

SI Prefixes

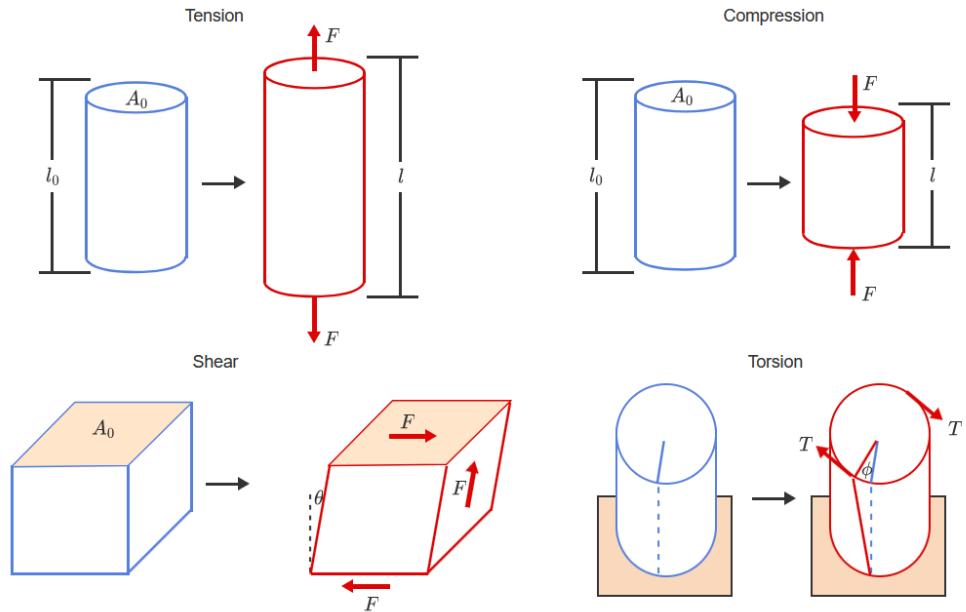
femto (f)	$= 10^{-15}$
pico (p)	$= 10^{-12}$
nano (n)	$= 10^{-9}$
micro (μ)	$= 10^{-6}$
milli (m)	$= 10^{-3}$
centi (c)	$= 10^{-2}$
deci (d)	$= 10^{-1}$
deca (da)	$= 10^1$
hecto (h)	$= 10^2$
kilo (k)	$= 10^3$
mega (M)	$= 10^6$
giga (G)	$= 10^9$
tera (T)	$= 10^{12}$
peta (P)	$= 10^{15}$

Note: 1Angstrom (\AA) = 10^{-10} m

Mechanical Properties

Stress and Strain

Types of loading



Stress (Force Normalized by Area)

Tensile and Compression Stress

$$\sigma = \frac{F}{A_o} \quad (\text{in units of pressure})$$

Shear Stress

$$\tau = \frac{F}{A_o}$$

Strain (Displacement Normalized by Original Length)

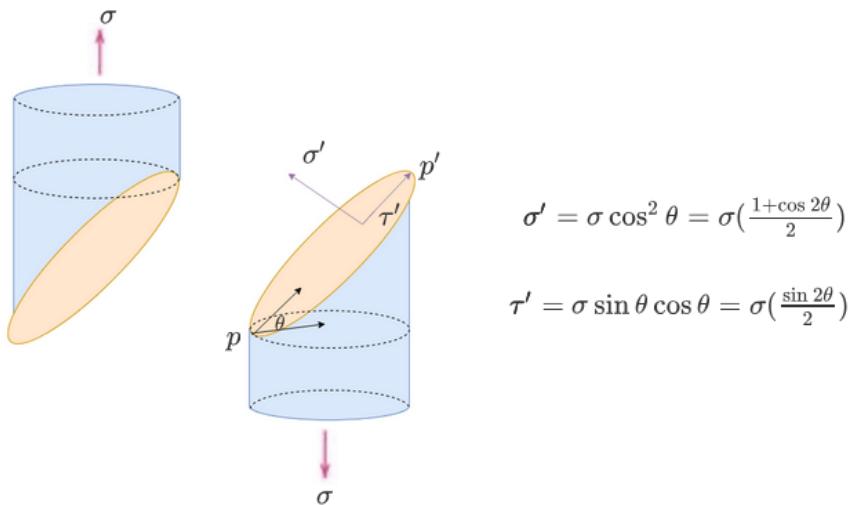
Tensile and Compression Strain

$$\epsilon = \frac{l - l_0}{l_0}$$

Shear Strain

$$\gamma = \tan \theta \quad (\theta \text{ is the shear angle})$$

Normal and Shear Stress Along an Angled Plane



Elastic Deformation

Relationship Between Stress and Strain

Tensile and Compression

$$\sigma = E\epsilon \quad (\text{E (GPa or psi) is the modulus of elasticity})$$

Shear

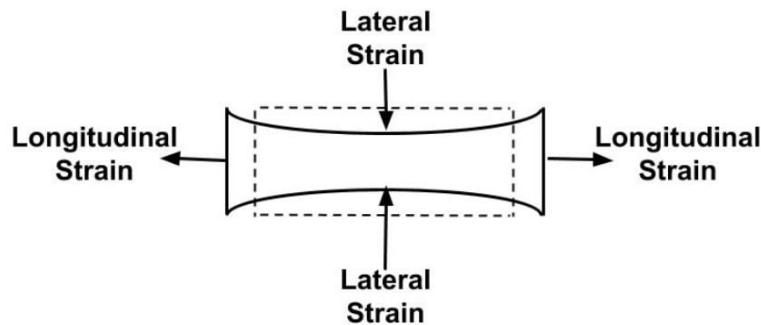
$$\tau = G\gamma \quad (G \text{ is the shear modulus})$$

Note: The modulus of elasticity (Young's modulus) is the slope of the stress - strain plot. (It describes a material's resistance to elastic deformation. Stiffer \Rightarrow higher E)

Anelasticity: time dependent elastic strain, where deformation and recovery is not instantaneous.

Viscoelastic behavior: materials (such as polymers) with significant anelasticity

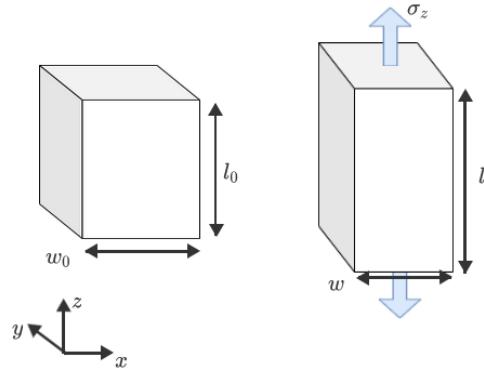
Poisson's Ratio



$$\nu = -\frac{\epsilon_{\text{lateral}}}{\epsilon_{\text{longitudinal}}}$$

Note: Lateral is perpendicular to the direction of loading and longitudinal is along the direction of loading

Example: Rectangular prism



If the applied stress is uniaxial (only along 1 axis) and the material is isotropic (constant properties regardless of direction), then for a σ_z , $\epsilon_x = \epsilon_y$

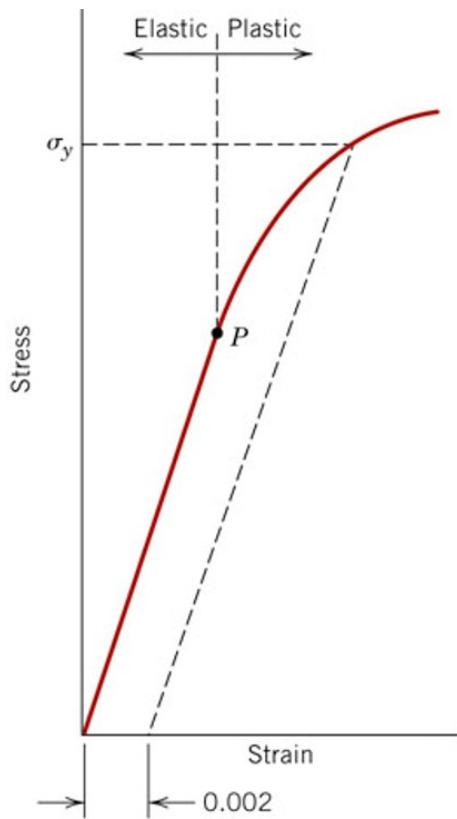
$$\nu = -\frac{\epsilon_x}{\epsilon_z} = -\frac{\epsilon_y}{\epsilon_z}$$

Relating modulus of elasticity, shear modulus and Poisson's ratio

$$E = 2G(1 + \nu)$$

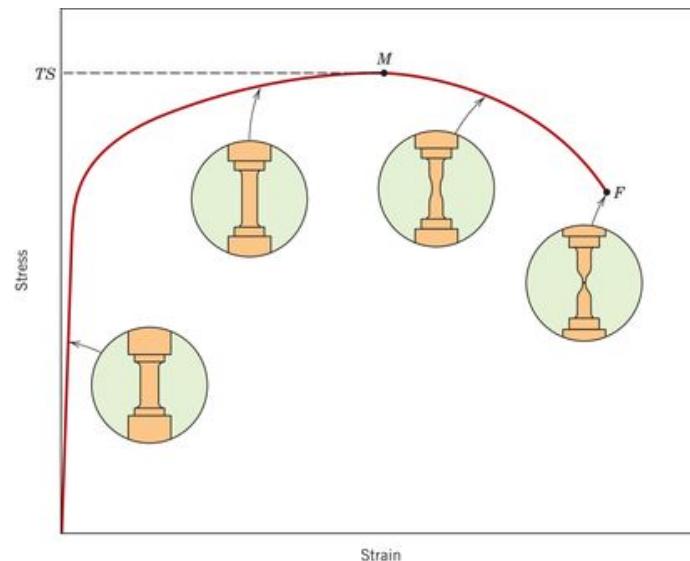
Note: Some materials (like foams) expand under tension so they have a negative Poisson's ratio, these materials are called **auxetics**.

Plastic Deformation



Point P is the **Proportional Limit** where the exact departure from linearity occurs and deformation becomes permanent.

Yield Stress (σ_y): stress at which noticeable strain has occurred (0.002)



Tensile Strength: Stress at the maximum point on the stress - strain plot. After this point, necking occurs and all deformation is focused at the neck until fracture (point F)

Ductility

As % elongation:

$$\%EL = \frac{l_f - l_0}{l_0} \times 100$$

As % reduction in area

$$\%RA = \frac{A_0 - A_f}{A_0} \times 100$$

l_f and A_f are length and cross-sectional area of sample at fracture respectively.

