned}
\hat{\rho}_1 &= \hat{I}_{1z} \\
&\downarrow (\pi/2)_{-y} \\
\hat{\rho}_2 &= -\hat{I}_{1x} \\
&\downarrow U_2 \\
&= -\hat{I}_{1x} \\
&\downarrow U_1 \\
&= -\hat{I}_{1x} \cos(\Omega_1 t_1) - \hat{I}_{1y} \sin(\Omega_1 t_1) \\
&\downarrow J_{12} \\
\hat{\rho}_3 &= -\cos(\Omega_1 t_1) [\cos(\pi J t_1) \hat{I}_{1x} + \sin(\pi J t_1) 2\hat{I}_{1y} \hat{I}_{2z}] \\
&\quad + \sin(\Omega_1 t_1) [-\cos(\pi J t_1) \hat{I}_{1y} + \sin(\pi J t_1) 2\hat{I}_{1x} \hat{I}_{2z}] \\
&\downarrow (\pi/2)_x \\
\hat{\rho}_4 &= -\cos(\Omega_1 t_1) [\cos(\pi J t_1) \hat{I}_{1x} - \sin(\pi J t_1) 2\hat{I}_{1z} \hat{I}_{2y}] \\
&\quad + \sin(\Omega_1 t_1) [-\cos(\pi J t_1) \hat{I}_{1z} - \underbrace{\sin(\pi J t_1) 2\hat{I}_{1x} \hat{I}_{2y}}_{\text{Double coherence}}] \\
\hat{\rho}_{4,\pm 2} &= -\sin(\Omega_1 t_1) \sin(\pi J t_1) 2\hat{I}_{1x} \hat{I}_{2y} \\
&\downarrow (\pi/2)_y \\
\hat{\rho}_{5,\cos} &= \sin(\Omega_1 t_1) \sin(\pi J t_1) 2\hat{I}_{1z} \hat{I}_{2y} \\
&= [\cos(\Omega_1 t_1 - \pi J t_1) - \cos(\Omega_1 t_1 + \pi J t_1)] \hat{I}_{1z} \hat{I}_{2y}
\end{aligned}$$

(iv) The product operators for the sine version pulse sequences

$$\begin{aligned}
\hat{\rho}_1 &= \hat{I}_{1z} \\
&\downarrow (\pi/2)_{-x} \\
\hat{\rho}_2 &= -\hat{I}_{1y} \\
&\downarrow U_2 \\
&= -\hat{I}_{1y} \\
&\downarrow U_1 \\
&= -\hat{I}_{1y} \cos(\Omega_1 t_1) - \hat{I}_{1x} \sin(\Omega_1 t_1) \\
&\downarrow J_{12} \\
\hat{\rho}_3 &= \cos(\Omega_1 t_1) [\cos(\pi J t_1) \hat{I}_{1y} - \sin(\pi J t_1) 2\hat{I}_{1x} \hat{I}_{2z}] \\
&\quad - \sin(\Omega_1 t_1) [\cos(\pi J t_1) \hat{I}_{1x} + \sin(\pi J t_1) 2\hat{I}_{1y} \hat{I}_{2z}] \\
&\downarrow (\pi/2)_x \\
\hat{\rho}_4 &= \cos(\Omega_1 t_1) [\cos(\pi J t_1) \hat{I}_{1z} + \underbrace{\sin(\pi J t_1) 2\hat{I}_{1x} \hat{I}_{2y}}_{\text{Double coherence}}] \\
&\quad - \sin(\Omega_1 t_1) [\cos(\pi J t_1) \hat{I}_{1x} - \sin(\pi J t_1) 2\hat{I}_{1z} \hat{I}_{2y}] \\
\hat{\rho}_{4,\pm 2} &= \cos(\Omega_1 t_1) \sin(\pi J t_1) 2\hat{I}_{1x} \hat{I}_{2y} \\
&\downarrow (\pi/2)_y \\
\hat{\rho}_{5,\sin} &= -\cos(\Omega_1 t_1) \sin(\pi J t_1) 2\hat{I}_{1z} \hat{I}_{2y} \\
&= [\sin(\Omega_1 t_1 - \pi J t_1) - \sin(\Omega_1 t_1 + \pi J t_1)] \hat{I}_{1z} \hat{I}_{2y}
\end{aligned}$$

Construction of phase-sensitive 2D peaks according to States Method

$$\begin{aligned}
\hat{\rho}_{States} &= \hat{\rho}_{\cos} + i \hat{\rho}_{\sin} \\
&= [\exp^{i(\Omega_1 - \pi J)t_1} - \exp^{i(\Omega_1 + \pi J)t_1}] \hat{I}_{1z} \hat{I}_{2y} \\
&= [\exp^{i(\Omega_1 - \pi J)t_1} - \exp^{i(\Omega_1 + \pi J)t_1}] \left[\underbrace{\frac{i}{4} \hat{I}_1^\alpha \hat{I}_2^-}_{\hat{\rho}_{\alpha-}} - \underbrace{\frac{i}{4} \hat{I}_1^\beta \hat{I}_2^-}_{\hat{\rho}_{\beta-}} + \dots \right]
\end{aligned}$$

Density operator $\hat{I}_{1z} \hat{I}_{2y}$ indicates an antiphase absorption peak, centered around frequency Ω_2^0 in the Ω_2 dimension. Therefore, equation signifies the doubly antiphase cross-peak, just like in the normal COSY scenario.