Resilience: capacity of a material to absorb energy when it is deformed elastically and unloaded (similar to spring potential energy)

Modulus of Resilience

$$U_r = \int_0^{\epsilon_{yield}} \sigma d\epsilon$$

Area under the stress - strain plot from 0 to yield point

Assuming a linear elastic region:

$$U_r = \frac{1}{2} \sigma_y \epsilon_y$$

Crystal Structures

Atomic Packing Factor

$$APF = \frac{\text{Volume of atoms in unit cell}}{\text{Total unit cell volume}}$$

Packing Fraction

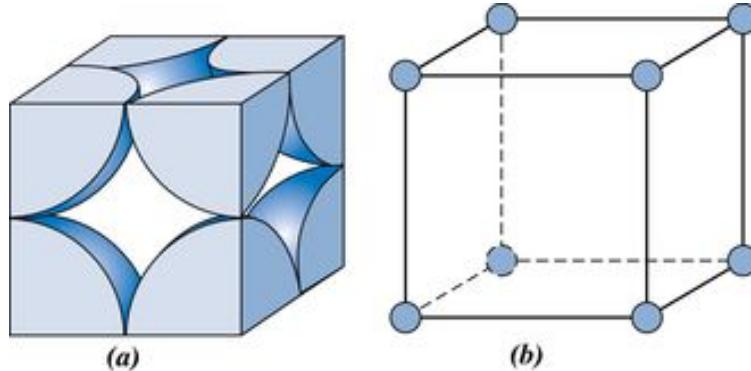
$$PF = \frac{\text{Total Cross Sectional Area of Atoms}}{\text{Total Area of Plane}}$$

Number of Atoms per Unit Cell

$$N = N_i + \frac{N_f}{2} + \frac{N_c}{8}$$

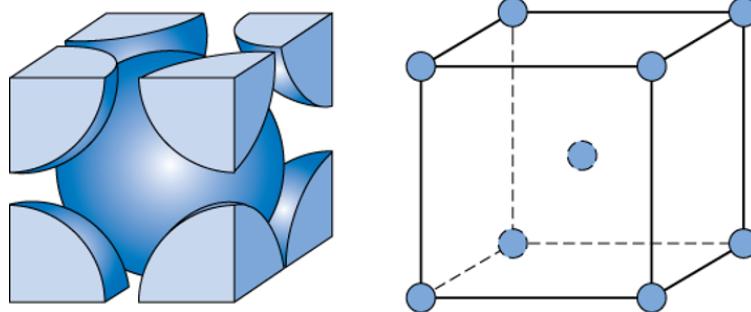
N_i are interior atoms, N_f are face atoms and N_c are corner atoms

Simple Cubic

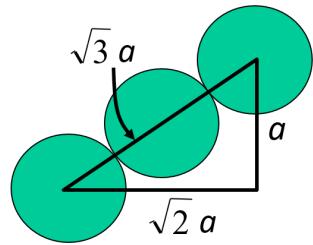


$$2R = a$$
$$APF = \frac{(\# \text{ atoms})(\text{volume/atom})}{(\text{volume/unit cell})}$$
$$APF = \frac{(1)(\frac{4}{3}\pi(a/2)^3)}{a^3}$$

Body Centered Cubic



Triangle formed along the main diagonal and face diagonal

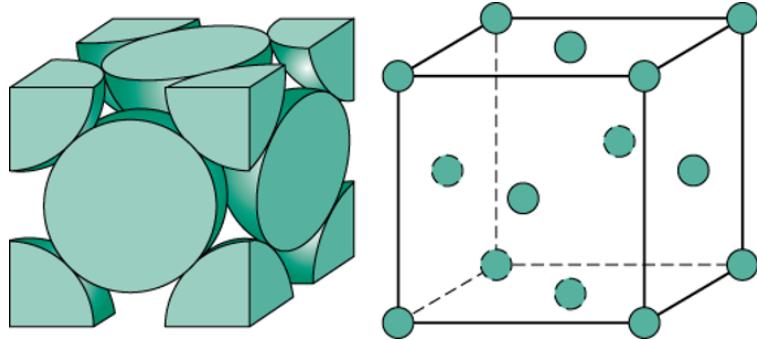


$$4R = \sqrt{3}a$$

$$APF = \frac{(\# \text{ atoms})(\text{volume/atom})}{(\text{volume/unit cell})}$$

$$APF = \frac{(2)(\frac{4}{3}\pi(\sqrt{3}/4a)^3)}{a^3}$$

Face Centered Cubic



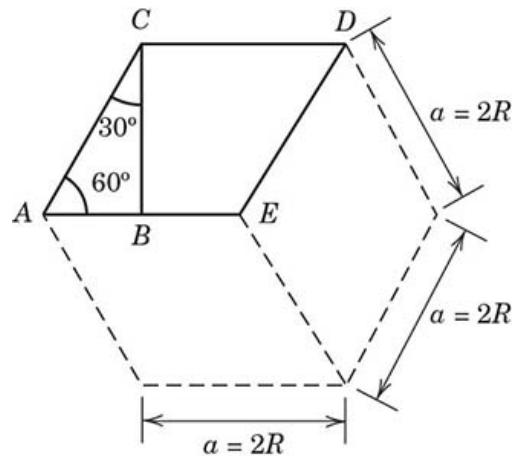
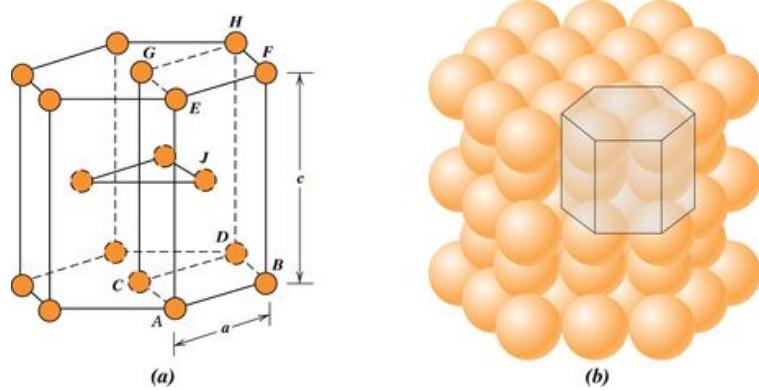
Along the face diagonal:

$$4R = \sqrt{2}a$$

$$APF = \frac{(\# \text{ atoms})(\text{volume/atom})}{(\text{volume/unit cell})}$$

$$APF = \frac{(4)(\frac{4}{3}\pi(\sqrt{2}/4a)^3)}{a^3}$$

Hexagonal Close Packed



Area of base hexagon is 3 parallelograms or 6 equilateral triangles:

$$\text{Area} = \frac{3a^2\sqrt{3}}{2}$$

Given height c:

$$\text{Volume of unit cell} = \frac{3a^3\sqrt{3}}{2}$$

Theoretical Density for Crystals

$$\rho = \frac{(\text{atoms/unit cell})(\text{g/mol})}{(\text{vol/unit cell})(\text{atoms/mol})} = (\text{g/vol})$$

$$\rho = \frac{nA}{V_c N_A}$$

Where:

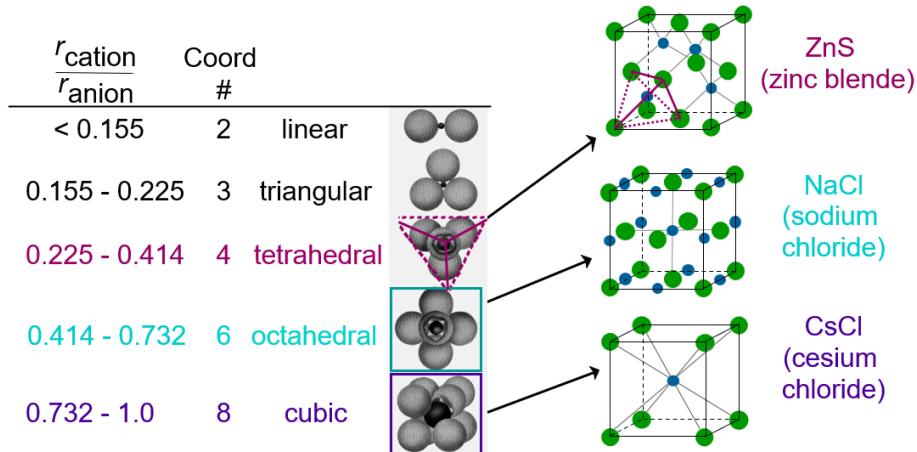
- n = # of atoms in unit cell
- A = atomic weight
- V_c = volume of unit cell
- N_a = Avogadro's number (6.022×10^{23} atoms/mol)

Ceramic Crystal Structures

Factors that determine crystal structure:

- Relative sizes of ions ($\frac{r_{\text{cation}}}{r_{\text{anion}}}$)
- Maintenance of charge neutrality (Net charge in ceramic is zero)

Note: As $\frac{r_{\text{cation}}}{r_{\text{anion}}}$ increases, so does coordination number



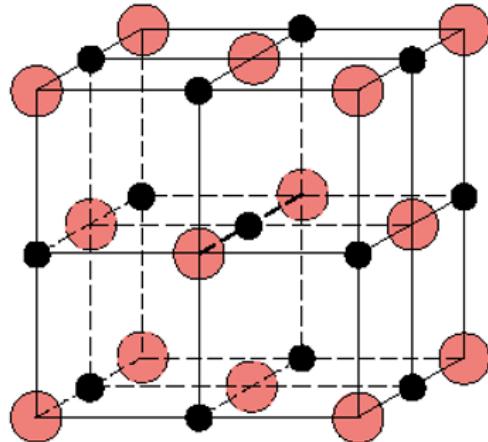
Theoretical Density for Ceramics

$$\rho = \frac{n'(\sum A_C + \sum A_A)}{V_c N_A}$$

Where:

- n' = # of Atoms per unit cell (For AX structures, this is equal for cations and anions)
- $\sum A_A$ = sum of cation molar mass
- $\sum A_C$ = sum of anion molar mass
- V_c = volume of unit cell
- N_A = Avogadro's number (6.022×10^{23} atoms/mol)

Rock Salt Structure

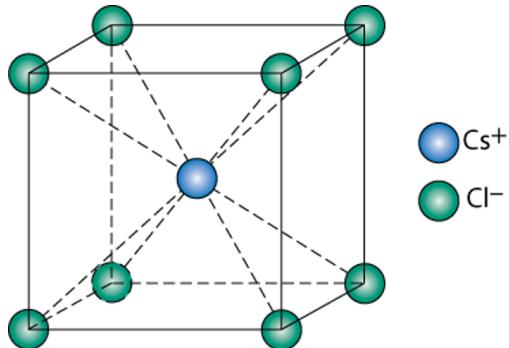


Note: Cations prefer octahedral sites (in black)

Along the edges:

$$2R_A + 2R_C = a$$

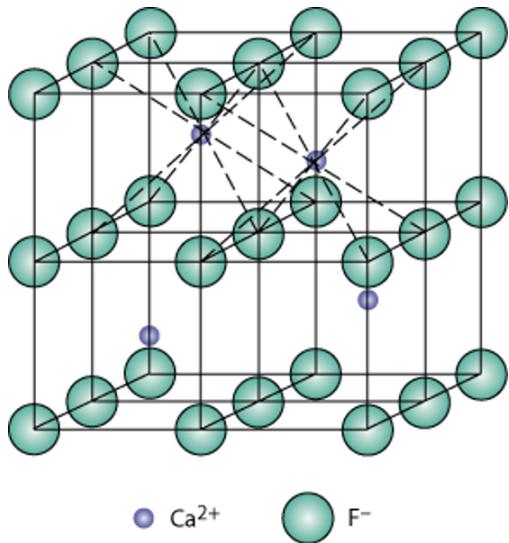
AX Crystal Structure (Cesium Chloride)



Note: Cations prefer cubic sites (Body Center, in blue)
Across the main diagonal:

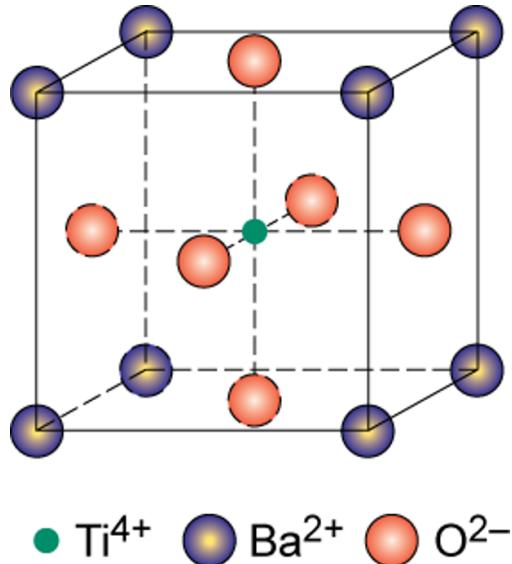
$$2R_A + 2R_C = \sqrt{3}a$$

AX₂ Crystal Structures (Flourite)



Note: Cations prefer cubic sites (Body Center, in blue)
There are half as many Ca^{2+} as F^- (for CaF_2)

ABX_3 Crystal Structure (Perovskite)



Point Coordinates

To find the coordinate indices (q, r, s), find the Cartesian coordinates and divide by the corresponding lattice parameter

$$(q, r, s) = \left(\frac{x}{a}, \frac{y}{b}, \frac{z}{c} \right)$$

Crystallographic Directions

How to define:

- Position vector to pass through origin
- Read off projections onto coordinate axes in terms of lattice parameters (a, b, c)
- Multiply through by common denominator
- Enclose in square brackets without commas, negatives go on top (ex: $[\bar{1}23]$)

How to read $[123]$:

- Divide by the common denominator used previously (say 6)
- Vector in Cartesian: $(1/6, 1/3, 1/2)$

Crystallographic Planes

How to define with Miller indices:

- Define any origin
- Read intercepts of the plane with the coordinate axes in terms of lattice parameters (a, b, c)
- Take reciprocals of intercepts
- Enclose in parentheses without commas, negatives go on top (ex: ($\bar{1}23$))

How to read:

- Take reciprocals of plane to identify intercepts

Linear and Planer Density

Linear Density

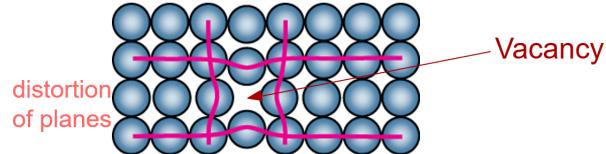
$$LD = \frac{\text{Number of atom centered on line}}{\text{Unit length of direction vector}}$$

Planer Density

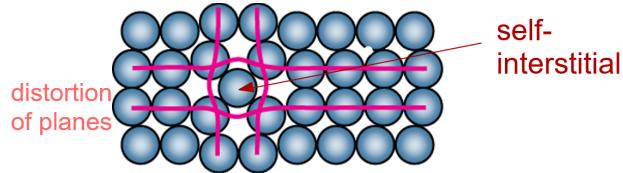
$$PD = \frac{\text{Number of atoms centered on plane}}{\text{Area of plane}}$$

Point Defects

- Vacancies: Missing atoms in the lattice



- Interstitials: Extra atoms positioned between lattice sites



Equilibrium Concentration of Vacancies

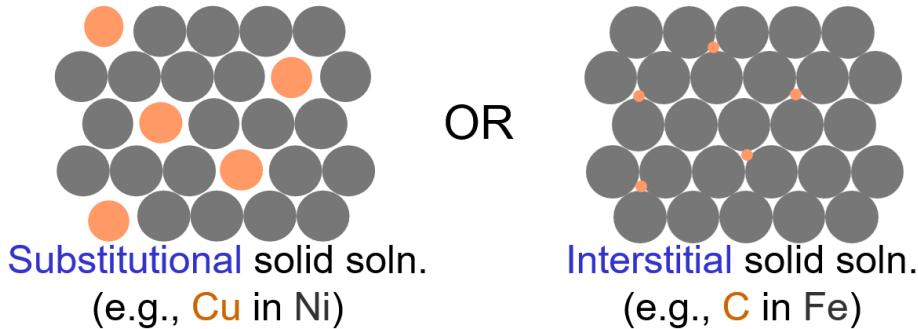
$$N_v = N e^{\left(\frac{-Q_v}{kT}\right)}$$

Where:

- N_v = number of vacancies
- N = total number of atomic sites (each lattice site can be occupied by an atom or be a vacancy)
- Q_v = activation energy required to form a vacancy (J or eV)
- k = Boltzmann's constant (8.617×10^{-5} eV/atom-K)
- T = temperature (K)

Alloys

Alloys are solid solutions of an impurity element (solute) dissolved in a base element (solvent). There are two types:



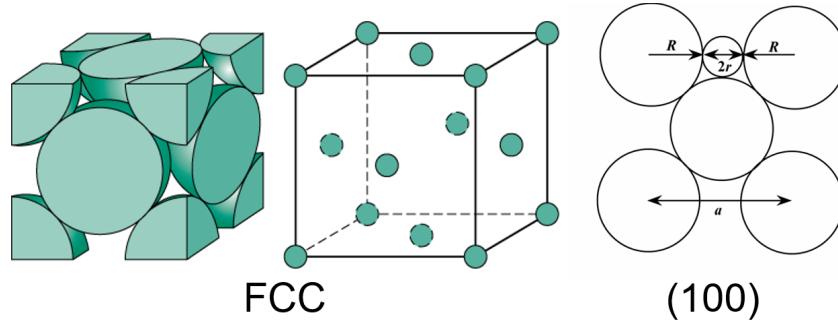
Hume-Rothery rule

Conditions for forming a solid solution:

- Atomic radii must be within 15% of each other ($\Delta r < 0.15\%$)
- Electronegativity values must be similar (i.e. Close together on the periodic table)
- Crystal structures of solvent and solute must be the same (for pure metals)
- Valencies should be similar (All else being equal, a metal will have a greater tendency to dissolve a metal of higher valency than one of lower valency)

Maximum Radius for an Interstitial Impurity

FCC

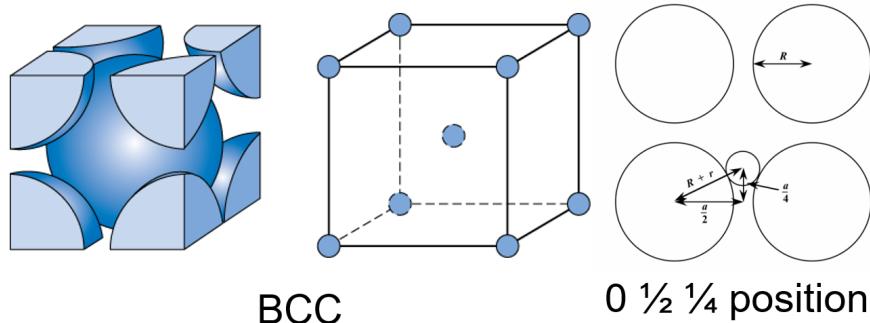


$$2r + 2R = a$$

$$r = \frac{a - 2R}{2}$$

$$r = \frac{2R\sqrt{2} - 2R}{2} = 0.41R$$

BCC



BCC

$0 \frac{1}{2} \frac{1}{4}$ position

$$\left(\frac{a}{2}\right)^2 + \left(\frac{a}{4}\right)^2 = (R + r)^2$$

$$\left(\frac{4R}{2\sqrt{3}}\right)^2 + \left(\frac{4R}{4\sqrt{3}}\right)^2 = R^2 + 2Rr + r^2$$

$$r^2 + 2Rr - 0.667R^2 = 0$$

$$\Rightarrow r = \frac{-2R \pm \sqrt{(2R)^2 - 4(1)(-0.667R^2)}}{2(1)}$$

$$r = 0.291R$$

Specification of composition

Weight Percent

$$C_1 = \frac{m_1}{m_1 + m_2} \times 100$$

Where

- C_1 = weight percent of element 1
- m_1 = mass of element 1
- m_2 = mass of element 2

Atomic Percent

$$C'_1 = \frac{n_{m1}}{n_{m1} + n_{m2}} \times 100$$

Where

- C'_1 = atomic percent of element 1
- n_{m1} = number of moles of element 1
- n_{m2} = number of moles of element 2

Note:

$$n_{m1} = \frac{m_1(\text{grams})}{A_1(\text{g/mol})}$$

Conversion between wt% and at%

$$C'_1 = \frac{C_1 A_2}{C_1 A_2 + C_2 A_1} \times 100$$

Where A_1 and A_2 are the atomic weights of elements 1 and 2 respectively

Diffusion

Diffusion: The general process where atoms move from one place to another, causing mass transport within or between materials.

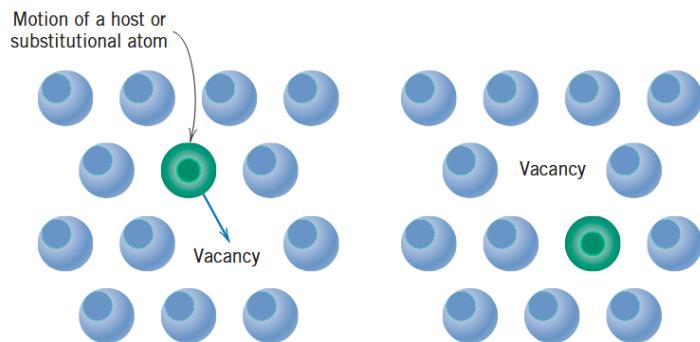
Interdiffusion: Occurs when atoms from two different metals migrate into each other, forming a region where the metals mix.

Self-diffusion: Happens in pure metals when atoms of the same kind exchange places, even though there is no concentration difference.

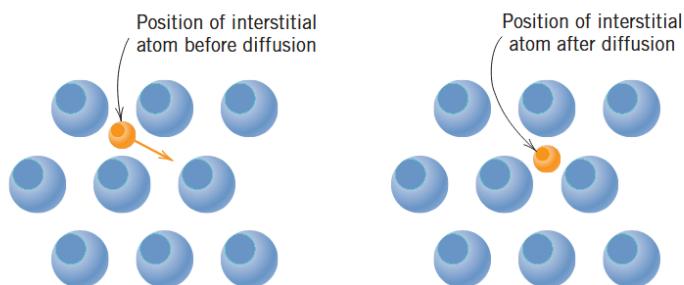
Diffusion Mechanisms

For diffusion to occur: (1) an adjacent site must be vacant, and (2) the atom must have enough energy to break bonds and distort the lattice. Only a fraction of atoms at a given temperature have sufficient vibrational energy for diffusion, and this fraction increases with temperature. For metals, two main models describe this atomic motion.

- **Vacancy Diffusion:** Atoms move into adjacent vacant lattice sites.



- **Interstitial Diffusion:** Smaller atoms move through the spaces between larger atoms.



Fick's First Law

We quantify how fast diffusion occurs by defining the rate of mass transfer or the **diffusion flux** (J). It is the amount of substance that diffuses through a unit area per unit time.

$$J = \frac{M}{At} \quad (\text{kg/m}^2 \cdot \text{s} \text{ or atoms/m}^2 \cdot \text{s})$$

Where:

- M = mass of diffusing substance (kg or atoms)
- A = cross-sectional area (m^2)
- t = time (s)

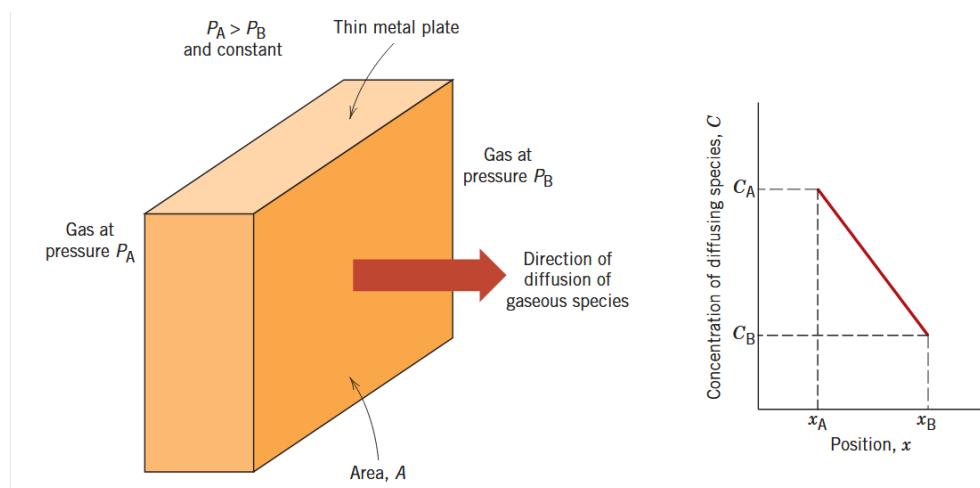
Fick's first law—diffusion flux for steady-state diffusion (in one direction)

$$J = -D \frac{dC}{dx}$$

Where:

- D = diffusion coefficient (m^2/s)
The negative sign indicates that the direction of diffusion is down the concentration gradient, from a high to a low concentration.
- $\frac{dC}{dx}$ = concentration gradient in the x direction.