(v) Using initial density operator \hat{I}_{2z}

For the full form of initial spin density operator $\hat{\rho}_1 \sim \hat{I}_{1z} + \hat{I}_{2z}$, if we start with \hat{I}_{2z} ,

$$\begin{aligned}
\hat{\rho}_1 &= \hat{I}_{2z} \\
&\downarrow (\pi/2)_{-y} \\
\hat{\rho}_2 &= -\hat{I}_{2x} \\
&\downarrow U_2 \\
&= -\hat{I}_{2x} \cos(\Omega_1 t_1) - \hat{I}_{2y} \sin(\Omega_1 t_1) \\
&\downarrow U_1 \\
&= -\hat{I}_{2x} \cos(\Omega_1 t_1) - \hat{I}_{2y} \sin(\Omega_1 t_1) \\
&\downarrow J_{12} \\
\hat{\rho}_3 &= -\cos(\Omega_1 t_1) [\cos(\pi J t_1) \hat{I}_{2x} + \sin(\pi J t_1) 2\hat{I}_{1z} \hat{I}_{2y}] \\
&\quad + \sin(\Omega_1 t_1) [-\cos(\pi J t_1) \hat{I}_{2y} + \sin(\pi J t_1) 2\hat{I}_{1z} \hat{I}_{2x}] \\
&\downarrow (\pi/2)_x \\
\hat{\rho}_4 &= -\cos(\Omega_1 t_1) [\cos(\pi J t_1) \hat{I}_{2x} - \sin(\pi J t_1) 2\hat{I}_{1y} \hat{I}_{2z}] \\
&\quad + \sin(\Omega_1 t_1) [-\cos(\pi J t_1) \hat{I}_{2z} - \underbrace{\sin(\pi J t_1) 2\hat{I}_{1y} \hat{I}_{2x}}_{\text{Double coherence}}] \\
\hat{\rho}_{4,\pm 2} &= -\sin(\Omega_1 t_1) \sin(\pi J t_1) 2\hat{I}_{1y} \hat{I}_{2x} \\
&\downarrow (\pi/2)_y \\
\hat{\rho}_{5,\cos} &= \sin(\Omega_1 t_1) \sin(\pi J t_1) 2\hat{I}_{1y} \hat{I}_{2z} \\
&= [\cos(\Omega_1 t_1 - \pi J t_1) - \cos(\Omega_1 t_1 + \pi J t_1)] \hat{I}_{1y} \hat{I}_{2z}
\end{aligned}$$

Similarly, for the sine version of pulse sequence,

$$\begin{aligned}
\hat{\rho}_1 &= \hat{I}_{2z} \\
&\downarrow_{(\pi/2)_{-x}} \\
\hat{\rho}_2 &= -\hat{I}_{2y} \\
&\downarrow_{U2} \\
&= \hat{I}_{2y} \cos(\Omega_1 t_1) - \hat{I}_{2x} \sin(\Omega_1 t_1) \\
&\downarrow_{U1} \\
&= \hat{I}_{2y} \cos(\Omega_1 t_1) - \hat{I}_{2x} \sin(\Omega_1 t_1) \\
&\downarrow_{J_{12}} \\
\hat{\rho}_3 &= \cos(\Omega_1 t_1) [\cos(\pi J t_1) \hat{I}_{2y} - \sin(\pi J t_1) 2 \hat{I}_{1z} \hat{I}_{2x}] \\
&\quad - \sin(\Omega_1 t_1) [\cos(\pi J t_1) \hat{I}_{2x} + \sin(\pi J t_1) 2 \hat{I}_{1z} \hat{I}_{2y}] \\
&\downarrow_{(\pi/2)_x} \\
\hat{\rho}_4 &= \cos(\Omega_1 t_1) [\cos(\pi J t_1) \hat{I}_{2z} + \underbrace{\sin(\pi J t_1) 2 \hat{I}_{1y} \hat{I}_{2x}}_{\text{Double coherence}}] \\
&\quad - \sin(\Omega_1 t_1) [\cos(\pi J t_1) \hat{I}_{2x} - \sin(\pi J t_1) 2 \hat{I}_{1y} \hat{I}_{2z}] \\
\hat{\rho}_{4,\pm 2} &= \cos(\Omega_1 t_1) \sin(\pi J t_1) 2 \hat{I}_{1y} \hat{I}_{2x} \\
&\downarrow_{(\pi/2)_y} \\
\hat{\rho}_{5,\sin} &= -\cos(\Omega_1 t_1) \sin(\pi J t_1) 2 \hat{I}_{1y} \hat{I}_{2z} \\
&= [\sin(\Omega_1 t_1 - \pi J t_1) - \sin(\Omega_1 t_1 + \pi J t_1)] \hat{I}_{1y} \hat{I}_{2z}
\end{aligned}$$

Overall,

$$\begin{aligned}
\hat{\rho}_{States} &= \hat{\rho}_{\cos} + i \hat{\rho}_{\sin} \\
&= [\exp^{i(\Omega_1 - \pi J)t_1} - \exp^{i(\Omega_1 + \pi J)t_1}] \hat{I}_{1y} \hat{I}_{2z} \\
&= [\exp^{i(\Omega_1 - \pi J)t_1} - \exp^{i(\Omega_1 + \pi J)t_1}] \left[\underbrace{\frac{i}{4} \hat{I}_1^- \hat{I}_2^\alpha}_{\hat{\rho}_{-\alpha}} - \underbrace{\frac{i}{4} \hat{I}_1^- \hat{I}_2^\beta}_{\hat{\rho}_{-\beta}} + \dots \right]
\end{aligned}$$

Density operator $\hat{I}_{1y} \hat{I}_{2z}$ indicates an antiphase absorption peak, centered around frequency Ω_1^0 in the Ω_2 dimension. Therefore, equation signifies the doubly antiphase diagonal peaks, which is different from the normal COSY scenario.

The form of DQF-COSY peaks