Fick's first law describes the diffusion of atoms through a thin metal plate when the concentrations (or pressures) on both surfaces are constant. In this case, a **steady-state diffusion** is achieved, where the diffusion flux remains constant over time and there is no net accumulation of the diffusing species within the plate.



When the concentration C is plotted against position x in a solid, the curve is called the **concentration profile**, and the **concentration gradient** is the slope at any point on this curve. For simplicity, the concentration profile is often assumed to be linear.

$$(\text{Concentration Gradient}) = \frac{dC}{dx} = \frac{\Delta C}{\Delta x} = \frac{C_A - C_B}{x_A - x_B}$$

The **concentration** C of a diffusing species is expressed as the mass of the species per unit volume of the solid, typically in units of kg/m³ or g/cm³.

In diffusion, the term **driving force** refers to what compels the process to occur. When diffusion follows Fick's law, the **concentration gradient** acts as the driving force.

Note: We can convert between weight percent and concentration using the density of the material.

$$C''_1 = \frac{C_1}{\frac{C_1}{\rho_1} + \frac{C_2}{\rho_2}} \times 10^3$$

Where:

- C_1 = weight percent of 1
- C_2 = weight percent of 2
- ρ_1 = density of 1 (g/cm³)
- ρ_2 = density of 2 (g/cm³)

Fick's second law—nonsteady-state diffusion

For nonsteady-state diffusion, the diffusion flux and the concentration gradient at some particular point in a solid vary with time, with a net accumulation or depletion of the diffusing species resulting.

Fick's second law—diffusion equation for nonsteady-state diffusion (in one direction)

$$\frac{\partial C}{\partial t} = D \frac{\partial^2 C}{\partial x^2}$$

Where:

- C = concentration of the diffusing species (kg/m^3 or g/cm^3)
- D = diffusion coefficient (m^2/s)

Solutions to this differential equation are possible when physically meaningful boundary conditions are specified.

Semi-infinite Solid

A practically important case is diffusion into a **semi-infinite solid** where the surface concentration is held constant, often from a gas phase with constant partial pressure. The assumptions are:

1. Initially, the solute atoms are uniformly distributed in the solid with concentration C_0 .
2. The surface is at $x = 0$ and x increases into the solid.
3. Time $t = 0$ is defined just before diffusion begins.

These conditions can be expressed as:

Initial condition:

$$t = 0, \quad C = C_0 \quad \text{for } 0 \leq x \leq \infty$$

Boundary conditions:

$$t > 0, \quad C = C_s \quad \text{at } x = 0 \quad (C_s \text{ is the constant surface concentration})$$
$$t > 0, \quad C = C_0 \quad \text{at } x = \infty$$

This results in the following solution:

$$\boxed{\frac{C_x - C_0}{C_s - C_0} = 1 - \operatorname{erf}\left(\frac{x}{2\sqrt{Dt}}\right)}$$

Where:

- C_x = concentration at distance x after time t
- C_0 = initial concentration inside solid
- C_s = surface concentration
- D = diffusion coefficient
- x = distance into solid

Note: The concentrations can be in any consistent units (wt%, at%, kg/m³, etc.) since they are in a ratio.

Target concentration (C_1):

If you want the solute in the alloy to reach a certain concentration C_1 , you can express it in a normalized form:

$$\frac{C_1 - C_0}{C_s - C_0} = \text{constant}$$

This ratio is constant for a given point in the material at a fixed time.

Relation to diffusion distance and time:

$$\frac{x}{2\sqrt{Dt}} = \text{constant} \quad \text{or equivalently} \quad \frac{x^2}{Dt} = \text{constant}$$

This is constant for fixed target concentration C_1 and initial and surface concentrations C_0 and C_s .

Note: **Carburizing** is the process by which the surface carbon concentration of a ferrous alloy is increased by diffusion from the surrounding environment. (C_0 is zero for pure iron)

Note: Values for the error function can be found for various $\frac{x}{2\sqrt{Dt}}$ inputs can be found can be found in tables. If the value does not match exactly, do a linear approx using the two closest values.

z	$erf(z)$	z	$erf(z)$	z	$erf(z)$
0	0	0.55	0.5633	1.3	0.9340
0.025	0.0282	0.60	0.6039	1.4	0.9523
0.05	0.0564	0.65	0.6420	1.5	0.9661
0.10	0.1125	0.70	0.6778	1.6	0.9763
0.15	0.1680	0.75	0.7112	1.7	0.9838
0.20	0.2227	0.80	0.7421	1.8	0.9891
0.25	0.2763	0.85	0.7707	1.9	0.9928
0.30	0.3286	0.90	0.7970	2.0	0.9953
0.35	0.3794	0.95	0.8209	2.2	0.9981
0.40	0.4284	1.0	0.8427	2.4	0.9993
0.45	0.4755	1.1	0.8802	2.6	0.9998
0.50	0.5205	1.2	0.9103	2.8	0.9999

Factors that influence diffusion

The diffusion coefficient D refers to the rate at which atoms diffuse. The diffusing species and the host material influence the diffusion coefficient.

Temperature Dependence

$$D = D_o e^{\left(-\frac{Q_d}{RT}\right)}$$

Where:

- D_o = a temperature-independent preexponential (m^2/s)
- Q_d = activation energy for diffusion (J/mol or eV/atom)
- R = universal gas constant (8.314 J/mol·K)
- T = absolute temperature (K)

Note: $K = {}^\circ C + 273$

<i>Diffusing Species</i>	<i>Host Metal</i>	$D_0 \text{ (m}^2/\text{s)}$	$Q_d \text{ (J/mol)}$
Interstitial Diffusion			
C ^b	Fe (α or BCC) ^a	1.1×10^{-6}	87,400
C ^c	Fe (γ or FCC) ^a	2.3×10^{-5}	148,000
N ^b	Fe (α or BCC) ^a	5.0×10^{-7}	77,000
N ^c	Fe (γ or FCC) ^a	9.1×10^{-5}	168,000
Self-Diffusion			
Fe ^c	Fe (α or BCC) ^a	2.8×10^{-4}	251,000
Fe ^c	Fe (γ or FCC) ^a	5.0×10^{-5}	284,000
Cu ^d	Cu (FCC)	2.5×10^{-5}	200,000
Al ^c	Al (FCC)	2.3×10^{-4}	144,000
Mg ^c	Mg (HCP)	1.5×10^{-4}	136,000
Zn ^c	Zn (HCP)	1.5×10^{-5}	94,000
Mo ^d	Mo (BCC)	1.8×10^{-4}	461,000
Ni ^d	Ni (FCC)	1.9×10^{-4}	285,000
Interdiffusion (Vacancy)			
Zn ^c	Cu (FCC)	2.4×10^{-5}	189,000
Cu ^c	Zn (HCP)	2.1×10^{-4}	124,000
Cu ^c	Al (FCC)	6.5×10^{-5}	136,000
Mg ^c	Al (FCC)	1.2×10^{-4}	130,000
Cu ^c	Ni (FCC)	2.7×10^{-5}	256,000
Ni ^d	Cu (FCC)	1.9×10^{-4}	230,000

Rewriting the equation:

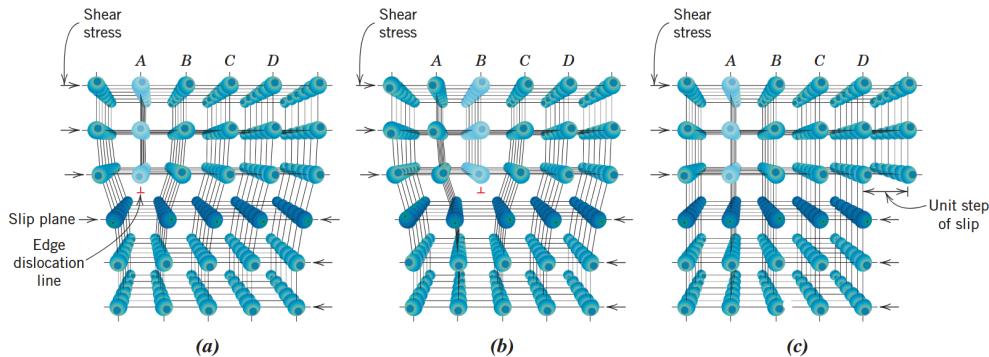
$$\ln D = \ln D_o - \frac{Q_d}{RT}$$

This is the equation of a straight line with slope $-\frac{Q_d}{R}$ and y-intercept $\ln D_o$. Thus, a plot of $\ln D$ versus $\frac{1}{T}$ yields a straight line, from which D_o and Q_d can be determined.

Line Defects

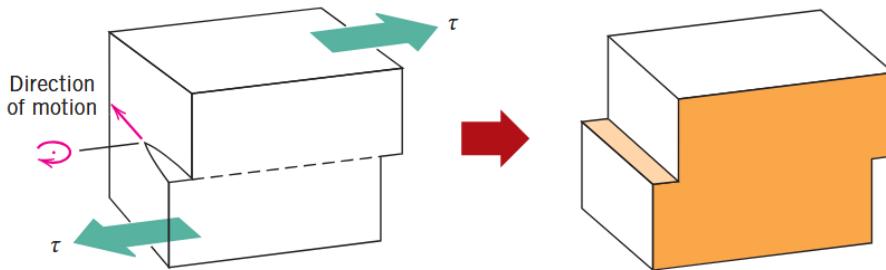
Dislocations are line defects. It occurs when there is slipping between crystal planes when dislocations move, and they cause permanent (plastic) deformation. There are two main types:

- Edge Dislocation



This occurs when an extra half plane of atoms is inserted in a crystal and moves when shear stress is applied. The direction of movement is parallel to the applied shear stress.

- Screw Dislocation



The direction of movement is perpendicular to the applied shear stress.

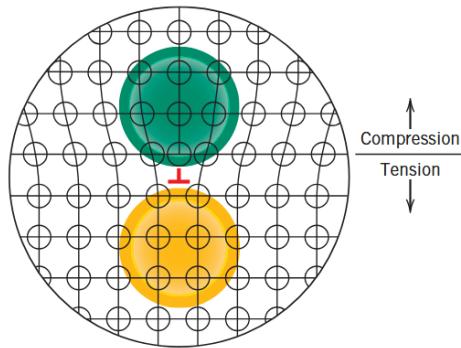
Plastic deformation occurs due to many dislocations. The process by which plastic deformation is produced by dislocation motion is termed **slip**; the crystallographic plane along which the dislocation line traverses is the **slip plane**.

All metals and alloys contain dislocations formed during solidification. The **solidification density** is the number of dislocations per unit volume in a material.

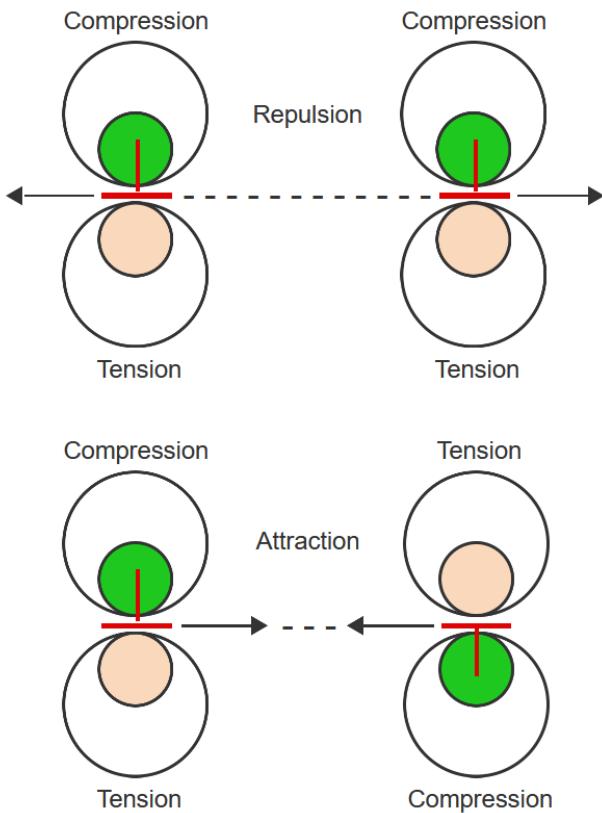
Characteristics of dislocations

Due to the presence of an extra half plane of atoms, there exists regions of compressive, tensile, and shear **lattice strains** around the dislocation line.

Ex:



These lattice strains interact in the following ways:



The attractive forces will result in the half planes merging into a complete plane, eliminating the dislocation. The repulsive forces will make it more difficult for dislocations to move. The number of dislocations increases dramatically during plastic deformation because of multiplication of existing dislocations, grain boundaries, internal defects, and surface irregularities.

Slip Systems

The preferred slip plane and slip direction are those with the densest atomic packing (highest planer density) and highest linear density respectively. The combination of slip plane and slip direction is called a **slip system**.

Ex: FCC $\{111\}\langle110\rangle$ slip system

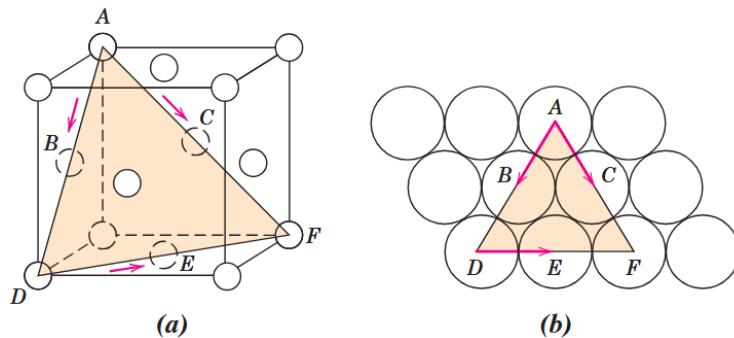


Table of Slip Systems

Metals	Slip Plane	Slip Direction	Number of Slip Systems
Face-Centered Cubic			
Cu, Al, Ni, Ag, Au	$\{111\}$	$\langle110\rangle$	12
Body-Centered Cubic			
α -Fe, W, Mo	$\{110\}$	$\langle111\rangle$	12
α -Fe, W	$\{211\}$	$\langle111\rangle$	12
α -Fe, K	$\{321\}$	$\langle111\rangle$	24
Hexagonal Close-Packed			
Cd, Zn, Mg, Ti, Be	$\{0001\}$	$\langle11\bar{2}0\rangle$	3
Ti, Mg, Zr	$\{10\bar{1}0\}$	$\langle11\bar{2}0\rangle$	3
Ti, Mg	$\{10\bar{1}1\}$	$\langle11\bar{2}0\rangle$	6

Note: More slip systems = more ductile material

The **Burgers vector** has the same direction as the dislocation motion (slip direction) and its magnitude is equal to the interatomic spacing in that direction (unit slip distance). Ex: Burgers vector for FCC and BCC

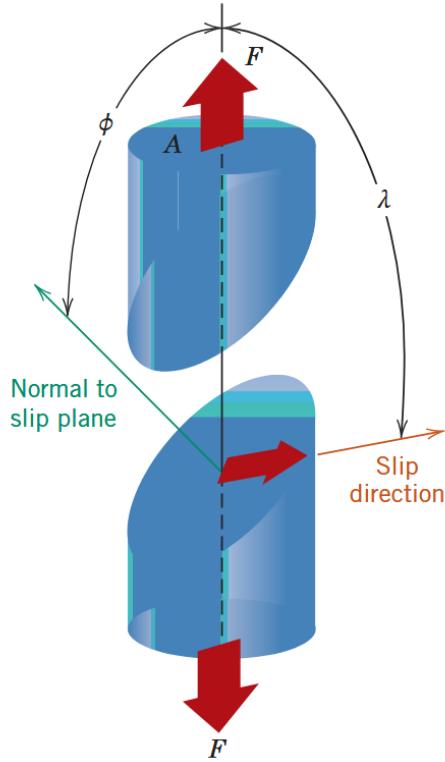
$$\vec{b}_{FCC} = \frac{a}{2}\langle110\rangle$$

$$\vec{b}_{BCC} = \frac{a}{2}\langle111\rangle$$

Slip in single crystals

For an applied shear stress σ , the resolved shear stress represents the magnitude of the shear stress acting on the slip plane in the slip direction. It is given by:

$$\tau_R = \sigma \cos \phi \cos \lambda$$



A metal single crystal has a number of different slip systems, each with its own value of resolved shear stresses. However, there will be one slip system with the maximum resolved shear stress, denoted $\tau_R(\max)$.

$$\tau_R(\max) = \sigma (\cos \phi \cos \lambda)_{\max}$$

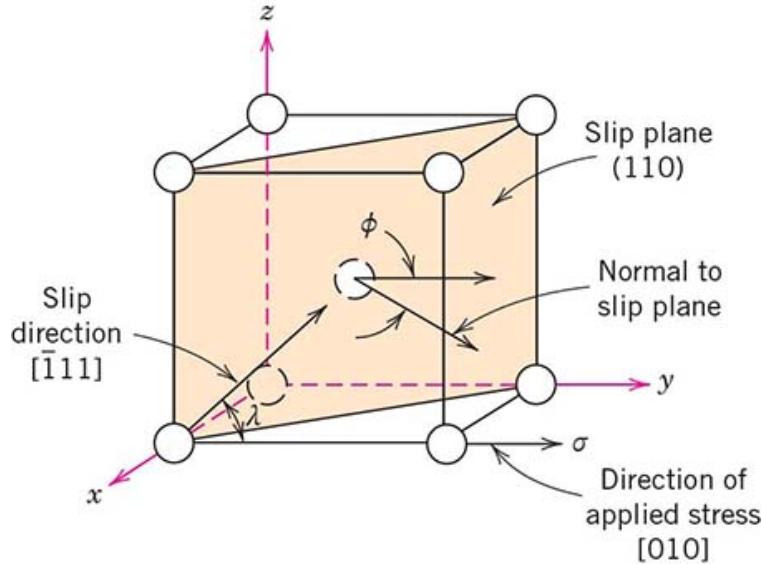
In this slip system with the maximum resolved shear stress, slip occurs when the resolved shear stress reaches a critical value (the critical resolved shear stress, τ_{CRSS}). This is the minimum shear stress required to initiate slip and thus determines when yielding will begin (when $\tau_R(\max) = \tau_{CRSS}$).

$$\sigma_y = \frac{\tau_{CRSS}}{(\cos \phi \cos \lambda)_{\max}} \quad (\text{yield stress})$$

This yield stress is minimized when $\phi = \lambda = 45^\circ$, maximizing $\cos \phi \cos \lambda = 0.5$. Thus,

$$\sigma_{y,\min} = 2\tau_{CRSS}$$

Ex: Consider a single crystal of BCC iron oriented such that a tensile stress is applied along a [010] direction. Compute the resolved shear stress along a (110) plane and in a [1̄11] direction when a tensile stress of 52 MPa (7500 psi) is applied.



Recall that we can find the angle by using the projection:

$$\cos \theta = \frac{\vec{a} \cdot \vec{b}}{|\vec{a}| |\vec{b}|}$$

Where:

- Applied tensile direction [010] $\Rightarrow \vec{l} = (0, 1, 0)$
- Slip plane (110) $\Rightarrow \vec{n} = (1, 1, 0)$ (normal vector to plane)
- Slip direction [1̄11] $\Rightarrow \vec{d} = (-1, 1, 1)$

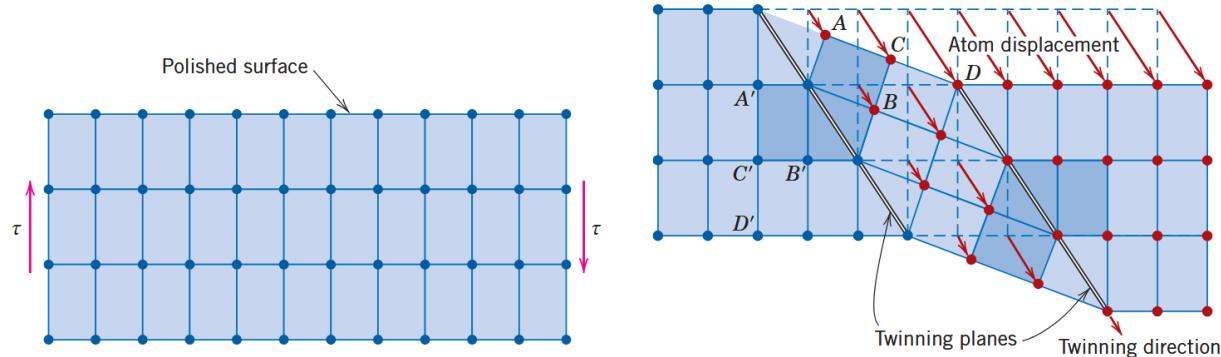
Thus,

$$\begin{aligned} \phi &= \cos^{-1} \left(\frac{\vec{n} \cdot \vec{l}}{|\vec{n}| |\vec{l}|} \right) & \lambda &= \cos^{-1} \left(\frac{\vec{d} \cdot \vec{l}}{|\vec{d}| |\vec{l}|} \right) \\ &= \cos^{-1} \left(\frac{(1, 1, 0) \cdot (0, 1, 0)}{\sqrt{1^2 + 1^2 + 0^2} \sqrt{0^2 + 1^2 + 0^2}} \right) & &= \cos^{-1} \left(\frac{(-1, 1, 1) \cdot (0, 1, 0)}{\sqrt{(-1)^2 + 1^2 + 1^2} \sqrt{0^2 + 1^2 + 0^2}} \right) \\ &= \cos^{-1} \left(\frac{1}{\sqrt{2}} \right) = 45^\circ & &= \cos^{-1} \left(\frac{1}{\sqrt{3}} \right) = 54.74^\circ \end{aligned}$$

$$\tau_R = (52 \text{ MPa}) \cos(45^\circ) \cos(54.74^\circ) = 21.3 \text{ MPa}$$

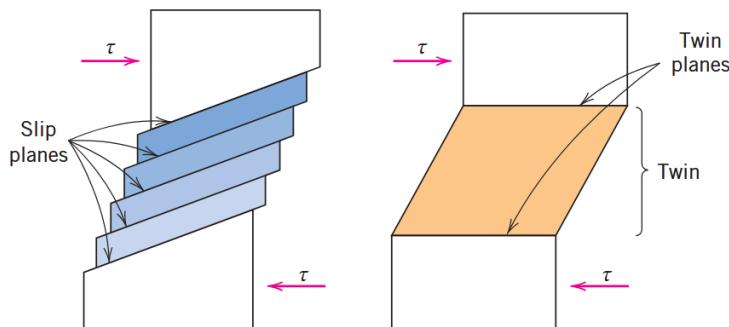
Deformation by Twinning

In addition to slip, plastic deformation in some metallic materials can occur by the formation of mechanical twins, or twinning. Twinning occurs when a shear force produces atomic displacements such that on one side of a plane (the twin boundary), atoms are located in mirror-image positions of atoms on the other side.



Twinning is different from slip in several ways:

- Twinning produces a reorientation of the crystal lattice, where slip crystallographic orientation is the same before and after deformation
- Slip occurs in distinct atomic spacing multiples, whereas the atomic displacement for twinning is less than the interatomic separation.



Mechanical twinning occurs in metals with BCC and HCP crystal structures at low temps and high rates of loading where there are fewer slip systems. Twinning is important because it may place new slip systems in orientations that are favorable relative to the stress axis such that the slip process can now take place.

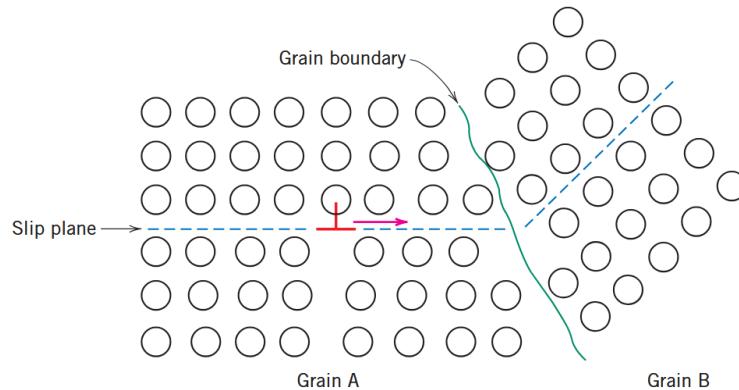
Strengthening by grain size reduction

The strength and hardness of metals depend on how easily dislocations can move. Plastic deformation occurs through dislocation motion, so if dislocations move freely, the metal is soft and weak. If their motion is restricted, more force is needed to deform the metal, making it harder and stronger. Thus, nearly all strengthening methods work by hindering dislocation movement.

Grain size affects a metal's mechanical properties because grain boundaries hinder dislocation motion. When a dislocation moves from one grain to another, it encounters two obstacles:

- The change in crystal orientation makes it harder for the dislocation to continue moving.
- The atomic disorder at the grain boundary disrupts slip-plane continuity.

As a result, smaller grains (more boundaries) strengthen the metal by restricting dislocation motion.



The relationship between yield strength and grain size is given by the **Hall-Petch equation**:

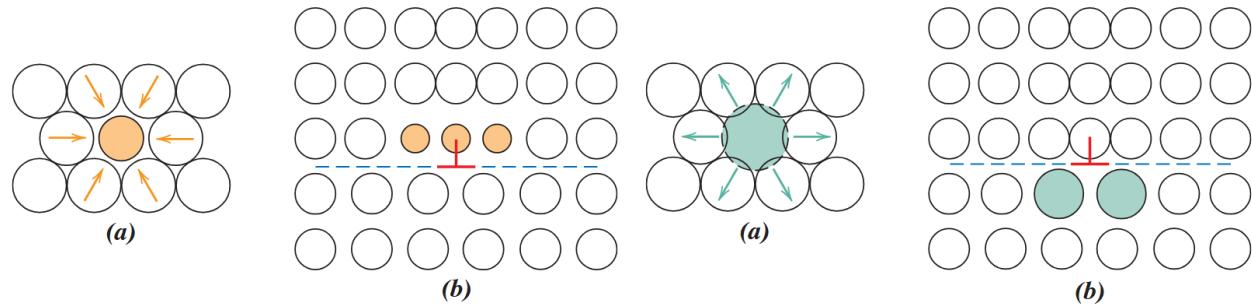
$$\sigma_y = \sigma_o + k_y d^{-1/2}$$

Where:

- σ_y = yield stress
- σ_o = friction stress (stress required to move dislocations, a material constant)
- k_y = strengthening coefficient (a material constant)
- d = average grain diameter

Solid-solution strengthening

By creating alloys with impurity atoms in either a substitutional or interstitial solid solution, the strength of a metal can be increased. This process is called **solid-solution strengthening**. High-purity metals are almost always softer and weaker than alloys composed of the same base metal. This is because the impurity atoms impose lattice strains that hinder dislocation motion.



Strain hardening

Strain hardening is the phenomenon by which a ductile metal becomes harder and stronger as it is plastically deformed. Sometimes it is also called work hardening, or **cold working**.

We can define the amount of plastic deformation as **percent cold work** instead of strain:

$$\%CW = \frac{A_o - A_d}{A_o} \times 100$$

Where:

- A_o = original cross-sectional area
- A_d = cross-sectional area after deformation

Recovery

When a polycrystalline metal is plastically deformed at low temperatures, its grains change shape, dislocation density increases, and it undergoes strain hardening. Some deformation energy is stored as strain energy, and properties such as electrical conductivity and corrosion resistance may change.

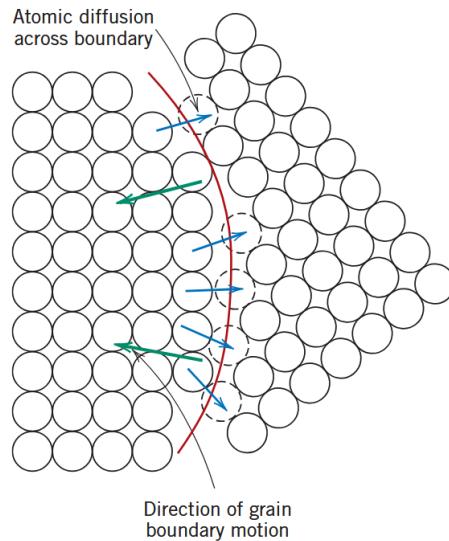
These effects can be reversed by **annealing**, a heat treatment involving **recovery** and **recrystallization** (sometimes followed by grain growth). During **recovery**, atomic diffusion allows dislocations to rearrange and reduce internal strain energy without external stress. This lowers dislocation density and restores properties such as electrical and thermal conductivity to their original states.

Recrystallization

Recrystallization replaces deformed grains with new, strain-free, equiaxed ones that have low dislocation density. This process, driven by stored strain energy, restores the metal's softness and ductility. It depends on time and temperature; the **recrystallization temperature** (temp for completion in 1 h) is typically $0.3\text{--}0.5 T_m$ for metals and up to $0.7 T_m$ for alloys (T_m is the melting temperature). Greater cold work lowers this temperature, while impurities raise it by slowing grain boundary motion.

Grain Growth

Grain growth occurs when strain-free grains enlarge at elevated temperatures, even without prior recovery or recrystallization. As grain boundaries migrate, large grains grow at the expense of smaller ones, reducing total boundary area and energy. The driving force is the decrease in grain boundary energy, and boundary motion occurs through short-range atomic diffusion.



The grain diameter varies with time according to the relationship:

$$d^n - d_o^n = Kt$$

Where:

- d = grain diameter at time t
- d_o = initial grain diameter
- n = grain growth exponent (typically 2)
- K = temperature-dependent rate constant

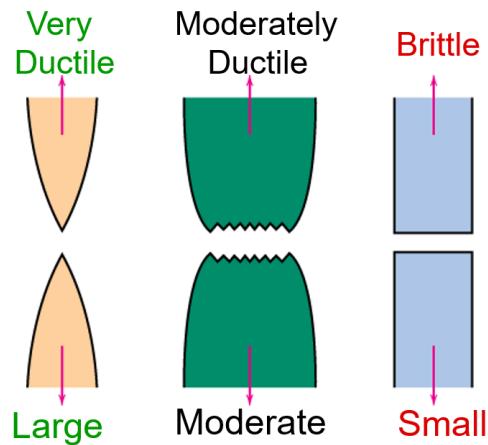
Failure

Fracture

Simple fracture is when a body breaks into two or more pieces under static stress. For metals, fracture can be classified as either **brittle** or **ductile**.

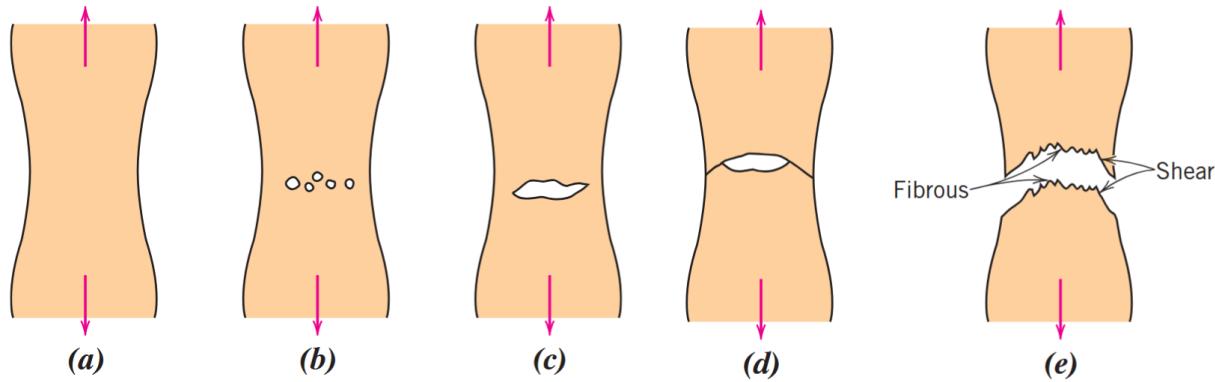
Brittle fracture occurs with little to no plastic deformation and is characterized by a rapid crack propagation.

Ductile fracture occurs with extensive plastic deformation and is characterized by significant necking before fracture.



Fracture involves two steps: crack formation and propagation.

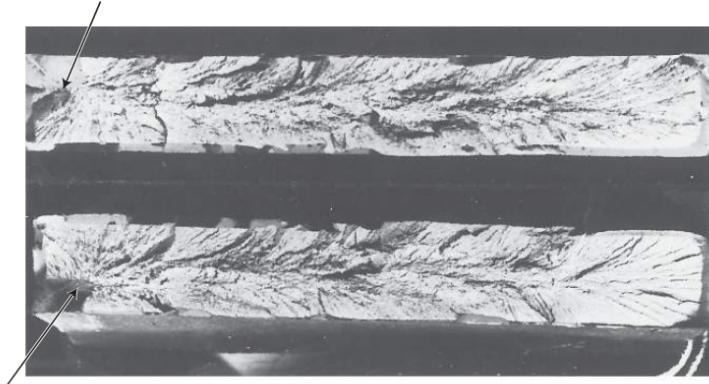
Ductile Fracture Mechanism



After necking begins, small cavities form in the center of the necked down region. The cavities grow and coalesce to form a single crack. The crack propagates toward the outer edge of the sample, resulting in fracture. Ductile materials show almost 100% reduction in area at fracture. Ductile fractures usually resemble a cup-and-cone shape (crack propagates at 45° to the applied load).

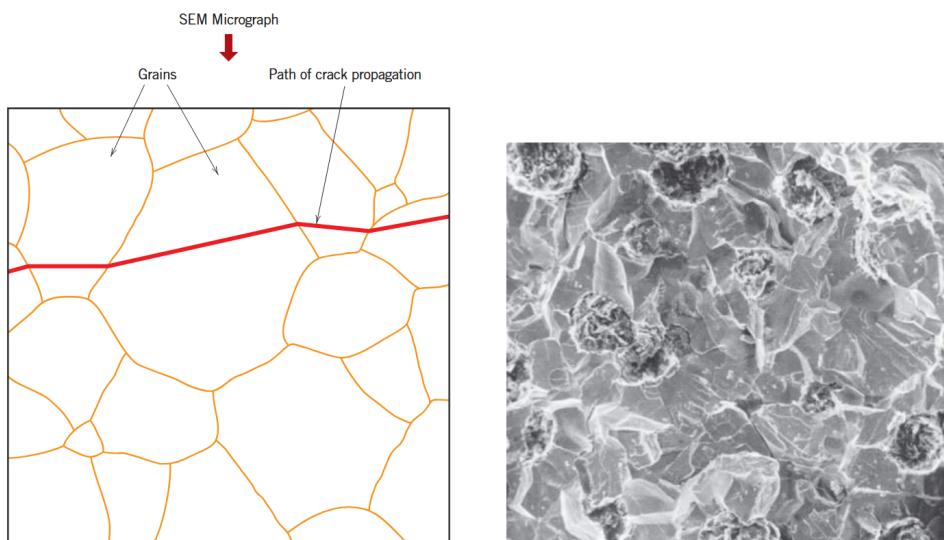
Brittle Fracture

In brittle fracture, cracks form in the center perpendicular to the direction of applied load. The crack then propagates rapidly. Brittle materials show very little reduction in area at fracture. The fracture surface is usually flat and perpendicular to the applied load. The fracture surface usually has V-shaped "chevron" markings that point in the direction of crack propagation.



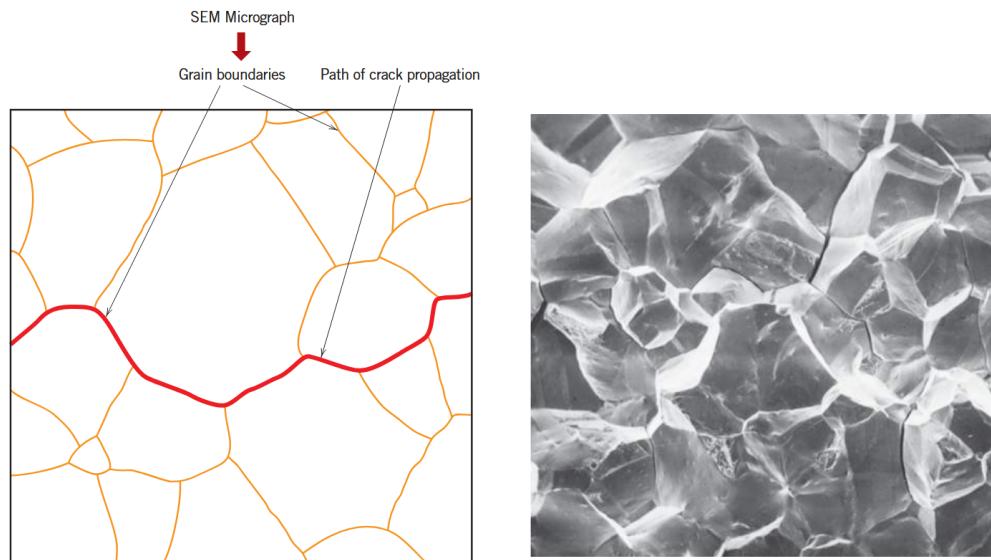
Transgranular Fracture

For most brittle materials, crack propagation corresponds to the breaking of atomic bonds along specific crystallographic planes, known as *cleavage*. When the fracture surface passes through grains, it is called **transgranular fracture**.



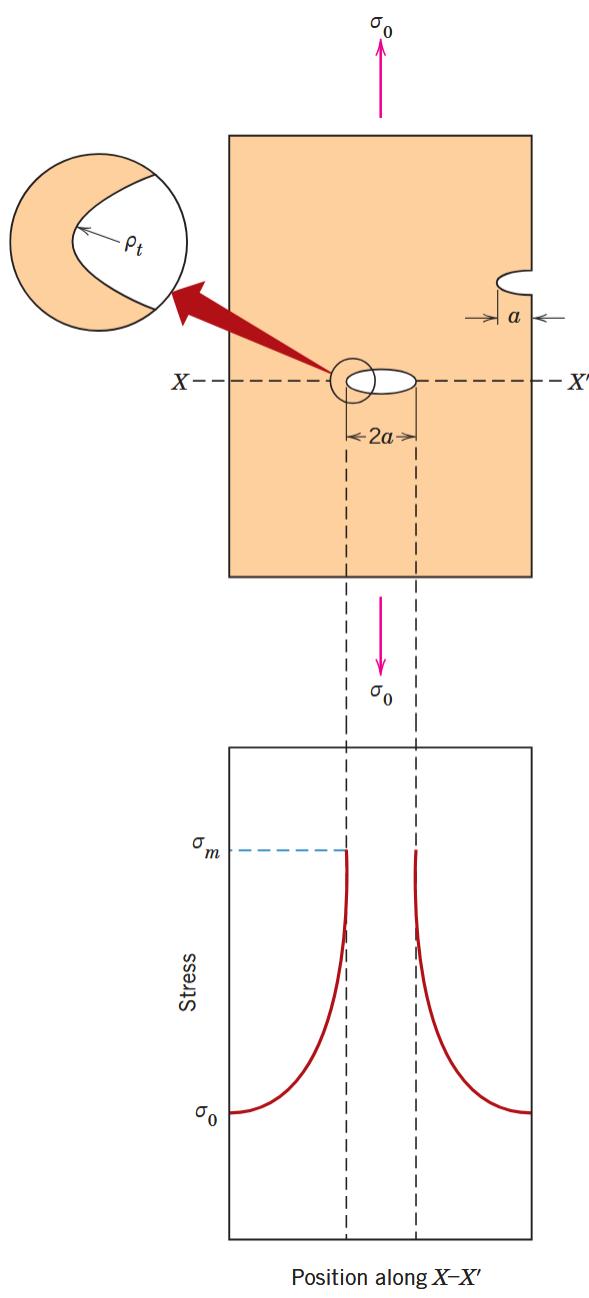
Intergranular Fracture

In some cases, cracks propagate along grain boundaries, resulting in **intergranular fracture**.



Principles of fracture mechanics

Stress concentration



The measured fracture strength of a material is often much lower than the theoretical strength calculated based on atomic bond strengths. This is due to the presence of microscopic flaws (cracks, voids, inclusions) that act as stress concentrators. The stress at the tip of a crack can be much higher than the applied stress, leading to premature fracture.

For internal cracks that are elliptical, the maximum stress at the crack tip is given by:

$$\sigma_m = 2\sigma_o \left(\frac{a}{\rho_t} \right)^{1/2} = K_t \sigma_o$$

$$K_t = \frac{\sigma_m}{\sigma_o}$$

Where

- σ_m = maximum stress at crack tip
- σ_o = applied stress
- a = length of a *surface* crack, or half of the length of an *internal* crack
- ρ_t = radius of curvature at crack tip
- K_t = stress concentration factor

Note: Sharper cracks (smaller ρ_t) lead to higher stress concentrations.

Note: Exterior fillets also act as stress concentrators (surface cracks).

Critical stress

Cracks propagate when the stress at the crack tip reaches a critical value, σ_c , (ie. $\sigma_m = \sigma_c$ or $K_t > K_c$). The critical stress for crack propagation is given by:

$$\sigma_c = \left(\frac{2E\gamma_s}{\pi a} \right)^{1/2}$$

Where:

- E = modulus of elasticity
- γ_s = specific surface energy
- a = length of a *surface* crack, or half of the length of an *internal* crack

For ductile fracture, replace γ_s with $\gamma_s + \gamma_p$, where γ_p is the plastic deformation energy per unit area.

Note: Larger cracks (larger a) lead to lower critical stresses for fracture.

Fracture Toughness

Fracture toughness, K_C , is a measure of a material's resistance to fracture in the presence of a crack. It represents the critical value of the stress intensity factor at which crack propagation becomes unstable. It is defined as:

$$K_C = Y\sigma_c\sqrt{\pi a}$$

Where:

- Y = geometric correction factor (dimensionless)
- σ_c = critical applied stress for crack propagation
- a = crack length (half-length for an internal crack)

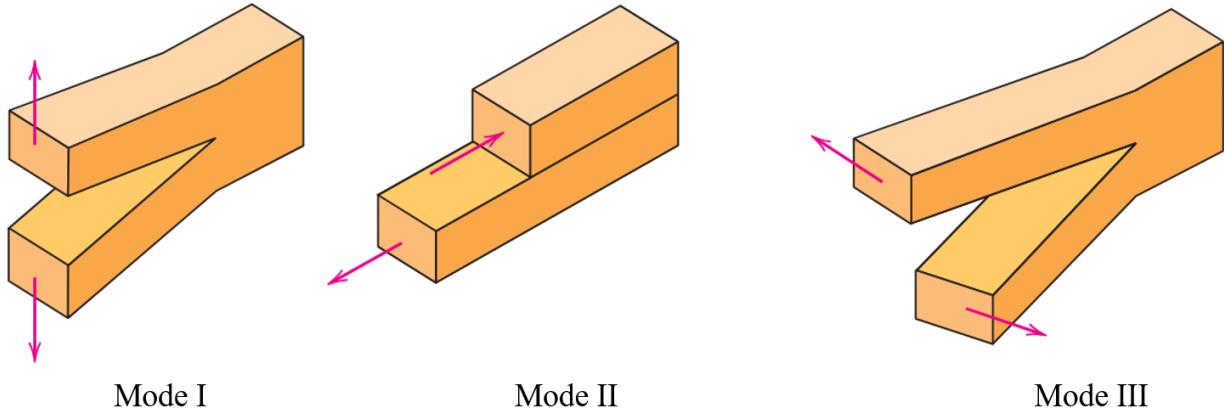
K_C has units of MPa·m^{1/2} or psi·in^{1/2}. In general, K_C may depend on specimen thickness and geometry. A higher K_C value indicates greater resistance to crack propagation.

Plane Strain Fracture Toughness

For specimens with thickness much greater than the crack dimensions and the crack-tip plastic zone, K_C becomes independent of thickness. Under these conditions, plane strain exists (no strain perpendicular to the front and back faces). The fracture toughness under plane strain conditions is denoted K_{IC} for Mode I crack opening.

$$K_{IC} = Y\sigma\sqrt{\pi a}$$

σ is the applied stress. Under valid plane strain conditions, K_{IC} is a true material property and represents the minimum fracture toughness of the material.



Design using fracture mechanics

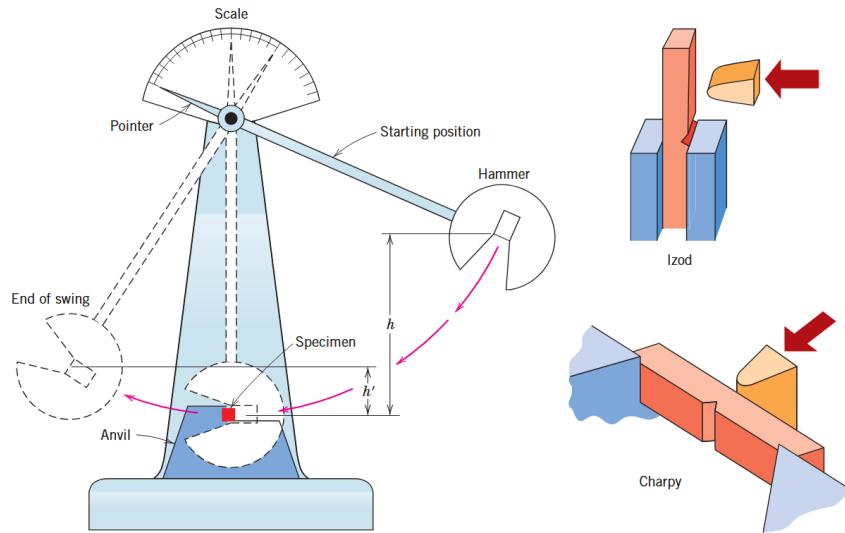
Typically when it comes to design, the fracture toughness K_C or K_{IC} is known for a material, and if we know either the applied stress σ or the crack length a , we can solve for the other variable to ensure that fracture does not occur. Thus, rearranging the equation we see that:

$$\sigma_c = \frac{K_{IC}}{Y\sqrt{\pi a}} \quad (\text{Max stress})$$

$$a_c = \frac{1}{\pi} \left(\frac{K_{IC}}{Y\sigma} \right)^2 \quad (\text{Max crack length/ defect size})$$

Fracture toughness testing

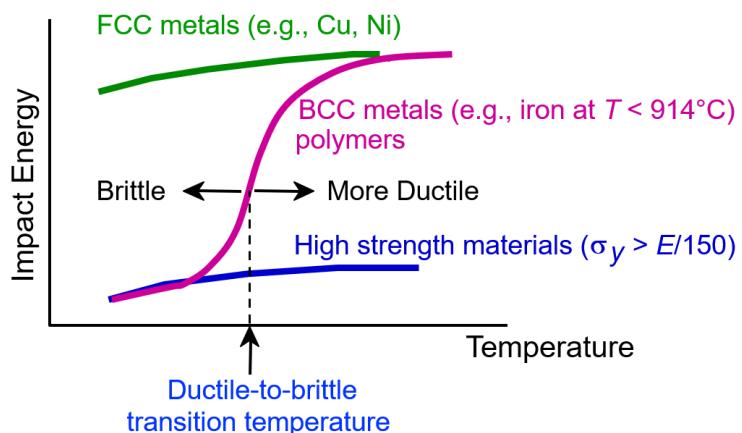
There are two tests, the **Charpy** and **Izod**, that measure the **impact energy** (energy absorbed during fracture) of a material. Both tests involve breaking a notched specimen with a swinging pendulum hammer. The difference is in how the specimen is held:



Ductile-to-brittle transition

Some materials exhibit a transition from ductile behavior at high temperatures to brittle behavior at low temperatures. This phenomenon is known as the ductile-to-brittle transition. The transition is characterized by a sharp drop in impact energy over a narrow temperature range.

The temperature at which the material changes from ductile to brittle fracture behavior is called the ductile-to-brittle transition temperature (DBTT). Above the DBTT, materials absorb a large amount of energy before fracturing (ductile). Below the DBTT, materials absorb very little energy (brittle). The Charpy V-notch impact test is commonly used to determine the ductile-to-brittle transition temperature (DBTT) of materials.



Cyclic stresses

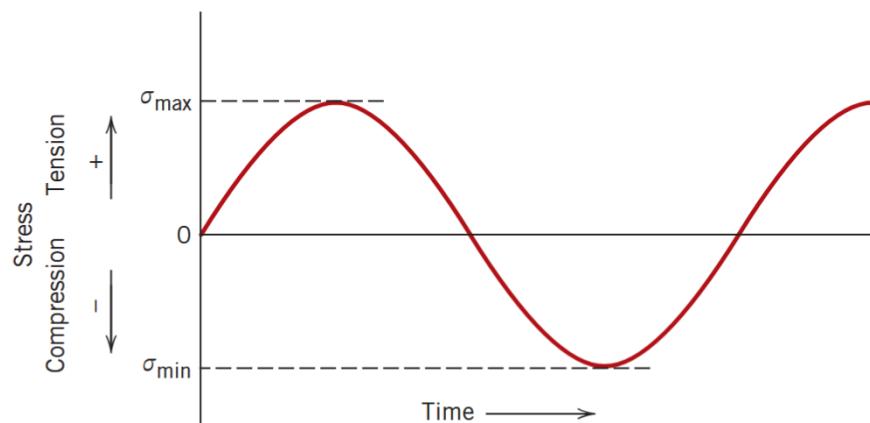
Fatigue is failure that occurs due to cyclic or fluctuating stresses, typically well below the static tensile or yield strength of the material. Most engineering failures occur by fatigue.

Fatigue failure proceeds in three stages:

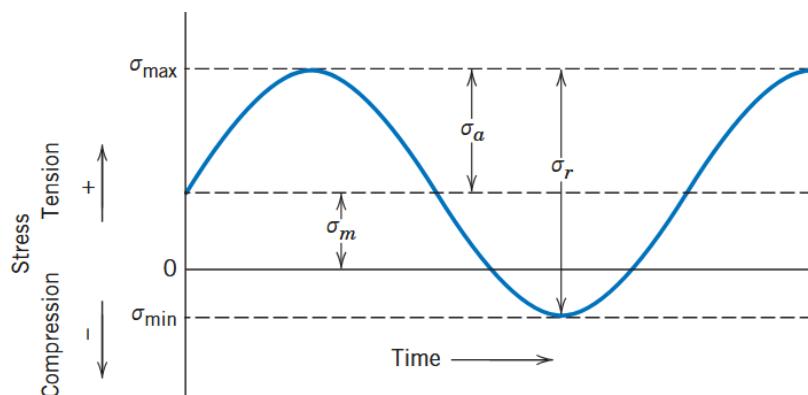
- Crack initiation (usually at surface defects, notches, scratches)
- Crack propagation (incremental crack advance per cycle)
- Final sudden fracture when remaining cross-section cannot support load

The applied stress may be axial (tension-compression), flexural (bending), or torsional (twisting) in nature. In general, three different fluctuating stress-time modes are possible.

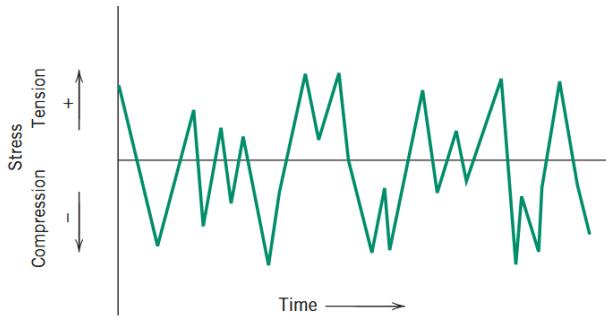
Reversed stress cycle, in which the stress alternates from a maximum tensile stress (+) to a maximum compressive stress (-) of equal magnitude:



Repeated stress cycle, in which maximum and minimum stresses are asymmetrical relative to the zero-stress level; mean stress σ_m , range of stress σ_r , and stress amplitude σ_a are indicated:



Random stress cycle:



We can characterize cyclic stresses using the following parameters:

- Mean stress for cyclic loading

$$\sigma_m = \frac{\sigma_{max} + \sigma_{min}}{2}$$

- Range of stress of cyclic loading

$$\sigma_r = \sigma_{max} - \sigma_{min}$$

- Stress amplitude of cyclic loading

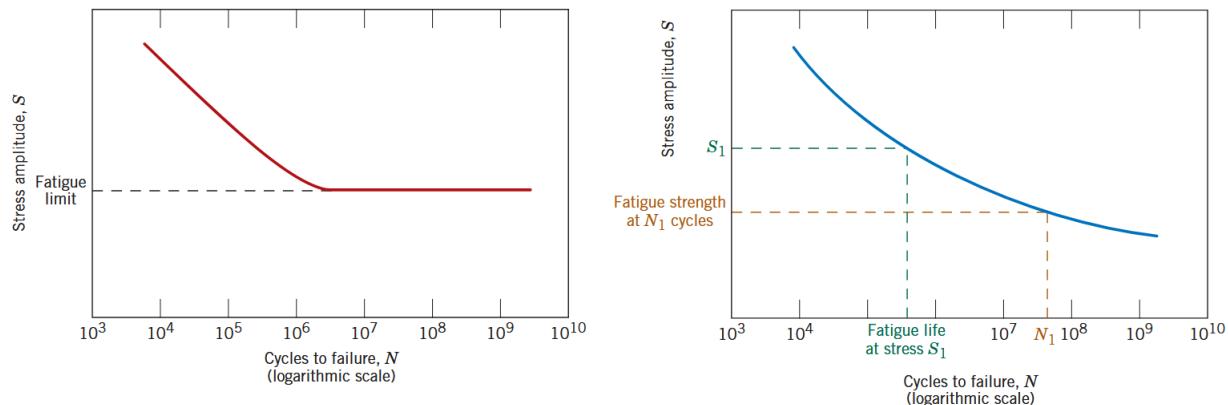
$$\sigma_a = \frac{\sigma_r}{2} = \frac{\sigma_{max} - \sigma_{min}}{2}$$

- Stress ratio of cyclic loading

$$R = \frac{\sigma_{min}}{\sigma_{max}}$$

The S-N curve

Fatigue behavior is commonly represented using an S-N curve, which plots stress amplitude vs. number of cycles to failure.



For some metals (primarily steels), the S–N curve levels off at a plateau called the: **fatigue limit**. Below this limit, the material can theoretically endure infinite cycles without failure. For other metals (such as aluminum and copper), there is no distinct fatigue limit; instead, the S–N curve continues to slope downward, indicating that failure will eventually occur at any stress level given enough cycles. For these metals, **fatigue strength** is defined as the stress amplitude at a specified number of cycles (**fatigue life**).

Fatigue Lifetime (Basquin's Law)

To determine the number of cycles to failure for a given stress amplitude, or to determine the stress amplitude for a desired fatigue life, we use the mathematical form of the S–N curve. For many metals in the high-cycle fatigue region, the S–N curve follows a power-law relationship known as **Basquin's Law**:

$$\sigma_a = AN^b$$

Where:

- σ_a = stress amplitude
- N = number of cycles to failure
- A = material constant (intercept)
- b = fatigue exponent (slope), typically negative

Taking the logarithm of both sides gives a linear relationship:

$$\log \sigma_a = \log A + b \log N$$

Thus, when plotted on a log–log graph, the S–N curve becomes a straight line.

Determining A and b From Data

Given two points on the S–N curve, (σ_{a1}, N_1) and (σ_{a2}, N_2) , we can compute the constants:

$$b = \frac{\log \sigma_{a2} - \log \sigma_{a1}}{\log N_2 - \log N_1}$$

$$A = \sigma_{a1} N_1^{-b}$$

Once A and b are known, we may solve for:

$$N = \left(\frac{\sigma_a}{A} \right)^{1/b} \quad \text{or} \quad \sigma_a = AN^b$$

Crack initiation and propagation

Cracks associated with fatigue failure almost always initiate (or nucleate) on the surface of a component at stress concentrators such as: surface scratches, sharp fillets, keyways, threads, dents, etc. Cyclic loading can also produce microscopic surface cracks from dislocation slip steps. Once a fatigue crack forms, it grows incrementally every cycle.

Fracture surfaces from fatigue failures often exhibit features such as

- **Benchmarks** (macroscopic): pattern formed by intermittent loading
- **Striations** (microscopic): each marks one cycle of crack growth

Phase Diagrams